

Enhanced Charge Carrier Separation in WO₃/BiVO₄ Photoanodes Achieved via Light Absorption in the BiVO₄ Layer

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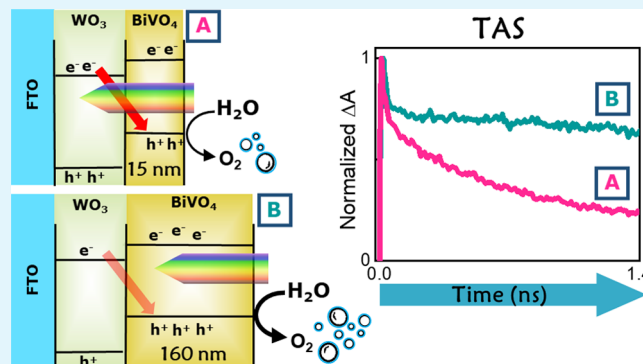
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Supporting Information

ABSTRACT: Photoelectrochemical (PEC) water splitting converts solar light and water into oxygen and energy-rich hydrogen. WO₃/BiVO₄ heterojunction photoanodes perform much better than the separate oxide components, though internal charge recombination undermines their PEC performance when both oxides absorb light. Here we exploit the BiVO₄ layer to sensitize WO₃ to visible light and shield it from direct photoexcitation to overcome this efficiency loss. PEC experiments and ultrafast transient absorption spectroscopy performed by frontside (through BiVO₄) or backside (through WO₃) irradiating photoanodes with different BiVO₄ layer thickness demonstrate that irradiation through BiVO₄ is beneficial for charge separation. Optimized electrodes irradiated through BiVO₄ show 40% higher photocurrent density compared to backside irradiation.

KEYWORDS: solar water oxidation, heterojunction, ultrafast transient absorption, photoactive layer thickness, filter effect



Bismuth vanadate, BiVO₄, is a promising semiconductor oxide employed in photoanodes for the oxygen evolution reaction in water-splitting devices.^{1,2} Its stability in contact with aqueous electrolytes,^{3,4} its good visible light-harvesting capability,⁵ and its simple preparation through cheap wet techniques⁶ point to this material as a possible component of future commercial photoelectrochemical (PEC) cells. Furthermore, in the last 15 years the efficiency of BiVO₄-based photoanodes (in terms of current density) rapidly grew from a few microamps per square centimeter in early reports to 4–5 mA cm⁻², with prolonged continuous operation of photoelectrodes modified with oxygen evolution cocatalysts.^{3,4,6–13} However, the fast charge recombination of BiVO₄-based electrodes still hampers the efficiency of this material.^{2,14}

A way to overcome this intrinsic flaw is to couple BiVO₄ with WO₃ in the WO₃/BiVO₄ heterojunction where visible light harvesting BiVO₄ sensitizes wider band gap WO₃.¹⁵ BiVO₄ photoanodes based on this heterojunction achieve the highest current densities among oxide-based photoanodes.^{16,17} The suitable band gap alignment between the two oxides, the efficient electron and hole transport in WO₃ and BiVO₄, respectively, and the spacial charge separation support the high performance of WO₃/BiVO₄ photoanodes.^{18–25}

In previous studies, we investigated the charge carrier dynamics in the WO₃/BiVO₄ system through transient absorption spectroscopy (TAS) with detection either in the visible range to observe the hole dynamics in BiVO₄^{18,26–28} or

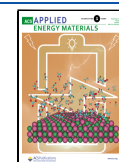
in the mid-infrared to follow the electron dynamics in WO₃ and BiVO₄.²⁹ We also identified wavelength-dependent processes by tuning the excitation wavelength across the WO₃ absorption edge (ca. 450 nm).^{18,26,29} Indeed, the type II band alignment between the two oxides (Figure 1A) allows distinct charge transfer processes leading to charge separation or recombination, depending on the excitation wavelength. Under visible light excitation of BiVO₄, electrons promoted in its conduction band (CB) flow into the energetically lower-lying CB of WO₃, while holes remain in the BiVO₄ valence band (VB). This electron transfer process (process A in Figure 1A) decreases charge recombination and leads to long-living charge carriers that are beneficial for PEC performance.²⁹

Conversely, irradiation at wavelengths below 450 nm leads to the excitation of both oxides and opens a detrimental recombination path between the electrons photopromoted in the CB of WO₃ and the holes in BiVO₄ (process B in Figure 1A). This process results in charge recombination on a ~200 ps time scale²⁶ and becomes more relevant with increasing WO₃ layer thickness.³⁰

