



Search for flavour-changing neutral currents in processes with one top quark and a photon using 81 fb^{-1} of pp collisions at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS experiment



The ATLAS Collaboration ^{*}

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ABSTRACT

A search for flavour-changing neutral current (FCNC) events via the coupling of a top quark, a photon, and an up or charm quark is presented using 81 fb^{-1} of proton–proton collision data taken at a centre-of-mass energy of 13 TeV with the ATLAS detector at the LHC. Events with a photon, an electron or muon, a b -tagged jet, and missing transverse momentum are selected. A neural network based on kinematic variables differentiates between events from signal and background processes. The data are consistent with the background-only hypothesis, and limits are set on the strength of the $tq\gamma$ coupling in an effective field theory. These are also interpreted as 95% CL upper limits on the cross section for FCNC $t\gamma$ production via a left-handed (right-handed) $tu\gamma$ coupling of 36 fb (78 fb) and on the branching ratio for $t \rightarrow \gamma u$ of 2.8×10^{-5} (6.1×10^{-5}). In addition, they are interpreted as 95% CL upper limits on the cross section for FCNC $t\gamma$ production via a left-handed (right-handed) $tc\gamma$ coupling of 40 fb (33 fb) and on the branching ratio for $t \rightarrow \gamma c$ of 22×10^{-5} (18×10^{-5}).

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1. Introduction

Flavour-changing neutral currents (FCNCs) are forbidden at tree level in the Standard Model (SM) and strongly suppressed at higher orders via the GIM mechanism [1]. Several extensions to the SM predict processes involving FCNCs. In particular, some of these models predict the branching ratios of top-quark decays via FCNC to be significantly larger [2] than those predicted by the SM, which are of the order of 10^{-14} [2]. Examples are R-parity-violating supersymmetric models [3] and models with two Higgs doublets [4]. Such models would allow the production of top quarks via FCNCs at a measurable rate.

This Letter presents a search for FCNCs in processes with a top quark (t) and a photon (γ) based on data collected with the ATLAS experiment at $\sqrt{s} = 13 \text{ TeV}$. This analysis is most sensitive to the production of a single top quark plus a photon, but also considers the decay of a pair-produced top quark into an up or charm quark (q) plus a photon. Tree-level Feynman diagrams for these processes are shown in Fig. 1, where in both cases, exactly one top quark decays via the SM-favoured tWb coupling. Compared to the SM production of a top quark and a photon, the FCNC processes result in higher photon transverse momenta on average.

FCNC contributions to the decay mode ($t \rightarrow q\gamma$) and the production mode ($q \rightarrow t\gamma$) can be expressed in terms of effective coupling parameters but also in terms of branching ratios and cross sections [5,6]. In the former case and following the notation in Ref. [7], the corresponding operators are $O_{uB}^{(ij)}$ and $O_{uW}^{(ij)}$, where $i \neq j$ are indices for the quark generation. In general, left-handed (LH) and right-handed (RH) couplings could exist, which result in different kinematic properties of the top-quark decay products, such as the transverse momentum of the charged lepton in semileptonic top-quark decays. The most stringent limits to date are limits on branching ratios of $\mathcal{B}(t \rightarrow u\gamma) < 1.3 \cdot 10^{-4}$ and $\mathcal{B}(t \rightarrow c\gamma) < 1.7 \cdot 10^{-3}$ set by the CMS Collaboration, assuming equal left- and right-handed couplings [8].

2. ATLAS detector

The ATLAS experiment [9] at the LHC is a multi-purpose particle detector with a forward-backward symmetric cylindrical geometry and a near 4π coverage in solid angle.¹ It consists of an inner tracking detector (ID) surrounded by a thin superconduct-

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms

^{*} E-mail address: atlas.publications@cern.ch.

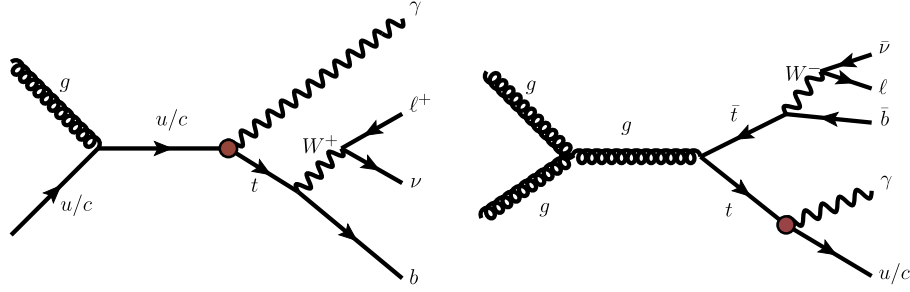


Fig. 1. Tree-level Feynman diagrams for top-quark production (left) and decay (right) via FCNCs. The $tq\gamma$ vertex, which is not present in the SM, is highlighted.

ing solenoid providing a 2 T axial magnetic field, electromagnetic and hadron calorimeters, and a muon spectrometer (MS). The ID covers the pseudorapidity range $|\eta| < 2.5$. It consists of silicon pixel, silicon microstrip, and transition radiation tracking detectors. Lead/liquid-argon (LAr) sampling calorimeters provide electromagnetic (EM) energy measurements with high granularity. A hadron (steel/scintillator-tile) calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The end-cap and forward regions are instrumented with LAr calorimeters for both EM and hadronic energy measurements up to $|\eta| = 4.9$. The MS surrounds the calorimeters and is based on three large air-core toroidal superconducting magnets with eight coils each. The field integral of the toroids ranges between 2.0 and 6.0 Tm across most of the detector. The MS includes a system of precision tracking chambers and fast detectors for triggering. A two-level trigger system is used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to at most nearly 100 kHz. This is followed by a software-based trigger that reduces the accepted event rate to 1 kHz on average depending on the data-taking conditions.

3. Analysis strategy

The search strategy selects events with a final state containing one prompt photon and the decay products of a top quark, namely an electron or a muon, a b -tagged jet, and missing transverse momentum, and estimates contributions from FCNC processes in this background-dominated dataset. A signal region (SR) is defined by loose requirements on the kinematic properties of the final-state objects, giving rise to a large acceptance for signal events in the production mode. With this selection, the search is most sensitive to FCNCs in this mode, but the decay mode is included in the analysis. The main background contribution stems from electrons misidentified as photons, primarily in top-quark–anti-top-quark events ($t\bar{t}$). These contributions and contributions from hadrons misidentified as photons (both labelled “fakes” in the following) are modelled by Monte Carlo (MC) simulations and scaled to data-driven estimates. Photons produced in association with a leptonically decaying W or Z boson are estimated in control regions (CRs) which do not overlap with the SR but are kinematically similar to it. The predictions for other small prompt-photon background processes are taken from MC simulation and include $t\bar{t}\gamma$ production, single-top quark production in association with a photon and the production of two massive gauge bosons with a photon. As the latter two processes result in a small contribution to the total background prediction, it is sufficient that in these processes prompt photons are generated by the parton-shower program. Signal and background events are distinguished using a neural network (NN).

Finally, the signal contribution is estimated in a profile likelihood fit to the NN output distributions in the SR and the CRs, in which each source of systematic uncertainty is modelled as a nuisance parameter.

4. Data and simulation

The proton–proton (pp) collision data analysed were recorded with the ATLAS detector from 2015 to 2017 at a centre-of-mass energy of $\sqrt{s} = 13$ TeV. The average number of interactions per bunch crossing was 13.4, 25.1, and 37.8 in 2015, 2016 and 2017, respectively. Events were selected by single-lepton triggers [10] and required to have at least one reconstructed primary vertex with at least three assigned tracks that have a transverse momentum greater than 400 MeV. After the application of data-quality requirements, the data sample corresponds to an integrated luminosity of 81 fb^{-1} . It is obtained using the LUCID-2 detector [11] for the primary luminosity measurements.

The data were modelled by MC simulations of the signal and background processes. After event generation, the response of the ATLAS detector was simulated using GEANT 4 [12] with a full detector model [13] or modelled by a fast simulation [14]. To account for additional pp collisions (pile-up), inelastic pp interactions were superimposed on the hard-scattering events and weighted according to the observed pile-up distribution. The pile-up events were simulated using PYTHIA 8.186 [15], with the A3 [16] set of tuned parameters (A3 tune).

The simulated signal samples were generated using MADGRAPH5_aMC@NLO 2.4.3 [17] with the TopFCNC model [6,18] at next-to-leading order (NLO) in QCD and the NNPDF3.0NLO [19] set of parton distribution functions (PDFs). The parton showering was done with PYTHIA 8.212 with the A14 tune set [20]. Simulated samples of SM $t\bar{t}$ and single-top-quark events were generated using POWHEG-Box [21–27] with the NNPDF3.0NLO PDF set. The parton showering, hadronisation, and the underlying event were modelled using PYTHIA 8.230. The top-quark mass m_{top} was set to 172.5 GeV in these samples, and the h_{damp} parameter that controls the transverse momentum of the first gluon emitted was set to 1.5 times m_{top} . Samples of $t\bar{t}\gamma$ events were generated as $2 \rightarrow 7$ process at leading order using MADGRAPH5_aMC@NLO 2.3.3 and the NNPDF2.3LO PDF [28] set and with the following fiducial photon criteria [29]: photon $p_T > 15$ GeV and $|\eta| < 5.0$, charged-lepton $p_T > 15$ GeV and $|\eta| < 5.0$, and $\Delta R < 0.2$ between the photon and any charged final-state particle. The cross sections for SM $t\bar{t}$ and single-top-quark production are scaled to the NNLO+NNLL predictions [30–33]. The leading-order cross section for $t\bar{t}\gamma$ production of 4.62 pb [29] is scaled to the NLO predictions [34] using a k -factor of 1.24. The NNLO+NNLL cross section for SM $t\bar{t}$ production is also used to normalise the signal in the decay mode, using the corresponding FCNC branching ratio. The cross sections for the signal in the production mode, however, are calculated at NLO with MADGRAPH5_aMC@NLO.

of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

For the study of systematic uncertainties in the modelling of processes involving top quarks, simulated $t\bar{t}$ samples were produced with POWHEG-Box + HERWIG 7.0.4 [35,36] and MADGRAPH5_aMC@NLO 2.6.0 plus PYTHIA 8.212. The MMHT2014LO [37] PDF set is used together with the H7-UE-MMHT [36] tune. An additional $t\bar{t}$ sample was produced with h_{damp} set to three times m_{top} and the factorisation and renormalisation scales set to half their nominal values using the A14 tune. For the tW process, a sample was produced with an alternative scheme for removing the overlap with $t\bar{t}$ production [38]. Moreover, single-top-quark samples were produced with POWHEG-Box + HERWIG 7.0.4 and MADGRAPH5_aMC@NLO 2.6.2 plus PYTHIA 8.212. The NNPDF2.3LO PDF set is used as well as the A14 tune.

Processes with one or two heavy gauge bosons, in particular the processes $W+\gamma$ +jets and $Z+\gamma$ +jets, were simulated using SHERPA 2.2.1 and 2.2.2 [39] with the matrix elements calculated using Comix [40] and OpenLoops [41]. All matrix elements were merged with the SHERPA parton showering [42] according to the ME+PS@NLO [43] prescription. The NNPDF3.0NNLO PDF set was used.

An overlap removal scheme was applied to remove double-counting of events stemming from photon radiation in samples in which a photon was not explicitly required in the final state [29]. This applies to the processes $t\bar{t}$, W +jets and Z +jets in order to remove the overlap with the $t\bar{t}\gamma$, $W+\gamma$ +jets and $Z+\gamma$ +jets samples.

Mismodelling of the photon p_T distribution is observed in the $W+\gamma$ +jet and $Z+\gamma$ CRs, which are defined in Section 5. The photon p_T spectrum in the $W+\gamma$ +jets and $Z+\gamma$ +jets processes was corrected by adjusting the MC prediction to the data in five p_T bins using a linear function that only changes the shape of the distribution and not its normalisation. This correction to the photon p_T improves the modelling of the NN output distribution in the CRs.

As discussed in Sections 6 and 7, the contribution of events with electrons and hadrons misidentified as photons is corrected using data. The contribution from processes with hadrons misidentified as leptons is estimated to be negligible.

5. Object and event selection

Electrons are reconstructed from clusters of energy deposits in the electromagnetic calorimeter cells with a matched ID track [44]. They are required to meet the *tight* identification criteria [44], and their tracks must point to the primary vertex. They must have a transverse momentum p_T larger than 27 GeV and $|\eta_{\text{cluster}}| < 2.47$, excluding $1.37 < |\eta_{\text{cluster}}| < 1.52$, where $|\eta_{\text{cluster}}|$ is the pseudorapidity of the electron's energy cluster. Muons are reconstructed by combining a track in the MS with a track in the ID [45]. They are required to meet the *medium* identification criteria [45] and must point to the primary vertex. They must have $p_T > 27$ GeV and $|\eta| < 2.5$. Isolated electrons and muons are selected by requiring the amount of energy in nearby energy deposits in the calorimeters and the scalar sum of the transverse momenta of nearby tracks in the ID to be small.

Photons are reconstructed from clusters of energy deposits in the electromagnetic calorimeter cells with no matched ID track (unconverted photons) or with one or two matched ID tracks that are compatible with the tracks from an electron or positron from a photon conversion (converted photons) [44]. They must have $p_T > 20$ GeV and $|\eta_{\text{cluster}}| < 2.37$, excluding $1.37 < |\eta_{\text{cluster}}| < 1.52$. They are required to meet the *tight* identification criteria for the shape of the shower in the electromagnetic calorimeter (shower shape) and for the energy deposited in the hadronic calorimeter [44]. Photons must be isolated from nearby energy deposits in the calorimeter and nearby tracks in the ID: the sum of the energy deposited (p_T of the tracks) within $\Delta R = 0.4$ ($\Delta R = 0.2$) of the photon

direction is required to be smaller than $0.022 \times p_T + 2.45$ GeV ($0.065 \times p_T$), excluding the photon energy deposition (tracks associated with the photon).

Jets are reconstructed from topological clusters [46,47] in the calorimeters with the anti- k_t algorithm [48] using FastJet [49] and a radius parameter of 0.4. Their energy is calibrated [50], and they must fulfil $p_T > 25$ GeV and $|\eta| < 2.5$. Jets with $p_T < 120$ GeV and $|\eta| < 2.4$ are required to pass a requirement on the jet-vertex-tagger (JVT) [51] to suppress pile-up jets. Jets are b -tagged with the MV2c10 algorithm [52], which uses a boosted decision tree with several b -tagging algorithms as input. The b -tagging efficiency for jets that originate from the hadronisation of b -quarks is 60% in simulated $t\bar{t}$ events. The b -tagging rejection² for jets that originate from the hadronisation of c -quarks (u -, d -, s -quarks or gluons) is 23 (1200).

The magnitude of the missing transverse momentum E_T^{miss} is reconstructed from the vector sum of the p_T of leptons, photons, and jets, as well as ID tracks that point to the primary vertex but are not associated with a reconstructed object (soft term) [53].

To avoid double-counting, objects are removed in the following order: electrons sharing a track with a muon; jets within $\Delta R = 0.2$ of an electron; electrons within $\Delta R = 0.4$ of a jet; jets within $\Delta R = 0.4$ of a muon if they have at most two associated tracks; muons within $\Delta R = 0.4$ of a jet; photons within $\Delta R = 0.4$ of an electron or muon; jets within $\Delta R = 0.4$ of a photon.

Scale factors (SFs) are used to correct the efficiencies in simulation in order to match the efficiencies measured in data for the electron [44] and muon [45] trigger, reconstruction, identification, and isolation criteria, as well as for the photon identification [44] and isolation requirements. SFs are also applied for the JVT requirement and for the b -tagging efficiencies for jets that originate from the hadronisation of b -quarks [54], c -quarks [52], and u -, d -, s -quarks or gluons [55].

The selected events have exactly one electron or muon, exactly one photon, exactly one b -tagged jet and no further jets, and $E_T^{\text{miss}} > 30$ GeV. This selection defines the SR with signal efficiencies³ for the production mode of 3.03% (2.45%) for the LH (RH) $tu\gamma$ coupling and of 3.79% (3.14%) for the LH (RH) $tc\gamma$ coupling. These efficiencies are defined with respect to signal events that include a leptonic decay of the W boson, i.e. a decay with an electron, muon or tau lepton. The efficiencies for the couplings that involve a c -quark are larger than those that involve a u -quark because of the difference in the kinematics that arises from the difference between the u - and c -quark PDFs. The difference between the efficiencies for the LH and the RH couplings is due to the kinematic distributions of the W boson's decay products, which differ depending on the handedness of the top quark. The efficiencies for the decay mode are 0.45% and 0.51% for the $tu\gamma$ and the $tc\gamma$ coupling, respectively, and are lower due to the requirement of not more than one jet in the final state, i.e. this analysis is optimised for the production mode. The absolute statistical uncertainties in the efficiencies are 0.03% or smaller. In the SR, 9557 data events are selected. The ratio of production-mode to decay-mode event yields is 4.2 (5.3) for the LH (RH) $tu\gamma$ coupling and it is 0.86 (0.68) for the LH (RH) $tc\gamma$ coupling, i.e. in the case of $tu\gamma$ coupling, the dominant signal process is indeed the production mode. However, in the case of the $tc\gamma$ coupling, the decay mode also plays an important role, because the production mode is suppressed by the c -quark PDF.

² The rejection is defined as the inverse of the efficiency.

³ Here, the signal efficiency includes the signal loss due to the limited acceptance of the detector.

Two CRs are defined for the $W+\gamma$ +jets and $Z+\gamma$ +jets processes, which are dominated by the respective background process and kinematically close to the SR. The $W+\gamma$ +jet CR is defined by the same criteria as the SR with two modifications: the jet must not be b -tagged and the lepton–photon invariant mass must be outside the range 60–100 GeV to suppress the contribution from Z +jets events with one electron that is misidentified as a photon. The $Z+\gamma$ CR is defined by requiring exactly one photon and exactly two leptons with the same flavour but opposite electric charge. No requirement is made in the $Z+\gamma$ CR on the number of jets or E_T^{miss} . In the $W+\gamma$ +jet CR, a total number of 127 864 events are observed, and in the $Z+\gamma$ CR, the total number is 85 347.

6. Data-driven estimate of electrons misidentified as photons

Electrons can be misidentified as photons, for example, if the track of the electron is not reconstructed or if the matching criteria between the track and cluster are not met. In particular, dileptonic $t\bar{t}$ events with at least one electron in the final state can enter the SR if an electron is misidentified as a photon. The probability for an electron to be misidentified as a photon, $f_{e\rightarrow\gamma}$, is measured from data and simulation following the methodology used previously [56], and a SF is derived that is applied to the simulation.

