### Title:

Rheological characterization of ethylcellulose-based melts for pharmaceutical applications

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### Running head:

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## 20 ABSTRACT

21 Rheological characterization of ethylcellulose (EC)-based melts intended for the production, via 22 micro-injection moulding (µIM), of oral capsular devices for prolonged release was carried out. 23 Neat EC, plasticized EC and plasticized EC containing solid particles of a release modifier (filler volume content in the melt around 30%) were examined by capillary and rotational rheometry tests. 24 25 Two release modifiers, differing in both chemical nature and particle geometry, were investigated. When studied by capillary rheometry, neat EC appeared at process temperatures as a highly viscous 26 27 melt with a shear-thinning characteristic that progressively diminished as the apparent shear rate 28 increased. Thus, EC as such could not successfully be processed via µIM. Plasticization, which 29 induces changes in the material microstructure, enhanced the shear-thinning characteristic of the 30 melt and reduced considerably its elastic properties. Marked wall slip effects were noticed in the 31 capillary flow of the plasticized EC-based melts, with or without release modifier particles. The 32 presence of these particles brought about an increase in viscosity, clearly highlighted by the 33 dynamic experiments at the rotational rheometer. However, it did not impair the material 34 processability. The thermal and rheological study undertaken would turn out a valid guideline for 35 the development of polymeric materials based on pharma-grade polymers with potential for new pharmaceutical applications of µIM. 36

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39 Keywords: rheology; micro-injection moulding; ethylcellulose; prolonged release

#### 41 INTRODUCTION

42 Well-established techniques used to process thermoplastic polymers, such as extrusion and injection moulding (IM), have gained over the years considerable interest in the pharmaceutical 43 44 field, as the number of papers recently published in the scientific literature and of drug products already on the market demonstrate [1-3]. In some applications, considering the weight, dimensions, 45 46 or dimension tolerances of molded parts, the manufacturing process should be regarded as micro-47 injection moulding (µIM) rather than IM. Several advantages have been highlighted with respect to 48 the use of IM techniques for pharmaceutical production. These are mainly relevant to the 49 improvement of manufacturing (cost-effective and eco-friendly processes, scalability, suitability for 50 continuous manufacturing) and technological/biopharmaceutical aspects (versatility of the product 51 design, patentability, achievement of solid molecular dispersions/solutions of the active ingredient) 52 [4]. Functional containers, namely gastroresistant and pulsatile-release capsular devices, were lately 53 prepared by µIM starting from polymers with peculiar properties, i.e. solubility at pH>5.5 and 54 erosion/dissolution after swelling, respectively [5-7]. For drug delivery purposes it would also be 55 interesting to develop new polymeric materials intended for the manufacturing by µIM of 56 prolonged-release capsular devices able to slowly release their contents over 12-24 h after oral 57 administration.

58 On the other hand, limitations may arise from processing pharma-grade materials, in 59 consideration of their rheological/thermomechanical behaviour in the melt state and of their thermal 60 or chemical stability. Moreover, issues may be associated with the active ingredient and the product 61 requirements, such as drug dose, release profile, stability and compliance. In this respect, a rational 62 design of new polymers, or of a new generation of traditional ones with improved thermal properties, has recently been attempted (e.g. Soluplus<sup>®</sup>, Affinisol<sup>™</sup>) [8]. Therefore, the need for a 63 64 thorough understanding of the thermal and rheological properties of pharma-grade polymeric melts 65 and of their impact on both processing conditions and physico-mechanical stability of the final products becomes of utmost importance [9]. 66

67 Based on such premises, the aim of this work was to investigate the rheological behaviour of 68 melts based on ethylcellulose (EC), chosen because it is the most widely employed polymer for the 69 preparation of prolonged drug delivery systems and, moreover, has already been processed in the 70 molten state [10-12]. Both neat and plasticized EC was studied. The plasticizer used was triethyl 71 citrate, the relevant effectiveness as a processing aid for EC-based materials being known from the 72 literature [13]. Besides, as preliminary experiments showed that these materials would fail to yield the pursued release performance, two permeability modifiers (a soluble and a swellable polymer) 73 74 were added to plasticized EC, based on their ability to create pores in the insoluble capsule wall. 75 The rheological behaviour of these systems, which at processing temperatures consisted of 76 suspensions of solid particles in the plasticized EC melt, was then evaluated. The characterization of 77 the EC-based melts was carried out by both capillary and rotational rheometry. Capillary rheometry 78 was employed to study the processability of the materials in view of their µIM, whereas rotational 79 rheometry experiments were performed to gain a more in-depth understanding of the viscoelastic 80 response.

81

## 82 MATERIALS AND METHODS

### 83 Materials

Ethylcellulose, EC (Ethocel<sup>™</sup> Std. 100 premium; Dow, United States); triethyl citrate, TEC
(Aldrich, Germany); KIR (Kollicoat<sup>®</sup> IR; BASF, Germany); low-substituted hydroxypropyl
cellulose, LHPC (L-HPC NBD 020; Shin-Etsu, Japan).

87

#### 88 Methods

89 The following EC-based materials were prepared:

90 - plasticized EC (pEC): EC powder was placed in a mortar and the liquid plasticizer, TEC (20% by

91 weight on the dry polymer), was manually added under hard mixing. The kneaded product was left

stand at least 12 h at room temperature (21±5°C, 55±5% RH) before it was ground by means of a
blade mill (La moulinette, Moulinex®, SEB, France).

94 - pEC filled with KIR or LHPC particles (pEC/KIR and pEC/LHPC, respectively): KIR powder ( $d_{50}$ 95 = 23µm;  $d_{90}$  = 55µm) or LHPC powder ( $d_{50}$  = 46µm;  $d_{90}$  = 129µm), previously desiccated at 40°C in 96 an oven for 24 hours, were mixed in a mortar with pEC in a 3:7 weight ratio. The mixed powders 97 were then stored into bags under vacuum, and characterized within one or two days.

98

#### 99 Thermal characterization

100 Thermal analysis on EC and pEC powders was performed on a calorimeter DSC1 STAR by 101 Mettler Toledo (Switzerland), using nitrogen as a purge gas (70 ml/min). Indium was used as a 102 calibration standard. Samples of about 10 mg were heated in aluminum crucibles from 30 to 200°C, 103 maintained at this temperature for 1 min, cooled down to 30°C and reheated up to 200°C. Both 104 heating and cooling steps were carried out at 5°C/min.