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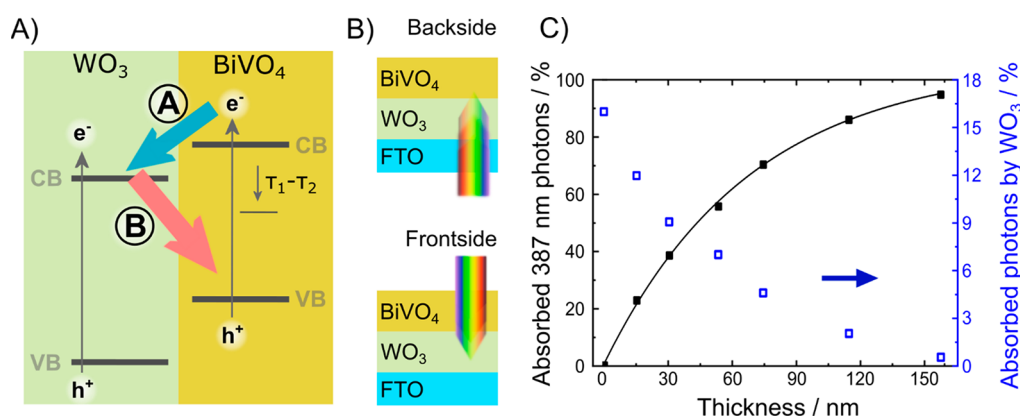


Figure 1. (A) Charge carrier transfer paths involved in the WO₃/BiVO₄ heterojunction: light absorption in BiVO₄ produces valence band (VB) holes h⁺ and conduction band (CB) electrons e⁻, which can recombine with holes (with time constants τ_1 and τ_2) or flow into WO₃ through process (A), leading to charge separation. When WO₃ is also photoexcited, electrons in its CB can recombine with holes in BiVO₄ through process (B). (B) Backside (through WO₃) and frontside (through BiVO₄) irradiation modes of the WO₃/BiVO₄ heterojunction photoanodes. (C) Percent amount of 387 nm photons (the pump wavelength used in TAS experiments) absorbed by the BiVO₄ layer (left ordinate) and by the WO₃ layer (right ordinate) under frontside excitation vs the thickness of the BiVO₄ layer in WO₃/BiVO₄ photoanodes.

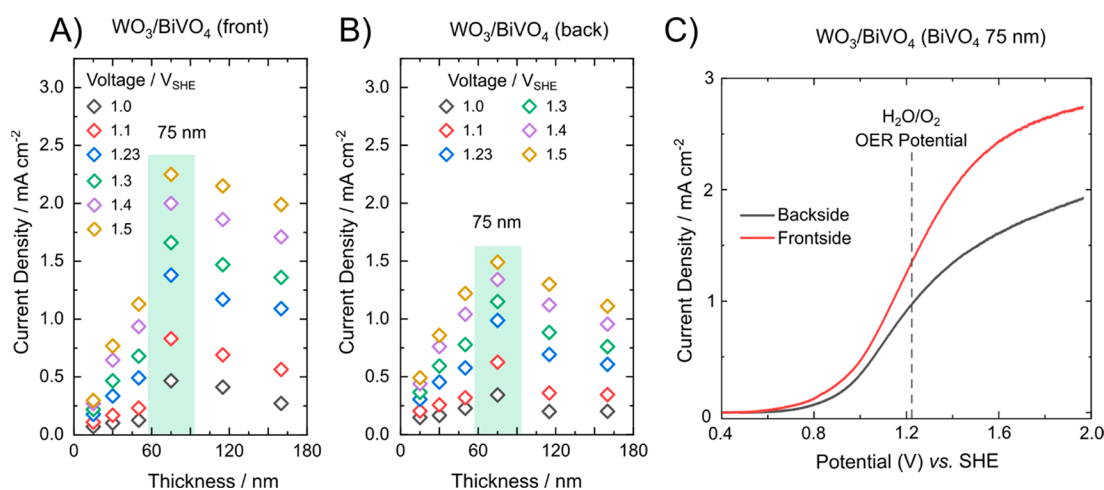


Figure 2. Current density generated under 1 sun simulated solar light (AM 1.5 G irradiation) at different applied potentials with the WO₃/BiVO₄ heterojunction electrodes in contact with a 0.5 M Na₂SO₄ aqueous solution under (A) frontside and (B) backside irradiation; (C) linear sweep voltammetry recorded with the best performing WO₃/BiVO₄ electrode (with a 75 nm thick BiVO₄ layer) under backside (black line) and frontside (red line) irradiation; the vertical dashed line indicates the standard oxygen evolution reaction potential.