Two regions are defined to measure $f_{e\rightarrow\gamma}$, called “electron fake regions” (EFR) in the following. The $Z \rightarrow e\gamma$ ($Z \rightarrow ee$) EFR is defined by requiring exactly one electron and one photon (exactly two electrons with opposite electric charge and no photons) with an electron–photon (dielectron) invariant mass in the range 60–120 GeV, a veto on the presence of jets, and $E_T^{\text{miss}} < 30$ GeV. Neither EFR overlaps with the SR nor the CRs. The $Z \rightarrow e\gamma$ EFR is rich in $Z \rightarrow ee$ events with one electron misidentified as a photon, and the $Z \rightarrow ee$ EFR is rich in $Z \rightarrow ee$ events.

In the $Z \rightarrow ee$ and the $Z \rightarrow e\gamma$ EFRs, the dielectron invariant mass or the electron–photon invariant mass, respectively, is fitted with analytic signal (for $Z \rightarrow ee$ with both electrons correctly identified or with one electron misidentified as a photon, respectively) and background functions. The signal function is a double-sided Crystal Ball function, and the background function is a fourth-order Bernstein polynomial. The integrals of the aforementioned fitted signal functions are divided in order to estimate $2f_{e\rightarrow\gamma}$, where the factor of two accounts for the two electrons in $Z \rightarrow ee$ events that may be misidentified as a photon. In the $Z \rightarrow e\gamma$ EFR, the expected contribution from $Z \rightarrow ee\gamma$ events, relative to the signal, is 8.8% and it is subtracted from the integral, because this process mainly contributes to events with prompt photons and in which one electron was not reconstructed or did not pass the identification or isolation criteria.

Systematic uncertainties from several sources are evaluated: the range of the invariant-mass fit is changed from 60–120 GeV to 65–115 GeV; the parameters of the signal function, except for its normalisation, are set to the values extracted from simulation and a Gaussian function is used for the background; instead of subtracting the expected relative contribution from $Z \rightarrow ee\gamma$ events in the $Z \rightarrow e\gamma$ EFR, the expected absolute contribution is subtracted. For each of these variations, a systematic uncertainty for $f_{e\rightarrow\gamma}$ is estimated as the deviation from the nominal value. The largest effect is due to the variation of the signal and background functions. The value of $f_{e\rightarrow\gamma}$ is $3.11\% \pm 0.01\%$ (stat.) $\pm 0.13\%$ (syst.). The SF for the simulation is 0.978 ± 0.004 (stat.) ± 0.040 (syst.).

The modelling of kinematic variables is checked in a validation region. The event selection for this validation region is similar to the SR selection, but a few changes are made in order to enhance the contributions from $Z \rightarrow ee$ events and dileptonic $t\bar{t}$ events with a misidentified electron, while ensuring no overlap with the SR, the CRs, or the EFRs. The validation region is defined by requir-

ing exactly one photon and one electron with an invariant mass in the range 70–110 GeV, at least one jet, and $E_T^{\text{miss}} < 30$ GeV. Satisfactory modelling of the kinematic variables is observed, but the relative uncertainty in $f_{e\rightarrow\gamma}$ is increased to 10% in order to cover the difference in the normalisation observed between data and the prediction.

7. Data-driven estimate of hadrons misidentified as photons

In some cases, hadrons can be misidentified as photons. For example, this can happen when a meson decays into two photons that are reconstructed as a single cluster in the electromagnetic calorimeter. Processes such as $t\bar{t}$ production can enter the SR if a high-energy hadron is misidentified as a photon. The number of events with misidentified hadrons is estimated from data, and a SF is applied to the simulation, defined as the estimated number of events in data divided by the predicted number in simulation. The SF is only used to correct the overall normalisation of this background, and the shapes of kinematic distributions are taken from simulations with the associated systematic uncertainties.

Three hadron fake regions (HFR) are defined by the same criteria as the SR but with modified photon criteria: $\text{HFR}^{\text{pass/fail}}$, $\text{HFR}^{\text{fail/pass}}$, and $\text{HFR}^{\text{fail/fail}}$. If the first index is “pass”, the photon has met the identification criteria defined in Section 5. If the first index is “fail”, the photon has failed to meet at least one of the criteria on the shower shapes that are calculated from the finely-segmented first layer of the electromagnetic calorimeter; however, it is required to meet all tight photon-identification criteria for the other shower variables. The second index represents whether the photon meets or fails to meet the isolation criterion.

Only the first-layer shower shapes are considered for the first index because these are mostly sensitive to the core of the shower and are expected to be only weakly correlated with the isolation variables, which are sensitive to the energy surrounding the photon. The number of SR events with misidentified hadrons is estimated as $N(\text{HFR}^{\text{pass/fail}}) \times N(\text{HFR}^{\text{fail/pass}}) / N(\text{HFR}^{\text{fail/fail}})$, where N is the number of observed events after subtracting both the expected number of events with misidentified electrons and the expected fraction of events with prompt photons (leakage). This estimate is additionally corrected for the non-zero correlation between the criteria for the first-layer shower shapes and the isolation criterion. The correction factor is determined using MC simulations and amounts to 0.85 with a statistical uncertainty of 0.14.

Systematic uncertainties from several sources are evaluated: the correction factor for the non-zero correlation is conservatively varied by $\pm 50\%$; the SF for misidentified electrons, used for the subtraction discussed above, is varied by one standard deviation up and down ($\pm 1\sigma$), and the larger of the two deviations is considered as a systematic uncertainty; and instead of subtracting the expected fraction of events with prompt photons, the expected prompt-photon contribution is subtracted in each HFR. For each of these variations, a systematic uncertainty is estimated as the resulting deviation from the nominal value. The largest effect is the variation of the correction factor. The SF for the simulations is 1.7 ± 0.3 (stat.) ± 1.0 (syst.).

8. Neural network for discrimination between signal and background

The signal is distinguished from the sum of the background processes by a fully connected feed-forward neural network (NN) with backpropagation, implemented in Keras [57] with the TensorFlow [58] back end. Separate NNs are trained for FCNC processes with a $t\bar{u}\gamma$ or a $t\bar{c}\gamma$ vertex and with LH or RH couplings. For the signal sample, the production mode was chosen, since the

kinematic differences between the production mode and the background are more pronounced than for the decay mode and thus lead to better discrimination between signal and background.

Ten variables are inputs to the NN: the p_T of the photon, the lepton and the jet; the charge of the lepton; E_T^{miss} ; the lepton–photon and lepton–jet invariant masses; the ΔR between the lepton and the photon, between the lepton and the jet, and between the jet and the photon. Additional kinematic variables were tested and did not improve the discrimination between signal and background. All variables are transformed using scikit-learn’s [59] RobustScaler.

The NN is trained with the Adam optimiser [60]. The MC samples are divided into 80% for training and 20% for testing, so that approximately 63 000 background MC events and approximately 10 000–13 000 signal MC events are available for the training—depending on the coupling. Two hidden layers with 11 nodes each are used and the hyperparameters of the NN are chosen from a series of tested values in a procedure with threefold cross validation.

9. Systematic uncertainties

Systematic effects may change the expected numbers of events from the signal and background processes and the shape of the fitted discriminants in the SR and in the CRs. These effects are evaluated by varying each source of systematic uncertainty by $\pm 1\sigma$ and considering the resulting difference from the nominal expectation as the uncertainty. For some sources, only one variation is available and the difference is symmetrised using the full difference. For sources with two variations, their effects are symmetrised using the average difference from the nominal prediction.

Uncertainties due to the theoretical cross sections are evaluated by varying them by $\pm 5.6\%$ for $t\bar{t}$ production [28,30,61–63], by $\pm 8\%$ for $t\bar{t}+\gamma$ production [34], by $^{+4.0\%}_{-3.4\%}$ ($^{+5.0\%}_{-4.5\%}$) for t -channel single- t (single- \bar{t}) production [31], by $^{+3.6\%}_{-3.1\%}$ ($^{+4.8\%}_{-4.3\%}$) for s -channel single- t (single- \bar{t}) production [32], by $\pm 5.3\%$ for tW production [33], by $\pm 5\%$ for W +jets and Z +jets production [64], and by $\pm 6\%$ for diboson production [65]. No cross-section uncertainty is considered for $W+\gamma$ +jets and $Z+\gamma$ +jets production, because their normalisations are determined in the fit.

Uncertainties due to the modelling of the signal are estimated by considering variations of the renormalisation and factorisation scales by factors of 2 and 0.5, but normalising the signal to the nominal cross section. In each bin of a distribution, the largest deviation among all variations is considered as an uncertainty (envelope). In addition, uncertainties due to the PDFs are estimated by following the PDF4LHC prescription for Run 2 [66].

Uncertainties due to the renormalisation and factorisation scales and from the PDFs of the background processes are estimated separately for each process following the same procedure as for the signal. For the $t\bar{t}$ and single-top processes, however, the scale variations are already included in the estimation of the uncertainty in the modelling of the initial-state radiation (see below). For the $W+\gamma$ +jets and the $Z+\gamma$ +jets processes, a correction is applied to the photon p_T spectrum, as described in Section 4. To account for the uncertainty due to the photon p_T correction, a conservative uncertainty is applied, for which the prediction with the correction applied is compared with the prediction without the correction.

For all background processes—except for $W+\gamma$ +jets and $Z+\gamma$ +jets production, for which the normalisation is estimated by a free parameter in the fit—an uncertainty of 2% in the integrated luminosity is included [67]. The uncertainty due to pile-up is determined by varying the average number of interactions by 9% in the simulation. The uncertainties due to the SFs for electrons and

Table 1

Expected number of events for the different background contributions in the SR and the two CRs after the fit including all uncertainties, as well as the observed number of events.

	SR	$W+\gamma$ +jet CR	$Z+\gamma$ CR
$e \rightarrow \gamma$ fake	4500 ± 400	8200 ± 1300	236 ± 32
$j \rightarrow \gamma$ fake	260 ± 200	2900 ± 2000	1300 ± 1000
$Z+\gamma$ +jets	780 ± 100	13400 ± 1300	81400 ± 1900
$W+\gamma$ +jets	2200 ± 400	101200 ± 2800	6 ± 2
Other prompt γ	1800 ± 400	1900 ± 500	2140 ± 200
Total background	9500 ± 220	127700 ± 3000	85100 ± 1600
Data	9557	127864	85347

hadrons that are misidentified as photons are determined as described in Sections 6 and 7.

To estimate the uncertainty due to the production of W and Z bosons together with b -quarks, the shape of the SR distribution in events with jets that originate from the hadronisation of a b -quark is used for events with jets that originate from other quarks or gluons (and vice versa) for the W +jets, $W+\gamma$ +jets, Z +jets, $Z+\gamma$ +jets, and diboson processes. Differences between the shapes of these backgrounds in association with b -quarks or with other quarks or gluons are small, however. An additional uncertainty in the normalisation of $W+\gamma$ production in association with b -quarks of 50% is assigned, covering observed differences between data and predictions in measurements of W and Z bosons in association with b -quarks [68–71].

To estimate the uncertainty due to the modelling of initial- and final-state radiation in $t\bar{t}$ and single-top-quark production, the effects of varying the A14 tune’s parameter values is evaluated. In addition, for $t\bar{t}$ production, the sample generated with $h_{\text{damp}} = 3m_{\text{top}}$, the factorisation and renormalisation scales set to half their nominal values, and a variation of the A14 tune’s parameter values is used to estimate the uncertainty due to the modelling of initial-state radiation. To estimate the uncertainty due to our choice of generator and shower programs for $t\bar{t}$ and single-top-quark production, the nominal MC samples, generated with POWHEG-BOX + PYTHIA 8, are replaced with samples generated with MADGRAPH5_aMC@NLO + PYTHIA 8 and with POWHEG-BOX + HERWIG 7. To evaluate the uncertainty due to the scheme for removing the overlap with $t\bar{t}$ production for the tW process, the nominal sample is compared with a sample produced with an alternative scheme [38].

For the triggering, reconstruction, identification, and calibration of the objects, the following systematic uncertainties are evaluated: electron and muon triggers, reconstruction, identification and isolation SFs [44,45]; photon identification [44] and isolation SFs; electron- and photon-energy and muon-momentum calibration and resolution [44,45]; jet energy scale (JES) [50] and jet energy resolution (JER) [72]; JVT SF; b -tagging SFs [52,55,73]; E_T^{miss} soft term [53].

10. Results

The normalisations of the signal contribution and the two contributions from $W/Z+\gamma$ +jets production are obtained from a simultaneous binned profile-likelihood fit to the NN-output distribution of the SR and $W+\gamma$ +jet CR as well as the photon- p_T distribution of the $Z+\gamma$ CR. The signal contribution scales the production- and decay-mode contributions consistently. Each source of systematic uncertainty is associated with a nuisance parameter. In Table 1, the expected number of events after a background-only fit to the SR and CRs for the case of the LH $tu\gamma$ coupling are shown, as well as the observed number of events. Fig. 2 shows the cor-

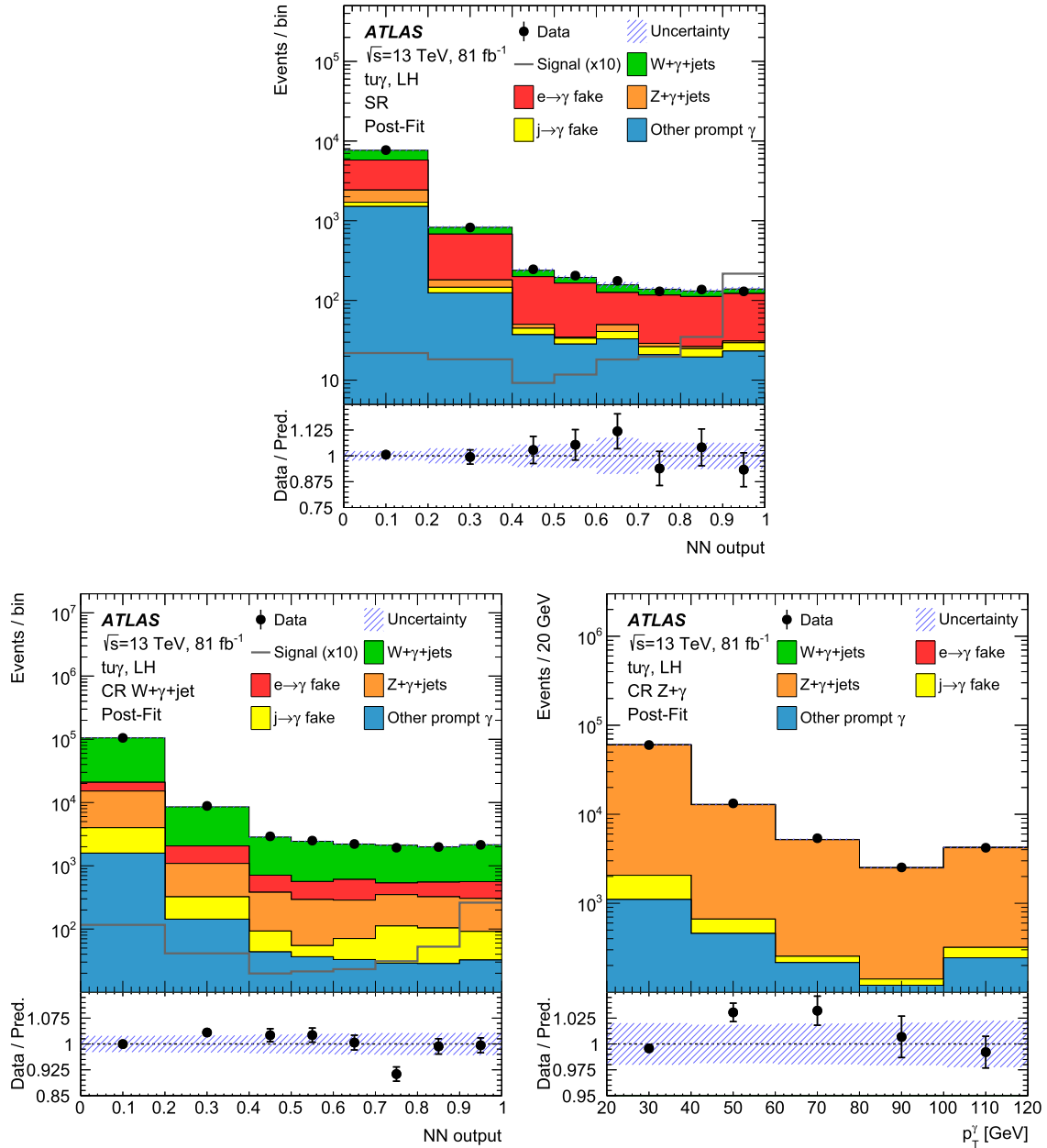


Fig. 2. Post-fit distributions of a background-only fit to the SR and the CRs of the NN output in the SR (top) and the $W+\gamma$ +jet CR (bottom left) and of the p_T -distribution of the $Z+\gamma$ CR (bottom right). The last bin of the distribution in the $Z+\gamma$ CR contains the overflow. In addition, in the SR and in the $W+\gamma$ +jet CR, the expected signal is overlaid for an effective coupling strength corresponding to the observed limit multiplied by a factor of ten. In the $Z+\gamma$ CR, the expected signal is not shown, because it is negligible.

responding post-fit distributions. The qualitative features of these distributions are similar for the other couplings studied.

The data and SM predictions agree within uncertainties and no significant FCNC contributions are observed. From the 95% confidence level (CL) limits on the signal contribution, derived using the CL_s method [74], the corresponding limits on the effective coupling parameters are calculated, and from these the limits on the production cross section and branching ratios are calculated. The background contributions from photons produced in association with a leptonically decaying W or Z boson are scaled by normalisation factors estimated to be 1.25 ± 0.09 and 1.12 ± 0.12 , respectively, from the fit for the LH $t\gamma$ coupling. The normalisation values determined in the fits for the other couplings are similar. The observed and expected 95% CL limits on the effective coupling strengths, the production cross section and the branching

ratio are summarised in Table 2 for different vertices and couplings. The sources of systematic uncertainty with the largest impact on the estimated signal contribution depend on the coupling studied. Among them are the jet energy resolution, the reweighting of the photon p_T , the factorisation and renormalisation scales, the choice of generator for the simulation of the $t\bar{t}$ and single-top processes, and the uncertainties due to the limited number of Monte Carlo events. The resulting limits on the strength of the effective operators are complementary to current limits on the single operators from a search for an FCNC tqZ coupling [75].

