

Observation of the $\gamma\gamma \rightarrow \tau\tau$ Process in Pb + Pb Collisions and Constraints on the τ -Lepton Anomalous Magnetic Moment with the ATLAS Detector

G. Aad *et al.*^{*}
(ATLAS Collaboration)



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This Letter reports the observation of τ -lepton-pair production in ultraperipheral lead-lead collisions $\text{Pb} + \text{Pb} \rightarrow \text{Pb}(\gamma\gamma \rightarrow \tau\tau)\text{Pb}$ and constraints on the τ -lepton anomalous magnetic moment a_τ . The dataset corresponds to an integrated luminosity of 1.44 nb^{-1} of LHC Pb + Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ recorded by the ATLAS experiment in 2018. Selected events contain one muon from a τ -lepton decay, an electron or charged-particle track(s) from the other τ -lepton decay, little additional central-detector activity, and no forward neutrons. The $\gamma\gamma \rightarrow \tau\tau$ process is observed in Pb + Pb collisions with a significance exceeding 5 standard deviations and a signal strength of $\mu_{\tau\tau} = 1.03^{+0.06}_{-0.05}$ assuming the standard model value for a_τ . To measure a_τ , a template fit to the muon transverse-momentum distribution from τ -lepton candidates is performed, using a dimuon ($\gamma\gamma \rightarrow \mu\mu$) control sample to constrain systematic uncertainties. The observed 95% confidence-level interval for a_τ is $-0.057 < a_\tau < 0.024$.

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Measurements of the anomalous magnetic moment, $a_\ell = \frac{1}{2}(g_\ell - 2)$, of charged leptons ℓ (electrons, muons, and τ leptons) are cornerstone tests of the standard model (SM) with unique sensitivity to beyond-the-SM (BSM) phenomena. The leading contribution to a_ℓ in the SM is the one-loop Schwinger term $\alpha_{\text{EM}}/2\pi \simeq 0.00116$ [1,2], where α_{EM} is the electromagnetic (EM) fine-structure constant. For the electron (muon), a_e (a_μ) is tested to parts per 10^{10} [3–8] (10^7 [9–11]) precision. Measurements of a_μ report tensions with the SM expectation [12–16], which may suggest BSM dynamics. Specific BSM scenarios such as supersymmetry [17] predict enhancements that scale quadratically with lepton mass m_ℓ , i.e., $\delta a_\ell \propto m_\ell^2$, resulting in a $(m_\tau/m_\mu)^2 \simeq 280$ times larger effect for τ leptons. However, the short τ -lepton lifetime precludes precise spin-precession measurements of a_τ to test the SM prediction of $a_\tau^{\text{SM}} = 0.001\,177\,21(5)$ [18] and potential BSM contributions.

Photon-induced events arise from interactions between the EM fields surrounding the beam particles at colliders. Observing photon-induced τ -lepton pairs ($\gamma\gamma \rightarrow \tau\tau$) predicted to occur at the Large Hadron Collider (LHC) [19–26] would open the way to hadron-collider probes of a_τ . Currently, the most precise single-experiment measurement

is $a_\tau = -0.018(17)$ by the DELPHI Collaboration [27,28] using $\gamma\gamma \rightarrow \tau\tau$ events at the Large Electron Positron (LEP) collider. The OPAL [29] and L3 [30] Collaborations also set constraints using radiative τ -lepton decays [31]. At the LHC, photon-induced dilepton production has only been measured in the dielectron (ee) and dimuon ($\mu\mu$) channels, using proton-proton (pp) [32–37] and lead-lead (Pb + Pb) collisions [38–43]. The $\tau\tau$ channel is challenging due to hadronic backgrounds and neutrinos in τ -lepton decays diluting visible final-state kinematics. This renders triggering and reconstruction more difficult, especially in high-luminosity pp collisions. Strategies to overcome these experimental obstacles using heavy-ion collisions were proposed in Refs. [44–46].

This Letter presents the observation of the $\text{Pb} + \text{Pb} \rightarrow \text{Pb}(\gamma\gamma \rightarrow \tau\tau)\text{Pb}$ process and measurement of a_τ using 1.44 nb^{-1} of Pb + Pb data recorded by ATLAS in 2018 at a nucleon-nucleon (NN) center-of-mass energy of $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$. The EM fields accompanying the ions coherently create photons that interact to produce τ -lepton pairs. The cross section is enhanced by Z^4 relative to pp collisions, where Z is the atomic number ($Z = 82$ for lead). The ions can remain intact, enabling selection of low-multiplicity events with one muon originating from one of the τ leptons, while the other τ -lepton decay is reconstructed as either an electron or one or three charged-particle tracks with low transverse momentum.

The ATLAS experiment [47–49] is a multipurpose particle detector with cylindrical geometry [50], comprising an inner-detector (ID) tracker, EM and hadronic calorimeters, and a muon spectrometer (MS). The zero-degree calorimeters (ZDCs) [51] are located at $z = \pm 140 \text{ m}$ from

^{*}Full author list given at the end of the article.