Based on these dynamics, we posited that an efficient heterojunction system needs to direct charges along process (A) and disfavor off-track routes such as process (B). Still, solar light includes photons energetic enough to excite WO₃ (~4% of the solar spectrum is at wavelengths below the absorption edge of bulk WO₃). Therefore, a portion of photogenerated charges in WO₃/BiVO₄ systems may be wasted through process (B). On the other hand, BiVO₄ efficiently absorbs light beyond the absorption edge of WO₃, up to 520 nm, allowing us to exploit a larger fraction of the solar spectrum. Therefore, in this work we pursue the strategy of using the BiVO₄ sensitizer to shield WO₃ from direct photoexcitation.

We assembled a series of heterojunction electrodes with a WO₃ scaffold layer of fixed thickness (ca. 150 nm) coated with BiVO₄ overlayers with different thickness (15–160 nm) to tune the amount of light absorbed by BiVO₄. First, a systematic PEC study allowed us to probe whether the irradiation mode (through WO₃ or BiVO₄, backside or frontside irradiation, respectively, Figure 1B) affects the overall PEC efficiency of the electrodes. Then, transient absorption spectroscopy (TAS)

with a pump in the UV region (387 nm) and detection in the visible range was employed to assess the effects of the irradiation mode on the lifetime of photogenerated holes in BiVO₄. These tests allowed us to evaluate the extent of charge recombination induced by process (B) and its impact on the PEC performance of the heterojunction photoanodes as a function of the BiVO₄ layer thickness.

The WO₃/BiVO₄ photoanodes were prepared through spin coating using fluorine-doped tin oxide (FTO) as the conductive glass substrate (see the Supporting Information). The heterojunction electrode with the thickest BiVO₄ layer almost entirely absorbs 387 nm photons, the pump wavelength in TAS experiments, Figure 1C. A series of control photoanodes consisting of pure BiVO₄ on FTO (without WO₃ layer) with variable BiVO₄ thickness was also prepared. The absorption spectra of the two electrode series are shown in Figures S1 and S2; the thickness of the BiVO₄ layer was estimated using the absorption coefficient at 420 nm, $\alpha_{40} = 6.7 \times 10^4 \text{ cm}^{-1}$. XRD analyses confirm the successful synthesis of WO₃ and BiVO₄ (Figure S3) and FESEM images