11. Conclusion

A search for flavour-changing neutral currents (FCNCs) in events with one top quark and a photon is presented using 81 fb^{-1} of

Table 2

Observed (expected) 95% CL limits on the effective coupling strengths for different vertices and couplings, the production cross section, and the branching ratio. For the former, the energy scale is assumed to be $\Lambda = 1$ TeV.

Observable	Vertex	Coupling	Obs.	Exp.
$ C_{uW}^{(13)*} + C_{uB}^{(13)*} $	$tu\gamma$	LH	0.19	$0.22_{-0.03}^{+0.04}$
$ C_{uW}^{(31)} + C_{uB}^{(31)} $	$tu\gamma$	RH	0.27	$0.27_{-0.04}^{+0.05}$
$ C_{uW}^{(23)*} + C_{uB}^{(23)*} $	$tc\gamma$	LH	0.52	$0.57_{-0.09}^{+0.11}$
$ C_{uW}^{(32)} + C_{uB}^{(32)} $	$tc\gamma$	RH	0.48	$0.59_{-0.09}^{+0.12}$
$\sigma(pp \rightarrow t\gamma)$ [fb]	$tu\gamma$	LH	36	52_{-14}^{+21}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$tu\gamma$	RH	78	75_{-21}^{+31}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$tc\gamma$	LH	40	49_{-14}^{+20}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$tc\gamma$	RH	33	52_{-14}^{+22}
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$tu\gamma$	LH	2.8	$4.0_{-1.1}^{+1.6}$
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$tu\gamma$	RH	6.1	$5.9_{-1.6}^{+2.4}$
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$tc\gamma$	LH	22	27_{-7}^{+11}
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$tc\gamma$	RH	18	28_{-8}^{+12}

$\sqrt{s} = 13$ TeV pp data collected with the ATLAS detector at the LHC. Events with a photon, an electron or muon, a b -tagged jet, and missing transverse momentum are selected. The contribution from events with electrons or hadrons that are misidentified as photons is estimated using data, and the two main background processes with a prompt photon are estimated in control regions. A neural network is used to distinguish the signal and background events, and the data are consistent with the background-only hypothesis. Limits are set on the strength of effective operators that introduce a left- or right-handed flavour-changing $tq\gamma$ coupling with an up-type quark q , on the production cross section for FCNC $t\gamma$ production, and on the branching ratio $t \rightarrow \gamma q$. The limits on the branching ratio and on the $t\gamma$ production cross section are the most stringent to date. The resulting limits on the strength of the effective operators are the most stringent limits obtained in searches for events with a $tq\gamma$ vertex.

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G. Aad¹⁰¹, B. Abbott¹²⁸, D.C. Abbott¹⁰², A. Abed Abud^{70a,70b}, K. Abeling⁵³, D.K. Abhayasinghe⁹³, S.H. Abidi¹⁶⁷, O.S. AbouZeid⁴⁰, N.L. Abraham¹⁵⁶, H. Abramowicz¹⁶¹, H. Abreu¹⁶⁰, Y. Abulaiti⁶, B.S. Acharya^{66a,66b,m}, B. Achkar⁵³, S. Adachi¹⁶³, L. Adam⁹⁹, C. Adam Bourdarios⁵, L. Adamczyk^{83a}, L. Adamek¹⁶⁷, J. Adelman¹²¹, M. Adersberger¹¹⁴, A. Adiguzel^{12c}, S. Adorni⁵⁴, T. Auyeub¹⁴⁴,

A.A. Affolder ¹⁴⁶, Y. Afik ¹⁶⁰, C. Agapopoulou ¹³², M.N. Agaras ³⁸, A. Aggarwal ¹¹⁹, C. Agheorghiesei ^{27c},
 J.A. Aguilar-Saavedra ^{140f,140a,ag}, F. Ahmadov ⁷⁹, W.S. Ahmed ¹⁰³, X. Ai ¹⁸, G. Aielli ^{73a,73b}, S. Akatsuka ⁸⁵,
 T.P.A. Åkesson ⁹⁶, E. Akilli ⁵⁴, A.V. Akimov ¹¹⁰, K. Al Khoury ¹³², G.L. Alberghi ^{23b,23a}, J. Albert ¹⁷⁶,
 M.J. Alconada Verzini ¹⁶¹, S. Alderweireldt ³⁶, M. Aleksa ³⁶, I.N. Aleksandrov ⁷⁹, C. Alexa ^{27b},
 D. Alexandre ¹⁹, T. Alexopoulos ¹⁰, A. Alfonsi ¹²⁰, F. Alfonsi ^{23b,23a}, M. Alhroob ¹²⁸, B. Ali ¹⁴²,
 G. Alimonti ^{68a}, J. Alison ³⁷, S.P. Alkire ¹⁴⁸, C. Allaire ¹³², B.M.M. Allbrooke ¹⁵⁶, B.W. Allen ¹³¹,
 P.P. Allport ²¹, A. Aloisio ^{69a,69b}, A. Alonso ⁴⁰, F. Alonso ⁸⁸, C. Alpigiani ¹⁴⁸, A.A. Alshehri ⁵⁷,
 M. Alvarez Estevez ⁹⁸, D. Álvarez Piqueras ¹⁷⁴, M.G. Alviggi ^{69a,69b}, Y. Amaral Coutinho ^{80b}, A. Ambler ¹⁰³,
 L. Ambroz ¹³⁵, C. Amelung ²⁶, D. Amidei ¹⁰⁵, S.P. Amor Dos Santos ^{140a}, S. Amoroso ⁴⁶, C.S. Amrouche ⁵⁴,
 F. An ⁷⁸, C. Anastopoulos ¹⁴⁹, N. Andari ¹⁴⁵, T. Andeen ¹¹, C.F. Anders ^{61b}, J.K. Anders ²⁰,
 A. Andreazza ^{68a,68b}, V. Andrei ^{61a}, C.R. Anelli ¹⁷⁶, S. Angelidakis ³⁸, A. Angerami ³⁹,
 A.V. Anisenkov ^{122b,122a}, A. Annovi ^{71a}, C. Antel ^{61a}, M.T. Anthony ¹⁴⁹, M. Antonelli ⁵¹, D.J.A. Antrim ¹⁷¹,
 F. Anulli ^{72a}, M. Aoki ⁸¹, J.A. Aparisi Pozo ¹⁷⁴, L. Aperio Bella ^{15a}, G. Arabidze ¹⁰⁶, J.P. Araque ^{140a},
 V. Araujo Ferraz ^{80b}, R. Araujo Pereira ^{80b}, C. Arcangeletti ⁵¹, A.T.H. Arce ⁴⁹, F.A. Arduh ⁸⁸, J-F. Arguin ¹⁰⁹,
 S. Argyropoulos ⁷⁷, J.-H. Arling ⁴⁶, A.J. Armbruster ³⁶, A. Armstrong ¹⁷¹, O. Arnaez ¹⁶⁷, H. Arnold ¹²⁰,
 Z.P. Arrubarrena Tame ¹¹⁴, A. Artamonov ^{111,*}, G. Artoni ¹³⁵, S. Artz ⁹⁹, S. Asai ¹⁶³, N. Asbah ⁵⁹,
 E.M. Asimakopoulou ¹⁷², L. Asquith ¹⁵⁶, J. Assahsah ^{35d}, K. Assamagan ²⁹, R. Astalos ^{28a}, R.J. Atkin ^{33a},
 M. Atkinson ¹⁷³, N.B. Atlay ¹⁹, H. Atmani ¹³², K. Augsten ¹⁴², G. Avolio ³⁶, R. Avramidou ^{60a}, M.K. Ayoub ^{15a},
 A.M. Azoulay ^{168b}, G. Azuelos ^{109,av}, H. Bachacou ¹⁴⁵, K. Bachas ^{67a,67b}, M. Backes ¹³⁵, F. Backman ^{45a,45b},
 P. Bagnaia ^{72a,72b}, M. Bahmani ⁸⁴, H. Bahrasemani ¹⁵², A.J. Bailey ¹⁷⁴, V.R. Bailey ¹⁷³, J.T. Baines ¹⁴⁴,
 M. Bajic ⁴⁰, C. Bakalis ¹⁰, O.K. Baker ¹⁸³, P.J. Bakker ¹²⁰, D. Bakshi Gupta ⁸, S. Balaji ¹⁵⁷,
 E.M. Baldin ^{122b,122a}, P. Balek ¹⁸⁰, F. Balli ¹⁴⁵, W.K. Balunas ¹³⁵, J. Balz ⁹⁹, E. Banas ⁸⁴, A. Bandyopadhyay ²⁴,
 Sw. Banerjee ^{181,i}, A.A.E. Bannoura ¹⁸², L. Barak ¹⁶¹, W.M. Barbe ³⁸, E.L. Barberio ¹⁰⁴, D. Barberis ^{55b,55a},
 M. Barbero ¹⁰¹, G. Barbour ⁹⁴, T. Barillari ¹¹⁵, M.-S. Barisits ³⁶, J. Barkeloo ¹³¹, T. Barklow ¹⁵³, R. Barnea ¹⁶⁰,
 S.L. Barnes ^{60c}, B.M. Barnett ¹⁴⁴, R.M. Barnett ¹⁸, Z. Barnovska-Blenessy ^{60a}, A. Baroncelli ^{60a}, G. Barone ²⁹,
 A.J. Barr ¹³⁵, L. Barranco Navarro ^{45a,45b}, F. Barreiro ⁹⁸, J. Barreiro Guimarães da Costa ^{15a}, S. Barsov ¹³⁸,
 R. Bartoldus ¹⁵³, G. Bartolini ¹⁰¹, A.E. Barton ⁸⁹, P. Bartos ^{28a}, A. Basalaeu ⁴⁶, A. Bassalat ^{132,ao},
 M.J. Basso ¹⁶⁷, R.L. Bates ⁵⁷, S. Batlamous ^{35e}, J.R. Batley ³², B. Batool ¹⁵¹, M. Battaglia ¹⁴⁶, M. Baucé ^{72a,72b},
 F. Bauer ¹⁴⁵, K.T. Bauer ¹⁷¹, H.S. Bawa ^{31,k}, J.B. Beacham ⁴⁹, T. Beau ¹³⁶, P.H. Beauchemin ¹⁷⁰, F. Becherer ⁵²,
 P. Bechtel ²⁴, H.C. Beck ⁵³, H.P. Beck ^{20,q}, K. Becker ⁵², M. Becker ⁹⁹, C. Becot ⁴⁶, A. Beddall ^{12d},
 A.J. Beddall ^{12a}, V.A. Bednyakov ⁷⁹, M. Bedognetti ¹²⁰, C.P. Bee ¹⁵⁵, T.A. Beermann ⁷⁶, M. Begalli ^{80b},
 M. Beger ²⁹, A. Behera ¹⁵⁵, J.K. Behr ⁴⁶, F. Beisiegel ²⁴, A.S. Bell ⁹⁴, G. Bella ¹⁶¹, L. Bellagamba ^{23b},
 A. Bellerive ³⁴, P. Bellos ⁹, K. Beloborodov ^{122b,122a}, K. Belotskiy ¹¹², N.L. Belyaev ¹¹², D. Benckekroun ^{35a},
 N. Benekos ¹⁰, Y. Benhammou ¹⁶¹, D.P. Benjamin ⁶, M. Benoit ⁵⁴, J.R. Bensinger ²⁶, S. Bentvelsen ¹²⁰,
 L. Beresford ¹³⁵, M. Beretta ⁵¹, D. Berge ⁴⁶, E. Bergeaas Kuutmann ¹⁷², N. Berger ⁵, B. Bergmann ¹⁴²,
 L.J. Bergsten ²⁶, J. Beringer ¹⁸, S. Berlendis ⁷, N.R. Bernard ¹⁰², G. Bernardi ¹³⁶, C. Bernius ¹⁵³, T. Berry ⁹³,
 P. Berta ⁹⁹, C. Bertella ^{15a}, I.A. Bertram ⁸⁹, O. Bessidskaia Bylund ¹⁸², N. Besson ¹⁴⁵, A. Bethani ¹⁰⁰,
 S. Bethke ¹¹⁵, A. Betti ²⁴, A.J. Bevan ⁹², J. Beyer ¹¹⁵, D.S. Bhattacharya ¹⁷⁷, R. Bi ¹³⁹, R.M. Bianchi ¹³⁹,
 O. Biebel ¹¹⁴, D. Biedermann ¹⁹, R. Bielski ³⁶, K. Bierwagen ⁹⁹, N.V. Biesuz ^{71a,71b}, M. Biglietti ^{74a},
 T.R.V. Billoud ¹⁰⁹, M. Bindi ⁵³, A. Bingul ^{12d}, C. Bini ^{72a,72b}, S. Biondi ^{23b,23a}, M. Birman ¹⁸⁰, T. Bisanz ⁵³,
 J.P. Biswal ¹⁶¹, D. Biswas ^{181,i}, A. Bitadze ¹⁰⁰, C. Bittrich ⁴⁸, K. Bjørke ¹³⁴, K.M. Black ²⁵, T. Blazek ^{28a},
 I. Bloch ⁴⁶, C. Blocker ²⁶, A. Blue ⁵⁷, U. Blumenschein ⁹², G.J. Bobbink ¹²⁰, V.S. Bobrovnikov ^{122b,122a},
 S.S. Bocchetta ⁹⁶, A. Bocci ⁴⁹, D. Boerner ⁴⁶, D. Bogavac ¹⁴, A.G. Bogdanchikov ^{122b,122a}, C. Bohm ^{45a},
 V. Boisvert ⁹³, P. Bokan ^{53,172}, T. Bold ^{83a}, A.S. Boldyrev ¹¹³, A.E. Bolz ^{61b}, M. Bomben ¹³⁶, M. Bona ⁹²,
 J.S. Bonilla ¹³¹, M. Boonekamp ¹⁴⁵, H.M. Borecka-Bielska ⁹⁰, A. Borisov ¹²³, G. Borissov ⁸⁹, J. Bortfeldt ³⁶,
 D. Bortoletto ¹³⁵, D. Boscherini ^{23b}, M. Bosman ¹⁴, J.D. Bossio Sola ¹⁰³, K. Bouaouda ^{35a}, J. Boudreau ¹³⁹,
 E.V. Bouhova-Thacker ⁸⁹, D. Boumediene ³⁸, S.K. Boutle ⁵⁷, A. Boveia ¹²⁶, J. Boyd ³⁶, D. Boye ^{33b,ap},
 I.R. Boyko ⁷⁹, A.J. Bozson ⁹³, J. Bracinik ²¹, N. Brahimi ¹⁰¹, G. Brandt ¹⁸², O. Brandt ³², F. Braren ⁴⁶,
 B. Brau ¹⁰², J.E. Brau ¹³¹, W.D. Breaden Madden ⁵⁷, K. Brendlinger ⁴⁶, L. Brenner ⁴⁶, R. Brenner ¹⁷²,
 S. Bressler ¹⁸⁰, B. Brickwedde ⁹⁹, D.L. Briglin ²¹, D. Britton ⁵⁷, D. Britzger ¹¹⁵, I. Brock ²⁴, R. Brock ¹⁰⁶,
 G. Brooijmans ³⁹, W.K. Brooks ^{147b}, E. Brost ¹²¹, J.H. Broughton ²¹, P.A. Bruckman de Renstrom ⁸⁴,
 D. Bruncko ^{28b}, A. Bruni ^{23b}, G. Bruni ^{23b}, L.S. Bruni ¹²⁰, S. Bruno ^{73a,73b}, B.H. Brunt ³², M. Bruschi ^{23b},