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the interaction point and detect neutral particles such as neutrons emitted from interacting nuclei. A two-level trigger system [52,53] was used to select events containing one muon with $p_T > 4$ GeV and, at most, 50 GeV (3 GeV) of transverse energy deposited in the whole (forward $3.2 < |\eta| < 4.9$) calorimeter [54]. An extensive software suite [55] is used in the reconstruction and analysis of real and simulated data, in detector operations and in the trigger and data acquisition systems. Standard data-quality requirements are imposed [56]. The average number of hadronic interactions per bunch crossing was 0.003.

Samples of simulated $\gamma\gamma \rightarrow \tau\tau$ signal events were produced at leading order in QED using the STARLIGHT2.0 [57] Monte Carlo (MC) generator, interfaced with TAUOLA [58,59] for τ -lepton decays. Final-state radiation (FSR) from the τ leptons and charged decay products of τ leptons was simulated using PYTHIA8.245 [60] and PHOTOS3.61 [61], respectively. One of the dominant background sources is the $\gamma\gamma \rightarrow \mu\mu$ process, and its contribution is estimated with the aid of MC samples generated using STARLIGHT; PYTHIA8 was used to model EM FSR from the muons. The photon-flux distribution in simulated $\gamma\gamma \rightarrow \tau\tau$ and $\gamma\gamma \rightarrow \mu\mu$ events was reweighted to that of SUPERCHIC3.05 [62], differentially in dilepton invariant mass and dilepton rapidity. In the STARLIGHT and SUPERCHIC simulations, no restriction on the Coulomb breakup of either nucleus was imposed. Dijet samples from photon-induced diquark production, $\gamma\gamma \rightarrow q\bar{q}$, were generated using PYTHIA8. Nondiffractive photonuclear events ($\gamma A \rightarrow X$) were simulated with STARLIGHT interfaced with DPMJET-III [63]. All MC samples were passed through a detailed detector simulation based on GEANT4 [64,65].

Charged-particle tracks reconstructed in the ID must satisfy $p_T > 100$ MeV, $|\eta| < 2.5$, transverse impact parameter $|d_0| < 1.5$ mm, and the “loose primary” track selection criterion [66–68]. Electrons must satisfy $p_T > 4$ GeV, $|\eta| < 2.47$ (excluding the calorimeter transition region, i.e., $|\eta| \notin [1.37, 1.52]$), $|d_0| < 0.5$ mm, and “loose” [69] likelihood-based identification criteria. Muons must satisfy $p_T > 4$ GeV, $|\eta| < 2.4$, $|d_0| < 0.3$ mm, and “low p_T ” [70] identification criteria. Small corrections, derived using tag-and-probe methods similar to those in Refs. [40,71,72], are applied to simulated reconstruction and trigger efficiencies of electrons and muons. Clusters of topologically connected calorimeter cells called topoclusters [73] must satisfy $|\eta| < 4.9$, $p_T > 0.1$ GeV ($p_T > 1$ GeV) for $2.5 < |\eta| < 4.9$ ($|\eta| < 2.5$), and the cell significance criteria for measured energies outlined in Ref. [74]. Topoclusters from calorimeter regions with an abnormal noise distribution are removed using a data-driven procedure based on analyzing η - ϕ distributions of topocluster activity for each calorimeter layer. Reconstructed photons must satisfy $E_T > 1.5$ GeV, $|\eta| < 2.37 \notin [1.37, 1.52]$, and dedicated identification criteria defined in Ref. [72].

Selected events must contain exactly one muon, which targets a muonic decay of one of the τ leptons while reducing backgrounds from $\gamma\gamma \rightarrow \mu\mu$ and $\gamma\gamma \rightarrow q\bar{q}$. Three signal regions (SRs) then categorize events by the decay signature of the other τ lepton. The μe -SR category additionally requires one electron and no additional tracks separated from the muon (electron) by $\Delta R_{\mu(e),\text{trk}} > 0.1$, which targets fully leptonic decays of both τ leptons. The different-flavor (μe) requirement suppresses same-flavor backgrounds dominated by $\gamma\gamma \rightarrow \mu\mu/ee$. The $\mu 1\text{T}$ -SR ($\mu 3\text{T}$ -SR) category requires exactly one track (three tracks) separated from the muon by $\Delta R_{\mu,\text{trk}} > 0.1$, which targets τ -lepton decays to one or three charged hadrons. The one-track requirement also captures leptonic τ -lepton decays that fail electron or muon reconstruction. The electric charges of the muon, electron, and tracks must sum to zero.

For both $\mu 1\text{T}$ -SR and $\mu 3\text{T}$ -SR, events must contain no additional muons satisfying looser criteria and no electrons to reject $\gamma\gamma \rightarrow \mu\mu/ee$ backgrounds. The looser requirements on muons comprise matched tracks in the ID and MS satisfying $p_T > 2$ GeV and $|\eta| < 2.5$. To suppress hadronic backgrounds such as photonuclear processes, there must be no topoclusters separated from the muon (track or three-track system [75]) by $\Delta R_{\text{clust},\mu} > 0.3$ ($\Delta R_{\text{clust},\text{trk}(s)} > 1.0$); this requirement is referred to as the topocluster veto. To further reduce photonuclear backgrounds, the acoplanarity between the muon and the track (three-track system) must satisfy $A_{\phi}^{\mu,\text{trk}(s)} \equiv 1 - |\Delta\phi_{\mu,\text{trk}(s)}|/\pi < 0.4(0.2)$. The signal has a narrower $A_{\phi}^{\mu,\text{trk}(s)}$ distribution in $\mu 3\text{T}$ -SR than $\mu 1\text{T}$ -SR, motivating the tighter requirement.

