

Nuclear Charge Radii of the Nickel Isotopes ^{58–68,70}Ni

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Collinear laser spectroscopy is performed on the nickel isotopes ^{58–68,70}Ni, using a time-resolved photon counting system. From the measured isotope shifts, nuclear charge radii R_c are extracted and compared to theoretical results. Three *ab initio* approaches all employ, among others, the chiral interaction NNLO_{sat}, which allows an assessment of their accuracy. We find agreement with experiment in differential radii $\delta\langle r_c^2 \rangle$ for all employed *ab initio* methods and interactions, while the absolute radii are consistent with data only for NNLO_{sat}. Within nuclear density functional theory, the Skyrme functional SV-min matches experiment more closely than the Fayans functional Fy(Δr , HFB).

Introduction. — The accurate description of rich physics phenomena encountered in atomic nuclei remains a formidable challenge for contemporary nuclear theory. The long-term goal of nuclear physics is thus to develop a universal framework to consistently describe atomic nuclei across the entire nuclear chart. Research in recent years has led to remarkable advances in nuclear many-body methods [1–7] as well as in the development of nuclear forces based on chiral effective field theory (EFT), rooted in symmetries of QCD and based on pion ex-

changes and short-ranged interactions [8–10]. A significant theoretical effort has been dedicated to the description of electromagnetic properties such as nuclear charge radii R_c . Since charge radii can be measured with high accuracy, they serve as robust benchmarks for nuclear theory. Presently, the region of medium- to heavy-mass nuclei constitutes the testing ground for developing the coherent theoretical nuclear framework. An important element of this endeavor is to connect *ab initio* models to nuclear density functional theory (DFT). In addition

to *ab initio* calculations, well-calibrated energy density functionals, such as the Fayans functional, are capable of a successful description of nuclear charge radii for multiple isotopic chains ranging from potassium ($Z = 19$) all the way to tin ($Z = 50$) [11–17].

In this Letter, we report nuclear charge radii of nickel isotopes (Ni, $Z = 28$) which, in terms of R_c , constitutes the last unexplored "magic" isotopic chain in this mass region. While the charge radius of ^{68}Ni was reported earlier [18], we here present additionally the results for $^{59,63,65-67,70}\text{Ni}$. The experimental data are compared with two DFT approaches as well as three independent *ab initio* methods based on chiral EFT interactions.

Experiment. — The experiment at ISOLDE/CERN has been described previously in [18]. Details on the general setup can be found in [19]. In brief, Ni isotopes were produced in a uranium carbide target bombarded with proton pulses of 1.4-GeV energy. Ions were formed by resonant laser ionization with RILIS [20] and accelerated in a first and a second beamtime to about 30 keV and 40 keV, respectively. Different ISOLDE targets were used with the aim to increase production and to suppress isobars, but they behaved comparably. After mass selection in a high-resolution mass separator, the ions were injected into the radio-frequency quadrupole (RFQ) ion beam cooler and buncher ISCOOL [21] where they were accumulated for typically 10 – 100 ms. After extraction as a short ion bunch, the ions were transported to the collinear laser spectroscopy beam line COLLAPS, where the beam was superimposed with a co-propagating laser beam. Bunching reduces the otherwise dominant background of scattered laser light compared to a continuous beam [22]. The ion beam energy was determined by the high-voltage applied to ISCOOL, which was recorded by a precision high-voltage divider. In the first beamtime a 30-kV divider was available, while a 50-kV divider was provided by PTB Braunschweig later on. This allowed independent voltage calibrations and the use of a higher beam energy, favorable for laser-spectroscopic resolution.

Laser spectroscopy on the neutral Ni atoms was performed after neutralization of the ions in a charge-exchange cell [23, 24] filled with potassium vapor. A frequency-doubled single-mode cw titanium-sapphire laser stabilized with a high-resolution wavemeter [25, 26] was used to excite the $3d^9 4s^3D_3 \rightarrow 3d^9 4p^3P_2$ transition at 352.45 nm. The wavemeter was calibrated regularly with a stabilized helium-neon laser. Fluorescence photons from spontaneous emission were detected by four photomultiplier tubes. All isotopes were measured alternating with the reference isotope ^{60}Ni to compensate for remaining long-term drifts in ion velocity or laser frequency.