105

## 106 **Capillary rheometry experiments**

107 The tests were performed on a capillary rheometer Rheologic 5000 by Instron CEAST (Italy). 108 In this instrument, the flow rate is imposed by a piston and the extrusion pressure monitored by a 109 melt pressure transducer (data recorded every 0.15 s) placed in the barrel at a distance of 15 mm 110 from the die entry plane; a schematic of a piston velocity-controlled capillary rheometer is 111 reported in Figure 1a. Preliminary trials were carried out on EC to find suitable test conditions. 112 Long residence times in the melt state promoted chemical degradation processes accompanied by 113 gas formation. As such processes were thermally activated, the maximum test temperature was set 114 at 180°C. 175°C was selected as the reference test temperature; additional tests were carried out at: 115 165, 170 and 180°C on EC; 150, 165 and 170°C on pEC; 150 and 160°C on pEC/KIR and pEC/LHPC. EC could not be examined under 165°C due to the high viscosity of the melt. Flat 116 entrance circular dies made from Vanadis 30<sup>TM</sup> steel were employed. A die with diameter, d, of 2 117

118 mm and length to diameter ratio, 1/d, of 10 was used with all the materials. Additional dies were 119 employed, at 175°C, for: i) the investigation of wall slip effects, on pEC-based systems; ii) the 120 determination of die entrance pressure losses, on EC and pEC. More in detail, the additional dies 121 were: with reference to i), a die having d = 1 mm and l/d = 10; with reference to ii), dies having d =122 2 mm and l/d = 5 and 20, for EC analysis, and a die having d = 1 mm and l/d = 20, for pEC analysis. A barrel with diameter, D, of 15 mm was employed. The apparent shear rate at the die wall,  $\dot{\gamma}_{app,w}$ , 123 ranged between 17.8 and 560 s<sup>-1</sup> (with five different levels of  $\dot{\gamma}_{app,w}$  per decade, equally spaced on 124 a logarithmic scale, examined from the lowest to the highest).  $\dot{\gamma}_{app,w}$  is related to the imposed flow 125 rate, Q, and the die diameter, d, as follows: 126

127 
$$\dot{\gamma}_{app,w} = \frac{32 \cdot Q}{\pi \cdot d^3}$$
 (1)

Shear rates lower than 17.8 s<sup>-1</sup> were not considered because time to the attainment of the steady-128 129 state would have been too long and degradation phenomena could occur. The tests at 175°C with l/d 130 = 10 and d = 1 mm or 2 mm dies were carried out at least three times to check repeatability. 131 In every test performed, the material, after being loaded into the instrument barrel, was subjected to a pre-heating step over 420 s, under a maximum compacting force of 300 N. 132 133 All tests with the 2 mm diameter dies were conducted in such a way that the various melts experienced the same pre-fixed flow-history, in consideration of the possibility that the materials 134 135 underwent residence time-controlled chemical processes while in the melt state. More in detail, the period of time during which a material was subjected to a specific  $\dot{\gamma}_{app,w}$  (step) was fixed. 136 Preliminary tests allowed to set the duration of each step ( $\Delta t_{step}$ ), chosen in such a way as to 137 guarantee the reaching of flow conditions which could be considered representative of a steady-138 139 state for every material, at the different test conditions (temperatures and, for EC, also die lengths).

140  $\Delta t_{step}$  varied between 12 and 20 s, at the highest and the lowest  $\dot{\gamma}_{app,w}$  explored, respectively. With

141 1/d fixed, in a 1 mm die test  $\Delta t_{step}$  had to be higher than in a 2 mm die test, at a given  $\dot{\gamma}_{app,w}$ , due to 142 the lower absolute time scale of the compression process experienced by the melt in the reservoir 143 (with a 1 mm die the fluid is forced to enter into the capillary with a flow rate that is one eighth of 144 the flow rate used with a 2 mm die).

145

## 146 Rotational rheometry experiments

147 Measurements of the linear viscoelastic properties were carried out using the strain-controlled 148 rheometer ARES G2 by TA Instruments (Unites States). Dynamic frequency sweep experiments 149 under low amplitude (2%) oscillations were carried out with a parallel-plate geometry; a schematic 150 of a strain-controlled rotational rheometer with parallel-plate geometry, for oscillatory shear 151 measurement, is reported in Figure 1b. The diameter of the plates was 25 mm, and the gap was 152 set between 1.3 and 1.8 mm. Preliminary experiments, carried out to identify the adequate test 153 conditions, put in evidence that the results were generally dependent on the overall duration of the 154 test (practically determined by the lower limit of the frequency range explored). This should be 155 attributed to chemical degradation processes, the impact of which is expected to increase with the 156 residence time. Thus, the lower limit of the frequency range explored was set at a relatively high 157 value (10 rad/s). Frequency ranged from 10 to 500 rad/s; ten different frequencies, equally spaced 158 on a logarithmic scale (except 500 rad/s), were examined from the lowest to the highest. The tests 159 were performed on EC at four different temperatures, equally spaced between 160 and 175°C, and 160 on the pEC-based systems, at six different temperatures, equally spaced between 150 and 175°C. 161 Each test was carried out three times at each temperature, to check repeatability.

162 The materials were placed on the rheometer in the form of disks, with diameter of 25 mm, 163 prepared *via* compression moulding on a P200E press by Collin (Germany). During the 164 compression operation, the material was maintained for 280 s under the pressure of 32 bar, at 183°C 165 for EC and 145°C for the pEC-based systems.

### 167 Scanning Electron Microscopy, SEM, analyses

SEM analyses were carried out both on KIR and LHPC powders, and on the fracture surfaces of the extrudates of the filled pEC-based materials, obtained with the tests at the capillary rheometer.

171 With reference to the analyses on the powders, these were carried out on gold coated powders 172 in a scanning electron microscope (model Zeiss Evo 50 EP) operated in high vacuum mode. 173 With reference to the analyses on the extrudates, for each material examined, a portion of the 174 extruded filament was collected at the exit from the rheometer die during an experiment at 175°C, with the 2 mm die with 1/d = 10, at a  $\dot{\gamma}_{app,w}$  of 100 s<sup>-1</sup>. The filament was then broken under liquid 175 176 nitrogen, in such a way to obtain a fracture surface almost perpendicular to the direction of the flow 177 in the capillary, and the cryo-fractured surface was observed by a scanning electron microscope 178 (model LEO EVO 40). The surface, prior to be analysed, was sputtered with a thin layer of gold.