N. Bruscinò¹³⁹, P. Bryant³⁷, L. Bryngemark⁹⁶, T. Buanes¹⁷, Q. Buat³⁶, P. Buchholz¹⁵¹, A.G. Buckley⁵⁷, I.A. Budagov⁷⁹, M.K. Bugge¹³⁴, F. Bühner⁵², O. Bulekov¹¹², T.J. Burch¹²¹, S. Burdin⁹⁰, C.D. Burgard¹²⁰, A.M. Burger¹²⁹, B. Burghgrave⁸, K. Burka^{83a}, J.T.P. Burr⁴⁶, C.D. Burton¹¹, J.C. Burzynski¹⁰², V. Büscher⁹⁹, E. Buschmann⁵³, P.J. Bussey⁵⁷, J.M. Butler²⁵, C.M. Buttar⁵⁷, J.M. Butterworth⁹⁴, P. Butti³⁶, W. Buttinger³⁶, C.J. Buxo Vazquez¹⁰⁶, A. Buzatu¹⁵⁸, A.R. Buzykaev^{122b,122a}, G. Cabras^{23b,23a}, S. Cabrera Urbán¹⁷⁴, D. Caforio⁵⁶, H. Cai¹⁷³, V.M.M. Cairo¹⁵³, O. Cakir^{4a}, N. Calace³⁶, P. Calafiura¹⁸, A. Calandri¹⁰¹, G. Calderini¹³⁶, P. Calfayan⁶⁵, G. Callea⁵⁷, L.P. Caloba^{80b}, S. Calvente Lopez⁹⁸, D. Calvet³⁸, S. Calvet³⁸, T.P. Calvet¹⁵⁵, M. Calvetti^{71a,71b}, R. Camacho Toro¹³⁶, S. Camarda³⁶, D. Camarero Munoz⁹⁸, P. Camarri^{73a,73b}, D. Cameron¹³⁴, R. Caminal Armadans¹⁰², C. Camincher³⁶, S. Campana³⁶, M. Campanelli⁹⁴, A. Camplani⁴⁰, A. Campoverde¹⁵¹, V. Canale^{69a,69b}, A. Canesse¹⁰³, M. Cano Bret^{60c}, J. Cantero¹²⁹, T. Cao¹⁶¹, Y. Cao¹⁷³, M.D.M. Capeans Garrido³⁶, M. Capua^{41b,41a}, R. Cardarelli^{73a}, F. Cardillo¹⁴⁹, G. Carducci^{41b,41a}, I. Carli¹⁴³, T. Carli³⁶, G. Carlino^{69a}, B.T. Carlson¹³⁹, L. Carminati^{68a,68b}, R.M.D. Carney^{45a,45b}, S. Caron¹¹⁹, E. Carquin^{147b}, S. Carrá⁴⁶, J.W.S. Carter¹⁶⁷, M.P. Casado^{14,d}, A.F. Casha¹⁶⁷, D.W. Casper¹⁷¹, R. Castelijin¹²⁰, F.L. Castillo¹⁷⁴, V. Castillo Gimenez¹⁷⁴, N.F. Castro^{140a,140e}, A. Catinaccio³⁶, J.R. Catmore¹³⁴, A. Cattai³⁶, J. Caudron²⁴, V. Cavaliere²⁹, E. Cavallaro¹⁴, M. Cavalli-Sforza¹⁴, V. Cavasinni^{71a,71b}, E. Celebi^{12b}, F. Ceradini^{74a,74b}, L. Cerda Alberich¹⁷⁴, K. Cerny¹³⁰, A.S. Cerqueira^{80a}, A. Cerri¹⁵⁶, L. Cerrito^{73a,73b}, F. Cerutti¹⁸, A. Cervelli^{23b,23a}, S.A. Cetin^{12b}, Z. Chadi^{35a}, D. Chakraborty¹²¹, S.K. Chan⁵⁹, W.S. Chan¹²⁰, W.Y. Chan⁹⁰, J.D. Chapman³², B. Chargeishvili^{159b}, D.G. Charlton²¹, T.P. Charman⁹², C.C. Chau³⁴, S. Che¹²⁶, S. Chekanov⁶, S.V. Chekulaev^{168a}, G.A. Chelkov^{79,au}, M.A. Chelstowska³⁶, B. Chen⁷⁸, C. Chen^{60a}, C.H. Chen⁷⁸, H. Chen²⁹, J. Chen^{60a}, J. Chen³⁹, S. Chen¹³⁷, S.J. Chen^{15c}, X. Chen^{15b,at}, Y. Chen⁸², Y-H. Chen⁴⁶, H.C. Cheng^{63a}, H.J. Cheng^{15a,15d}, A. Cheplakov⁷⁹, E. Cheremushkina¹²³, R. Cherkaoui El Moursli^{35e}, E. Cheu⁷, K. Cheung⁶⁴, T.J.A. Chevaléris¹⁴⁵, L. Chevalier¹⁴⁵, V. Chiarella⁵¹, G. Chiarelli^{71a}, G. Chiodini^{67a}, A.S. Chisholm²¹, A. Chitan^{27b}, I. Chiu¹⁶³, Y.H. Chiu¹⁷⁶, M.V. Chizhov⁷⁹, K. Choi⁶⁵, A.R. Chomont^{72a,72b}, S. Chouridou¹⁶², Y.S. Chow¹²⁰, M.C. Chu^{63a}, X. Chu^{15a}, J. Chudoba¹⁴¹, A.J. Chuinard¹⁰³, J.J. Chwastowski⁸⁴, L. Chytka¹³⁰, D. Cieri¹¹⁵, K.M. Ciesla⁸⁴, D. Cinca⁴⁷, V. Cindro⁹¹, I.A. Cioară^{27b}, A. Ciocio¹⁸, F. Ciotto^{69a,69b}, Z.H. Citron¹⁸⁰, M. Citterio^{68a}, D.A. Ciubotaru^{27b}, B.M. Ciungu¹⁶⁷, A. Clark⁵⁴, M.R. Clark³⁹, P.J. Clark⁵⁰, C. Clement^{45a,45b}, Y. Coadou¹⁰¹, M. Cobl^{66a,66c}, A. Coccaro^{55b}, J. Cochran⁷⁸, H. Cohen¹⁶¹, A.E.C. Coimbra³⁶, L. Colasurdo¹¹⁹, B. Cole³⁹, A.P. Colijn¹²⁰, J. Collot⁵⁸, P. Conde Muiño^{140a,e}, E. Coniavitis⁵², S.H. Connell^{33b}, I.A. Connelly⁵⁷, S. Constantinescu^{27b}, F. Conventi^{69a,aw}, A.M. Cooper-Sarkar¹³⁵, F. Cormier¹⁷⁵, K.J.R. Cormier¹⁶⁷, L.D. Corpe⁹⁴, M. Corradi^{72a,72b}, E.E. Corrigan⁹⁶, F. Corriveau^{103,ac}, A. Cortes-Gonzalez³⁶, M.J. Costa¹⁷⁴, F. Costanza⁵, D. Costanzo¹⁴⁹, G. Cowan⁹³, J.W. Cowley³², J. Crane¹⁰⁰, K. Cranmer¹²⁴, S.J. Crawley⁵⁷, R.A. Creager¹³⁷, S. Crépe-Renaudin⁵⁸, F. Crescioli¹³⁶, M. Cristinziani²⁴, V. Croft¹²⁰, G. Crosetti^{41b,41a}, A. Cueto⁵, T. Cuhadar Donszelmann¹⁴⁹, A.R. Cukierman¹⁵³, S. Czekierda⁸⁴, P. Czodrowski³⁶, M.J. Da Cunha Sargedas De Sousa^{60b}, J.V. Da Fonseca Pinto^{80b}, C. Da Via¹⁰⁰, W. Dabrowski^{83a}, T. Dado^{28a}, S. Dahbi^{35e}, T. Dai¹⁰⁵, C. Dallapiccola¹⁰², M. Dam⁴⁰, G. D'amen²⁹, V. D'Amico^{74a,74b}, J. Damp⁹⁹, J.R. Dandoy¹³⁷, M.F. Daneri³⁰, N.P. Dang¹⁸¹, N.D. Dann¹⁰⁰, M. Danninger¹⁷⁵, V. Dao³⁶, G. Darbo^{55b}, O. Dartsis⁵, A. Dattagupta¹³¹, T. Daubney⁴⁶, S. D'Auria^{68a,68b}, W. Davey²⁴, C. David⁴⁶, T. Davidek¹⁴³, D.R. Davis⁴⁹, I. Dawson¹⁴⁹, K. De⁸, R. De Asmundis^{69a}, M. De Beurs¹²⁰, S. De Castro^{23b,23a}, S. De Cecco^{72a,72b}, N. De Groot¹¹⁹, P. de Jong¹²⁰, H. De la Torre¹⁰⁶, A. De Maria^{15c}, D. De Pedis^{72a}, A. De Salvo^{72a}, U. De Sanctis^{73a,73b}, M. De Santis^{73a,73b}, A. De Santo¹⁵⁶, K. De Vasconcelos Corga¹⁰¹, J.B. De Vivie De Regie¹³², C. Debenedetti¹⁴⁶, D.V. Dedovich⁷⁹, A.M. Deiana⁴², M. Del Gaudio^{41b,41a}, J. Del Peso⁹⁸, Y. Delabat Diaz⁴⁶, D. Delgove¹³², F. Deliot^{145,p}, C.M. Delitzsch⁷, M. Della Pietra^{69a,69b}, D. Della Volpe⁵⁴, A. Dell'Acqua³⁶, L. Dell'Asta^{73a,73b}, M. Delmastro⁵, C. Delporte¹³², P.A. Delsart⁵⁸, D.A. DeMarco¹⁶⁷, S. Demers¹⁸³, M. Demichev⁷⁹, G. Demontigny¹⁰⁹, S.P. Denisov¹²³, D. Denysiuk¹²⁰, L. D'Eramo¹³⁶, D. Derendarz⁸⁴, J.E. Derkaoui^{35d}, F. Derue¹³⁶, P. Dervan⁹⁰, K. Desch²⁴, C. Deterre⁴⁶, K. Dette¹⁶⁷, C. Deutsch²⁴, M.R. Devesa³⁰, P.O. Deviveiros³⁶, A. Dewhurst¹⁴⁴, F.A. Di Bello⁵⁴, A. Di Ciaccio^{73a,73b}, L. Di Ciaccio⁵, W.K. Di Clemente¹³⁷, C. Di Donato^{69a,69b}, A. Di Girolamo³⁶, G. Di Gregorio^{71a,71b}, B. Di Micco^{74a,74b}, R. Di Nardo¹⁰², K.F. Di Petrillo⁵⁹, R. Di Sipio¹⁶⁷, D. Di Valentino³⁴, C. Diaconu¹⁰¹, F.A. Dias⁴⁰, T. Dias Do Vale^{140a}, M.A. Diaz^{147a}, J. Dickinson¹⁸, E.B. Diehl¹⁰⁵, J. Dietrich¹⁹, S. Díez Cornell⁴⁶,

A. Dimitrievska¹⁸, W. Ding^{15b}, J. Dingfelder²⁴, F. Dittus³⁶, F. Djama¹⁰¹, T. Djobava^{159b}, J.I. Djuvsland¹⁷,
 M.A.B. Do Vale^{80c}, M. Dobre^{27b}, D. Dodsworth²⁶, C. Doglioni⁹⁶, J. Dolejsi¹⁴³, Z. Dolezal¹⁴³,
 M. Donadelli^{80d}, B. Dong^{60c}, J. Donini³⁸, A. D'onofrio⁹², M. D'Onofrio⁹⁰, J. Dopke¹⁴⁴, A. Doria^{69a},
 M.T. Dova⁸⁸, A.T. Doyle⁵⁷, E. Drechsler¹⁵², E. Dreyer¹⁵², T. Dreyer⁵³, A.S. Drobac¹⁷⁰, D. Du^{60b},
 Y. Duan^{60b}, F. Dubinin¹¹⁰, M. Dubovsky^{28a}, A. Dubreuil⁵⁴, E. Duchovni¹⁸⁰, G. Duckeck¹¹⁴,
 A. Ducourthial¹³⁶, O.A. Ducu¹⁰⁹, D. Duda¹¹⁵, A. Dudarev³⁶, A.C. Dudder⁹⁹, E.M. Duffield¹⁸, L. Dufлот¹³²,
 M. Dührssen³⁶, C. Dülsen¹⁸², M. Dumancic¹⁸⁰, A.E. Dumitriu^{27b}, A.K. Duncan⁵⁷, M. Dunford^{61a},
 A. Duperrin¹⁰¹, H. Duran Yildiz^{4a}, M. Düren⁵⁶, A. Durglishvili^{159b}, D. Duschinger⁴⁸, B. Dutta⁴⁶,
 D. Duvnjak¹, G.I. Dyckes¹³⁷, M. Dyndal³⁶, S. Dysch¹⁰⁰, B.S. Dziejczak⁸⁴, K.M. Ecker¹¹⁵, R.C. Edgar¹⁰⁵,
 M.G. Eggleston⁴⁹, T. Eifert³⁶, G. Eigen¹⁷, K. Einsweiler¹⁸, T. Ekelof¹⁷², H. El Jarrari^{35e}, M. El Kacimi^{35c},
 R. El Kosseifi¹⁰¹, V. Ellajosyula¹⁷², M. Ellert¹⁷², F. Ellinghaus¹⁸², A.A. Elliot⁹², N. Ellis³⁶,
 J. Elmsheuser²⁹, M. Elsing³⁶, D. Emelianov¹⁴⁴, A. Emerman³⁹, Y. Enari¹⁶³, M.B. Epland⁴⁹,
 J. Erdmann⁴⁷, A. Ereditato²⁰, M. Errenst³⁶, M. Escalier¹³², C. Escobar¹⁷⁴, O. Estrada Pastor¹⁷⁴,
 E. Etzion¹⁶¹, H. Evans⁶⁵, A. Ezhilov¹³⁸, F. Fabbri⁵⁷, L. Fabbri^{23b,23a}, V. Fabiani¹¹⁹, G. Facini⁹⁴,
 R.M. Faisca Rodrigues Pereira^{140a}, R.M. Fakhruddinov¹²³, S. Falciano^{72a}, P.J. Falke⁵, S. Falke⁵,
 J. Faltova¹⁴³, Y. Fang^{15a}, Y. Fang^{15a}, G. Fanourakis⁴⁴, M. Fanti^{68a,68b}, M. Faraj^{66a,66c}, A. Farbin⁸,
 A. Farilla^{74a}, E.M. Farina^{70a,70b}, T. Farooque¹⁰⁶, S. Farrell¹⁸, S.M. Farrington⁵⁰, P. Farthouat³⁶, F. Fassi^{35e},
 P. Fassnacht³⁶, D. Fassouliotis⁹, M. Fauci Giannelli⁵⁰, W.J. Fawcett³², L. Fayard¹³², O.L. Fedin^{138,n},
 W. Fedorko¹⁷⁵, M. Feickert⁴², L. Feligioni¹⁰¹, A. Fell¹⁴⁹, C. Feng^{60b}, E.J. Feng³⁶, M. Feng⁴⁹,
 M.J. Fenton⁵⁷, A.B. Fenyuk¹²³, J. Ferrando⁴⁶, A. Ferrante¹⁷³, A. Ferrari¹⁷², P. Ferrari¹²⁰, R. Ferrari^{70a},
 D.E. Ferreira de Lima^{61b}, A. Ferrer¹⁷⁴, D. Ferrere⁵⁴, C. Ferretti¹⁰⁵, F. Fiedler⁹⁹, A. Filipčič⁹¹,
 F. Filthaut¹¹⁹, K.D. Finelli²⁵, M.C.N. Fiolhais^{140a}, L. Fiorini¹⁷⁴, F. Fischer¹¹⁴, W.C. Fisher¹⁰⁶, I. Fleck¹⁵¹,
 P. Fleischmann¹⁰⁵, R.R.M. Fletcher¹³⁷, T. Flick¹⁸², B.M. Flierl¹¹⁴, L.F. Flores¹³⁷, L.R. Flores Castillo^{63a},
 F.M. Follega^{75a,75b}, N. Fomin¹⁷, J.H. Foo¹⁶⁷, G.T. Forcolin^{75a,75b}, A. Formica¹⁴⁵, F.A. Förster¹⁴,
 A.C. Forti¹⁰⁰, A.G. Foster²¹, M.G. Foti¹³⁵, D. Fournier¹³², H. Fox⁸⁹, P. Francavilla^{71a,71b},
 S. Francescato^{72a,72b}, M. Franchini^{23b,23a}, S. Franchino^{61a}, D. Francis³⁶, L. Franconi²⁰, M. Franklin⁵⁹,
 A.N. Fray⁹², P.M. Freeman²¹, B. Freund¹⁰⁹, W.S. Freund^{80b}, E.M. Freundlich⁴⁷, D.C. Frizzell¹²⁸,
 D. Froidevaux³⁶, J.A. Frost¹³⁵, C. Fukunaga¹⁶⁴, E. Fullana Torregrosa¹⁷⁴, E. Fumagalli^{55b,55a},
 T. Fusayasu¹¹⁶, J. Fuster¹⁷⁴, A. Gabrielli^{23b,23a}, A. Gabrielli¹⁸, G.P. Gach^{83a}, S. Gadatsch⁵⁴, P. Gadow¹¹⁵,
 G. Gagliardi^{55b,55a}, L.G. Gagnon¹⁰⁹, C. Galea^{27b}, B. Galhardo^{140a}, G.E. Gallardo¹³⁵, E.J. Gallas¹³⁵,
 B.J. Gallop¹⁴⁴, G. Galster⁴⁰, R. Gamboa Goni⁹², K.K. Gan¹²⁶, S. Ganguly¹⁸⁰, J. Gao^{60a}, Y. Gao⁵⁰,
 Y.S. Gao^{31,k}, C. García¹⁷⁴, J.E. García Navarro¹⁷⁴, J.A. García Pascual^{15a}, C. Garcia-Argos⁵²,
 M. Garcia-Sciveres¹⁸, R.W. Gardner³⁷, N. Garelli¹⁵³, S. Gargiulo⁵², V. Garonne¹³⁴, A. Gaudiello^{55b,55a},
 G. Gaudio^{70a}, I.L. Gavrilenko¹¹⁰, A. Gavrilyuk¹¹¹, C. Gay¹⁷⁵, G. Gaycken⁴⁶, E.N. Gazis¹⁰, A.A. Geanta^{27b},
 C.M. Gee¹⁴⁶, C.N.P. Gee¹⁴⁴, J. Geisen⁵³, M. Geisen⁹⁹, M.P. Geisler^{61a}, C. Gemme^{55b}, M.H. Genest⁵⁸,
 C. Geng¹⁰⁵, S. Gentile^{72a,72b}, S. George⁹³, T. Gerialis⁴⁴, L.O. Gerlach⁵³, P. Gessinger-Befurt⁹⁹,
 G. Gessner⁴⁷, S. Ghasemi¹⁵¹, M. Ghasemi Bostanabad¹⁷⁶, M. Ghneimat²⁴, A. Ghosh¹³², A. Ghosh⁷⁷,
 B. Giacobbe^{23b}, S. Giagu^{72a,72b}, N. Giangiacomi^{23b,23a}, P. Giannetti^{71a}, A. Giannini^{69a,69b}, G. Giannini¹⁴,
 S.M. Gibson⁹³, M. Gignac¹⁴⁶, D. Gillberg³⁴, G. Gilles¹⁸², D.M. Gingrich^{3,av}, M.P. Giordani^{66a,66c},
 F.M. Giorgi^{23b}, P.F. Giraud¹⁴⁵, G. Giugliarelli^{66a,66c}, D. Giugni^{68a}, F. Giuli^{73a,73b}, S. Gkaitatzis¹⁶²,
 I. Gkialas^{9,g}, E.L. Gkougkousis¹⁴, P. Gkoutoumis¹⁰, L.K. Gladilin¹¹³, C. Glasman⁹⁸, J. Glatzer¹⁴,
 P.C.F. Glaysher⁴⁶, A. Glazov⁴⁶, G.R. Gledhill¹³¹, M. Goblirsch-Kolb²⁶, D. Godin¹⁰⁹, S. Goldfarb¹⁰⁴,
 T. Golling⁵⁴, D. Golubkov¹²³, A. Gomes^{140a,140b}, R. Goncalves Gama⁵³, R. Gonçalves^{140a,140b}, G. Gonella⁵²,
 L. Gonella²¹, A. Gongadze⁷⁹, F. Gonnella²¹, J.L. Gonski⁵⁹, S. González de la Hoz¹⁷⁴,
 S. Gonzalez-Sevilla⁵⁴, G.R. Gonzalvo Rodriguez¹⁷⁴, L. Goossens³⁶, P.A. Gorbounov¹¹¹, H.A. Gordon²⁹,
 B. Gorini³⁶, E. Gorini^{67a,67b}, A. Gorišek⁹¹, A.T. Goshaw⁴⁹, M.I. Gostkin⁷⁹, C.A. Gottardo¹¹⁹,
 M. Gouighri^{35b}, D. Goujdami^{35c}, A.G. Goussiou¹⁴⁸, N. Govender^{33b}, C. Goy⁵, E. Gozani¹⁶⁰,
 I. Grabowska-Bold^{83a}, E.C. Graham⁹⁰, J. Gramling¹⁷¹, E. Gramstad¹³⁴, S. Grancagnolo¹⁹, M. Grandi¹⁵⁶,
 V. Gratchev¹³⁸, P.M. Gravila^{27f}, F.G. Gravili^{67a,67b}, C. Gray⁵⁷, H.M. Gray¹⁸, C. Greife²⁴, K. Gregersen⁹⁶,
 I.M. Gregor⁴⁶, P. Grenier¹⁵³, K. Grevtsov⁴⁶, C. Grieco¹⁴, N.A. Grieser¹²⁸, J. Griffiths⁸, A.A. Grillo¹⁴⁶,
 K. Grimm^{31,j}, S. Grinstein^{14,x}, J.-F. Grivaz¹³², S. Groh⁹⁹, E. Gross¹⁸⁰, J. Grosse-Knetter⁵³, Z.J. Grout⁹⁴,
 C. Grud¹⁰⁵, A. Grummer¹¹⁸, L. Guan¹⁰⁵, W. Guan¹⁸¹, J. Guenther³⁶, A. Guerguichon¹³²,