For the present work, a new data acquisition system called "TILDA" [27] was employed for the first time at COLLAPS. It is based on photon tagging with reference to ISCOOL's release trigger [28] and relaxes the need for

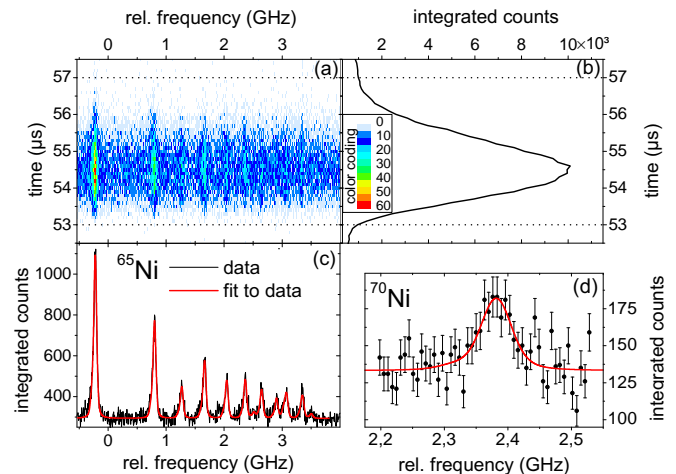


FIG. 1. Frequency-time spectrum of a ^{65}Ni resonance (a), the temporal ion-bunch structure (b) and the laser-spectroscopic resonance spectrum (c). In (d), a resonance of ^{70}Ni is shown. See text for details.

hard-wired gates set during a beamtime. Comparable schemes have previously been employed at other laser experiments with bunched ion beams [29–33]. A typical spectrum recorded with TILDA is shown for ^{65}Ni in Fig. 1(a): The x -axis represents the laser frequency calculated from the scanning voltage at the charge exchange cell, while the y -axis is the time elapsed since the RFQ extraction pulse was recorded. The color represents the number of photons detected within a 100-ns interval during 900 extractions from the RFQ. The time structure of the ion bunch is shown in Fig. 1(b), where the counts at a specified time are integrated over all frequencies. Similarly, summing all counts at a fixed frequency within the (adjustable) time interval between 53 and 57 μs reveals the resonance spectrum of the isotope in Fig. 1(c).

According to an analysis with ISOLTRAP's multi-reflection time-of-flight mass spectrometer [34], the beam of the most exotic isotope ^{70}Ni was dominated by the isobar ^{70}Ga with a ratio of of $\approx 1 : 10^4$. The large amount of isobaric ions can cause an overfilling of ISCOOL and a corresponding shift in beam-energy due to the ions' space-charge potential, which can degrade the accuracy of the spectroscopic measurements. Moreover, non-resonant light emitted by the unwanted ions after collisional excitation or neutralization in the charge exchange cell will reduce the sensitivity for ^{70}Ni . To suppress ^{70}Ga , we took advantage of the different target-release properties of the two elements: the beam gate at ISOLDE, allowing the ions to be transported to the experiments, stayed closed during the first 2 s after the proton impact. Then, most of the more volatile ^{70}Ga had been released from the target while the remaining fraction of ^{70}Ni ($T_{1/2} = 6\text{ s}$) was accumulated for 1.2 s in the RFQ and then sent as a single bunch to COLLAPS, before the next proton pulse arrived. A ^{70}Ni resonance

is shown in Fig. 1(d).