179

### 180 RESULTS AND DISCUSSION

### 181 Thermal response

182 Figure 2 shows the thermograms of EC and pEC obtained from the first and the second 183 heating scan, respectively. EC exhibits an endothermic peak, close to 190°C, whereas a similar 184 endothermic peak was not observed in pEC heating scans. This endothermic peak, which was 185 observed in both heating scans and thus related to a reversible thermal process, may be ascribed to 186 the melting of high order (crystalline) domains, as suggested by Davidivich-Pinhas et al. [14] with respect to the same EC grade. This would indicate that EC has a semi-crystalline nature that it 187 188 maintains up to around 190°C. Interestingly, based on the absence of this endothermic peak in the 189 pEC trace, it could be inferred that plasticization with TEC makes the polymer amorphous. The 190 glass transition of EC was clearly detected in the second heating scan only (glass transition 191 temperature, T<sub>g</sub>, around 130°C), because the thermal response in the first heating scan was 192 complicated by the very broad peak around 100°C probably relevant to water evaporation. A similar broad peak was found also in the first heating thermogram of pEC, the glass transition of whichcould not be identified even in the second heating scan.

195

### 196 **Capillary rheometry experiments**

#### 197 *Extrudate appearance*

198 Following rheometry tests on the EC melt, defect-free (at a macroscopic scale) extrudates were always obtained, except at the highest  $\dot{\gamma}_{app,w}$  explored (that is 560 s<sup>-1</sup>). At this apparent shear 199 200 rate, the presence of bubbles of size comparable to the extrudate cross-section was noticed, 201 irrespective of the working temperature. The formation of these bubbles, observed for the melt that stayed longer in the reservoir of the rheometer ( $\dot{\gamma}_{app,w}$  of 560 s<sup>-1</sup> corresponds to the highest 202 residence time in the rheometer), might be explained as the consequence of chemical degradation 203 204 processes [15]. This was supported by results obtained from infrared spectroscopy analyses (data 205 not reported).

pEC melt showed two distinct flow regimes within the range of  $\dot{\gamma}_{app,w}$  explored in the tests 206 carried out between 165 and 175°C: *i*) a stable steady flow regime, at low levels of  $\dot{\gamma}_{app,w}$ , and *ii*) 207 an unstable flow regime characterized by the occurrence of melt fracture phenomena, at high levels 208 of  $\dot{\gamma}_{app,w}$ . The transition from the stable to the unstable regime was detected from the extrudate 209 aspect, whereas no effects on the pressure signal trend occurred. The extrudate from the former 210 211 regime appeared undistorted, appreciably transparent and with a relatively smooth surface. In the 212 unstable regime, a highly deformed and semi-transparent extrudate, characterized by regular wavy distortions with a typical length around 3 mm, was obtained. Interestingly, the level of  $\dot{\gamma}_{app,w}$  at 213 the onset of melt fracture decreased by lowering the test temperature, from 560 s<sup>-1</sup> at 175°C to 178 214 s<sup>-1</sup> at 165°C. Extrudates at 150°C were different as compared to those obtained at higher 215

216 temperatures (distortions were in general more irregular and less marked) and appeared deformed in 217 the whole range of  $\dot{\gamma}_{app,w}$  explored.

With pEC/KIR and pEC/LHPC melts, at the highest  $\dot{\gamma}_{app,w}$  explored (that is 560 s<sup>-1</sup>), large 218 219 amplitude fluctuations in the pressure signal were recorded and acceptable steady-state extrusion 220 pressure data could not be measured. Fluctuations could not be ascribed to the occurrence of melt 221 fracture phenomena, since no flow instability could be recognized, and the processes at the basis of 222 their origin are not clear. Due to the presence of the KIR and LHPC particles, the extrudate 223 appeared opaque and with a considerably rough surface. Figures 3a and b show SEM images of 224 these particles. KIR particles have pseudo-spherical shape and smooth surface, whereas LHPC 225 particles appear as flakes or short fibres with rough surface. Figures 3c and d show higher 226 magnification SEM images of the fracture surfaces obtained from the pEC/KIR and pEC/LHPC 227 extrudate, respectively. It is worth underlining that both KIR and LHPC particles maintained their 228 solid state at the test temperatures. LHPC particles appear better wetted by the polymer with respect 229 to KIR ones. SEM images of the same surfaces at lower magnification (data not reported) showed 230 that both types of particles are uniformly dispersed within the polymeric matrix.

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#### 232 Apparent viscosity functions

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Figures **4**a-d show the steady-state apparent shear viscosity,  $\eta_{app}$ , *vs* apparent shear rate at the die wall,  $\dot{\gamma}_{app,w}$ , curves obtained for EC and the pEC-based melts at different temperatures and with the die having d = 2 mm and l/d = 10. At a given  $\dot{\gamma}_{app,w}$ ,  $\eta_{app}$  is evaluated as the ratio between the apparent shear stress at the die wall,  $\tau_{app,w}$ , and  $\dot{\gamma}_{app,w}$ , with  $\tau_{app,w}$  calculated from the steadystate extrusion pressure,  $\Delta P$ , uncorrected for the die entrance pressure drop,  $\Delta P_{ent}$ , as:

238 
$$\tau_{app,w} = \frac{\Delta P}{4 \cdot l/d}$$
 (2)

Data at 175°C are the average of three replicated measures (bars in figure represent the standard deviation; when not visible, they are shorter than the symbol height). Working with EC at 165°C,  $\dot{\gamma}_{app,w}$  higher than 178 s<sup>-1</sup> could not be determined since extrusion pressure exceeded the maximum capacity of the pressure transducer (50 MPa). Data indicated with open or thinner symbols and connected by dotted segments should be regarded as critical since melt fracture occurred (for pEC) or acceptable values of steady-state extrusion pressure could not be detected (for

EC and filled pEC-based melts, at 
$$\gamma_{app,w} = 560 \text{ s}^{-1}$$
)

246 The pronounced effect of the test temperature on the viscosity of EC melt is clearly shown in 247 Figure **4**a. Under the assumption that the relationship between  $\Delta P_{ent}$ , and the true shear stress at the 248 die wall,  $\tau_w$ , is temperature invariant [16], the time-temperature superposition principle was

tentatively applied to the  $\tau_{app,w}$  data obtained. Each logarithmic plot of  $\tau_{app,w} vs \dot{\gamma}_{app,w}$  referred to a specific temperature, T, was horizontally shifted to overlap as much as possible the reference curve at  $T_0 = 175^{\circ}$ C. Such a procedure permitted to evaluate, for each temperature, the shift factor,  $a_T$ , at constant  $\tau_{app,w}$ , averaged over the shear stress range explored. Good results with respect to the overlapping were generally obtained, and the shift factor turned out practically independent from the stress level within the whole shear stress range explored. This appears rather surprising taking into account the complex evolving micro-structure of the EC melt.