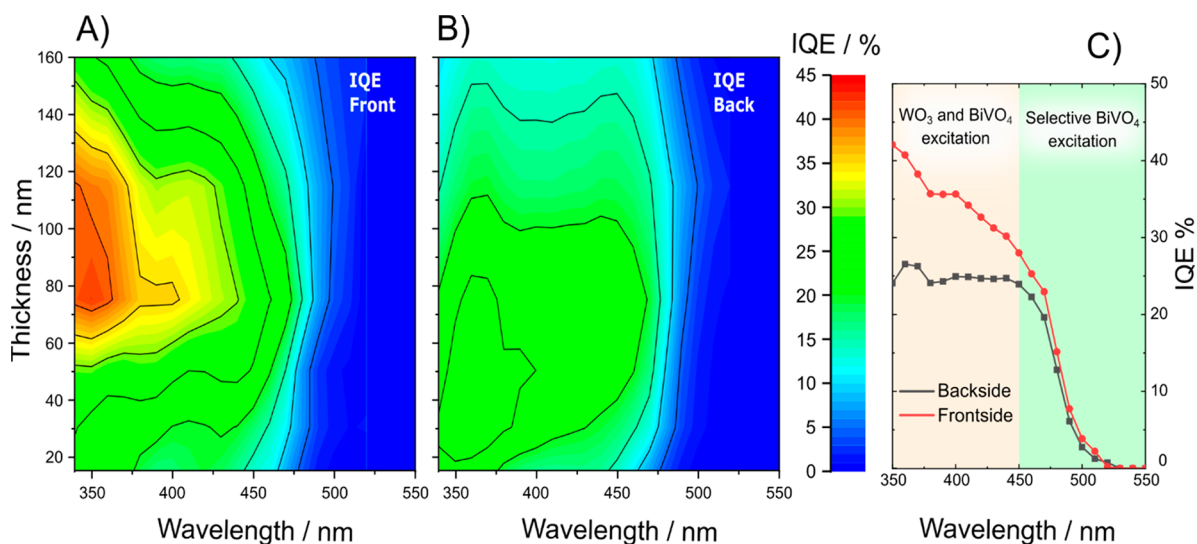


Figure 3. Internal quantum efficiency (IQE) 3D contour plots (color scale) vs the incident wavelength and the BiVO₄ film thickness of the WO₃/BiVO₄ electrodes, obtained in contact with a 0.5 M Na₂SO₄ aqueous solution at 1.23 V_{SHE} under (A) frontside or (B) backside irradiation; (C) IQE plots under backside and frontside irradiation obtained with the best performing electrode (with a 75 nm thick BiVO₄ layer).

demonstrate the uniform coating of the photoanodes (Figure S4).

In order to explore the shielding hypothesis, we carried out PEC experiments on the electrodes. Figure 2A, B shows the photocurrent density generated with the WO₃/BiVO₄ electrodes under simulated solar light irradiation in 0.5 M Na₂SO₄ solution under back- and frontside irradiation at different applied potentials. The linear sweep voltammetry plots are reported in Figures S5 and S6. As a general trend, all heterojunction photoanodes outperform control pure BiVO₄ electrodes (see Figure S7). Furthermore, the better light exploitation achieved with increasing the BiVO₄ layer thickness drives the photocurrent increase under both irradiation conditions up to a 75 nm thick BiVO₄ layer.

Under frontside simulated solar light irradiation (Figure 2A), the heterojunction photoanodes generate considerably higher photocurrent than in backside mode (Figure 2B). The best performing electrode with a 75 nm BiVO₄ layer thickness (Figure 2C), when irradiated frontside shows a ca. 40% increase in the current density, from 1.0 to 1.38 mA cm⁻², with respect to backside irradiation, at the formal H₂O/O₂ redox potential of 1.23 V vs the standard hydrogen electrode (V_{SHE}).

We used single-wavelength efficiency measurements to gather further information on this PEC performance increase. Specifically, internal quantum efficiency (IQE, Figure 3 and Figures S8 and S9), measuring the efficiency of absorbed photons, was calculated from the incident photon to current efficiency (IPCE, see Figures S10 and S11) recorded with the WO₃/BiVO₄ electrodes in contact with a 0.5 M Na₂SO₄ solution at 1.23 V_{SHE}. Figure 3A, B shows the IQE vs BiVO₄ thickness contour plots measured under frontside and backside irradiation.

Under frontside irradiation, the IQE reaches the highest values for 70–130 nm thick BiVO₄ layers in the WO₃/BiVO₄ heterojunction, as evidenced by the red/yellow island appearing in Figure 3A as opposed to the green plot obtained under backside irradiation, which leads to lower IQE values (Figure 3B). Notably, the largest IQE enhancement under frontside irradiation occurs below 450 nm, where WO₃ absorbs light, and for BiVO₄ layers thicker than 50 nm, which absorb a

substantial fraction of incident light. The IQE in this irradiation mode maintains above 30% up to 450 nm for the best-performing electrodes, while it is seldom above 25% under backside irradiation, see for example the IQE traces for the WO₃/BiVO₄ electrode with 75 nm thick BiVO₄ in Figure 3C.

At the same time, the IQE curves are similar under the two irradiation modes in the 450–520 nm range because only BiVO₄ absorbs light at these wavelengths and process @ (Figure 1A) is predominantly active. Thus, these wavelength-dependent PEC analyses indicate that the performance of the WO₃/BiVO₄ photoanodes benefits from avoiding WO₃ excitation. This condition occurs by selectively exciting BiVO₄ under frontside irradiation, i.e., by shielding WO₃ with BiVO₄, and in both irradiation modes under excitation at wavelengths above the WO₃ absorption onset (Figure 3C).

We then investigated the effect of WO₃ shielding on the lifetime of photogenerated holes in the BiVO₄ layer of the WO₃/BiVO₄ system. Previous work ascribed the transient absorption ΔA signal at 470 nm to trapped holes in BiVO₄, based on experiments in the presence of hole scavengers.^{18,21,31} TAS proved an essential tool for studying charge carrier dynamics and diffusion in photocatalysis and photovoltaics.^{32–35} Therefore, TAS with detection at 470 nm was here employed to investigate the dynamics of photogenerated holes in both pure BiVO₄ and WO₃/BiVO₄ electrode series upon backside and frontside excitation at 387 nm.