Analysis. — Isotope shifts $\delta\nu^{60,A} = \nu^A - \nu^{60}$ for all isotopes were calculated from their respective center frequency ν^A with respect to the center frequency ν^{60} of the reference isotope ^{60}Ni . Both beamtimes were analyzed individually and a linear displacement in their isotope shifts was corrected by introducing a correction to the main acceleration voltage within the uncertainty of the corresponding voltage dividers. The main acceleration voltage of 30 kV (1st beamtime) was reduced by 3.5 V and the 40 kV (2nd beamtime) was increased by 2.5 V in the analysis. A still remaining scatter in the isotope shifts of individual isotopes as obtained in the two beamtimes could not be explained by their statistical uncertainties. However, this variation was not systematic and could not be traced back to definite reasons individually. Therefore, an additional statistical uncertainty was added to all isotopes, such that the scatter appeared statistically reasonable, i.e., the χ^2_{red} calculated from the deviations between the final isotope shifts of the two beamtimes and their average was reduced to 1. Results are listed in Table I. The changes in mean-square nuclear charge radii $\delta\langle r_c^2 \rangle^{60,A} \equiv \langle r_c^2 \rangle^A - \langle r_c^2 \rangle^{60}$ are obtained using the field-shift factor $F = -783(94)$ MHz/fm² and the mass-shift factor $M_{\alpha=396} = 950(5)$ GHz u, as explained in [18]. These values are in excellent agreement with independent measurements reported in [35]. Negligible deviations from our values in [18] arise from a correction in the analysis code but lead only to insignificant changes of $\delta\langle r_c^2 \rangle$ values. The uncertainties of $\delta\langle r_c^2 \rangle^{60,A}$ are dominated by the correlated error based on the uncertainty of F . The absolute charge radii $R_c \equiv \langle r_c^2 \rangle^{1/2}$ are obtained from $\delta\langle r_c^2 \rangle^{60,A}$ by utilising $R_c(^{60}\text{Ni}) = 3.806(2)$ fm [36].

Theory. — *Ab initio* approaches compute the mean-square charge radius $\langle r_c^2 \rangle$ starting from the calculated point-proton mean-square radius $\langle r_p^2 \rangle$,

$$\langle r_c^2 \rangle = \langle r_p^2 \rangle + \langle R_p^2 \rangle + \frac{N}{Z} \langle R_n^2 \rangle + \langle r^2 \rangle_{\text{so}} + \frac{3\hbar^2}{4m_p^2 c^2}, \quad (1)$$

where $\langle R_p^2 \rangle$ and $\langle R_n^2 \rangle$ are the mean-square charge radii of the proton and the neutron respectively, $\langle r^2 \rangle_{\text{so}}$ denotes a spin-orbit correction [1, 37] and the last term corresponds to the relativistic Darwin-Foldy correction [38], with m_p being the proton mass [39]. The intrinsic (i.e. with respect to the center of mass) squared charge radius operator [1, 40, 41] is employed for $\langle r_p^2 \rangle$ in all calculations. In the present work, the values of $\langle R_p^2 \rangle = 0.709$ fm² [42, 43] and $\langle R_n^2 \rangle = -0.106$ fm² [44] were used.

We employ the following two- plus three-nucleon (3N) interactions from chiral EFT: (i) NNLO_{sat} [45], which gives a good description of charge radii in light and mid-mass isotopes but somewhat underbinds finite nuclei [1, 2, 7, 12, 15, 46, 47]; (ii) 1.8/2.0(EM) [6, 48, 49],

and (iii) $NN+3N(\text{lnl})$ [46], which reproduce ground-state and excitation energies throughout the medium- and heavy mass region, but generally underpredict absolute charge radii [6, 7, 50]. The present work addresses a long sequence of charge radii along the Ni isotopic chain for the first time with three *ab initio* techniques, using these three nuclear interactions. This provides a new, stringent accuracy benchmark of state-of-the-art methods which implement different computational schemes. Importantly, a thorough evaluation of theoretical uncertainties is carried out for each many-body technique, as briefly described in the following.

The self-consistent Green's function (SCGF) approach [51–53] is a full-space correlation-expansion method applicable to the description of medium-mass nuclei [7, 46, 54–58]. Nickel isotopes were recently addressed in Ref. [46] where the calculation of radii, however, was not optimized and lacked theoretical uncertainties. Here, we present nickel charge radii with a full analysis of basis convergence and an assessment of associated theoretical errors. To this end, SCGF calculations are performed in the Gorkov ADC(2) scheme [59, 60] using a spherical harmonic-oscillator basis including up to 14 major shells ($e_{\text{max}} \equiv \max(2n + l) = 13$), with matrix elements of three-body operators further restricted to $e_{3\text{max}} = 16$. Theoretical errors comprise uncertainties arising from both many-body and model-space truncations. The former are estimated from differences between ADC(2) and ADC(3) [61, 62] calculations, available for