J.G.R. Guerrero Rojas ¹⁷⁴, F. Guescini ¹¹⁵, D. Guest ¹⁷¹, R. Gugel ⁵², T. Guillemin ⁵, S. Guindon ³⁶, U. Gul ⁵⁷, J. Guo ^{60c}, W. Guo ¹⁰⁵, Y. Guo ^{60a,r}, Z. Guo ¹⁰¹, R. Gupta ⁴⁶, S. Gurbuz ^{12c}, G. Gustavino ¹²⁸, M. Guth ⁵², P. Gutierrez ¹²⁸, C. Gutsche ⁹⁴, C. Guyot ¹⁴⁵, C. Gwenlan ¹³⁵, C.B. Gwilliam ⁹⁰, A. Haas ¹²⁴, C. Haber ¹⁸, H.K. Hadavand ⁸, N. Haddad ^{35e}, A. Hadeef ^{60a}, S. Hageböck ³⁶, M. Haleem ¹⁷⁷, J. Haley ¹²⁹, G. Halladjian ¹⁰⁶, G.D. Hallowell ¹⁰¹, K. Hamacher ¹⁸², P. Hamal ¹³⁰, K. Hamano ¹⁷⁶, H. Hamdaoui ^{35e}, G.N. Hamity ¹⁴⁹, K. Han ^{60a,ai}, L. Han ^{60a}, S. Han ^{15a,15d}, Y.F. Han ¹⁶⁷, K. Hanagaki ^{81,v}, M. Hance ¹⁴⁶, D.M. Handl ¹¹⁴, B. Haney ¹³⁷, R. Hankache ¹³⁶, E. Hansen ⁹⁶, J.B. Hansen ⁴⁰, J.D. Hansen ⁴⁰, M.C. Hansen ²⁴, P.H. Hansen ⁴⁰, E.C. Hanson ¹⁰⁰, K. Hara ¹⁶⁹, T. Harenberg ¹⁸², S. Harkusha ¹⁰⁷, P.F. Harrison ¹⁷⁸, N.M. Hartmann ¹¹⁴, Y. Hasegawa ¹⁵⁰, A. Hasib ⁵⁰, S. Hassani ¹⁴⁵, S. Haug ²⁰, R. Hauser ¹⁰⁶, L.B. Havener ³⁹, M. Havranek ¹⁴², C.M. Hawkes ²¹, R.J. Hawkins ³⁶, D. Hayden ¹⁰⁶, C. Hayes ¹⁵⁵, R.L. Hayes ¹⁷⁵, C.P. Hays ¹³⁵, J.M. Hays ⁹², H.S. Hayward ⁹⁰, S.J. Haywood ¹⁴⁴, F. He ^{60a}, M.P. Heath ⁵⁰, V. Hedberg ⁹⁶, L. Heelan ⁸, S. Heer ²⁴, K.K. Heidegger ⁵², W.D. Heidorn ⁷⁸, J. Heilman ³⁴, S. Heim ⁴⁶, T. Heim ¹⁸, B. Heinemann ^{46,aq}, J.J. Heinrich ¹³¹, L. Heinrich ³⁶, C. Heinz ⁵⁶, J. Hejbal ¹⁴¹, L. Helary ^{61b}, A. Held ¹⁷⁵, S. Hellesund ¹³⁴, C.M. Helling ¹⁴⁶, S. Hellman ^{45a,45b}, C. Helsens ³⁶, R.C.W. Henderson ⁸⁹, Y. Heng ¹⁸¹, S. Henkelmann ¹⁷⁵, A.M. Henriques Correia ³⁶, G.H. Herbert ¹⁹, H. Herde ²⁶, V. Herget ¹⁷⁷, Y. Hernández Jiménez ^{33c}, H. Herr ⁹⁹, M.G. Herrmann ¹¹⁴, T. Herrmann ⁴⁸, G. Herten ⁵², R. Hertenberger ¹¹⁴, L. Hervas ³⁶, T.C. Herwig ¹³⁷, G.G. Hesketh ⁹⁴, N.P. Hessey ^{168a}, A. Higashida ¹⁶³, S. Higashino ⁸¹, E. Higón-Rodríguez ¹⁷⁴, K. Hildebrand ³⁷, E. Hill ¹⁷⁶, J.C. Hill ³², K.K. Hill ²⁹, K.H. Hiller ⁴⁶, S.J. Hillier ²¹, M. Hils ⁴⁸, I. Hinchliffe ¹⁸, F. Hinterkeuser ²⁴, M. Hirose ¹³³, S. Hirose ⁵², D. Hirschbuehl ¹⁸², B. Hiti ⁹¹, O. Hladik ¹⁴¹, D.R. Hlaluku ^{33c}, X. Hoad ⁵⁰, J. Hobbs ¹⁵⁵, N. Hod ¹⁸⁰, M.C. Hodgkinson ¹⁴⁹, A. Hoecker ³⁶, F. Hoenig ¹¹⁴, D. Hohn ⁵², D. Hohov ¹³², T.R. Holmes ³⁷, M. Holzbock ¹¹⁴, L.B.A.H. Hommels ³², S. Honda ¹⁶⁹, T.M. Hong ¹³⁹, A. Hönle ¹¹⁵, B.H. Hooberman ¹⁷³, W.H. Hopkins ⁶, Y. Horii ¹¹⁷, P. Horn ⁴⁸, L.A. Horyn ³⁷, S. Hou ¹⁵⁸, A. Hoummada ^{35a}, J. Howarth ¹⁰⁰, J. Hoya ⁸⁸, M. Hrabovsky ¹³⁰, J. Hrdinka ⁷⁶, I. Hristova ¹⁹, J. Hrivnac ¹³², A. Hrynevich ¹⁰⁸, T. Hryn'ova ⁵, P.J. Hsu ⁶⁴, S.-C. Hsu ¹⁴⁸, Q. Hu ²⁹, S. Hu ^{60c}, Y.F. Hu ^{15a}, D.P. Huang ⁹⁴, Y. Huang ^{60a}, Y. Huang ^{15a}, Z. Hubacek ¹⁴², F. Hubaut ¹⁰¹, M. Huebner ²⁴, F. Huegging ²⁴, T.B. Huffman ¹³⁵, M. Huhtinen ³⁶, R.F.H. Hunter ³⁴, P. Huo ¹⁵⁵, A.M. Hupe ³⁴, N. Huseynov ^{79,ae}, J. Huston ¹⁰⁶, J. Huth ⁵⁹, R. Hyneman ¹⁰⁵, S. Hyrych ^{28a}, G. Iacobucci ⁵⁴, G. Iakovidis ²⁹, I. Ibragimov ¹⁵¹, L. Iconomidou-Fayard ¹³², Z. Idrissi ^{35e}, P.I. Iengo ³⁶, R. Ignazzi ⁴⁰, O. Igonkina ^{120,z,*}, R. Iguchi ¹⁶³, T. Iizawa ⁵⁴, Y. Ikegami ⁸¹, M. Ikeno ⁸¹, D. Iliadis ¹⁶², N. Ilic ^{119,s}, F. Iltzsche ⁴⁸, G. Introzzi ^{70a,70b}, M. Iodice ^{74a}, K. Iordanidou ^{168a}, V. Ippolito ^{72a,72b}, M.F. Isacson ¹⁷², M. Ishino ¹⁶³, W. Islam ¹²⁹, C. Issever ¹³⁵, S. Istin ¹⁶⁰, F. Ito ¹⁶⁹, J.M. Iturbe Ponce ^{63a}, R. Iuppa ^{75a,75b}, A. Ivina ¹⁸⁰, H. Iwasaki ⁸¹, J.M. Izen ⁴³, V. Izzo ^{69a}, P. Jacka ¹⁴¹, P. Jackson ¹, R.M. Jacobs ²⁴, B.P. Jaeger ¹⁵², V. Jain ², G. Jäkel ¹⁸², K.B. Jakobi ⁹⁹, K. Jakobs ⁵², S. Jakobsen ⁷⁶, T. Jakoubek ¹⁴¹, J. Jamieson ⁵⁷, K.W. Janas ^{83a}, R. Jansky ⁵⁴, J. Janssen ²⁴, M. Janus ⁵³, P.A. Janus ^{83a}, G. Jarlskog ⁹⁶, N. Javadov ^{79,ae}, T. Javůrek ³⁶, M. Javurkova ⁵², F. Jeanneau ¹⁴⁵, L. Jeanty ¹³¹, J. Jejelava ^{159a,af}, A. Jelinskas ¹⁷⁸, P. Jenni ^{52,a}, J. Jeong ⁴⁶, N. Jeong ⁴⁶, S. Jézéquel ⁵, H. Ji ¹⁸¹, J. Jia ¹⁵⁵, H. Jiang ⁷⁸, Y. Jiang ^{60a}, Z. Jiang ^{153,o}, S. Jiggins ⁵², F.A. Jimenez Morales ³⁸, J. Jimenez Pena ¹¹⁵, S. Jin ^{15c}, A. Jinaru ^{27b}, O. Jinnouchi ¹⁶⁵, H. Jivan ^{33c}, P. Johansson ¹⁴⁹, K.A. Johns ⁷, C.A. Johnson ⁶⁵, K. Jon-And ^{45a,45b}, R.W.L. Jones ⁸⁹, S.D. Jones ¹⁵⁶, S. Jones ⁷, T.J. Jones ⁹⁰, J. Jongmanns ^{61a}, P.M. Jorge ^{140a}, J. Jovicevic ³⁶, X. Ju ¹⁸, J.J. Junggeburth ¹¹⁵, A. Juste Rozas ^{14,x}, A. Kaczmarska ⁸⁴, M. Kado ^{72a,72b}, H. Kagan ¹²⁶, M. Kagan ¹⁵³, C. Kahra ⁹⁹, T. Kaji ¹⁷⁹, E. Kajomovitz ¹⁶⁰, C.W. Kalderon ⁹⁶, A. Kaluza ⁹⁹, A. Kamenshchikov ¹²³, M. Kaneda ¹⁶³, L. Kanjir ⁹¹, Y. Kano ¹⁶³, V.A. Kantserov ¹¹², J. Kanzaki ⁸¹, L.S. Kaplan ¹⁸¹, D. Kar ^{33c}, K. Karava ¹³⁵, M.J. Kareem ^{168b}, S.N. Karpov ⁷⁹, Z.M. Karpova ⁷⁹, V. Kartvelishvili ⁸⁹, A.N. Karyukhin ¹²³, L. Kashif ¹⁸¹, R.D. Kass ¹²⁶, A. Kastanas ^{45a,45b}, C. Kato ^{60d,60c}, J. Katzy ⁴⁶, K. Kawade ¹⁵⁰, K. Kawagoe ⁸⁷, T. Kawaguchi ¹¹⁷, T. Kawamoto ¹⁶³, G. Kawamura ⁵³, E.F. Kay ¹⁷⁶, V.F. Kazanin ^{122b,122a}, R. Keeler ¹⁷⁶, R. Kehoe ⁴², J.S. Keller ³⁴, E. Kellermann ⁹⁶, D. Kelsey ¹⁵⁶, J.J. Kempster ²¹, J. Kendrick ²¹, O. Kepka ¹⁴¹, S. Kersten ¹⁸², B.P. Kerševan ⁹¹, S. Ketabchi Haghighat ¹⁶⁷, M. Khader ¹⁷³, F. Khalil-Zada ¹³, M. Khandoga ¹⁴⁵, A. Khanov ¹²⁹, A.G. Kharlamov ^{122b,122a}, T. Kharlamova ^{122b,122a}, E.E. Khoda ¹⁷⁵, A. Khodinov ¹⁶⁶, T.J. Khoo ⁵⁴, E. Khramov ⁷⁹, J. Khubua ^{159b}, S. Kido ⁸², M. Kiehn ⁵⁴, C.R. Kilby ⁹³, Y.K. Kim ³⁷, N. Kimura ⁹⁴, O.M. Kind ¹⁹, B.T. King ^{90,*}, D. Kirchmeier ⁴⁸, J. Kirk ¹⁴⁴, A.E. Kiryunin ¹¹⁵, T. Kishimoto ¹⁶³, D.P. Kisliuk ¹⁶⁷, V. Kitali ⁴⁶, O. Kivernyk ⁵, T. Klapdor-Kleingrothaus ⁵², M. Klassen ^{61a}, M.H. Klein ¹⁰⁵, M. Klein ⁹⁰, U. Klein ⁹⁰, K. Kleinknecht ⁹⁹, P. Klimek ¹²¹, A. Klimentov ²⁹, T. Klingl ²⁴, T. Klioutchnikova ³⁶, F.F. Klitzner ¹¹⁴, P. Kluit ¹²⁰, S. Kluth ¹¹⁵,