The ΔA signals recorded with pure BiVO₄ electrodes were analyzed first. For this system, similar transient dynamics were obtained in the two irradiation modes. Figure S12 reports representative transient absorption spectra, while Figure S13 shows the transient decay ΔA profiles at 470 nm, which were fitted according to a biexponential decay model (eq 1).

$$\Delta A = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} + \Delta A_0 \quad (1)$$

In this equation, τ_1 and τ_2 are the lifetimes of the faster and slower decay processes typical of BiVO₄, respectively, A_1 and A_2 are the weighted coefficients that represent the contribution of each of the two processes to the overall decay and ΔA_0 is the offset (set at zero in the fitting).²¹ The fitting parameters for the BiVO₄ electrodes (Table S1) are in line with literature

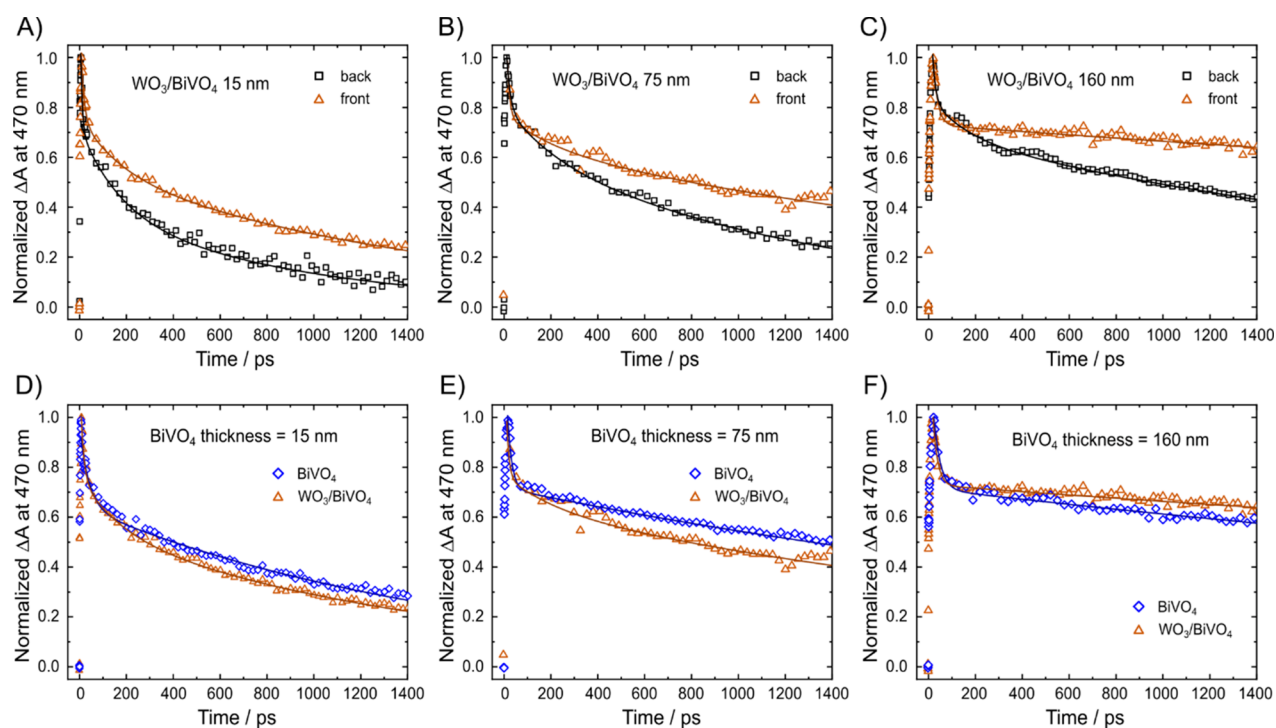


Figure 4. Comparison between the normalized TAS decay signals in (A–C) $\text{WO}_3/\text{BiVO}_4$ electrodes with different BiVO_4 thicknesses recorded under backside (through WO_3 , black squares) or frontside (through BiVO_4 , red triangles) irradiation and in (D–F) $\text{WO}_3/\text{BiVO}_4$ (red triangles) and pure BiVO_4 films (blue diamonds) with the same BiVO_4 thickness recorded under frontside irradiation. The solid lines are the fitting curves according to eq 1 (BiVO_4 films) or eq 2 ($\text{WO}_3/\text{BiVO}_4$ films). Excitation pump at 387 nm, TAS signal monitored at 470 nm.

