

# BCS-BEC Crossover in a Two Dimensional Fermi Gas

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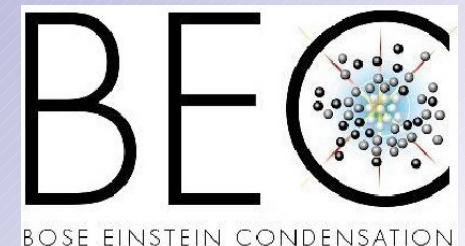
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INO-BEC, Trento, Italy

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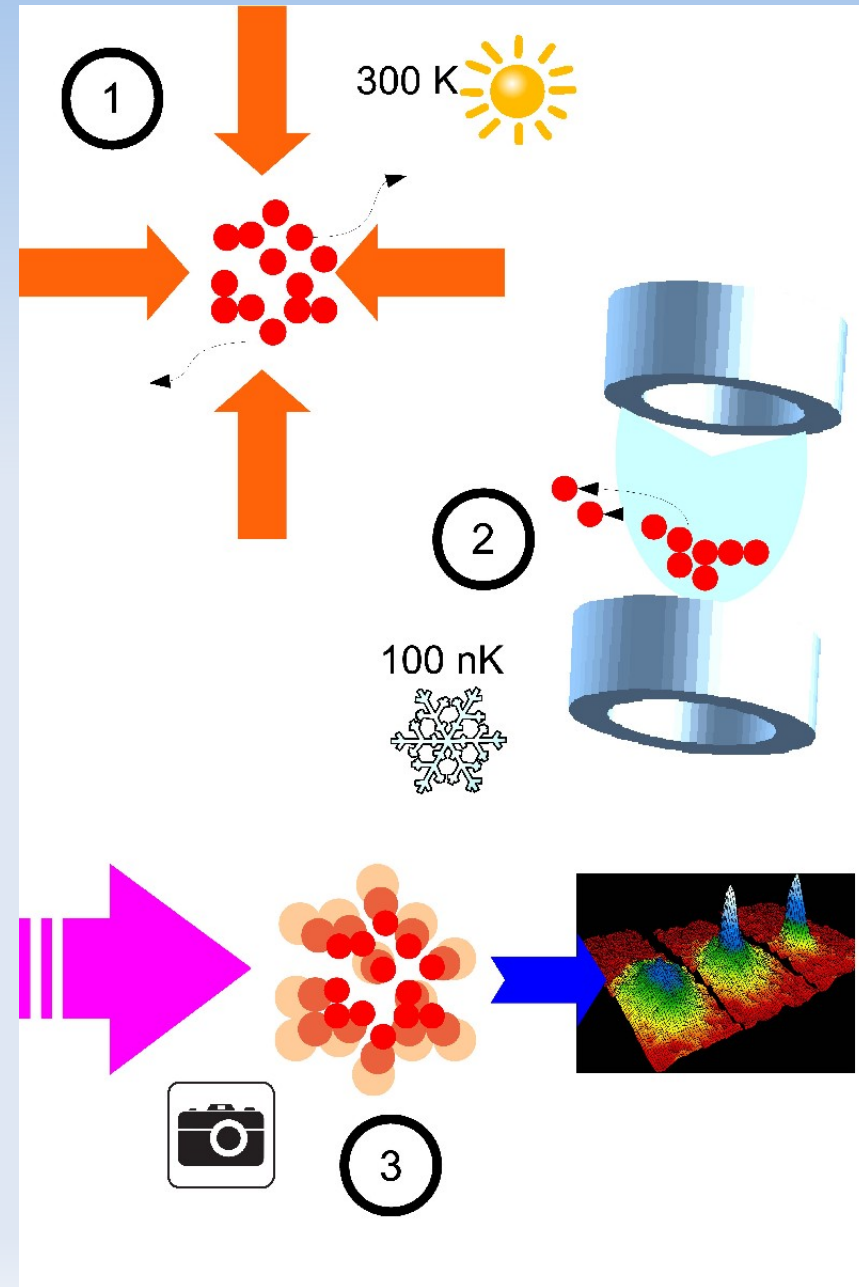