E. Kneringer⁷⁶, E.B.F.G. Knoops¹⁰¹, A. Knue⁵², D. Kobayashi⁸⁷, T. Kobayashi¹⁶³, M. Kobel⁴⁸, M. Kocian¹⁵³, P. Kodys¹⁴³, P.T. Koenig²⁴, T. Koffas³⁴, N.M. Köhler³⁶, T. Koi¹⁵³, M. Kolb^{61b}, I. Koletsou⁵, T. Komarek¹³⁰, T. Kondo⁸¹, N. Kondrashova^{60c}, K. Köneke⁵², A.C. König¹¹⁹, T. Kono¹²⁵, R. Konoplich^{124,al}, V. Konstantinides⁹⁴, N. Konstantinidis⁹⁴, B. Konya⁹⁶, R. Kopeliansky⁶⁵, S. Koperny^{83a}, K. Korcyl⁸⁴, K. Kordas¹⁶², G. Koren¹⁶¹, A. Korn⁹⁴, I. Korolkov¹⁴, E.V. Korolkova¹⁴⁹, N. Korotkova¹¹³, O. Kortner¹¹⁵, S. Kortner¹¹⁵, T. Kosek¹⁴³, V.V. Kostyukhin¹⁶⁶, A. Kotsokechagia¹³², A. Kotwal⁴⁹, A. Koulouris¹⁰, A. Kourkoumeli-Charalampidi^{70a,70b}, C. Kourkoumelis⁹, E. Kourlitis¹⁴⁹, V. Kouskoura²⁹, A.B. Kowalewska⁸⁴, R. Kowalewski¹⁷⁶, C. Kozakai¹⁶³, W. Kozanecki¹⁴⁵, A.S. Kozhin¹²³, V.A. Kramarenko¹¹³, G. Kramberger⁹¹, D. Krasnopevtsev^{60a}, M.W. Krasny¹³⁶, A. Krasznahorkay³⁶, D. Krauss¹¹⁵, J.A. Kremer^{83a}, J. Kretzschmar⁹⁰, P. Krieger¹⁶⁷, F. Krieter¹¹⁴, A. Krishnan^{61b}, K. Krizka¹⁸, K. Kroeninger⁴⁷, H. Kroha¹¹⁵, J. Kroll¹⁴¹, J. Kroll¹³⁷, K.S. Krowpman¹⁰⁶, J. Krstic¹⁶, U. Kruchonak⁷⁹, H. Krüger²⁴, N. Krumnack⁷⁸, M.C. Kruse⁴⁹, J.A. Krzysiak⁸⁴, T. Kubota¹⁰⁴, O. Kuchinskaia¹⁶⁶, S. Kuday^{4b}, J.T. Kuechler⁴⁶, S. Kuehn³⁶, A. Kugel^{61a}, T. Kuhl⁴⁶, V. Kukhtin⁷⁹, R. Kukla¹⁰¹, Y. Kulchitsky^{107,ah}, S. Kuleshov^{147b}, Y.P. Kulinich¹⁷³, M. Kuna⁵⁸, T. Kunigo⁸⁵, A. Kupco¹⁴¹, T. Kupfer⁴⁷, O. Kuprash⁵², H. Kurashige⁸², L.L. Kurchaninov^{168a}, Y.A. Kurochkin¹⁰⁷, A. Kurova¹¹², M.G. Kurth^{15a,15d}, E.S. Kuwertz³⁶, M. Kuze¹⁶⁵, A.K. Kvam¹⁴⁸, J. Kvita¹³⁰, T. Kwan¹⁰³, A. La Rosa¹¹⁵, L. La Rotonda^{41b,41a}, F. La Ruffa^{41b,41a}, C. Lacasta¹⁷⁴, F. Lacava^{72a,72b}, D.P.J. Lack¹⁰⁰, H. Lacker¹⁹, D. Lacour¹³⁶, E. Ladygin⁷⁹, R. Lafaye⁵, B. Laforge¹³⁶, T. Lagouri^{33c}, S. Lai⁵³, S. Lammers⁶⁵, W. Lampl⁷, C. Lampoudis¹⁶², E. Lançon²⁹, U. Landgraf⁵², M.P.J. Landon⁹², M.C. Lanfermann⁵⁴, V.S. Lang⁴⁶, J.C. Lange⁵³, R.J. Langenberg³⁶, A.J. Lankford¹⁷¹, F. Lanni²⁹, K. Lantzsck²⁴, A. Lanza^{70a}, A. Lapertosa^{55b,55a}, S. Laplace¹³⁶, J.F. Laporte¹⁴⁵, T. Lari^{68a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁶, T.S. Lau^{63a}, A. Laudrain¹³², A. Laurier³⁴, M. Lavorgna^{69a,69b}, S.D. Lawlor⁹³, M. Lazzaroni^{68a,68b}, B. Le¹⁰⁴, E. Le Guirriec¹⁰¹, M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁸, A.C.A. Lee⁹⁴, C.A. Lee²⁹, G.R. Lee¹⁷, L. Lee⁵⁹, S.C. Lee¹⁵⁸, S.J. Lee³⁴, S. Lee⁷⁸, B. Lefebvre^{168a}, M. Lefebvre¹⁷⁶, F. Legger¹¹⁴, C. Leggett¹⁸, K. Lehmann¹⁵², N. Lehmann¹⁸², G. Lehmann Miotto³⁶, W.A. Leight⁴⁶, A. Leisos^{162,w}, M.A.L. Leite^{80d}, C.E. Leitgeb¹¹⁴, R. Leitner¹⁴³, D. Lellouch^{180,*}, K.J.C. Leney⁴², T. Lenz²⁴, B. Lenzi³⁶, R. Leone⁷, S. Leone^{71a}, C. Leonidopoulos⁵⁰, A. Leopold¹³⁶, G. Lerner¹⁵⁶, C. Leroy¹⁰⁹, R. Les¹⁶⁷, C.G. Lester³², M. Levchenko¹³⁸, J. Levêque⁵, D. Levin¹⁰⁵, L.J. Levinson¹⁸⁰, D.J. Lewis²¹, B. Li^{15b}, B. Li¹⁰⁵, C-Q. Li^{60a}, F. Li^{60c}, H. Li^{60a}, H. Li^{60b}, J. Li^{60c}, K. Li¹⁵³, L. Li^{60c}, M. Li^{15a}, Q. Li^{15a,15d}, Q.Y. Li^{60a}, S. Li^{60d,60c}, X. Li⁴⁶, Y. Li⁴⁶, Z. Li^{60b}, Z. Liang^{15a}, B. Liberti^{73a}, A. Liblong¹⁶⁷, K. Lie^{63c}, C.Y. Lin³², K. Lin¹⁰⁶, T.H. Lin⁹⁹, R.A. Linck⁶⁵, J.H. Lindon²¹, A.L. Lioni⁵⁴, E. Lipeles¹³⁷, A. Lipniacka¹⁷, M. Lisovyi^{61b}, T.M. Liss^{173,as}, A. Lister¹⁷⁵, A.M. Litke¹⁴⁶, J.D. Little⁸, B. Liu⁷⁸, B.L. Liu⁶, H.B. Liu²⁹, H. Liu¹⁰⁵, J.B. Liu^{60a}, J.K.K. Liu¹³⁵, K. Liu¹³⁶, M. Liu^{60a}, P. Liu¹⁸, Y. Liu^{15a,15d}, Y.L. Liu¹⁰⁵, Y.W. Liu^{60a}, M. Livan^{70a,70b}, A. Lleres⁵⁸, J. Llorente Merino¹⁵², S.L. Lloyd⁹², C.Y. Lo^{63b}, F. Lo Sterzo⁴², E.M. Lobodzinska⁴⁶, P. Loch⁷, S. Loffredo^{73a,73b}, T. Lohse¹⁹, K. Lohwasser¹⁴⁹, M. Lokajicek¹⁴¹, J.D. Long¹⁷³, R.E. Long⁸⁹, L. Longo³⁶, K.A. Looper¹²⁶, J.A. Lopez^{147b}, I. Lopez Paz¹⁰⁰, A. Lopez Solis¹⁴⁹, J. Lorenz¹¹⁴, N. Lorenzo Martinez⁵, M. Losada²², P.J. Lösel¹¹⁴, A. Lösle⁵², X. Lou⁴⁶, X. Lou^{15a}, A. Lounis¹³², J. Love⁶, P.A. Love⁸⁹, J.J. Lozano Bahilo¹⁷⁴, M. Lu^{60a}, Y.J. Lu⁶⁴, H.J. Lubatti¹⁴⁸, C. Luci^{72a,72b}, A. Lucotte⁵⁸, C. Luedtke⁵², F. Luehring⁶⁵, I. Luise¹³⁶, L. Luminari^{72a}, B. Lund-Jensen¹⁵⁴, M.S. Lutz¹⁰², D. Lynn²⁹, R. Lysak¹⁴¹, E. Lytken⁹⁶, F. Lyu^{15a}, V. Lyubushkin⁷⁹, T. Lyubushkina⁷⁹, H. Ma²⁹, L.L. Ma^{60b}, Y. Ma^{60b}, G. Maccarrone⁵¹, A. Macchiolo¹¹⁵, C.M. Macdonald¹⁴⁹, J. Machado Miguens¹³⁷, D. Madaffari¹⁷⁴, R. Madar³⁸, W.F. Mader⁴⁸, N. Madysa⁴⁸, J. Maeda⁸², S. Maeland¹⁷, T. Maeno²⁹, M. Maerker⁴⁸, A.S. Maevskiy¹¹³, V. Magerl⁵², N. Magini⁷⁸, D.J. Mahon³⁹, C. Maidantchik^{80b}, T. Maier¹¹⁴, A. Maio^{140a,140b,140d}, O. Majersky^{28a}, S. Majewski¹³¹, Y. Makida⁸¹, N. Makovec¹³², B. Malaescu¹³⁶, Pa. Malecki⁸⁴, V.P. Maleev¹³⁸, F. Malek⁵⁸, U. Mallik⁷⁷, D. Malon⁶, C. Malone³², S. Maltezos¹⁰, S. Malyukov³⁶, J. Mamuzic¹⁷⁴, G. Mancini⁵¹, I. Mandić⁹¹, L. Manhaes de Andrade Filho^{80a}, I.M. Maniatis¹⁶², J. Manjarres Ramos⁴⁸, K.H. Mankinen⁹⁶, A. Mann¹¹⁴, A. Manousos⁷⁶, B. Mansoulie¹⁴⁵, I. Manthos¹⁶², S. Manzoni¹²⁰, A. Marantis¹⁶², G. Marceca³⁰, L. Marchese¹³⁵, G. Marchiori¹³⁶, M. Marcisovsky¹⁴¹, L. Marcocchia^{73a,73b}, C. Marcon⁹⁶, C.A. Marin Tobon³⁶, M. Marjanovic¹²⁸, Z. Marshall¹⁸, M.U.F. Martensson¹⁷², S. Marti-Garcia¹⁷⁴, C.B. Martin¹²⁶, T.A. Martin¹⁷⁸, V.J. Martin⁵⁰, B. Martin dit Latour¹⁷, L. Martinelli^{74a,74b}, M. Martinez^{14,x}, V.I. Martinez Outschoorn¹⁰², S. Martin-Haugh¹⁴⁴, V.S. Martoiu^{27b}, A.C. Martyniuk⁹⁴, A. Marzin³⁶, S.R. Maschek¹¹⁵, L. Masetti⁹⁹, T. Mashimo¹⁶³, R. Mashinistov¹¹⁰, J. Masik¹⁰⁰, A.L. Maslennikov^{122b,122a},

L. Massa ^{73a,73b}, P. Massarotti ^{69a,69b}, P. Mastrandrea ^{71a,71b}, A. Mastroberardino ^{41b,41a}, T. Masubuchi ¹⁶³, D. Matakias ¹⁰, A. Matic ¹¹⁴, P. Mättig ²⁴, J. Maurer ^{27b}, B. Maček ⁹¹, D.A. Maximov ^{122b,122a}, R. Mazini ¹⁵⁸, I. Maznas ¹⁶², S.M. Mazza ¹⁴⁶, S.P. Mc Kee ¹⁰⁵, T.G. McCarthy ¹¹⁵, W.P. McCormack ¹⁸, E.F. McDonald ¹⁰⁴, J.A. Mcfayden ³⁶, G. Mchedlidze ^{159b}, M.A. McKay ⁴², K.D. McLean ¹⁷⁶, S.J. McMahon ¹⁴⁴, P.C. McNamara ¹⁰⁴, C.J. McNicol ¹⁷⁸, R.A. McPherson ^{176,ac}, J.E. Mdhului ^{33c}, Z.A. Meadows ¹⁰², S. Meehan ³⁶, T. Megy ⁵², S. Mehlhase ¹¹⁴, A. Mehta ⁹⁰, T. Meideck ⁵⁸, B. Meirose ⁴³, D. Melini ¹⁷⁴, B.R. Mellado Garcia ^{33c}, J.D. Mellenthin ⁵³, M. Melo ^{28a}, F. Meloni ⁴⁶, A. Melzer ²⁴, S.B. Menary ¹⁰⁰, E.D. Mendes Gouveia ^{140a,140e}, L. Meng ³⁶, X.T. Meng ¹⁰⁵, S. Menke ¹¹⁵, E. Meoni ^{41b,41a}, S. Mergelmeyer ¹⁹, S.A.M. Merkt ¹³⁹, C. Merlassino ²⁰, P. Mermod ⁵⁴, L. Merola ^{69a,69b}, C. Meroni ^{68a}, O. Meshkov ^{113,110}, J.K.R. Meshreki ¹⁵¹, A. Messina ^{72a,72b}, J. Metcalfe ⁶, A.S. Mete ¹⁷¹, C. Meyer ⁶⁵, J. Meyer ¹⁶⁰, J-P. Meyer ¹⁴⁵, H. Meyer Zu Theenhausen ^{61a}, F. Miano ¹⁵⁶, M. Michetti ¹⁹, R.P. Middleton ¹⁴⁴, L. Mijović ⁵⁰, G. Mikenberg ¹⁸⁰, M. Mikestikova ¹⁴¹, M. Mikuž ⁹¹, H. Mildner ¹⁴⁹, M. Milesi ¹⁰⁴, A. Milic ¹⁶⁷, D.A. Millar ⁹², D.W. Miller ³⁷, A. Milov ¹⁸⁰, D.A. Milstead ^{45a,45b}, R.A. Mina ^{153,o}, A.A. Minaenko ¹²³, M. Miñano Moya ¹⁷⁴, I.A. Minashvili ^{159b}, A.I. Mincer ¹²⁴, B. Mindur ^{83a}, M. Mineev ⁷⁹, Y. Minegishi ¹⁶³, L.M. Mir ¹⁴, A. Mirto ^{67a,67b}, K.P. Mistry ¹³⁷, T. Mitani ¹⁷⁹, J. Mitrevski ¹¹⁴, V.A. Mitsou ¹⁷⁴, M. Mittal ^{60c}, O. Miu ¹⁶⁷, A. Miucci ²⁰, P.S. Miyagawa ¹⁴⁹, A. Mizukami ⁸¹, J.U. Mjörnmark ⁹⁶, T. Mkrtchyan ¹⁸⁴, M. Mlynarikova ¹⁴³, T. Moa ^{45a,45b}, K. Mochizuki ¹⁰⁹, P. Mogg ⁵², S. Mohapatra ³⁹, R. Moles-Valls ²⁴, M.C. Mondragon ¹⁰⁶, K. Mönig ⁴⁶, J. Monk ⁴⁰, E. Monnier ¹⁰¹, A. Montalbano ¹⁵², J. Montejo Berlingen ³⁶, M. Montella ⁹⁴, F. Monticelli ⁸⁸, S. Monzani ^{68a}, N. Morange ¹³², D. Moreno ²², M. Moreno Llácer ³⁶, C. Moreno Martinez ¹⁴, P. Morettini ^{55b}, M. Morgenstern ¹²⁰, S. Morgenstern ⁴⁸, D. Mori ¹⁵², M. Morii ⁵⁹, M. Morinaga ¹⁷⁹, V. Morisbak ¹³⁴, A.K. Morley ³⁶, G. Mornacchi ³⁶, A.P. Morris ⁹⁴, L. Morvaj ¹⁵⁵, P. Moschovakos ³⁶, B. Moser ¹²⁰, M. Mosidze ^{159b}, T. Moskalets ¹⁴⁵, H.J. Moss ¹⁴⁹, J. Moss ^{31,l}, E.J.W. Moyses ¹⁰², S. Muanza ¹⁰¹, J. Mueller ¹³⁹, R.S.P. Mueller ¹¹⁴, D. Muenstermann ⁸⁹, G.A. Mullier ⁹⁶, J.L. Munoz Martinez ¹⁴, F.J. Munoz Sanchez ¹⁰⁰, P. Murin ^{28b}, W.J. Murray ^{178,144}, A. Murrone ^{68a,68b}, M. Muškinja ¹⁸, C. Mwewa ^{33a}, A.G. Myagkov ^{123,am}, J. Myers ¹³¹, M. Myska ¹⁴², B.P. Nachman ¹⁸, O. Nackenhorst ⁴⁷, A.Nag Nag ⁴⁸, K. Nagai ¹³⁵, K. Nagano ⁸¹, Y. Nagasaka ⁶², M. Nagel ⁵², J.L. Nagle ²⁹, E. Nagy ¹⁰¹, A.M. Nairz ³⁶, Y. Nakahama ¹¹⁷, K. Nakamura ⁸¹, T. Nakamura ¹⁶³, I. Nakano ¹²⁷, H. Nanjo ¹³³, F. Napolitano ^{61a}, R.F. Naranjo Garcia ⁴⁶, R. Narayan ⁴², I. Naryshkin ¹³⁸, T. Naumann ⁴⁶, G. Navarro ²², P.Y. Nechaeva ¹¹⁰, F. Nechansky ⁴⁶, T.J. Neep ²¹, A. Negri ^{70a,70b}, M. Negrini ^{23b}, C. Nellist ⁵³, M.E. Nelson ¹³⁵, S. Nemecek ¹⁴¹, P. Nemethy ¹²⁴, M. Nessi ^{36,c}, M.S. Neubauer ¹⁷³, M. Neumann ¹⁸², P.R. Newman ²¹, Y.S. Ng ¹⁹, Y.W.Y. Ng ¹⁷¹, B. Ngair ^{35e}, H.D.N. Nguyen ¹⁰¹, T. Nguyen Manh ¹⁰⁹, E. Nibigira ³⁸, R.B. Nickerson ¹³⁵, R. Nicolaidou ¹⁴⁵, D.S. Nielsen ⁴⁰, J. Nielsen ¹⁴⁶, N. Nikiforou ¹¹, V. Nikolaenko ^{123,am}, I. Nikolic-Audit ¹³⁶, K. Nikolopoulos ²¹, P. Nilsson ²⁹, H.R. Nindhito ⁵⁴, Y. Ninomiya ⁸¹, A. Nisati ^{72a}, N. Nishu ^{60c}, R. Nisius ¹¹⁵, I. Nitsche ⁴⁷, T. Nitta ¹⁷⁹, T. Nobe ¹⁶³, Y. Noguchi ⁸⁵, I. Nomidis ¹³⁶, M.A. Nomura ²⁹, M. Nordberg ³⁶, N. Norjoharuddeen ¹³⁵, T. Novak ⁹¹, O. Novgorodova ⁴⁸, R. Novotny ¹⁴², L. Nozka ¹³⁰, K. Ntekas ¹⁷¹, E. Nurse ⁹⁴, F.G. Oakham ^{34,av}, H. Oberlack ¹¹⁵, J. Ocariz ¹³⁶, A. Ochi ⁸², I. Ochoa ³⁹, J.P. Ochoa-Ricoux ^{147a}, K. O'Connor ²⁶, S. Oda ⁸⁷, S. Odaka ⁸¹, S. Oerdek ⁵³, A. Ogrodnik ^{83a}, A. Oh ¹⁰⁰, S.H. Oh ⁴⁹, C.C. Ohm ¹⁵⁴, H. Oide ¹⁶⁵, M.L. Ojeda ¹⁶⁷, H. Okawa ¹⁶⁹, Y. Okazaki ⁸⁵, Y. Okumura ¹⁶³, T. Okuyama ⁸¹, A. Olariu ^{27b}, L.F. Oleiro Seabra ^{140a}, S.A. Olivares Pino ^{147a}, D. Oliveira Damazio ²⁹, J.L. Oliver ¹, M.J.R. Olsson ¹⁷¹, A. Olszewski ⁸⁴, J. Olszowska ⁸⁴, D.C. O'Neil ¹⁵², A.P. O'Neill ¹³⁵, A. Onofre ^{140a,140e}, P.U.E. Onyisi ¹¹, H. Oppen ¹³⁴, M.J. Oreglia ³⁷, G.E. Orellana ⁸⁸, D. Orestano ^{74a,74b}, N. Orlando ¹⁴, R.S. Orr ¹⁶⁷, V. O'Shea ⁵⁷, R. Ospanov ^{60a}, G. Otero y Garzon ³⁰, H. Otono ⁸⁷, P.S. Ott ^{61a}, M. Ouchrif ^{35d}, J. Ouellette ²⁹, F. Ould-Saada ¹³⁴, A. Ouraou ¹⁴⁵, Q. Ouyang ^{15a}, M. Owen ⁵⁷, R.E. Owen ²¹, V.E. Ozcan ^{12c}, N. Ozturk ⁸, J. Pacalt ¹³⁰, H.A. Pacey ³², K. Pachal ⁴⁹, A. Pacheco Pages ¹⁴, C. Padilla Aranda ¹⁴, S. Pagan Griso ¹⁸, M. Paganini ¹⁸³, G. Palacino ⁶⁵, S. Palazzo ⁵⁰, S. Palestini ³⁶, M. Palka ^{83b}, D. Pallin ³⁸, I. Panagoulas ¹⁰, C.E. Pandini ³⁶, J.G. Panduro Vazquez ⁹³, P. Pani ⁴⁶, G. Panizzo ^{66a,66c}, L. Paolozzi ⁵⁴, C. Papadatos ¹⁰⁹, K. Papageorgiou ^{9,g}, S. Parajuli ⁴³, A. Paramonov ⁶, D. Paredes Hernandez ^{63b}, S.R. Paredes Saenz ¹³⁵, B. Parida ¹⁶⁶, T.H. Park ¹⁶⁷, A.J. Parker ³¹, M.A. Parker ³², F. Parodi ^{55b,55a}, E.W.P. Parrish ¹²¹, J.A. Parsons ³⁹, U. Parzefall ⁵², L. Pascual Dominguez ¹³⁶, V.R. Pascuzzi ¹⁶⁷, J.M.P. Pasner ¹⁴⁶, F. Pasquali ¹²⁰, E. Pasqualucci ^{72a}, S. Passaggio ^{55b}, F. Pastore ⁹³, P. Pasuwan ^{45a,45b}, S. Patariaia ⁹⁹, J.R. Pater ¹⁰⁰, A. Pathak ^{181,i}, T. Pauly ³⁶, B. Pearson ¹¹⁵, M. Pedersen ¹³⁴, L. Pedraza Diaz ¹¹⁹, R. Pedro ^{140a}, T. Peiffer ⁵³, S.V. Peleganchuk ^{122b,122a},