For $\mu 1\text{T}$ -SR, the p_T of the muon-track pair must satisfy $p_T^{\mu,\text{trk}} > 1$ GeV to reject p_T -balanced backgrounds, such as $\gamma\gamma \rightarrow \mu\mu$. To further reduce the $\gamma\gamma \rightarrow \mu\mu\gamma$ background, the p_T of the muon, track, and photon (topocluster) system must fulfill $p_T^{\mu,\text{trk},\gamma(\text{clust})} > 1$ GeV for events containing a photon (topocluster) within $\Delta R_{\gamma(\text{clust}),\text{trk}} = 1$ of the track. If there are multiple nearby photons (topoclusters), the highest- p_T photon (topocluster) is used. The topoclusters considered here must have $p_T > 2$ GeV and not be track matched; these criteria avoid track-induced topoclusters from, e.g., charged-pion energy deposits. Topoclusters within $\Delta R_{\text{clust},\text{trk}} = 0.1$ of a track with $p_T > 0.7$ GeV extrapolated to the calorimeter are considered track matched. Low-multiplicity events in minimum-bias data are used to correct for the bending due to the magnetic field, such that the $\Delta R_{\text{clust},\text{trk}}$ distribution of topoclusters associated with a track peaks at zero. Tracks with $p_T < 0.7$ GeV are not considered in the track-cluster matching, as they typically do not deposit significant energy in the calorimeter.

For $\mu 3\text{T}$ -SR, the three-track system mass must fulfill $m_{3\text{trk}} < 1.7$ GeV, assuming each track has the charged-pion mass of 140 MeV. This requirement retains three-prong hadronic τ -lepton decays and suppresses background from exclusive ρ^0 mesons ($\gamma A \rightarrow \rho^0 \rightarrow \pi\pi$) produced

simultaneously with $\gamma\gamma \rightarrow \mu\mu$ events. In this background process, neither muon is correlated with the $\pi\pi$ system and $m_{\pi\pi\mu}$ typically exceeds 1.7 GeV for a muon with p_T of several GeV.

To constrain the $\gamma\gamma \rightarrow \mu\mu$ background, a control region (CR) of dimuon events called 2μ -CR is defined. It requires exactly two muons with invariant mass above 11 GeV to suppress quarkonia [$\Upsilon(nS) \rightarrow \mu\mu$] backgrounds and no additional tracks separated from the muons by $\Delta R_{\mu,\text{trk}} > 0.1$.

Events must additionally not have ZDC energies satisfying $E_{\text{ZDC}} > 1$ TeV on each side, mainly to suppress photonuclear backgrounds where ion dissociation typically occurs. This class of events with no forward neutrons detected ($0n0n$) corresponds to the absence of Coulomb breakup of either nucleus. Such breakup typically proceeds through the giant dipole resonance and induces the emission of one or more neutrons [40]. This requirement also fully suppresses lepton-pair production in which the initial photon emission results in the dissociation of one or both nuclei [40]. The SRs and 2μ -CR are all statistically independent. Since the extra forward neutron emissions are not simulated, the $\gamma\gamma \rightarrow \ell\ell$ MC samples are corrected using data-driven probabilities for the $0n0n$ event topology, which are found to be between 0.4 and 0.7. These are extracted from 2μ -CR without the E_{ZDC} requirement, differentially in dilepton invariant mass and dilepton rapidity.

The dominant sources of background after event selection are radiative dimuon ($\gamma\gamma \rightarrow \mu\mu\gamma$) and photonuclear events with low central-detector activity.

The $\gamma\gamma \rightarrow \mu\mu\gamma$ background is estimated with the aid of MC samples. This process enters $\mu 1\text{T-SR}$ when FSR photons substantially modify the dimuon kinematics so as to mimic the signal kinematics and enters $\mu 3\text{T-SR}$ and $\mu e\text{-SR}$ primarily when photons convert to e^+e^- in detector material. To improve the modeling of high- p_T ($p_T^\gamma \gtrsim p_T^\mu$) photon emissions for t - and u -channel $\gamma\gamma \rightarrow \mu\mu$ processes, an additional $\gamma\gamma \rightarrow \mu\mu\gamma$ MC sample generated using MADGRAPH5_AMC@NLO [76] with the photon flux reweighted to SUPERCHIC is used instead of STARLIGHT+PYTHIA8 if a leading photon has $p_T^\gamma > 2$ GeV. Comparing the $\gamma\gamma \rightarrow \mu\mu(\gamma)$ simulated events with data in 2μ -CR shows reasonable data-to-MC agreement in differential distributions and normalization to within 5%. The simulated $\gamma\gamma \rightarrow \mu\mu(\gamma)$ event yield in 2μ -CR is 15% lower than data if STARLIGHT photon-flux calculations are used due to known limitations of STARLIGHT [40,77]. Before the fit to data, this method estimates 70, 6.5, and 2.8 $\gamma\gamma \rightarrow \mu\mu(\gamma)$ events enter $\mu 1\text{T-SR}$, $\mu 3\text{T-SR}$, and $\mu e\text{-SR}$, respectively.