# Outline

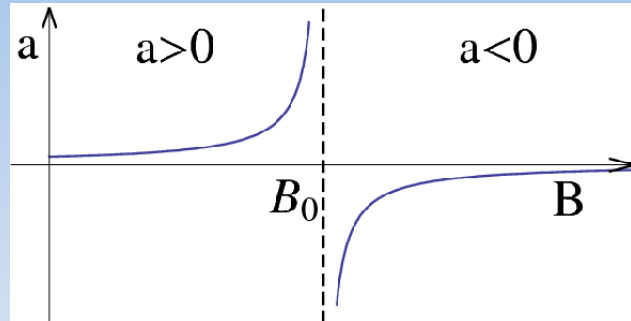
- Ultracold Fermi gases and BCS-BEC crossover
- Quantum Monte Carlo techniques
- BCS-BEC crossover in two dimensions

# Ultracold Fermi gases

- Short range interaction  $R$
- Diluteness  $R \ll l$
- Low temperature  $R \ll \lambda_T$
- 1) Laser cooling of alkali atoms ( ${}^6\text{Li}, {}^{40}\text{K} \dots$ )
- 2) Magneto-optical cooling and trapping (two hyperfine states)
- Quantum degeneracy  $n \lambda_T^d > 1$
- 3) Imaging: density and spin density

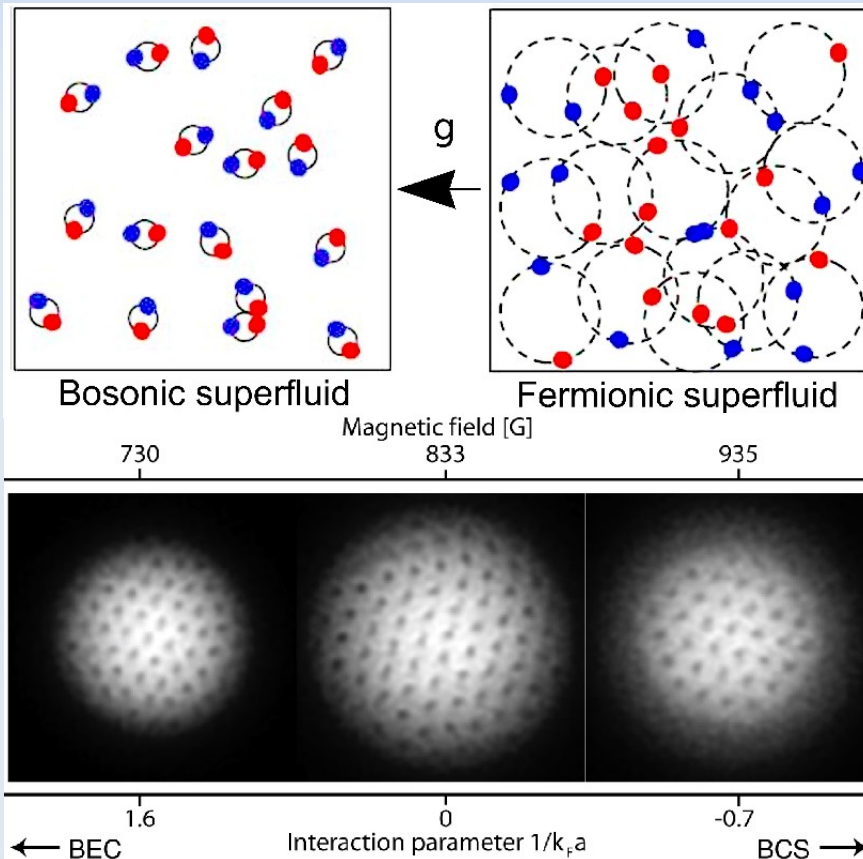


# BCS-BEC crossover



2 species of fermions with attractive interaction at  $T=0$  (Feshbach resonance)

- Weak coupling:  $E/N \sim \alpha \varepsilon_F$   
Cooper instability
- Strong coupling:  $E/N \sim \frac{\varepsilon_b}{2}$   
Condensate of dimers
- Crossover:  
no phase transition and no small parameter for perturbation theory  
→ QMC



Ketterle group, MIT (3D system)