O. Penc¹⁴¹, H. Peng^{60a}, B.S. Peralva^{80a}, M.M. Perego¹³², A.P. Pereira Peixoto^{140a}, D.V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{68a,68b}, H. Pernegger³⁶, S. Perrella^{69a,69b}, K. Peters⁴⁶, R.F.Y. Peters¹⁰⁰, B.A. Petersen³⁶, T.C. Petersen⁴⁰, E. Petit¹⁰¹, A. Petridis¹, C. Petridou¹⁶², P. Petroff¹³², M. Petrov¹³⁵, F. Petrucci^{74a,74b}, M. Pettee¹⁸³, N.E. Pettersson¹⁰², K. Petukhova¹⁴³, A. Peyaud¹⁴⁵, R. Pezoa^{147b}, L. Pezzotti^{70a,70b}, T. Pham¹⁰⁴, F.H. Phillips¹⁰⁶, P.W. Phillips¹⁴⁴, M.W. Phipps¹⁷³, G. Piacquadio¹⁵⁵, E. Pianori¹⁸, A. Picazio¹⁰², R.H. Pickles¹⁰⁰, R. Piegaia³⁰, D. Pietreanu^{27b}, H.P. Pikhartova⁹³, J.E. Pilcher³⁷, A.D. Pilkington¹⁰⁰, M. Pinamonti^{73a,73b}, J.L. Pinfold³, M. Pitt¹⁶¹, L. Pizzimento^{73a,73b}, M.-A. Pleier²⁹, V. Pleskot¹⁴³, E. Plotnikova⁷⁹, P. Podberezko^{122b,122a}, R. Poettgen⁹⁶, R. Poggi⁵⁴, L. Poggioli¹³², I. Pogrebnyak¹⁰⁶, D. Pohl²⁴, I. Pokharel⁵³, G. Polesello^{70a}, A. Poley¹⁸, A. Policicchio^{72a,72b}, R. Polifka¹⁴³, A. Polini^{23b}, C.S. Pollard⁴⁶, V. Polychronakos²⁹, D. Ponomarenko¹¹², L. Pontecorvo³⁶, S. Popa^{27a}, G.A. Popeneciu^{27d}, L. Portales⁵, D.M. Portillo Quintero⁵⁸, S. Pospisil¹⁴², K. Potamianos⁴⁶, I.N. Potrap⁷⁹, C.J. Potter³², H. Potti¹¹, T. Poulsen⁹⁶, J. Poveda³⁶, T.D. Powell¹⁴⁹, G. Pownall⁴⁶, M.E. Pozo Astigarraga³⁶, P. Pralavorio¹⁰¹, S. Prell⁷⁸, D. Price¹⁰⁰, M. Primavera^{67a}, S. Prince¹⁰³, M.L. Proffitt¹⁴⁸, N. Proklova¹¹², K. Prokofiev^{63c}, F. Prokoshin⁷⁹, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{83a}, D. Pudzha¹³⁸, A. Puri¹⁷³, P. Puzo¹³², J. Qian¹⁰⁵, Y. Qin¹⁰⁰, A. Quadt⁵³, M. Queitsch-Maitland⁴⁶, A. Qureshi¹, M. Racko^{28a}, P. Rados¹⁰⁴, F. Ragusa^{68a,68b}, G. Rahal⁹⁷, J.A. Raine⁵⁴, S. Rajagopalan²⁹, A. Ramirez Morales⁹², K. Ran^{15a,15d}, T. Rashid¹³², S. Raspopov⁵, D.M. Rauch⁴⁶, F. Rauscher¹¹⁴, S. Rave⁹⁹, B. Ravina¹⁴⁹, I. Ravinovich¹⁸⁰, J.H. Rawling¹⁰⁰, M. Raymond³⁶, A.L. Read¹³⁴, N.P. Readioff⁵⁸, M. Reale^{67a,67b}, D.M. Rebuffi^{70a,70b}, A. Redelbach¹⁷⁷, G. Redlinger²⁹, K. Reeves⁴³, L. Rehnisch¹⁹, J. Reichert¹³⁷, D. Reikher¹⁶¹, A. Reiss⁹⁹, A. Rej¹⁵¹, C. Rembser³⁶, M. Renda^{27b}, M. Rescigno^{72a}, S. Resconi^{68a}, E.D. Resseguie¹³⁷, S. Rettie¹⁷⁵, E. Reynolds²¹, O.L. Rezanova^{122b,122a}, P. Reznicek¹⁴³, E. Ricci^{75a,75b}, R. Richter¹¹⁵, S. Richter⁴⁶, E. Richter-Was^{83b}, O. Ricken²⁴, M. Ridel¹³⁶, P. Rieck¹¹⁵, C.J. Riegel¹⁸², O. Rifki⁴⁶, M. Rijssenbeek¹⁵⁵, A. Rimoldi^{70a,70b}, M. Rimoldi⁴⁶, L. Rinaldi^{23b}, G. Ripellino¹⁵⁴, I. Riu¹⁴, J.C. Rivera Vergara¹⁷⁶, F. Rizatdinova¹²⁹, E. Rizvi⁹², C. Rizzi³⁶, R.T. Roberts¹⁰⁰, S.H. Robertson^{103,ac}, M. Robin⁴⁶, D. Robinson³², J.E.M. Robinson⁴⁶, C.M. Robles Gajardo^{147b}, A. Robson⁵⁷, A. Rocchi^{73a,73b}, E. Rocco⁹⁹, C. Roda^{71a,71b}, S. Rodriguez Bosca¹⁷⁴, A. Rodriguez Perez¹⁴, D. Rodriguez Rodriguez¹⁷⁴, A.M. Rodríguez Vera^{168b}, S. Roe³⁶, O. Röhne¹³⁴, R. Röhrig¹¹⁵, R.A. Rojas^{147b}, C.P.A. Roland⁶⁵, J. Roloff⁵⁹, A. Romaniouk¹¹², M. Romano^{23b,23a}, N. Rompotis⁹⁰, M. Ronzani¹²⁴, L. Roos¹³⁶, S. Rosati^{72a}, K. Rosbach⁵², G. Rosin¹⁰², B.J. Rosser¹³⁷, E. Rossi⁴⁶, E. Rossi^{74a,74b}, E. Rossi^{69a,69b}, L.P. Rossi^{55b}, L. Rossini^{68a,68b}, R. Rosten¹⁴, M. Rotaru^{27b}, J. Rothberg¹⁴⁸, D. Rousseau¹³², G. Rovelli^{70a,70b}, A. Roy¹¹, D. Roy^{33c}, A. Rozanov¹⁰¹, Y. Rozen¹⁶⁰, X. Ruan^{33c}, F. Rubbo¹⁵³, F. Rühr⁵², A. Ruiz-Martinez¹⁷⁴, A. Rummler³⁶, Z. Rurikova⁵², N.A. Rusakovich⁷⁹, H.L. Russell¹⁰³, L. Rustige^{38,47}, J.P. Rutherford⁷, E.M. Rüttinger¹⁴⁹, Y.F. Ryabov¹³⁸, M. Rybar³⁹, G. Rybkin¹³², E.B. Rye¹³⁴, A. Ryzhov¹²³, P. Sabatini⁵³, G. Sabato¹²⁰, S. Sacerdoti¹³², H.F.W. Sadrozinski¹⁴⁶, R. Sadykov⁷⁹, F. Safai Tehrani^{72a}, B. Safarzadeh Samani¹⁵⁶, P. Saha¹²¹, S. Saha¹⁰³, M. Sahinsoy^{61a}, A. Sahu¹⁸², M. Saimpert⁴⁶, M. Saito¹⁶³, T. Saito¹⁶³, H. Sakamoto¹⁶³, A. Sakharov^{124,al}, D. Salamani⁵⁴, G. Salamanna^{74a,74b}, J.E. Salazar Loyola^{147b}, P.H. Sales De Bruin¹⁷², A. Salmikov¹⁵³, J. Salt¹⁷⁴, D. Salvatore^{41b,41a}, F. Salvatore¹⁵⁶, A. Salvucci^{63a,63b,63c}, A. Salzburger³⁶, J. Samarati³⁶, D. Sammel⁵², D. Sampsonidis¹⁶², D. Sampsonidou¹⁶², J. Sánchez¹⁷⁴, A. Sanchez Pineda^{66a,66c}, H. Sandaker¹³⁴, C.O. Sander⁴⁶, I.G. Sanderswood⁸⁹, M. Sandhoff¹⁸², C. Sandoval²², D.P.C. Sankey¹⁴⁴, M. Sannino^{55b,55a}, Y. Sano¹¹⁷, A. Sansoni⁵¹, C. Santoni³⁸, H. Santos^{140a,140b}, S.N. Santpur¹⁸, A. Santra¹⁷⁴, A. Saprnov⁷⁹, J.G. Saraiva^{140a,140d}, O. Sasaki⁸¹, K. Sato¹⁶⁹, F. Sauerburger⁵², E. Sauvan⁵, P. Savard^{167,av}, N. Savic¹¹⁵, R. Sawada¹⁶³, C. Sawyer¹⁴⁴, L. Sawyer^{95,aj}, C. Sbarra^{23b}, A. Sbrizzi^{23a}, T. Scanlon⁹⁴, J. Schaarschmidt¹⁴⁸, P. Schacht¹¹⁵, B.M. Schachtner¹¹⁴, D. Schaefer³⁷, L. Schaefer¹³⁷, J. Schaeffer⁹⁹, S. Schaepe³⁶, U. Schäfer⁹⁹, A.C. Schaffer¹³², D. Schaile¹¹⁴, R.D. Schamberger¹⁵⁵, N. Scharmberg¹⁰⁰, V.A. Schegelsky¹³⁸, D. Scheirich¹⁴³, F. Schenck¹⁹, M. Schernau¹⁷¹, C. Schiavi^{55b,55a}, S. Schier¹⁴⁶, L.K. Schildgen²⁴, Z.M. Schillaci²⁶, E.J. Schioppa³⁶, M. Schioppa^{41b,41a}, K.E. Schleicher⁵², S. Schlenker³⁶, K.R. Schmidt-Sommerfeld¹¹⁵, K. Schmieden³⁶, C. Schmitt⁹⁹, S. Schmitt⁴⁶, S. Schmitz⁹⁹, J.C. Schmoeckel⁴⁶, U. Schnoor⁵², L. Schoeffel¹⁴⁵, A. Schoening^{61b}, P.G. Scholer⁵², E. Schopf¹³⁵, M. Schott⁹⁹, J.F.P. Schouwenberg¹¹⁹, J. Schovancova³⁶, S. Schramm⁵⁴, F. Schroeder¹⁸², A. Schulte⁹⁹, H-C. Schultz-Coulon^{61a}, M. Schumacher⁵², B.A. Schumm¹⁴⁶, Ph. Schune¹⁴⁵, A. Schwartzman¹⁵³, T.A. Schwarz¹⁰⁵, Ph. Schwemling¹⁴⁵, R. Schwienhorst¹⁰⁶, A. Sciandra¹⁴⁶, G. Sciolla²⁶, M. Scodeggio⁴⁶,

M. Scornajenghi ^{41b,41a}, F. Scuri ^{71a}, F. Scutti ¹⁰⁴, L.M. Scyboz ¹¹⁵, C.D. Sebastiani ^{72a,72b}, P. Seema ¹⁹,
 S.C. Seidel ¹¹⁸, A. Seiden ¹⁴⁶, B.D. Seidlitz ²⁹, T. Seiss ³⁷, J.M. Seixas ^{80b}, G. Sekhniaidze ^{69a}, K. Sekhon ¹⁰⁵,
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 A. Sharma ¹³⁵, A.S. Sharma ¹, P.B. Shatalov ¹¹¹, K. Shaw ¹⁵⁶, S.M. Shaw ¹⁰⁰, A. Shcherbakova ¹³⁸,
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 J. Shlomi ¹⁸⁰, A. Shmeleva ¹¹⁰, M.J. Shochet ³⁷, J. Shojaii ¹⁰⁴, D.R. Shope ¹²⁸, S. Shrestha ¹²⁶, E.M. Shrif ^{33c},
 E. Shulga ¹⁸⁰, P. Sicho ¹⁴¹, A.M. Sickles ¹⁷³, P.E. Sidebo ¹⁵⁴, E. Sideras Haddad ^{33c}, O. Sidiropoulou ³⁶,
 A. Sidoti ^{23b,23a}, F. Siegert ⁴⁸, Dj. Sijacki ¹⁶, M. Jr. Silva ¹⁸¹, M.V. Silva Oliveira ^{80a}, S.B. Silverstein ^{45a},
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 M. Slawinska ⁸⁴, K. Sliwa ¹⁷⁰, R. Slovak ¹⁴³, V. Smakhtin ¹⁸⁰, B.H. Smart ¹⁴⁴, J. Smiesko ^{28a}, N. Smirnov ¹¹²,
 S.Yu. Smirnov ¹¹², Y. Smirnov ¹¹², L.N. Smirnova ^{113,t}, O. Smirnova ⁹⁶, J.W. Smith ⁵³, M. Smizanska ⁸⁹,
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 Y. Takubo ⁸¹, M. Talby ¹⁰¹, A.A. Talyshev ^{122b,122a}, N.M. Tamir ¹⁶¹, J. Tanaka ¹⁶³, M. Tanaka ¹⁶⁵,
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 J. Terron ⁹⁸, S. Terzo ¹⁴, M. Testa ⁵¹, R.J. Teuscher ^{167,ac}, S.J. Thais ¹⁸³, T. Theveneaux-Pelzer ⁴⁶, F. Thiele ⁴⁰,
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 E. Thomson ¹³⁷, E.J. Thorpe ⁹², R.E. Ticse Torres ⁵³, V.O. Tikhomirov ^{110,an}, Yu.A. Tikhonov ^{122b,122a},
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K. Uno ¹⁶³, J. Urban ^{28b}, P. Urquijo ¹⁰⁴, G. Usai ⁸, Z. Uysal ^{12d}, L. Vacavant ¹⁰¹, V. Vacek ¹⁴², B. Vachon ¹⁰³, K.O.H. Vadla ¹³⁴, A. Vaidya ⁹⁴, C. Valderanis ¹¹⁴, E. Valdes Santurio ^{45a,45b}, M. Valente ⁵⁴, S. Valentini ^{23b,23a}, A. Valero ¹⁷⁴, L. Valéry ⁴⁶, R.A. Vallance ²¹, A. Vallier ³⁶, J.A. Valls Ferrer ¹⁷⁴, T.R. Van Daalen ¹⁴, P. Van Gemmeren ⁶, I. Van Vulpen ¹²⁰, M. Vanadia ^{73a,73b}, W. Vandelli ³⁶, E.R. Vandewall ¹²⁹, A. Vaniachine ¹⁶⁶, D. Vannicola ^{72a,72b}, R. Vari ^{72a}, E.W. Varnes ⁷, C. Varni ^{55b,55a}, T. Varol ¹⁵⁸, D. Varouchas ¹³², K.E. Varvell ¹⁵⁷, M.E. Vasile ^{27b}, G.A. Vasquez ¹⁷⁶, J.G. Vasquez ¹⁸³, F. Vazeille ³⁸, D. Vazquez Furelos ¹⁴, T. Vazquez Schroeder ³⁶, J. Veatch ⁵³, V. Vecchio ^{74a,74b}, M.J. Veen ¹²⁰, L.M. Veloce ¹⁶⁷, F. Veloso ^{140a,140c}, S. Veneziano ^{72a}, A. Ventura ^{67a,67b}, N. Venturi ³⁶, A. Verbytskyi ¹¹⁵, V. Vercesi ^{70a}, M. Verducci ^{71a,71b}, C.M. Vergel Infante ⁷⁸, C. Vergis ²⁴, W. Verkerke ¹²⁰, A.T. Vermeulen ¹²⁰, J.C. Vermeulen ¹²⁰, M.C. Vetterli ^{152.av}, N. Viaux Maira ^{147b}, M. Vicente Barreto Pinto ⁵⁴, T. Vickey ¹⁴⁹, O.E. Vickey Boeriu ¹⁴⁹, G.H.A. Viehhauser ¹³⁵, L. Vigani ^{61b}, M. Villa ^{23b,23a}, M. Villaplana Perez ^{68a,68b}, E. Vilucchi ⁵¹, M.G. Vincker ³⁴, G.S. Virdee ²¹, A. Vishwakarma ⁴⁶, C. Vittori ^{23b,23a}, I. Vivarelli ¹⁵⁶, M. Vogel ¹⁸², P. Vokac ¹⁴², S.E. von Buddenbrock ^{33c}, E. Von Toerne ²⁴, V. Vorobel ¹⁴³, K. Vorobev ¹¹², M. Vos ¹⁷⁴, J.H. Vosseveld ⁹⁰, M. Vozak ¹⁰⁰, N. Vranjes ¹⁶, M. Vranjes Milosavljevic ¹⁶, V. Vrba ¹⁴², M. Vreeswijk ¹²⁰, T. Šfiligoj ⁹¹, R. Vuillermet ³⁶, I. Vukotic ³⁷, T. Ženiš ^{28a}, L. Živković ¹⁶, P. Wagner ²⁴, W. Wagner ¹⁸², J. Wagner-Kuhr ¹¹⁴, S. Wahdan ¹⁸², H. Wahlberg ⁸⁸, V.M. Walbrecht ¹¹⁵, J. Walder ⁸⁹, R. Walker ¹¹⁴, S.D. Walker ⁹³, W. Walkowiak ¹⁵¹, V. Wallangen ^{45a,45b}, A.M. Wang ⁵⁹, C. Wang ^{60c}, C. Wang ^{60b}, F. Wang ¹⁸¹, H. Wang ¹⁸, H. Wang ³, J. Wang ^{63a}, J. Wang ¹⁵⁷, J. Wang ^{61b}, P. Wang ⁴², Q. Wang ¹²⁸, R.-J. Wang ⁹⁹, R. Wang ^{60a}, R. Wang ⁶, S.M. Wang ¹⁵⁸, W.T. Wang ^{60a}, W. Wang ^{15c,ad}, W.X. Wang ^{60a,ad}, Y. Wang ^{60a,ak}, Z. Wang ^{60c}, C. Wanotayaroj ⁴⁶, A. Warburton ¹⁰³, C.P. Ward ³², D.R. Wardrope ⁹⁴, N. Warrack ⁵⁷, A. Washbrook ⁵⁰, A.T. Watson ²¹, M.F. Watson ²¹, G. Watts ¹⁴⁸, B.M. Waugh ⁹⁴, A.F. Webb ¹¹, S. Webb ⁹⁹, C. Weber ¹⁸³, M.S. Weber ²⁰, S.A. Weber ³⁴, S.M. Weber ^{61a}, A.R. Weidberg ¹³⁵, J. Weingarten ⁴⁷, M. Weirich ⁹⁹, C. Weiser ⁵², P.S. Wells ³⁶, T. Wenaus ²⁹, T. Wengler ³⁶, S. Wenig ³⁶, N. Wermes ²⁴, M.D. Werner ⁷⁸, M. Wessels ^{61a}, T.D. Weston ²⁰, K. Whalen ¹³¹, N.L. Whallon ¹⁴⁸, A.M. Wharton ⁸⁹, A.S. White ¹⁰⁵, A. White ⁸, M.J. White ¹, D. Whiteson ¹⁷¹, B.W. Whitmore ⁸⁹, W. Wiedenmann ¹⁸¹, M. Wielers ¹⁴⁴, N. Wieseotte ⁹⁹, C. Wiglesworth ⁴⁰, L.A.M. Wiik-Fuchs ⁵², F. Wilk ¹⁰⁰, H.G. Wilkens ³⁶, L.J. Wilkins ⁹³, H.H. Williams ¹³⁷, S. Williams ³², C. Willis ¹⁰⁶, S. Willocq ¹⁰², J.A. Wilson ²¹, I. Wingerter-Seez ⁵, E. Winkels ¹⁵⁶, F. Winklmeier ¹³¹, O.J. Winston ¹⁵⁶, B.T. Winter ⁵², M. Wittgen ¹⁵³, M. Wobisch ⁹⁵, A. Wolf ⁹⁹, T.M.H. Wolf ¹²⁰, R. Wolff ¹⁰¹, R.W. Wölker ¹³⁵, J. Wollrath ⁵², M.W. Wolter ⁸⁴, H. Wolters ^{140a,140c}, V.W.S. Wong ¹⁷⁵, N.L. Woods ¹⁴⁶, S.D. Worm ²¹, B.K. Wosiek ⁸⁴, K.W. Woźniak ⁸⁴, K. Wraight ⁵⁷, S.L. Wu ¹⁸¹, X. Wu ⁵⁴, Y. Wu ^{60a}, T.R. Wyatt ¹⁰⁰, B.M. Wynne ⁵⁰, S. Xella ⁴⁰, Z. Xi ¹⁰⁵, L. Xia ¹⁷⁸, X. Xiao ¹⁰⁵, I. Xiotidis ¹⁵⁶, D. Xu ^{15a}, H. Xu ^{60a,b}, L. Xu ²⁹, T. Xu ¹⁴⁵, W. Xu ¹⁰⁵, Z. Xu ^{60b}, Z. Xu ¹⁵³, B. Yabsley ¹⁵⁷, S. Yacoob ^{33a}, K. Yajima ¹³³, D.P. Yallup ⁹⁴, D. Yamaguchi ¹⁶⁵, Y. Yamaguchi ¹⁶⁵, A. Yamamoto ⁸¹, M. Yamatani ¹⁶³, T. Yamazaki ¹⁶³, Y. Yamazaki ⁸², Z. Yan ²⁵, H.J. Yang ^{60c,60d}, H.T. Yang ¹⁸, S. Yang ⁷⁷, X. Yang ^{60b,58}, Y. Yang ¹⁶³, W.-M. Yao ¹⁸, Y.C. Yap ⁴⁶, Y. Yasu ⁸¹, E. Yatsenko ^{60c,60d}, J. Ye ⁴², S. Ye ²⁹, I. Yeletsikh ⁷⁹, M.R. Yexley ⁸⁹, E. Yigitbasi ²⁵, K. Yorita ¹⁷⁹, K. Yoshihara ¹³⁷, C.J.S. Young ³⁶, C. Young ¹⁵³, J. Yu ⁷⁸, R. Yuan ^{60b,h}, X. Yue ^{61a}, S.P.Y. Yuen ²⁴, M. Zaazoua ^{35e}, B. Zabinski ⁸⁴, G. Zacharis ¹⁰, E. Zaffaroni ⁵⁴, J. Zahreddine ¹³⁶, A.M. Zaitsev ^{123,am}, T. Zakareishvili ^{159b}, N. Zakharchuk ³⁴, S. Zambito ⁵⁹, D. Zanzi ³⁶, D.R. Zaripovas ⁵⁷, S.V. Zeiβner ⁴⁷, C. Zeitnitz ¹⁸², G. Zemaityte ¹³⁵, J.C. Zeng ¹⁷³, O. Zenin ¹²³, D. Zerwas ¹³², M. Zgubič ¹³⁵, D.F. Zhang ^{15b}, G. Zhang ^{15b}, H. Zhang ^{15c}, J. Zhang ⁶, L. Zhang ^{15c}, L. Zhang ^{60a}, M. Zhang ¹⁷³, R. Zhang ²⁴, X. Zhang ^{60b}, Y. Zhang ^{15a,15d}, Z. Zhang ^{63a}, Z. Zhang ¹³², P. Zhao ⁴⁹, Y. Zhao ^{60b}, Z. Zhao ^{60a}, A. Zhemchugov ⁷⁹, Z. Zheng ¹⁰⁵, D. Zhong ¹⁷³, B. Zhou ¹⁰⁵, C. Zhou ¹⁸¹, M.S. Zhou ^{15a,15d}, M. Zhou ¹⁵⁵, N. Zhou ^{60c}, Y. Zhou ⁷, C.G. Zhu ^{60b}, C. Zhu ^{15a}, H.L. Zhu ^{60a}, H. Zhu ^{15a}, J. Zhu ¹⁰⁵, Y. Zhu ^{60a}, X. Zhuang ^{15a}, K. Zhukov ¹¹⁰, V. Zhulanov ^{122b,122a}, D. Zieminska ⁶⁵, N.I. Zimine ⁷⁹, S. Zimmermann ⁵², Z. Zinonos ¹¹⁵, M. Ziolkowski ¹⁵¹, G. Zobernig ¹⁸¹, A. Zoccoli ^{23b,23a}, K. Zoch ⁵³, T.G. Zorbas ¹⁴⁹, R. Zou ³⁷, L. Zwalinski ³⁶