Diffraction photonuclear backgrounds with low particle activity are estimated using fully data-driven methods. Dedicated CRs are introduced, called $\mu 2\text{T-CR}$ ($\mu 4\text{T-CR}$), which apply the same selection as $\mu 1\text{T-SR}$ ($\mu 3\text{T-SR}$), but

require an additional track satisfying $p_T < 0.5$ GeV. Furthermore, the $E_{\text{ZDC}} < 1$ TeV requirement is removed on either side to enrich the sample with events from photonuclear processes. To suppress the $\gamma\gamma \rightarrow \tau\tau$ signal contamination in $\mu 2\text{T-CR}$, the two-track system mass must fulfill $m_{2\text{trk}} > 1$ GeV; if $m_{2\text{trk}} < 1$ GeV, the acoplanarity of the muon and highest- p_T track is required to exceed 0.2. The event yields in CRs are extrapolated to SRs by loosening the veto on topoclusters not matched to the muon or tracks from $n_{\text{TC}}^{\text{unmatch}} = 0$ to $n_{\text{TC}}^{\text{unmatch}} \leq 8$, both in CRs and SRs. The $\mu 2\text{T-CR}$ ($\mu 4\text{T-CR}$) templates for $n_{\text{TC}}^{\text{unmatch}}$ distributions are normalized to the event yield in $\mu 1\text{T-SR}$ ($\mu 3\text{T-SR}$) in the region $4 \leq n_{\text{TC}}^{\text{unmatch}} \leq 8$. In this region, the signal and dimuon background contributions are found to be negligible, and events exhibit properties that suggest nonexclusive diffractive production, such as a small or no rapidity gap [78]. As the additional track in $\mu 2\text{T-CR}$ and $\mu 4\text{T-CR}$ is soft ($p_T < 0.5$ GeV), its possible correlation with topocluster activity is very small. This method estimates that 13 (2.8) photonuclear events enter $\mu 1\text{T-SR}$ ($\mu 3\text{T-SR}$); photonuclear events are expected to be negligible in $\mu e\text{-SR}$.

Other sources of background are predicted to be negligible in the SRs. Nondiffractive photonuclear interactions are estimated using the STARLIGHT+DPMJET-III sample and are found to be negligible in all three SRs. The PYTHIA8 simulation of $\gamma\gamma \rightarrow q\bar{q}$ estimates dijet backgrounds contribute less than 0.3 events in both $\mu 1\text{T-SR}$ and $\mu 3\text{T-SR}$. Similarly, the contribution from resolved $\gamma\gamma$ interactions, as estimated with PYTHIA8 [79], is found to be negligible. Exclusive ρ^0 -meson production with simultaneous $\gamma\gamma \rightarrow \mu\mu$ production in $\mu 3\text{T-SR}$ is studied using a data-driven method. Template distributions for this process are built from events with two muons and two additional charged-particle tracks. The acoplanarities of both the muon pair and the track pair must be below 0.05. All events in these templates are found to have $m_{3\text{trk}} > 1.7$ GeV, so this background is expected to be negligible in $\mu 3\text{T-SR}$.

Systematic uncertainties affecting the measurement arise from the reconstruction of leptons, photons, charged-particle tracks and topoclusters, the signal and background modeling, and integrated luminosity.

Uncertainties in the muon momentum scale and resolution follow those in Ref. [80]. The analysis includes uncertainties in the data-to-MC correction factors applied to simulated samples for the muon trigger and reconstruction efficiencies. Uncertainties in the reconstruction, identification, and energy calibration of electrons and photons are evaluated in accord with Ref. [72]. The uncertainty in the inclusive track reconstruction efficiency is dominated by the uncertainty in the amount of ID material [67]. This uncertainty is applied in the simulation by randomly removing tracks with a p_T - and η -dependent probability corresponding to the material uncertainty. Uncertainties in the topocluster reconstruction efficiency and energy calibration are