# Diffusion Monte Carlo (DMC)

Schroedinger equation in imaginary time

$$-\frac{\partial}{\partial \tau} \Psi = H \Psi$$

$$\Psi(x, \tau) = e^{-H \tau} \Psi_T(x)$$

Eigenfunctions  $\psi_n(x, \tau) = e^{-E_n \tau} \psi_n(x)$

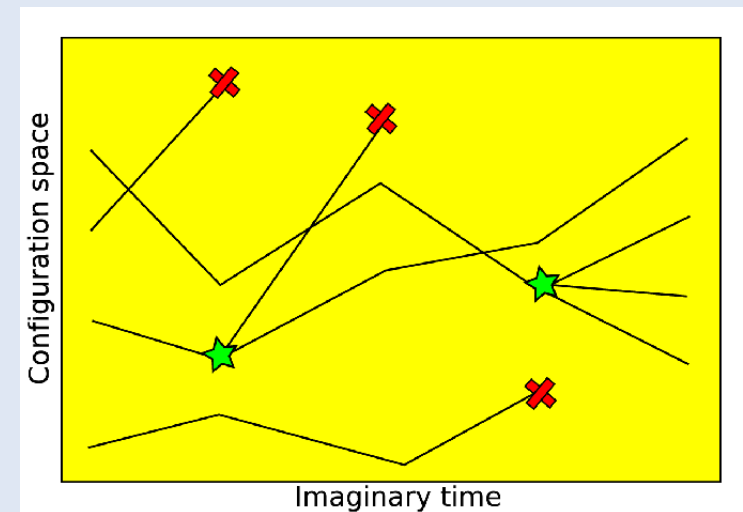
$$\text{Energy } \langle H \rangle = \frac{\langle \psi_0 | H | \Psi_T \rangle}{\langle \psi_0 | \Psi_T \rangle} = \frac{\int \psi_0(x) \Psi_T(x) \frac{\langle x | H | \Psi_T \rangle}{\Psi_T(x)} dx}{\int \psi_0(x) \Psi_T(x) dx} = \int g(x) E_L(x) dx$$

$\Psi_T$  Trial (guide) wavefunction

*Basic Algorithm (small time-step expansion)*

Random walk + Drift

Branching

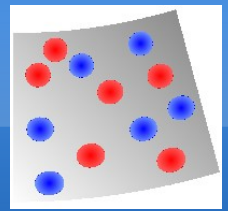


**Bosons** : *exact (ground state)*

**Fermions** : *Fixed Node approx. (variational principle)*  $\Psi(x) \Psi_T(x) \geq 0$

# BCS-BEC crossover in 2D

G. Bertaina and S. Giorgini, Phys. Rev. Lett. 106, 110403 (2011).

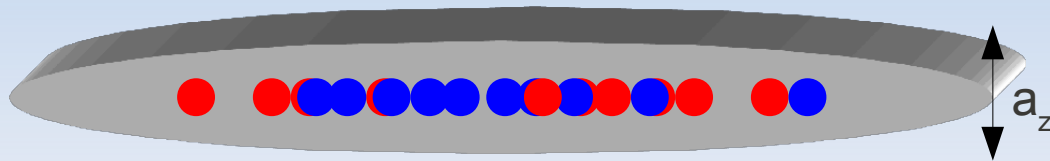


Optically trapped gas: Quasi-2D  $\epsilon_F \ll \hbar \omega_z$

Mapping: trapped 3D  $\rightarrow$  Pure 2D

Experiments (2010): Turlapov, Koehl

$$a_{2D} \propto a_z \exp\left(-\sqrt{\frac{\pi}{2}} \frac{a_z}{a_{3D}}\right)$$



Bound state  $\epsilon_b \propto -\frac{\hbar^2}{m a_{2D}^2}$

Model hamiltonian: Square well interaction (in universal regime  $nR^d \ll 1$ )

**Trial Wavefunctions (used to fix the nodal surface in DMC):**

- Weak coupling:

$$\Psi_{JS}(x) = J_{\uparrow\downarrow} D_{\uparrow}(N_{\uparrow}) D_{\downarrow}(N_{\downarrow})$$

Jastrow factor:

$$J_{\uparrow\downarrow} = \prod_{ii'} f_{\uparrow\downarrow}(x_{ii'}) \quad \text{f: two-body problem}$$