TABLE I. Measured isotope shifts $\delta\nu^{60,A}$ relative to ^{60}Ni with statistical uncertainties in parentheses and systematic uncertainties in square brackets. Values for the stable isotopes and ^{68}Ni are those from [18]. The statistical uncertainty includes variations between the two beamtimes that are partially of systematic but uncorrelated origin and change statistically from isotope to isotope, while the systematic uncertainty is restricted to the correlated uncertainty caused by the high-voltage measurement. The extracted change in mean-square charge radius $\delta\langle r_c^2 \rangle^{60,A}$ and the total charge radii R_c are listed with the total uncertainties. Please note that there are diminutive corrections (flips in the last digit) in $\delta\langle r_c^2 \rangle^{60,A}$ compared to [18] caused by a correction in the analysis code.

| A | $\delta\nu^{60,A}/\text{MHz}$ | $\delta\langle r_c^2 \rangle^{60,A}/\text{fm}^2$ | R_c/fm |
|-----|-------------------------------|--|-----------------|
| 58 | -509.1(25)[42] | -0.275(8) | 3.770(2) |
| 59 | -214.3(27)[22] | -0.180(9) | 3.782(2) |
| 60 | 0.0 | 0.0 | 3.806(2) |
| 61 | 280.8(27)[20] | 0.082(5) | 3.817(2) |
| 62 | 503.9(25)[39] | 0.223(5) | 3.835(2) |
| 63 | 784.9(26)[57] | 0.277(8) | 3.842(2) |
| 64 | 1027.2(25)[77] | 0.367(10) | 3.854(2) |
| 65 | 1317.5(26)[94] | 0.385(18) | 3.856(3) |
| 66 | 1526.8(26)[113] | 0.493(17) | 3.870(3) |
| 67 | 1796.6(26)[130] | 0.514(25) | 3.873(3) |
| 68 | 1992.3(27)[147] | 0.619(24) | 3.886(3) |
| 70 | 2377.2(49)[181] | 0.806(24) | 3.910(3) |

closed-shell isotopes. The latter are evaluated from a range of oscillator frequencies, $\hbar\Omega$, within 2 MeV from the optimal values.

The valence-space in-medium similarity renormalization group (VS-IMSRG) method [63–69] decouples a valence-space Hamiltonian and consistent operators from the full-space problem via an approximate unitary transformation. To obtain charge radii, we first decouple the core and valence-space intrinsic proton mean-squared radius operator and then apply Eq. (1). We use the IMSRG(2) approximation where induced operators are truncated at the two-body level and the ensemble normal ordering procedure [67, 70], which captures the physics of 3N forces between valence particles. We take the neutron $p_{3/2}, p_{1/2}, f_{5/2}, g_{9/2}$ valence space outside a ^{56}Ni core, decouple a valence-space Hamiltonian for each isotope studied and diagonalize with the KSHELL code [71] to obtain expectation values for the intrinsic proton mean-square radius operator. While model-space uncertainties are obtained analogously to SCGF, errors due to the many-body method cannot be estimated currently [72]. The $e_{\text{max}}/e_{3\text{max}}$ and basis choices are as in Ref. [6].

The coupled-cluster method performs a similarity transformation of the Hamiltonian and decouples a reference state from its n -particle- n -hole (np - nh) excitations [73–76]. This method was used to compute the structure of doubly-magic nuclei and their neighbours [2, 3, 18, 77–79] and can also be extended to open-shell nuclei [80]. Our calculations for nickel isotopes employ a single-particle basis of up to 13 harmonic oscillator shells with a frequency $\hbar\Omega = 16$ MeV; matrix elements of three-nucleon forces are truncated at $e_{3\text{max}} = 16$. We start from an axially symmetric Hartree-Fock reference, normal-order the resulting Hamiltonian with respect to this state, and truncate it at the two-body level [81, 82]. The ensuing coupled-cluster calculations employ the CCSD approximation, i.e., 1p-1h and 2p-2h excitations of the reference are fully decoupled. While this captures (only) about 90% of the correlation energy, the omission of 3p-3h excitations has a much smaller effect on radii and introduces an estimated 1% uncertainty. Uncertainties from the finite model space are estimated from the difference between calculations in 11 and 13 harmonic oscillator shells. Overall, we estimate coupled-cluster uncertainties on R_c to be +2%/-1%.