256 An Arrhenius type dependence on temperature was observed for a<sub>T</sub>, and a viscous flow 257 activation energy,  $E_{0.s}$ , of  $\approx 240$  kJ/mol was obtained (see [16] for Arrhenius equation) – the 258 subscript "s" is used in  $E_{0,s}$  to indicate that the activation energy is derived from steady-state flow 259 properties. By considering that, for an ordinary industrial polymer melt, the viscous flow activation 260 energy derived from fully developed flow tests is of the order of a few tens of kJ/mol and typically 261 keeps smaller than 100 kJ/mol [17,18], the value of  $E_{0,s}$  obtained for EC seems very high. The 262 pronounced sensitivity of apparent viscosity to temperature could be related to the thermal response 263 of EC. Due to the tendency to undergo chemical degradation processes at high temperatures, the test 264 temperatures have been kept relatively low, just below the endothermic peak ascribed to the melting 265 of the crystalline domains (see *Thermal response* section). At the test temperatures, EC melt is expected to have a heterogeneous microstructure with high order regions dispersed within an 266 267 amorphous matrix and, therefore, it can be argued that small temperature variations can induce 268 substantial changes in the material microstructure and in polymer chains mobility. These 269 temperature-induced microstructural modifications are likely to be reflected in the viscous response 270 obtained in the capillary rheometer, which varied widely with temperature. A contribution of 271 thermally activated chemical processes occurring during the rheometry test, which could reduce the 272 length of the polymeric chains and, therefore, the viscosity of the material, could not be ruled out. It 273 is worth noting that the phenomena described above do not explain the "apparent thermorheological 274 simplicity" observed for EC. A comprehensive analysis of the processes that govern the temperature 275 dependence of the viscosity here observed would require an extensive work aimed at studying the 276 chemical-physical processes undergone by EC during a capillary rheometer test, and it could be a 277 matter of further studies.

From a<sub>T</sub> data, the  $\eta_{app} vs \dot{\gamma}_{app,w}$  master curve of EC was constructed at T<sub>0</sub> = 175°C (data at 278 560 s<sup>-1</sup> have been removed). In Figure 4a, shifted data-points [( $\dot{\gamma}_{app,w} \cdot a_T; \eta_{app}/a_T$ ) where  $a_T$  is 279 280 temperature shift factor] are indicated with grey symbols, and dashed line is sketched as a guideline 281 for the eye. The EC melt shows a non-Newtonian behaviour, apparently shear-thinning, but 282 different from that of an ordinary polymer melt. It is characterized by a threshold level of apparent 283 viscosity,  $\eta_{app,\infty}$  ( $\eta_{app,\infty} \approx 670 \text{ Pa} \cdot \text{s}$ ), approached at high shear rates. This behaviour recalls that of 284 Bingham fluids [19], even if in this case the existence of a yield stress cannot be firmly stated on the 285 basis of the data obtained. The presence of a threshold at a relatively high level of viscosity clearly 286 indicates that EC can be very hardly processed via µIM, at least at process temperatures compatible 287 with its chemical stability. This was confirmed by explorative  $\mu$ IM tests, performed with a  $\mu$ IM 288 press (BabyPlast 6/10P; Cronoplast S.L., Barcelona, Spain) equipped with a disk-shaped screening

mould (the details of machine and mould used are reported in [7]). At process temperatures
comparable with those used in the capillary rheometer tests, EC could not be processed if unplasticized.

For comparison purposes, the flow curve (master curve at  $T_0 = 175^{\circ}$ C) of EC is also reported 292 in Figures 4b-d relevant to pEC-based melts. From Figure 4b it can be observed that the 293 294 plasticization with TEC makes the rheological response of pEC consistent with that of an ordinary 295 polymer melt. The trend of the flow curve does not indicate the existence of a threshold level of apparent viscosity and, at the highest  $\dot{\gamma}_{app,w}$  explored,  $\eta_{app}$  of pEC is remarkably lower than that of 296 297 EC (at 175°C). This suggests for the pEC melt the possibility to be processed via µIM. By 298 comparing the flow curves of pEC/KIR and pEC/LHPC, reported in Figures 4c and d, respectively, 299 with that of pEC, it emerges that the presence of a secondary phase in pEC, either KIR or LHPC particles, promotes an increase in  $\eta_{app}$ , at a given level of  $\dot{\gamma}_{app,w}$ , without impairing the shear-300 thinning nature of the response, and therefore the potential for the material to be processed at high 301 302 shear rates. Also these assumptions were confirmed by preliminary µIM trials performed as above 303 described.

304 As shown in Figures 4b-d, the test temperature leaves the apparent viscosity function of pECbased melts practically unchanged. This seems to suggest that the flow of these materials is not 305 306 governed by the shear viscosity of the melt, and that a more complex flow takes place. Moreover, for unfilled pEC the test temperature was already shown to influence the level of  $\dot{\gamma}_{app,w}$  at the 307 308 onset of melt fracture phenomena (see Extrudate appearance section), which are generally related 309 to a complex equilibrium between adhesion and slip of the polymer melt at the wall metal interface 310 [20,21]. As polymer/wall interface processes can play a role in defining the capillary flow of the 311 pEC-based melts, the occurrence of wall slip effects in the capillary flow of these melts was 312 verified.