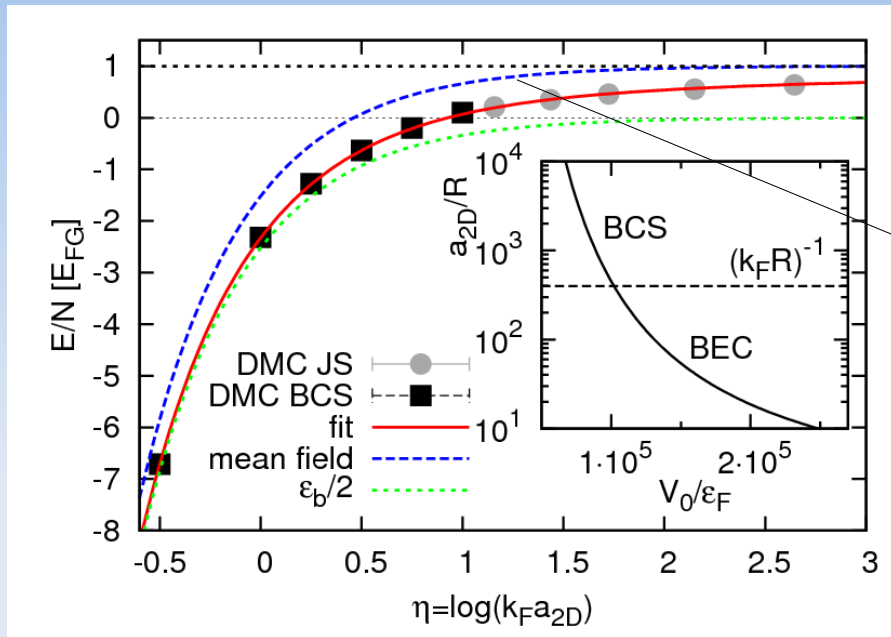
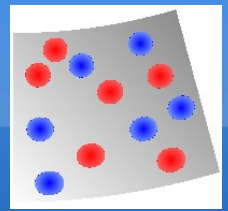
- Strong coupling:

$$\Psi_{BCS} = A [\phi(x_{11'}) \dots \phi(x_{N_{\uparrow}N_{\downarrow}})]$$

$\phi$ : bound state

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Energy per particle

Crossover parameter:  $\eta = \log(k_F a_{2D})$

BCS selfconsistent theory (Randeria) misses interaction between bosons

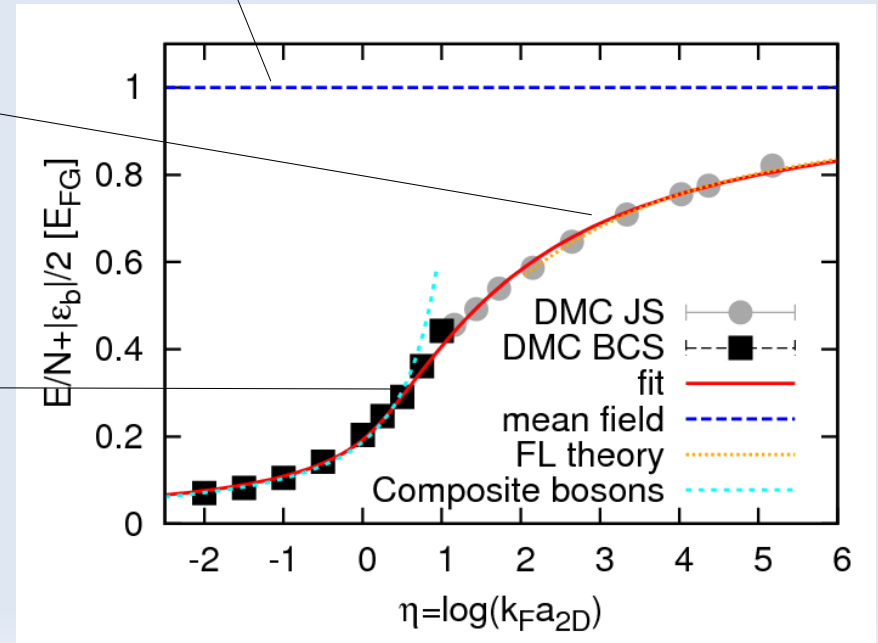
$$\frac{E}{N} = E_{FG} - \frac{|\epsilon_b|}{2}$$

Fermi liquid energy functional (small gap)

$$\frac{E}{N} = E_{FG} \left( 1 - \frac{1}{\eta} + \frac{A}{\eta^2} \right)$$

Composite bosons:  $a_{dd} \sim 0.55(4) a_{2D}$  (Petrov)