¹ Department of Physics, University of Adelaide, Adelaide, Australia² Physics Department, SUNY Albany, Albany, NY, United States of America³ Department of Physics, University of Alberta, Edmonton, AB, Canada⁴ (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey⁵ LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France⁶ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States of America⁷ Department of Physics, University of Arizona, Tucson, AZ, United States of America⁸ Department of Physics, University of Texas at Arlington, Arlington, TX, United States of America⁹ Physics Department, National and Kapodistrian University of Athens, Athens, Greece

- ¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece
- ¹¹ Department of Physics, University of Texas at Austin, Austin, TX, United States of America
- ¹² ^(a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; ^(b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; ^(c) Department of Physics, Bogazici University, Istanbul; ^(d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
- ¹³ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
- ¹⁴ Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain
- ¹⁵ ^(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; ^(b) Physics Department, Tsinghua University, Beijing; ^(c) Department of Physics, Nanjing University, Nanjing;
- ^(d) University of Chinese Academy of Science (UCAS), Beijing, China
- ¹⁶ Institute of Physics, University of Belgrade, Belgrade, Serbia
- ¹⁷ Department for Physics and Technology, University of Bergen, Bergen, Norway
- ¹⁸ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States of America
- ¹⁹ Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany
- ²⁰ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
- ²¹ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- ²² Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia
- ²³ ^(a) INFN Bologna and Università di Bologna, Dipartimento di Fisica; ^(b) INFN Sezione di Bologna, Italy
- ²⁴ Physikalisches Institut, Universität Bonn, Bonn, Germany
- ²⁵ Department of Physics, Boston University, Boston, MA, United States of America
- ²⁶ Department of Physics, Brandeis University, Waltham, MA, United States of America
- ²⁷ ^(a) Transilvania University of Brasov, Brasov; ^(b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; ^(c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; ^(d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; ^(e) University Politehnica Bucharest, Bucharest; ^(f) West University in Timisoara, Timisoara, Romania
- ²⁸ ^(a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; ^(b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
- ²⁹ Physics Department, Brookhaven National Laboratory, Upton, NY, United States of America
- ³⁰ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- ³¹ California State University, CA, United States of America
- ³² Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- ³³ ^(a) Department of Physics, University of Cape Town, Cape Town; ^(b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; ^(c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³⁴ Department of Physics, Carleton University, Ottawa, ON, Canada
- ³⁵ ^(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies – Université Hassan II, Casablanca; ^(b) Faculté des Sciences, Université Ibn-Tofail, Kénitra; ^(c) Faculté des Sciences Semailia, Université Cadi Ayyad, LPHEA, Marrakech; ^(d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; ^(e) Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ³⁶ CERN, Geneva, Switzerland
- ³⁷ Enrico Fermi Institute, University of Chicago, Chicago, IL, United States of America
- ³⁸ LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ³⁹ Nevis Laboratory, Columbia University, Irvington, NY, United States of America
- ⁴⁰ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ⁴¹ ^(a) Dipartimento di Fisica, Università della Calabria, Rende; ^(b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴² Physics Department, Southern Methodist University, Dallas, TX, United States of America
- ⁴³ Physics Department, University of Texas at Dallas, Richardson, TX, United States of America
- ⁴⁴ National Centre for Scientific Research "Demokritos", Agia Paraskevi, Greece
- ⁴⁵ ^(a) Department of Physics, Stockholm University; ^(b) Oskar Klein Centre, Stockholm, Sweden
- ⁴⁶ Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁷ Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴⁸ Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁴⁹ Department of Physics, Duke University, Durham, NC, United States of America
- ⁵⁰ SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁵¹ INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵² Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ⁵³ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁵⁴ Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ⁵⁵ ^(a) Dipartimento di Fisica, Università di Genova, Genova; ^(b) INFN Sezione di Genova, Italy
- ⁵⁶ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁷ SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁵⁸ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁵⁹ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States of America
- ⁶⁰ ^(a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; ^(b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; ^(c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; ^(d) Tsung-Dao Lee Institute, Shanghai, China
- ⁶¹ ^(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ⁶² Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶³ ^(a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, University of Hong Kong, Hong Kong; ^(c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶⁴ Department of Physics, National Tsing Hua University, Hsinchu, Taiwan
- ⁶⁵ Department of Physics, Indiana University, Bloomington, IN, United States of America
- ⁶⁶ ^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine, Italy
- ⁶⁷ ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- ⁶⁸ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁶⁹ ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- ⁷⁰ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ⁷¹ ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ⁷² ^(a) INFN Sezione di Roma; ^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- ⁷³ ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ⁷⁴ ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- ⁷⁵ ^(a) INFN-TIFPA; ^(b) Università degli Studi di Trento, Trento, Italy
- ⁷⁶ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁷⁷ University of Iowa, Iowa City, IA, United States of America

- 78 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States of America
- 79 Joint Institute for Nuclear Research, Dubna, Russia
- 80 (a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; (b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro;
- (c) Universidade Federal de São João del Rei (UFSJ), São João del Rei; (d) Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
- 81 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- 82 Graduate School of Science, Kobe University, Kobe, Japan
- 83 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
- 84 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- 85 Faculty of Science, Kyoto University, Kyoto, Japan
- 86 Kyoto University of Education, Kyoto, Japan
- 87 Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
- 88 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- 89 Physics Department, Lancaster University, Lancaster, United Kingdom
- 90 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- 91 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
- 92 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- 93 Department of Physics, Royal Holloway University of London, Egham, United Kingdom
- 94 Department of Physics and Astronomy, University College London, London, United Kingdom
- 95 Louisiana Tech University, Ruston, LA, United States of America
- 96 Fysiska institutionen, Lunds universitet, Lund, Sweden
- 97 Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
- 98 Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
- 99 Institut für Physik, Universität Mainz, Mainz, Germany
- 100 School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- 101 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- 102 Department of Physics, University of Massachusetts, Amherst, MA, United States of America
- 103 Department of Physics, McGill University, Montreal, QC, Canada
- 104 School of Physics, University of Melbourne, Victoria, Australia
- 105 Department of Physics, University of Michigan, Ann Arbor, MI, United States of America
- 106 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States of America
- 107 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- 108 Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
- 109 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- 110 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
- 111 Institute for Theoretical and Experimental Physics of the National Research Centre Kurchatov Institute, Moscow, Russia
- 112 National Research Nuclear University MEPhI, Moscow, Russia
- 113 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- 114 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- 115 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- 116 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 117 Graduate School of Science and Kobayashi–Maskawa Institute, Nagoya University, Nagoya, Japan
- 118 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States of America
- 119 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- 120 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- 121 Department of Physics, Northern Illinois University, DeKalb, IL, United States of America
- 122 (a) Budker Institute of Nuclear Physics and NSU, SB RAS, Novosibirsk; (b) Novosibirsk State University Novosibirsk, Russia
- 123 Institute for High Energy Physics of the National Research Centre Kurchatov Institute, Protvino, Russia
- 124 Department of Physics, New York University, New York, NY, United States of America
- 125 Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan
- 126 Ohio State University, Columbus, OH, United States of America
- 127 Faculty of Science, Okayama University, Okayama, Japan
- 128 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States of America
- 129 Department of Physics, Oklahoma State University, Stillwater, OK, United States of America
- 130 Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic
- 131 Center for High Energy Physics, University of Oregon, Eugene, OR, United States of America
- 132 LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
- 133 Graduate School of Science, Osaka University, Osaka, Japan
- 134 Department of Physics, University of Oslo, Oslo, Norway
- 135 Department of Physics, Oxford University, Oxford, United Kingdom
- 136 LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France
- 137 Department of Physics, University of Pennsylvania, Philadelphia, PA, United States of America
- 138 Konstantinov Nuclear Physics Institute of National Research Centre "Kurchatov Institute", PNPI, St. Petersburg, Russia
- 139 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States of America
- 140 (a) Laboratório de Instrumentação e Física Experimental de Partículas – LIP; (b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; (c) Departamento de Física, Universidade de Coimbra, Coimbra; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; (e) Departamento de Física, Universidade do Minho, Braga; (f) Universidad de Granada, Granada (Spain); (g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica; (h) Av. Rovisco Pais, 1 1049-001 Lisbon, Portugal, Portugal
- 141 Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
- 142 Czech Technical University in Prague, Prague, Czech Republic
- 143 Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- 144 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- 145 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
- 146 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States of America
- 147 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- 148 Department of Physics, University of Washington, Seattle, WA, United States of America
- 149 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- 150 Department of Physics, Shinshu University, Nagano, Japan
- 151 Department Physik, Universität Siegen, Siegen, Germany

- ¹⁵² Department of Physics, Simon Fraser University, Burnaby, BC, Canada
¹⁵³ SLAC National Accelerator Laboratory, Stanford, CA, United States of America
¹⁵⁴ Physics Department, Royal Institute of Technology, Stockholm, Sweden
¹⁵⁵ Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, United States of America
¹⁵⁶ Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
¹⁵⁷ School of Physics, University of Sydney, Sydney, Australia
¹⁵⁸ Institute of Physics, Academia Sinica, Taipei, Taiwan
¹⁵⁹ ^(a) E. Andronikashvili Institute of Physics, Iv. Javakhsishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
¹⁶⁰ Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
¹⁶¹ Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
¹⁶² Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
¹⁶³ International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
¹⁶⁴ Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
¹⁶⁵ Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
¹⁶⁶ Tomsk State University, Tomsk, Russia
¹⁶⁷ Department of Physics, University of Toronto, Toronto, ON, Canada
¹⁶⁸ ^(a) TRIUMF, Vancouver, BC; ^(b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
¹⁶⁹ Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
¹⁷⁰ Department of Physics and Astronomy, Tufts University, Medford, MA, United States of America
¹⁷¹ Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States of America
¹⁷² Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
¹⁷³ Department of Physics, University of Illinois, Urbana, IL, United States of America
¹⁷⁴ Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia – CSIC, Valencia, Spain
¹⁷⁵ Department of Physics, University of British Columbia, Vancouver, BC, Canada
¹⁷⁶ Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
¹⁷⁷ Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany
¹⁷⁸ Department of Physics, University of Warwick, Coventry, United Kingdom
¹⁷⁹ Waseda University, Tokyo, Japan
¹⁸⁰ Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel
¹⁸¹ Department of Physics, University of Wisconsin, Madison, WI, United States of America
¹⁸² Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
¹⁸³ Department of Physics, Yale University, New Haven, CT, United States of America
¹⁸⁴ Yerevan Physics Institute, Yerevan, Armenia

- ^a Also at CERN, Geneva; Switzerland.
^b Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
^c Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
^d Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
^e Also at Departamento de Física, Instituto Superior Técnico, Universidade de Lisboa, Lisboa; Portugal.
^f Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah; United Arab Emirates.
^g Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
^h Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
ⁱ Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.
^j Also at Department of Physics, California State University, East Bay; United States of America.
^k Also at Department of Physics, California State University, Fresno; United States of America.
^l Also at Department of Physics, California State University, Sacramento; United States of America.
^m Also at Department of Physics, King's College London, London; United Kingdom.
ⁿ Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.
^o Also at Department of Physics, Stanford University, Stanford CA; United States of America.
^p Also at Department of Physics, University of Adelaide, Adelaide; Australia.
^q Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
^r Also at Department of Physics, University of Michigan, Ann Arbor MI; United States of America.
^s Also at Department of Physics, University of Toronto, Toronto ON; Canada.
^t Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow; Russia.
^u Also at Giresun University, Faculty of Engineering, Giresun; Turkey.
^v Also at Graduate School of Science, Osaka University, Osaka; Japan.
^w Also at Hellenic Open University, Patras; Greece.
^x Also at Institutio Catalana de Recerca i Estudis Avançats, ICREA, Barcelona; Spain.
^y Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
^z Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.
^{aa} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia; Bulgaria.
^{ab} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.
^{ac} Also at Institute of Particle Physics (IPP); Canada.
^{ad} Also at Institute of Physics, Academia Sinica, Taipei; Taiwan.
^{ae} Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
^{af} Also at Institute of Theoretical Physics, Ilia State University, Tbilisi; Georgia.
^{ag} Also at Instituto de Física Teórica, IFT-UAM/CSIC, Madrid; Spain.
^{ah} Also at Joint Institute for Nuclear Research, Dubna; Russia.
^{ai} Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.
^{aj} Also at Louisiana Tech University, Ruston LA; United States of America.
^{ak} Also at LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris; France.
^{al} Also at Manhattan College, New York NY; United States of America.
^{am} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.
^{an} Also at National Research Nuclear University MEPhI, Moscow; Russia.
^{ao} Also at Physics Department, An-Najah National University, Nablus; Palestine.
^{ap} Also at Physics Dept, University of South Africa, Pretoria; South Africa.

^{aq} Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.

^{ar} Also at School of Physics, Sun Yat-sen University, Guangzhou; China.

^{as} Also at The City College of New York, New York NY; United States of America.

^{at} Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.

^{au} Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

^{av} Also at TRIUMF, Vancouver BC; Canada.

^{aw} Also at Università di Napoli Parthenope, Napoli; Italy.

* Deceased.