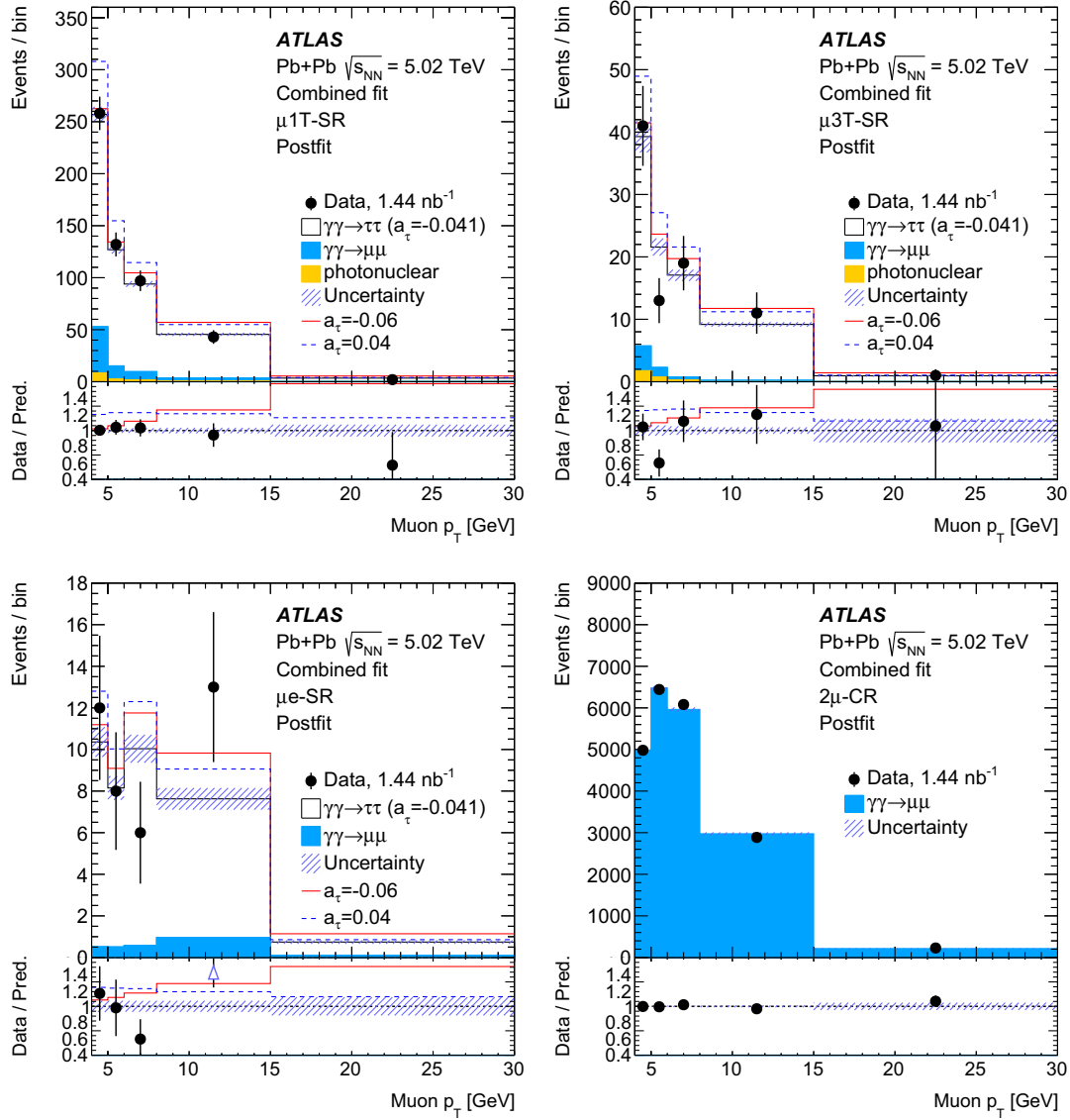


FIG. 1. Muon transverse-momentum distributions in the (top left) $\mu 1T$ -SR, (top right) $\mu 3T$ -SR, (bottom left) μe -SR, and (bottom right) 2μ -CR categories. Black markers denote data and stacked histograms indicate the different components contributing to the regions. Postfit distributions are shown with the signal contribution corresponding to the best-fit a_τ value ($a_\tau = -0.041$). For comparison, signal contributions with alternative a_τ values are shown as solid red ($a_\tau = -0.06$) or dashed blue ($a_\tau = 0.04$) lines. The bottom panel shows the ratio of the data to postfit predictions. Vertical bars denote uncertainties from the finite number of data events. Hatched bands represent $\pm 1\sigma$ systematic uncertainties of the prediction with the constraints from the fit applied.

estimated using $\gamma\gamma \rightarrow ee$ events where one of the electrons emits a hard bremsstrahlung photon due to its interaction with detector material [72].

Uncertainties in the photonuclear background evaluation are estimated by repeating the procedure with alternative requirements for CRs. These resemble the $\mu 1T$ -SR ($\mu 3T$ -SR) selection except that the track (three-track system) has the same electric charge as the muon candidate. The difference between the photonuclear background contribution evaluated with alternative and nominal CRs defines the uncertainties, affecting both the normalization and differential distributions. Uncertainties in modeling the photon

flux are estimated by using the STARLIGHT MC samples without reweighting to SUPERCHIC. This affects the normalization and differential distributions of the signal and $\gamma\gamma \rightarrow \mu\mu$ background. Uncertainties in modeling τ -lepton decays are estimated using PYTHIA8 [81] as an alternative MC simulation to TAUOLA. The effect of τ -lepton spin correlations in TAUOLA is implemented using helicity amplitudes from the $\gamma^* \rightarrow \tau\tau$ process. This modeling is therefore cross-checked by comparing signal events simulated using two versions of PYTHIA8: v8.245 uses helicity amplitudes from the $\gamma^* \rightarrow \tau\tau$ process, whereas v8.305 uses the $\gamma\gamma \rightarrow \tau\tau$ elementary process. The difference between the