The fourth theory considered is nuclear DFT [83]. Here, we focus on non-relativistic energy density functionals (EDF) and employ two EDF parametrizations, namely SV-min [84] as representative of the widely used Skyrme functionals and Fy(Δr , HFB) as the recent example of a Fayans functional [13]. Both have the basic structure in common and are calibrated with the same fitting strategy to the same large body of nuclear ground state data (energy, radii, surface thickness, ...) as described in [84]. The Fayans functional Fy(Δr , HFB) differs in that it contains additional gradient terms in surface and pair-

ing energies [85, 86] and that isotopic shifts of charge radii in the calcium chain were added to the optimization data set. The rms charge radii are computed directly from the nuclear charge form factor. The latter is obtained from folding the proton and neutron densities with the intrinsic charge and magnetic densities of the nucleons, for details see Ref. [87]. The calculations are done with codes allowing for deformed ground states, for SV-min with SkyAx [88], and for Fy(Δr , HFB) with a version of HFBTHO [89] extended to Fayans EDF. Results for spherical nuclei have been counter-checked with our spherical BCS/HFB code, the one which was used for the calibration of both functionals [13, 84]. DFT parametrizations carry statistical uncertainties [84] as well as systematic errors related to principle limitations of the model [90].

Discussion. — Theoretical and experimental nuclear charge radii are compared in Fig. 2. Charge radii R_c provide a comparison on the absolute scale, while the differential charge radii $\delta\langle r_c^2 \rangle$ probe local variations in the nuclear charge distribution more closely, since various theoretical uncertainties cancel in $\delta\langle r_c^2 \rangle$. For instance, the errors on R_c in DFT contain a sizable, nearly constant offset along the chain reflecting a certain vibrational softness for all Ni isotopes. These vibrational corrections enhance total radii and are thus predominantly positive, but are greatly reduced in $\delta\langle r_c^2 \rangle$.

For SV-min based DFT as well as for all *ab initio* calculations based on NNLO_{sat}, the overall agreement with experiment is very good. For both, R_c and $\delta\langle r_c^2 \rangle$, the experimental values are within, or very close to, the theoretical error band, which is of the order of $\approx 1\%$. The same holds for the differential radii $\delta\langle r_c^2 \rangle$ when considering *ab initio* results for the other employed nuclear interactions, see Fig. 2(d), while those deviate notably from experiment in the absolute charge radii R_c , as shown in Fig. 2(c). This is in line with the expectation from previous work [7, 15].

Within the same nuclear many-body method, calculations of R_c with NNLO_{sat} disagree with the results of both 1.8/2.0(EM) and $NN+3N$ (Inl). This illustrates the sensitivity of R_c on the accurate encoding of the relevant physics for medium-mass nuclei into nuclear forces [45, 49]. On the other hand, a comprehensive assessment of uncertainties due to a many-body method itself remains a challenge. Employing the same nuclear interaction in conjunction with different many-body methods is one way to evaluate many-body uncertainties. As shown in Fig. 2(a) and 2(b), the results of SCGF, VS-IMSRG, and coupled cluster theory, all utilising NNLO_{sat}, agree with each other within the theoretical uncertainties, thus, providing strong evidence for the accuracy of the methods. Small differences can be seen for ^{56}Ni where uncertainties of SCGF and VS-IMSRG do not overlap. Note that the error bars in VS-IMSRG account for model-space uncertainties only. We have confirmed that the latter are consistent in size across different meth-