Table 1. Fitting Parameters of the TAS Dataset Collected with the $\text{WO}_3/\text{BiVO}_4$ Electrodes under Backside Excitation at 387 nm^a

BiVO_4 thickness (nm)	A_1 (%)	τ_1 (ps)	A_r (%)	τ_r (ps)	A_2 (%)	τ_2 (ns)
15	30 ± 1	7.4 ± 0.8	32 ± 3	170 ± 22	38 ± 3	0.98 ± 0.07
30	26 ± 2	6.2 ± 1.1	34 ± 3	189 ± 27	40 ± 3	1.9 ± 0.2
50	24 ± 2	7.8 ± 1.2	19 ± 2	156 ± 30	57 ± 2	1.99 ± 0.11
75	22 ± 1	15.8 ± 1.8	17 ± 2	215 ± 48	61 ± 3	1.44 ± 0.07
115	23 ± 2	22 ± 3	19 ± 2	210 ± 25	58.0 ± 0.3	3.57 ± 0.11
160	18 ± 1	13.1 ± 1.7	15 ± 1	159 ± 26	67.4 ± 0.9	3.05 ± 0.11

^aData fitted according to eq 2.

reports on pure BiVO_4 . Regardless of the BiVO_4 thickness, A_1 and A_2 account of ~30 and 70% of the hole decay, respectively. The fast decay lifetime, which is associated with the recombination of trapped holes in BiVO_4 with photopromoted free electrons, is independent of the BiVO_4 layer thickness (τ_1 , ~20 ps), because all electrodes are excited at the same pump wavelength (i.e., with the same energy excess with respect to the BiVO_4 CB).^{21,29} On the other hand, τ_2 , which is ascribed to the recombination of trapped holes with trapped electrons, increases from ~1 to 6.5 ns with increasing BiVO_4 layer thickness as more holes get trapped in bulk sites.^{28,36,37}

The decay signal of photoproduced holes in the BiVO_4 layer of the $\text{WO}_3/\text{BiVO}_4$ electrodes series recorded at 470 nm under backside and frontside irradiation are reported in Figure 4 and Figures S14 and S15. Under frontside excitation, the ΔA signals decay slower than under backside excitation (Figure 4A–C). Indeed, under backside irradiation mode a significant fraction of 387 nm photons is absorbed by WO_3 , leading to photoexcitation of electrons into its CB (the individual WO_3 layer absorbs ca. 16% of 387 photons, Figure 1C). Therefore, many photoproduced charge carriers recombine through

process @ (Figure 1A). This additional recombination channel leads to the abrupt ΔA drop observed during the first 400 ps following backside photoexcitation (Figure 4A). Furthermore, under frontside irradiation the ΔA signal recorded with the $\text{WO}_3/\text{BiVO}_4$ heterojunctions becomes progressively slower and comparable with those recorded with pure BiVO_4 (Figure 4D–F and Figure S13).

In previous studies on the $\text{WO}_3/\text{BiVO}_4$ heterojunction, we evaluated the contribution of process @ to the overall ΔA decay signal by fitting the ΔA decay traces including an additional decay component in eq 1, to take into account also process @.²⁶ Here, we used the same approach to assess its contribution to the charge carrier dynamics in the two irradiation modes and fitted the ΔA decay with eq 2.

$$\Delta A = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} + A_r e^{-t/\tau_r} + \Delta A_0 \quad (2)$$

where τ_r accounts for the additional recombination process and A_r is its weighted contribution.