313 Wall slip effects were studied by the application of the Mooney technique [22]. This approach 314 assumes that the fluid is homogeneous and the slip velocity at the die wall, v<sub>s</sub>, is a function of the wall shear stress,  $\tau_w$ , only. The apparent shear rate at the die wall,  $\dot{\gamma}_{app,w}$ , which the polymer melt 315 is subjected to during capillary flow at a given  $\tau_w$ , results from: (i.) a purely viscous bulk 316 contribution,  $\dot{\gamma}_{app,w,v}$ , that would only depend on the intrinsic behaviour of the material, and (ii.) a

wall slip contribution,  $\dot{\gamma}_{app,w,s}$ , that depends on v<sub>s</sub> and die diameter, d, according to Eq. 3: 318

317

319 
$$\dot{\gamma}_{app,w,s} = \frac{8}{d} \cdot v_s(\tau_w)$$
 (3)

Tests with different die diameters are carried out and  $\tau_w vs \dot{\gamma}_{app,w}$  curves are obtained.  $v_s$  and 320  $\dot{\gamma}_{app,w,v}$  can thus be determined from the slope and the intercept, respectively, of the least square 321 regression straight line that interpolates the  $\dot{\gamma}_{app,w}$  vs 1/d data (Mooney plot), at a given level of  $\tau_w$ . 322 In this work, under the assumption that  $\Delta P_{ent}$ , do not significantly change with the die diameter 323 324 (supported by the experimental observation that for pEC the levels of  $\Delta P_{ent}$  measured with the 325 highest contraction ratio, D/d, available are very small, as shown below, in the Entrance pressure *losses* section), extrusion pressure data uncorrected for  $\Delta P_{ent}$  were used, and  $\tau_{app,w}$  rather than  $\tau_w$ 326 327 data were employed. By considering that slip phenomena can depend in principle also on the 328 extrusion pressure [23], the dies for wall slip analysis were chosen in such a way that similar levels 329 of steady-state extrusion pressure were reached in the experiments. More specifically, data from 330 tests carried out, at 175°C, with dies having different diameter (d = 1 and 2 mm) but same l/d (l/d = 10) were used. In Figure 5 the  $\tau_{app,w} vs \dot{\gamma}_{app,w}$  curves obtained with the two different dies for pEC, 331 pEC/KIR and pEC/LHPC are compared (bars in figure represent the standard deviation; when not 332 visible, they are shorter than the symbol height). The data obtained with the 2 mm die at 560 s<sup>-1</sup> 333 were not taken into account; for pEC, also the datum obtained with the 1 mm die at 560 s<sup>-1</sup> was 334

335 omitted since melt fracture was observed. The curve obtained with the 2 mm die is above the 1 mm die curve over the whole range of  $\dot{\gamma}_{app,w}$ , independent of the material, thus indicating that slip 336 337 effects occur. It is worth noting that, contrary to what generally observed for melts based on 338 ordinary polymeric materials exhibiting wall slip in capillary flow [23], at the highest levels of  $\tau_{app,w}$ 339 (that is of  $\Delta P$ ) the curves obtained with the different dies tend to get closer. For pEC, they even overlap at the highest  $\dot{\gamma}_{app,w}$ , indicating that, for this material, high levels of  $\Delta P$ , close to the onset 340 of melt fracture phenomena, lead to slipping suppression. Considering that (i.)  $\dot{\gamma}_{app,w}$  data from the 341 342 curves obtained with the two different die diameters, at a given  $\tau_{app,w}$ , refer to different residence 343 times in the rheometer (longer for the 1 mm die), and that (ii.) the linearity of the Mooney plot, corresponding to a given  $\tau_{app,w}$ , could not be proven since only two diameters were used, only slip 344 345 velocity was determined and referred to as an "apparent slip velocity" (indicated with v<sub>s,app</sub>). In Figure 6,  $v_{s,app}$  is reported as a function of  $\tau_{app,w}$  for pEC and the filled pEC-based systems, in a 346 347 double-logarithmic plot. For pEC, the downward arrow indicates the tendency of slipping 348 phenomena to be suppressed at high levels of  $\tau_{app,w}$ . As the dependence of  $v_{s,app}$  on  $\tau_{app,w}$  within the explored range of  $\tau_{app,w}$  appeared of power-law-type (linear in a double-logarithmic plot), as for 349 350 ordinary polymer melts, with particular reference to polyolefins [23,24], data were accordingly 351 fitted. In Figure 6, power-law fitting curve (dashed line) is traced for each series of data (for pEC 352 the point at  $\tau_{app,w} = 76770$  Pa, indicated with an open symbol, was not included in the fitting). The 353 order of magnitude of the slip velocity obtained for the pEC-melts is comparable with that of 354 polyolefin-based melts, but the levels of shear stress at which wall slip phenomena occur are 355 generally lower [23,25]. With respect to the unfilled pEC, the presence of a secondary phase, in 356 pEC/KIR and pEC/LHPC, promotes a decrease in  $v_{s,app}$ , at a given level of  $\tau_{app,w}$ , but enlarges the 357  $\tau_{app,w}$  range of slip, allowing high values of  $v_{s,app}$  to be reached. This latter effect is due to the 358 different sensitivity of  $v_{s,app}$  on  $\tau_{app,w}$ , exhibited by pEC and the filled pEC-based melts, here

359 represented by the slope of the power-law fitting line (that is 2.0 for pEC, 3.4 and 4.0 for pEC/KIR 360 and pEC/LHPC, respectively). Most of the flow occurs by wall slip, and the slip contribution to the flow rate decreases by increasing the shear stress level; at the lowest level of  $\tau_{app,w}$ , for the capillary 361 362 with the highest surface/volume ratio (that is the 1 mm die), it is approximately the 90, 80 and 70% of the imposed flow rate for pEC, pEC/KIR and pEC/LHPC melt, respectively. It can be thought 363 364 that the pronounced wall slip effects of pEC-based melts are due to the presence of the plasticizer, 365 TEC, which, during the capillary flow, could migrate at the die wall. This hypothesis should be 366 confirmed with specific analyses, which could be a matter of further studies.

367

#### 368 Entrance pressure losses

Figures 7a and b show Bagley plots constructed from the steady-state extrusion pressure,  $\Delta P$ , data measured at various  $\dot{\gamma}_{app,w}$ , at 175°C, for EC and pEC, respectively; for EC,  $\Delta P$  data were measured in tests carried out with dies having diameter of 2 mm (data at 560 s<sup>-1</sup> were not included in the analysis) whereas, for pEC, in tests with 1 mm dies. At 1/d = 10, each datum is the average of three values measured in replicated experiments (standard deviation bars indicated).