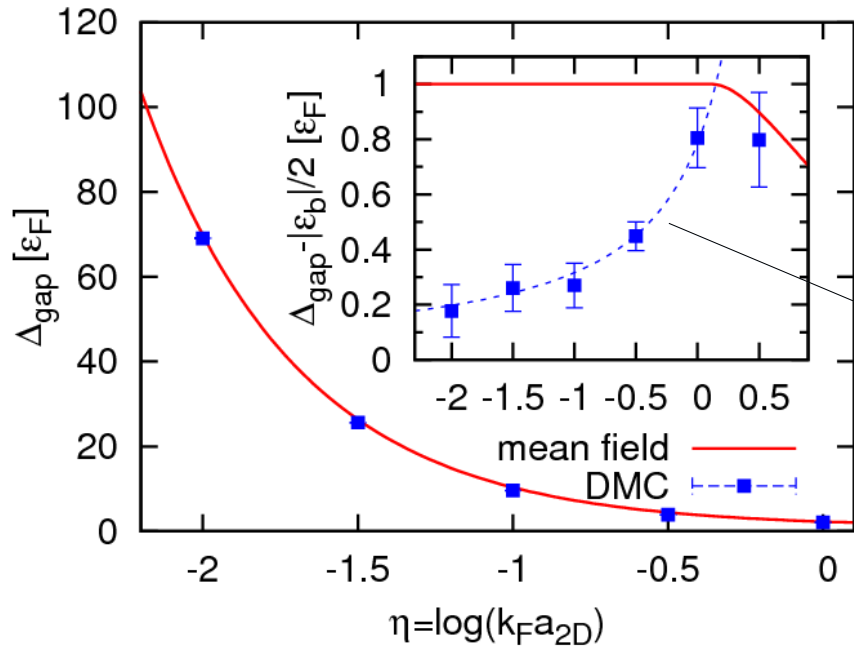
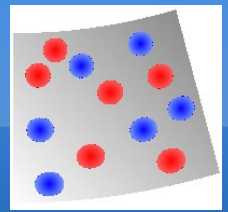
$$\frac{E}{N} = -\frac{\epsilon_b}{2} + \frac{\pi \hbar^2 n_d}{m_d} \frac{1}{\log \frac{1}{n_d a_{dd}^2}} \left[ 1 + 2^{nd} \text{ order} \right]$$





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Gap in the spectrum

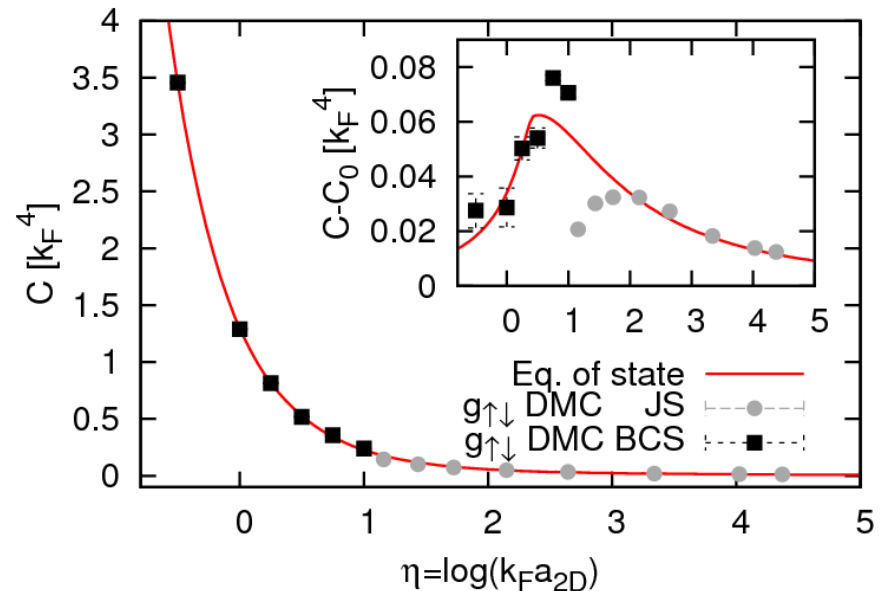
$$E(N_{\uparrow} + 1, N_{\downarrow}) = E(N_{\uparrow}, N_{\downarrow}) + \mu_{\uparrow} + \Delta_{gap}$$

Interpretation:  
One fermion in a BEC  
 $a_{ad} \sim 1.7(1) a_{2D}$

Short range interactions:  
Contact parameter (Tan)

$$\frac{dE}{d \log(k_F a_{2D})} \propto C$$

$$g_{\uparrow\downarrow}^{(2)}(r) \underset{r \rightarrow 0}{\propto} C \log(r/a_{2D})^2$$





# Conclusions

QMC shows absolute relevance of beyond mean-field corrections

The nodal surface of the used wavefunctions contains the relevant  $T=0$  physics

# Outlook

Polaron-molecule problem in 2D  
Phase separation