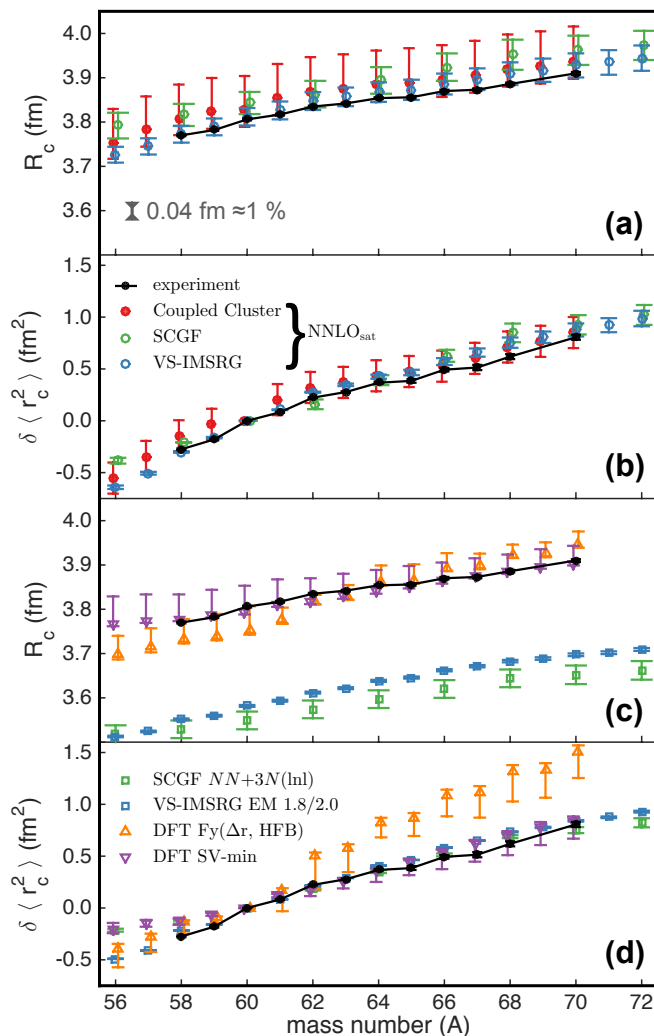


FIG. 2. Nuclear charge radii R_c and differentials $\delta \langle r_c^2 \rangle^{60,A}$ of Ni isotopes with respect to ^{60}Ni as reference. Experimental data are compared to theoretical results. See text for details.

ods.

With respect to nuclear charge radii, the Fayans functional has been very successful in describing an odd-even staggering as well as characteristic kinks typically found at shell closures [12, 13, 15–17]. In contrast, DFT utilizing Skyrme functionals such as SV-min generally fails to reproduce both. However, compared to the large odd-even staggering in Ca [12, 13] or the sizeable kink at $N = 82$ in Sn [17], charge radii along the measured Ni isotopes do not exhibit these features very prominently. Interestingly, the SV-min follows in this case the experimental trend more closely compared to Fy(Δr , HFB), see Fig. 2(d). On closer inspection, analogous conclusions for the mid-shell region also hold for the charge radii of Cu [15] and Sn isotopes [17]. A potential deficiency of the present Fayans functional could be its lack of an isovector component in its pairing part [91]. Hence, future efforts in Fayans-based DFT will focus on pinning down

the (presently unused) isovector term in the pairing functional, see Ref. [92].

Summary — Collinear laser spectroscopy of short-lived nickel isotopes $^{58-68,70}\text{Ni}$ was performed. The extracted nuclear mean-square charge radii R_c benchmark theoretical work applying density functional theory as well as three *ab initio* methods. When the same chiral EFT-based nuclear potential NNLO_{sat} is utilized in all *ab initio* calculations, their results show excellent consistency and they agree well with experiment. Calculations exploiting other nuclear interactions perform equally well for $\delta \langle r_c^2 \rangle$, but struggle in reproducing the absolute radii. Interestingly, in the absence of prominent features such as unusually-large odd-even staggering or kinks in R_c , which have been successfully described by Fayans-based functionals, Skyrme-based DFT yields results closer to experiment. Overall, this comparative work combining experiment, density functional theory and *ab initio* calculations establishes a theoretical accuracy of $\sim 1\%$ for the description of nuclear charge radii in the Ni region.

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