In Bagley plot, the experimental (l/d;  $\Delta P$ ) data at a given  $\dot{\gamma}_{app,w}$  typically lie on a straight 374 line, and  $\Delta P_{ent}$  value is determined by extrapolating to zero the least square regression line that 375 interpolates them. From Figure 7a, it clearly emerges that, for EC, at each level of  $\dot{\gamma}_{app,w}$  explored, 376 377 the dependence of  $\Delta P$  on the die length to diameter ratio, 1/d, is not linear. This curvature of the Bagley plots cannot be related to the effect of pressure on viscosity, nor to slipping phenomena, 378 379 since these factors would induce an upward shift of curves [23]; it rather might be ascribed, at least for the highest levels of  $\dot{\gamma}_{app,w}$ , to viscous heating effects, mainly occurring in the longest die (that 380 381 is the die with l = 40 mm), where they are expected to be more marked [26,27]. In support of this 382 discussion is the observation that EC viscosity is very sensitive to temperature, and a small

temperature increase produced by viscous heating can give rise to a pronounced decrease in viscosity (and therefore in  $\Delta P$ ). In accordance with these results, for the determination of  $\Delta P_{ent}$  of EC, only data obtained with the two shorter dies (with 1/d = 5 and 10), and indicated in Figure 7a wih full symbols, were used.

387 For the pEC melt, too small  $\Delta P_{ent}$  values were obtained by applying the Bagley procedure to data from tests performed with 2 mm dies. When very small,  $\Delta P_{ent}$  values can turn out comparable 388 389 with the statistical error induced by the extrapolation procedure, thus becoming not reliable. 390 Therefore, it is necessary to resort to a test flow geometry that could induce higher level of  $\Delta P_{ent}$  to be reached. By considering that at a given  $\dot{\gamma}_{app,w}$ , the higher the reservoir to die diameter ratio, 391 D/d, the higher is  $\Delta P_{ent}$  (with D fixed) [28], 1 mm dies were used for  $\Delta P_{ent}$  evaluation, and more 392 reliable  $\Delta P_{ent}$  data could be obtained except for the lowest  $\dot{\gamma}_{app,w}$ . In consideration of the fact that 393 394 pEC shows wall slip phenomena that can be pressure dependent, for the determination of  $\Delta P_{ent}$  only 395 data obtained with the two shortest dies available (with l/d = 10 and 20) were used (Figure 7b). For both EC and pEC, at each  $\dot{\gamma}_{app,w}$  examined, the value of  $\Delta P_{ent}$  was determined as the 396 397 intercept of the straight line traced on the two (l/d;  $\Delta P$ ) valid points (dashed lines in Figures 7a and b), and, in turn, the true shear stress at the die wall,  $\tau_w$ , was evaluated. Figure 7 clearly shows that 398 399 the levels of  $\Delta P_{ent}$  of EC are considerably higher than those of pEC (for EC, at 316 s<sup>-1</sup>,  $\Delta P_{ent}$  is higher than 5 MPa, whereas for pEC, at 560 s<sup>-1</sup> it does not even reach 1 MPa). In Figure 8, the ratio 400 between  $\Delta P_{ent}$  and  $\tau_w$  is plotted as a function of  $\dot{\gamma}_{app,w}$ , for both EC and pEC, and power-law fitting 401 402 curve (dashed line) is traced for each series of data. It clearly emerges that, in the whole range of  $\dot{\gamma}_{app,w}$ , the level of  $\Delta P_{ent}/\tau_w$  obtained for EC is largely higher than that obtained for pEC. By 403 404 considering that, quoting Jerman and Baird [29], "values of  $\Delta P_{ent}/\tau_w$  are thought to represent the 405 elasticity of a polymer melt", this result indicates that plasticization with TEC considerably reduces 406 the elastic character of EC in the melt state. This is further confirmed by the swell behaviour

407 exhibited by the two extrudates: the diameter of the EC extrudate (evaluated after material 408 solidification on samples obtained from tests carried out at 175°C, with the 2 mm die with 1/d = 10) 409 was larger than that of the pEC one. The difference between the two materials increased by 410 increasing the level of  $\dot{\gamma}_{app,w}$  undergone by the melt in the capillary flow (at 316 s<sup>-1</sup>, the diameter 411 of EC extrudate was nearly twice that of pEC). The pronounced elastic character exhibited by the 412 EC melt is probably another reflection, at a macroscopic scale, of its peculiar microstructure.

413

#### 414 **Rotational rheometry experiments**

415 Figures 9a and b show the absolute value of the complex viscosity,  $|\eta^*|$ , and the dynamic 416 storage modulus, G', vs frequency,  $\omega$ , master curves of the various materials, constructed from the 417 data obtained at the various temperatures examined (reference temperature,  $T_0 = 175^{\circ}$ C). Each 418 datum reported in the graphs is the average of three values measured in replicated experiments (bars 419 representing the standard deviation are not reported in the graphs as they were generally shorter 420 than the symbol height). Differently from that observed with the capillary rheometer tests, at the 421 rotational rheometer also for the pEC-based melts a clear effect of temperature was observed. This 422 is probably due to the fact that the shear stress levels reached in these latter tests, remarkably lower 423 than those reached in the capillary experiments, are insufficient to promote slip phenomena. The 424 successful attainment of a master curve for each system examined by shifting the rheological data 425 with respect to frequency, on the basis of the temperature shift factors, a<sub>T</sub>, demonstrates the validity 426 of the time-temperature superposition principle [already demonstrated for the EC melt by referring 427 to the capillary rheometer results (see *Capillary rheometry experiments*, *Apparent viscosity* 428 *functions* section)]. An Arrhenius type dependence on temperature was observed for  $a_T$ , and a flow 429 activation energy,  $E_{0,d}$ , was obtained for each melt – the subscript "d" is used in  $E_{0,d}$  to indicate that 430 this flow activation energy is obtained from flow properties measured in dynamic tests. E<sub>0,d</sub> values

- 431 are reported in Table I. For the EC melt, an appreciable difference between the flow activation
- 432 energy calculated from capillary and rotational rhometer tests was obtained. This difference should

be related to the different mechanical histories experienced by the polymer melt in the two rheometers.  $E_{0,d}$  of pEC is largely lower than that of EC. In consideration of the differences between the microstructures of these two melts (see *Thermal response* section), this result seems to support the idea that the high sensitivity to temperature exhibited by EC might be ascribed to the presence of high order domains in its melt microstructure, that were not found in the pEC melt. Finally, the incorporation of either KIR or LHPC particles in the pEC induces an increase in  $E_{0,d}$ .