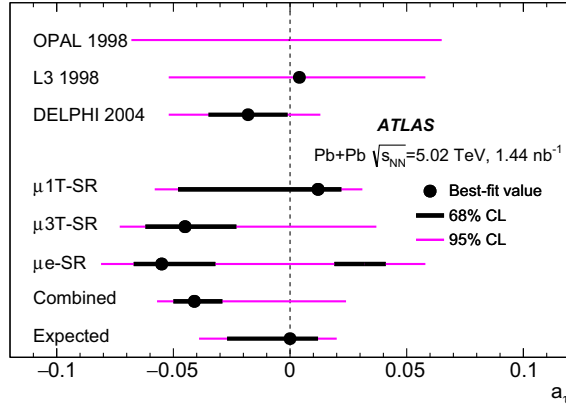


FIG. 2. Measurements of a_τ from fits to individual signal regions (including the dimuon control region) and from the combined fit. These are compared with existing measurements from the OPAL [29], L3 [30], and DELPHI [27] experiments at LEP. A point denotes the best-fit a_τ value for each measurement if available, while thick black (thin magenta) lines show 68% CL (95% CL) intervals. The expected interval from the ATLAS combined fit is also shown.

two implementations is found to be negligible, and no further systematic uncertainty is assigned.

The uncertainty in the integrated luminosity is 1.9%, obtained with the LUCID-2 detector [82] using methods similar to Ref. [83] for the primary luminosity measurements.

After applying the event selection, a total of 532, 85, and 39 data events are observed, compared with 84 ± 19 , 9 ± 3 , and 2.8 ± 0.7 expected background events in $\mu 1T$ -SR, $\mu 3T$ -SR, and μe -SR, respectively. The background-only hypothesis is rejected with significance exceeding 5σ , establishing the observation of the $\gamma\gamma \rightarrow \tau\tau$ process at ATLAS. The signal significance is highest in $\mu 1T$ -SR, while μe -SR has the largest signal-to-background ratio. The prefit signal-plus-background hypothesis predicts 543 ± 111 , 93 ± 20 , and 35 ± 8 events in $\mu 1T$ -SR, $\mu 3T$ -SR, and μe -SR, respectively, which is compatible with the observed data. The signal strength $\mu_{\tau\tau}$, defined as the ratio of the observed signal yield to the SM expectation assuming the SM value for a_τ , is measured using a profile-likelihood fit [84,85] to be $\mu_{\tau\tau} = 1.03^{+0.06}_{-0.05}(\text{tot}) = 1.03^{+0.05}_{-0.05}(\text{stat})^{+0.03}_{-0.03}(\text{syst})$. The fit uses the p_T^μ distribution in the three SRs and 2μ -CR with $\mu_{\tau\tau}$ being the only parameter of interest.

Approximately 80 nuisance parameters representing the systematic uncertainties are included in the fit. Many systematic uncertainties are correlated between the SRs and 2μ -CR, so their impact on the measurement precision is minimized since they are constrained by 2μ -CR. The dominant prefit contribution is the photon-flux uncertainty, which mainly affects the signal yield (by approximately 20%), with a significantly smaller impact on the signal shape found upon decorrelation from the normalization component. After the fit, the photon-flux uncertainty

becomes subdominant and luminosity uncertainty becomes negligible relative to other sources. The leading contributions to the total systematic uncertainty are the estimation of the muon trigger efficiency, τ -lepton decay modeling, and track reconstruction efficiency.

To measure a_τ , an alternative fit is performed where a_τ is the only free parameter using the p_T^μ distribution in the three SRs and 2μ -CR; p_T^μ is chosen because of its high sensitivity to a_τ [46]. Simulated signal samples with various a_τ values are employed. In the nominal sample, a_τ is set to its SM value. Signal templates for alternative a_τ hypotheses are obtained by reweighting the nominal sample in three dimensions, differentially in $\tau\tau$ invariant mass, $\tau\tau$ rapidity, and rapidity difference between the two τ leptons, according to calculations from Ref. [46]. These calculations parametrize the $\tau\tau\gamma$ coupling by $F_1(q^2)\gamma^\mu + F_2(q^2)(i/2m_\tau)\sigma^{\mu\nu}q_\nu$, where q_ν is the photon four-momentum, $\sigma^{\mu\nu} = i[\gamma^\mu, \gamma^\nu]/2$ the spin tensor, and the form factors satisfy $F_1(q^2 \rightarrow 0) = 1$ and $F_2(q^2 \rightarrow 0) = a_\tau$. A similar parametrization was used in previous LEP measurements [27,29,30], which exploits the near-zero virtuality of initial-state photons. A total of 14 templates for different a_τ values are created to model the dependence of the p_T^μ distribution on a_τ in the three SRs.

Figure 1 shows the p_T^μ distributions of the four analysis regions for the data and postfit expectation. The fit describes the data well.

The best-fit value of a_τ is $a_\tau = -0.041$, with the corresponding 68% and 95% confidence level (CL) intervals being $(-0.050, -0.029)$ and $(-0.057, 0.024)$, respectively. The higher-than-expected observed yields lead to the highly asymmetric 95% CL interval. This arises from the nearly quadratic signal cross section dependence on a_τ , caused by the interference of the SM and BSM amplitudes [29,30,46]. The expected 95% CL interval is $-0.039 < a_\tau < 0.020$. The impact of systematic uncertainties on the final results is small relative to statistical uncertainties. Figure 2 shows the a_τ measurement alongside previous results obtained at LEP. The precision of this measurement is similar to the most precise single-experiment measurement by the DELPHI Collaboration.