We first fitted the dynamics recorded in the $\text{WO}_3/\text{BiVO}_4$ electrodes under backside excitation. The fitting parameters are reported in Table 1. In this configuration, the WO_3 layer is

Table 2. Fitting Parameters of the TAS Dataset Collected with the WO₃/BiVO₄ Electrodes under Frontside Excitation at 387 nm^a

BiVO ₄ thickness (nm)	A ₁ (%)	τ ₁ (ps)	A _r (%)	τ _r (ps)	A ₂ (%)	τ ₂ (ns)
15	25 ± 1	13 ± 1	20 ± 1	170	55 ± 1	1.55 ± 0.03
30	23 ± 3	15 ± 2	27 ± 3	189	50 ± 3	2.8 ± 0.4
50	21 ± 3	19 ± 7	13 ± 5	156	66 ± 2	3.0 ± 0.3
75	23 ± 3	10 ± 3	12 ± 4	215	65 ± 2	3.1 ± 0.4
115	20 ± 1	24 ± 4	4 ± 2	210	76.3 ± 0.8	5.4 ± 0.3
160	27 ± 1	21 ± 2	-	-	73.1 ± 0.3	10.6 ± 1.2

^aData fitted according to eq 2 using the τ_r values reported in Table 1.

irradiated directly and absorbs the same amount of light in all electrodes. Therefore, we expect that process ⑥ has a comparable effect on the charge carrier dynamics at each BiVO₄ layer thickness. Indeed, this process accounts for ca. 23 ± 8% of the holes decay, with a time constant τ_r of ~200 ps (Table 1).

Under frontside irradiation, the BiVO₄ layer in the heterojunction electrodes shields WO₃ from light absorption as the BiVO₄ layer thickness increases. Indeed, the percent amount of incident 387 nm photons absorbed by the WO₃ underlayer in the coupled system progressively decreases (Figure 1C). Therefore, we sought to quantify the shielding effect of the BiVO₄ layer in decreasing the extent of process ⑥ in the WO₃/BiVO₄ electrodes. By assuming that process ⑥ operates with its intrinsic time constant τ_r regardless of the excitation mode, we fitted the decay dynamics recorded under frontside irradiation by employing the τ_r previously extracted from the TAS signals recorded upon excitation in backside mode (Table 1). Because of the reduced amount of charge carriers generated in WO₃, the weight of process ⑥ in terms of the A_r parameter (Table 2) decreases with increasing the BiVO₄ layer thickness. Additionally, as fewer charge carriers undergo process ⑥, A₂ increases, suggesting that a larger number of photogenerated charge carriers recombines through the slower process. A 160 nm thick BiVO₄ layer almost entirely absorbs the pump (Figure 1C), preventing WO₃ excitation. Due to the lack of photoexcited electrons in the CB of WO₃, the electrons photogenerated in the BiVO₄ CB can only recombine with trapped holes in BiVO₄, or flow into WO₃ CB via process ④, resulting in better charge carrier separation. Consequently, the holes photogenerated in the BiVO₄ layer of the WO₃/BiVO₄ heterojunction live longer than those in the individual 160 nm thick BiVO₄ electrode (Figure 4F). This condition is akin to selective BiVO₄ excitation in WO₃/BiVO₄ at wavelengths beyond WO₃ absorption edge, which we previously observed extending the hole lifetimes compared to individual BiVO₄.^{26,29}

Thus, TAS and PEC experiments suggest that light absorption by the BiVO₄ layer in WO₃/BiVO₄ electrodes selectively suppresses process ⑥ and promotes process ④, which leads to an increase of trapped hole lifetime in BiVO₄. However, despite the promise of long-living holes in the heterojunction electrode with a 160 nm thick BiVO₄ layer, it performs poorly compared to the most active heterojunction with the 75 nm thick BiVO₄ layer. This latter electrode possesses an optimal matching between (i) WO₃ sensitization to the visible light, (ii) photogenerated charge separation at the heterojunction, and (iii) efficient charge extraction toward the external circuit and the electrolyte. Indeed, thinner BiVO₄ layers limit the electrode performance due to the low visible light absorption, while thicker films may suffer from a greater

charge recombination owing to hole accumulation in the BiVO₄ bulk under *operando* conditions.

In conclusion, we identified a shielding strategy to suppress the internal charge recombination occurring in the WO₃/BiVO₄ heterojunction due to WO₃ excitation. Optimized light absorption in BiVO₄ layers considerably suppresses this recombination channel. The best performing electrode tested in this work shows a 40% increase in the PEC performance under frontside irradiation compared to backside irradiation. Furthermore, these findings suggest that methods to suppress undesired wavelength-dependent recombination processes and optimize charge transport and surface catalysis are required to design efficient photoelectrodes based on type-II heterojunctions.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsaem.2c02597>.

Experimental details, absorption spectra, XRD and FESEM analyses, linear sweep voltammetry tests, IQE and IPCE plots, TAS decay profiles, and fitting parameters (PDF)

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Notes

The authors declare no competing financial interest.

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