439 From Figure 9 it emerges that the viscoelastic behaviour of the pEC-based melts is that 440 typical of homogeneous polymer melts [30]. With respect to the neat pEC melt, the presence of a 441 secondary phase in pEC/KIR and pEC/LHPC melts induces a significant increase in both G' and  $|\eta^*|$ , within the explored range of frequency. This result is in agreement with the effects reported in 442 443 literature for polymer melts filled with rigid particles [19,31], indicating that KIR and LHPC 444 particles enhance the capacity of the pEC melts to store elastic energy and dissipate mechanical 445 energy. Interestingly, LHPC particles appear the most effective in promoting an increase in both dynamic properties. By considering for pEC melt a density of 1.01 g/cm<sup>3</sup> (measured by melt flow 446 447 rate tests carried out by a Modular Melt Flow Indexer by Ceast-Instron, at 170°C, with an extruding 448 mass of 10 kg, according to ISO 1133-1:2011), and for KIR and LHPC particles density values of 449 0.94 and 1.29 g/cm<sup>3</sup>, respectively (evaluated from volume and mass data collected from three 450 compacts prepared by means of a hydraulic press equipped with flat punches), 32 and 25% of actual 451 filler volume contents in pEC/KIR and pEC/LHPC melts were estimated. This result indicates that 452 the higher effectiveness exhibited by LHPC particles cannot be ascribed to a higher volume content 453 of rigid particles in the melt, but it is intrinsically associated with the nature of these particles, 454 which differ from those of KIR in geometry, surface roughness and chemical composition. 455 Interestingly, by considering that LHPC particles (i.) promote an increase in G' and  $|\eta^*|$  that is practically independent of the frequency, and (ii.) leave the out-of-phase angle,  $\delta$  (tan[ $\delta$ ]=G"/G'), 456 457 practically unchanged (data of  $\delta$  not shown here), with respect to the unfilled melt, the viscoelastic 458 response of pEC/LHPC seems to be governed by the matrix. By contrast, for pEC/KIR melt, the

459 frequency dependence of the dynamic properties and the values of the out-of-phase angle obtained, 460 which are different from those of the matrix, suggest that KIR particles embedded in the polymer 461 melt actively contribute to its viscoelastic response.

462 The viscoelastic behaviour of the EC melt differs considerably from that of the pEC-based 463 melts and this might be ascribed to its particular microstructure. In Figure 10, the shear viscosity,  $\eta$ , obtained from capillary rheometer tests is compared with the complex viscosity,  $|\eta^*|$ , obtained from 464 465 rotational rheometer tests (viscosity data measured at 175°C). n data are determined from data of true shear stress (corrected for  $\Delta P_{ent}$ ),  $\tau_w$ , and true shear rate (calculated according to the 466 Rabinowitsch–Weissenberg method [30]),  $\dot{\gamma}$ , both evaluated at the die wall, and are plotted as a 467 function of  $\dot{\gamma}$ . It emerges that the Cox-Merz rule [30], according to which an approximate 468 agreement exists between  $|\eta^*(\omega)|$  and the steady-state shear viscosity,  $\eta(\dot{\gamma})$ , if the angular 469 frequency,  $\omega$  in [rad/s], is set equal to the shear rate,  $\dot{\gamma}$  in [s<sup>-1</sup>], is not verified. In fact, within the 470 range of  $\dot{\gamma}$  explored, the value of  $|\eta^*|$  is roughly 2.5 times that of  $\eta$ . Even by replotting the complex 471 472 viscosity data as a function of the maximum (or "effective") shear rate experienced in the dynamic 473 tests, that is  $\omega$  multiplied by the amplitude of the oscillation (0.02 in these experiments), as 474 proposed for filled melts by Doraiswamy et al. [32], steady-state and complex viscosity data do not 475 overlap. The mechanical history experienced by the melt in the two types of rheometer is different 476 with respect to: (i.) the type of flow (a steady-state pressure-induced shear flow in the capillary 477 rheometer, and a dynamic drag-induced shear flow under low amplitude oscillation in the rotational 478 rheometer); (ii.) the residence time (a capillary rheometer test is roughly two times longer than a 479 rotational rheometer test). In addition, the material tested at the rotational rheometer underwent a 480 compression moulding operation prior to be tested. These differences are likely to be at the basis of 481 the differences between viscosity and flow activation energy data obtained with the two types of 482 test. More specifically, with reference to the viscosity, it is reasonable to think that data coming 483 from the capillary rheometer refer to a melt characterized by a higher degree of microstructural

484 orientation. In spite of the difference between the values of  $\eta$  and  $|\eta^*|$ , the dependence on the timevariable (either  $\dot{\gamma}$  or  $\omega$ ) observed for the two different sets of viscosity data is similar, and may 485 suggest that the shear-thinning characteristic of the melt is being progressively reduced as  $\dot{\gamma}$  (or  $\omega$ ) 486 487 increases. Interestingly, the trend of the viscosity function obtained from the rotational rheometer 488 tests does not suggest the existence of a threshold viscosity for EC melt. This result seems to 489 indicate that the threshold noted for the apparent viscosity ( $\eta_{app,\infty}$  in Figure 4a), coming from 490 capillary rheometer tests, originates in part from the high levels of die entrance pressure losses that this melt exhibits at high  $\dot{\gamma}_{app,w}$  (it is worth recalling that  $\eta_{app}$  data were evaluated from data of 491 apparent shear stress,  $\tau_{app,w}$ , uncorrected for  $\Delta P_{ent}$ ). 492

493

#### 494 **CONCLUSION**

In this work, a rheological characterization of ethylcellulose (EC)-based melts intended to be
used for the production *via* µIM of prolonged-release capsular devices was carried out. Neat EC,
EC plasticized with TEC (pEC), and pEC containing particles of either a water soluble polymer
(KIR) or a hydrophilic swellable polymer (LHPC), added to enhance the permeability of the ECbased materials in view of their final use, were examined by both capillary and rotational rheometry
tests.