In summary, τ -lepton-pair production in ultraperipheral heavy-ion collisions, $\text{Pb} + \text{Pb} \rightarrow \text{Pb}(\gamma\gamma \rightarrow \tau\tau)\text{Pb}$, is observed by ATLAS with a significance exceeding 5σ in 1.44 nb^{-1} of $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ data at the LHC. The observed event yield is compatible with that expected from the SM prediction within uncertainties. The events are used to set constraints on the τ -lepton anomalous magnetic moment, corresponding to $-0.057 < a_\tau < 0.024$ at 95% CL. The measurement precision is limited by statistical uncertainties. This result introduces the use of hadron-collider data to test electromagnetic properties of the τ lepton, and the results are competitive with existing lepton-collider constraints.

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 H. M. Borecka-Bielska¹⁰⁷, L. S. Borgna⁹⁵, G. Borissov⁹⁰, D. Bortoletto¹²⁵, D. Boscherini^{23b}, M. Bosman¹³,
 J. D. Bossio Sola³⁶, K. Bouaouda^{35a}, J. Boudreau¹²⁸, E. V. Bouhova-Thacker⁹⁰, D. Boumediene⁴⁰, R. Bouquet⁵

- A. Boveia¹¹⁸, J. Boyd³⁶, D. Boye²⁹, I. R. Boyko³⁸, J. Bracinik²⁰, N. Brahimi^{62d}, G. Brandt¹⁶⁹, O. Brandt³², F. Braren⁴⁸, B. Brau¹⁰², J. E. Brau¹²², W. D. Breaden Madden⁵⁹, K. Brendlinger⁴⁸, R. Brenner¹⁶⁷, L. Brenner³⁶, R. Brenner¹⁵⁹, S. Bressler¹⁶⁷, B. Brickwedde⁹⁹, D. Britton⁵⁹, D. Britzger¹⁰⁹, I. Brock²⁴, G. Brooijmans⁴¹, W. K. Brooks^{136f}, E. Brost²⁹, P. A. Bruckman de Renstrom⁸⁵, B. Brüers⁴⁸, D. Bruncko^{28b,a}, A. Bruni^{23b}, G. Bruni^{23b}, M. Bruschi^{23b}, N. Bruscino^{74a,74b}, L. Bryngemark¹⁴², T. Buanes¹⁶, Q. Buat¹³⁷, P. Buchholz¹⁴⁰, A. G. Buckley⁵⁹, I. A. Budagov^{38,a}, M. K. Bugge¹²⁴, O. Bulekov³⁷, B. A. Bullard⁶¹, S. Burdin⁹¹, C. D. Burgard⁴⁸, A. M. Burger⁴⁰, B. Burghgrave⁸, J. T. P. Burr³², C. D. Burton¹¹, J. C. Burzynski¹⁴¹, E. L. Busch⁴¹, V. Büscher⁹⁹, P. J. Bussey⁵⁹, J. M. Butler²⁵, C. M. Buttar⁵⁹, J. M. Butterworth⁹⁵, W. Buttinger¹³³, C. J. Buxo Vazquez¹⁰⁶, A. R. Buzykaev³⁷, G. Cabras^{23b}, S. Cabrera Urbán¹⁶¹, D. Caforio⁵⁸, H. Cai¹²⁸, Y. Cai^{14a,14d}, V. M. M. Cairo³⁶, O. Cakir^{3a}, N. Calace³⁶, P. Calafiura^{17a}, G. Calderini¹²⁶, P. Calfayan⁶⁷, G. Callea⁵⁹, L. P. Caloba^{81b}, D. Calvet⁴⁰, S. Calvet⁴⁰, T. P. Calvet¹⁰¹, M. Calvetti^{73a,73b}, R. Camacho Toro¹²⁶, S. Camarda³⁶, D. Camarero Munoz⁹⁸, P. Camarri^{75a,75b}, M. T. Camerlingo^{76a,76b}, D. Cameron¹²⁴, C. Camincher¹⁶³, M. Campanelli⁹⁵, A. Camplani⁴², V. Canale^{71a,71b}, A. Canesse¹⁰³, M. Cano Bret⁷⁹, J. Cantero¹⁶¹, Y. Cao¹⁶⁰, F. Capocasa²⁶, M. Capua^{43b,43a}, A. Carbone^{70a,70b}, R. Cardarelli^{75a}, J. C. J. Cardenas⁸, F. Cardillo¹⁶¹, T. Carli³⁶, G. Carlino^{71a}, B. T. Carlson^{128,j}, E. M. Carlson^{163,155a}, L. Carminati^{70a,70b}, M. Carnesale^{74a,74b}, S. Caron¹¹², E. Carquin^{136f}, S. Carrá^{70a,70b}, G. Carratta^{23b,23a}, F. Carrio