501 Capillary rheometry tests showed that this peculiar melt, which is expected to have a 502 heterogeneous microstructure with high order regions dispersed within an amorphous matrix, 503 exhibits a non-Newtonian, apparently shear-thinning behaviour characterized by a threshold in the 504 steady-state apparent shear viscosity,  $\eta_{app}$ , approached at high flow rates. The existence of this  $\eta_{app}$ 505 threshold indicates that neat EC cannot be practically processed *via* µIM. From both capillary and 506 rotational rheometry experiments EC melt shows a very high sensitivity of viscosity to temperature. 507 Capillary rheometry tests suggest that plasticization with TEC makes EC (pEC) processable 508 *via*  $\mu$ IM. For the pEC melt, which has a homogeneous amorphous structure, the threshold in the  $\eta_{app}$  509 was not observed, and a  $\eta_{app}$  function consistent with that of an ordinary polymer melt was 510 obtained. Plasticization with TEC considerably reduces the elastic character of EC, and fosters the 511 occurrence of pronounced wall slip phenomena in the capillary flow. Also for the pEC melts filled 512 with the permeability modifiers, either KIR or LHPC particles (filler volume content around 30%), 513 most of the capillary flow occurs by wall slip. The presence of these particles in the pEC melt 514 induces an increase in  $\eta_{app}$ , which however does not compromise the material shear-thinning behaviour, and therefore its processability. Rotational rheometry experiments clearly indicate that, 515 516 in the frequency range explored, both LHPC and KIR particles induce an increase in both the 517 complex viscosity,  $|\eta^*|$ , and the dynamic storage modulus, G', with LHPC particles being the most 518 effective at promoting this increase. KIR particles seem to actively contribute to the frequency 519 dependency of the dynamic properties, whereas the frequency response of the pEC melt filled with 520 LHPC particles seems to be governed by the matrix.

521 The study performed on EC turned out to be a valuable guideline for the development of
 522 polymeric materials with potential for new pharmaceutical applications of μIM.

523

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637	Dear Editor,
638	
639	we would like to thank the Reviewers for appreciating our work and providing helpful
640	comments.
641	Please find below an outline explaining how such comments have been addressed.
642	
643	Reviewer #1
644	The current manuscript is a lucid read and clearly explains each phenomenon observed while using
645	rotational and capillary rheometry experiments and extending the results and confirming the same
646	to understand the processability of the polymers during thermal extrusion. I recommend to accept
647	the manuscript in its current form.
648	
649	We kindly thank the reviewer for his comments.
650	
651	Reviewer #2
652	Authors are reporting a study of rheological characterization of molten polymer systems using
653	capillary and rotational rheometers. It would be helpful for an average reader if authors could add
654	a schematic figure comparing the measurement principle and sample placement with these two
655	methods.
656	
657	We followed the suggestion of the reviewer by introducing a new Figure 1 in the manuscript
658	showing the measurement principle and sample placement in both capillary and rotational
659	rheometers. The main text of the manuscript has been changed accordingly (in the revised
660	manuscript Figure 1 has been mentioned at lines 110 and 149, and all the other figures have been

661 renumbered).

Material	E <sub>0,d</sub> [kJ/mol]
EC	335
pEC	120
pEC/KIR	140
pEC/LHPC	150

Table I – Flow activation energies,  $E_{0,d}$ , of melts from dynamic tests at the rotational rheometer.

#### **Figure captions**

Figure 1 – Schematic of (a) a capillary rheometer (piston velocity-controlled) and (b) a rotational rheometer with parallel-plate geometry (strain-controlled, for oscillatory shear measurement). With reference to (a):  $v_p$ , piston velocity; P, resulting pressure in the reservoir. With reference to (b):  $\theta$ , angular displacement;  $\omega$ , oscillation frequency; M, resulting torque.

Figure 2 - DSC traces of EC and pEC obtained from the first and the second heating scans (thermograms vertically shifted for clarity).

Figure **3** – SEM images of (a) KIR and (b) LHPC particles, and of cryo-fracture surfaces obtained from the extrudates of (c) pEC/KIR and (d) pEC/LHPC.

Figure **4** – Steady-state apparent shear viscosity,  $\eta_{app}$ , *vs* apparent shear rate at die wall,  $\dot{\gamma}_{app,w}$ , for (a) EC, (b) pEC, (c) pEC/KIR, (d) pEC/LHPC, from capillary rheometry tests at: 180°C, *circle* (a); 175°C, *square* (a-d); 170°C, *triangle* (a,b); 165°C, *rhomb* (a,b); 160°C, "×" (c,d); 150°C, *asterisk* (b-d).

Figure 5 – Apparent shear stress at the die wall,  $\tau_{app,w}$ , as a function of the apparent shear rate at the die wall,  $\dot{\gamma}_{app,w}$ , for pEC (*square*), pEC/KIR (*circle*), pEC/LHPC (*triangle*), from capillary rheometry tests at 175°C, with d = 2 mm (full symbols) and d = 1 mm (open symbols).

Figure **6** – "Apparent wall slip velocity",  $v_{s,app}$ , as a function of the apparent shear stress at the die wall,  $\tau_{app,w}$ , for pEC (*square*), pEC/KIR (*circle*), pEC/LHPC (*triangle*), from capillary rheometry tests at 175°C.

Figure 7 – Bagley plots for (a) EC and (b) pEC, from capillary rheometry tests carried out with dies having d = 2 mm for EC (a) and d = 1 mm for pEC (b), at 175°C. For EC (a), data-points at l/d = 20 (open symbols) not used for  $\Delta P_{ent}$  determination.

Figure 8 – Ratio between the die entrance pressure drops,  $\Delta P_{ent}$ , and the true shear stress at the die wall,  $\tau_w$ , as a function of the apparent shear rate at the die wall,  $\dot{\gamma}_{app,w}$ , for EC (*rhomb*) and pEC (*square*), from capillary rheometry tests at 175°C.

Figure **9** – (a) Absolute value of the complex viscosity,  $|\eta^*|$ , and (b) dynamic storage modulus, G', *vs* frequency,  $\omega$ , master curves (T<sub>0</sub> = 175°C) of EC (*rhomb*), pEC (*square*), pEC/KIR (*circle*) and pEC/LHPC (*triangle*), from rotational rheometry experiments.

Figure 10 – Steady-state shear viscosity,  $\eta$ , *vs* shear rate (at the die wall),  $\dot{\gamma}$ , curve from capillary rheometry tests (open symbols) compared with the absolute value of the complex viscosity,  $|\eta^*|$ , *vs* frequency,  $\omega$ , master curve from rotational rheometry tests (full symbols) for EC, at 175°C.



