Argos^{33g}, J. W. S. Carter¹⁵⁴, T. M. Carter⁵², M. P. Casado^{13,k}, A. F. Casha¹⁵⁴, E. G. Castiglia¹⁷⁰, F. L. Castillo^{63a}, L. Castillo Garcia¹³, V. Castillo Gimenez¹⁶¹, N. F. Castro^{129a,129e}, A. Catinaccio³⁶, J. R. Catmore¹²⁴, V. Cavaliere²⁹, N. Cavalli^{23b,23a}, V. Cavasinni^{73a,73b}, E. Celebi^{21a}, F. Celli¹²⁵, M. S. Centonze^{69a,69b}, K. Cerny¹²¹, A. S. Cerqueira^{81a}, A. Cerri¹⁴⁵, L. Cerrito^{75a,75b}, F. Cerutti^{17a}, A. Cervelli^{23b}, S. A. Cetin^{21d}, Z. Chadi^{35a}, D. Chakraborty¹¹⁴, M. Chala^{129f}, J. Chan¹⁶⁸, W. S. Chan¹¹³, W. Y. Chan¹⁵², J. D. Chapman³², B. Chargeishvili^{148b}, D. G. Charlton²⁰, T. P. Charman⁹³, M. Chatterjee¹⁹, S. Chekanov⁶, S. V. Chekulaev^{155a}, G. A. Chelkov^{38,l}, A. Chen¹⁰⁵, B. Chen¹⁵⁰, B. Chen¹⁶³, C. Chen^{62a}, H. Chen^{14c}, H. Chen²⁹, J. Chen^{62c}, J. Chen²⁶, S. Chen¹⁵², S. J. Chen^{14c}, X. Chen^{62c}, X. Chen^{14b,m}, Y. Chen^{62a}, C. L. Cheng¹⁶⁸, H. C. Cheng^{64a}, A. Cheplakov³⁸, E. Cheremushkina⁴⁸, E. Cherepanova¹¹³, R. Cherkaoui El Moursli^{35e}, E. Cheu⁷, K. Cheung⁶⁵, L. Chevalier¹³⁴, V. Chiarella⁵³, G. Chiarelli^{73a}, N. Chiedde¹⁰¹, G. Chiodini^{69a}, A. S. Chisholm²⁰, A. Chitan^{27b}, Y. H. Chiu¹⁶³, M. V. Chizhov³⁸, K. Choi¹¹, A. R. Chomont^{74a,74b}, Y. Chou¹⁰², E. Y. S. Chow¹¹³, T. Chowdhury^{33g}, L. D. Christopher^{33g}, K. L. Chu^{64a}, M. C. Chu^{64a}, X. Chu^{14a,14d}, J. Chudoba¹³⁰, J. J. Chwastowski⁸⁵, D. Cieri¹⁰⁹, K. M. Ciesla^{84a}, V. Cindro⁹², A. Ciocio^{17a}, F. Ciotto^{71a,71b}, Z. H. Citron^{167,n}, M. Citterio^{70a}, D. A. Ciubotaru^{27b}, B. M. Ciungu¹⁵⁴, A. Clark⁵⁶, P. J. Clark⁵², J. M. Clavijo⁴⁸, S. E. Clawson¹⁰⁰, C. Clement^{47a,47b}, J. Clercx⁴⁸, L. Clissa^{23b,23a}, Y. Coadou¹⁰¹, M. Cobal^{68a,68c}, A. Coccaro^{57b}, R. F. Coelho Barrue^{129a}, R. Coelho Lopes De Sa¹⁰², S. Coelli^{70a}, H. Cohen¹⁵⁰, A. E. C. Coimbra^{70a,70b}, B. Cole⁴¹, J. Collot⁶⁰, P. Conde Muño^{129a,129g}, M. P. Connell^{33c}, S. H. Connell^{33c}, I. A. Connelly⁵⁹, E. I. Conroy¹²⁵, F. Conventi^{71a,o}, H. G. Cooke²⁰, A. M. Cooper-Sarkar¹²⁵, F. Cormier¹⁶², L. D. Corpe³⁶, M. Corradi^{74a,74b}, E. E. Corrigan⁹⁷, F. Corriveau^{103,p}, A. Cortes-Gonzalez¹⁸, M. J. Costa¹⁶¹, F. Costanza⁴, D. Costanzo¹³⁸, B. M. Cote¹¹⁸, G. Cowan⁹⁴, J. W. Cowley³², K. Cranmer¹¹⁶, S. Crépe-Renaudin⁶⁰, F. Crescioli¹²⁶, M. Cristinziani¹⁴⁰, M. Cristoforetti^{77a,77b,q}, V. Croft¹⁵⁷, G. Crosetti^{43b,43a}, A. Cueto³⁶, T. Cuhadar Donszelmann¹⁵⁸, H. Cui^{14a,14d}, Z. Cui⁷, A. R. Cukierman¹⁴², W. R. Cunningham⁵⁹, F. Curcio^{43b,43a}, P. Czodrowski³⁶, M. M. Czurylo^{63b}, M. J. Da Cunha Sargedas De Sousa^{62a}, J. V. Da Fonseca Pinto^{81b}, C. Da Via¹⁰⁰, W. Dabrowski^{84a}, T. Dado⁴⁹, S. Dahbi^{33g}, T. Dai¹⁰⁵, C. Dallapiccola¹⁰², M. Dam⁴², G. D'amen²⁹, V. D'Amico^{76a,76b}, J. Damp⁹⁹, J. R. Dandoy¹²⁷, M. F. Daneri³⁰, M. Danninger¹⁴¹, V. Dao³⁶, G. Darbo^{57b}, S. Darmora⁶, S. J. Das^{29,h}, A. Dattagupta¹²², S. D'Auria^{70a,70b}, C. David^{155b}, T. Davidek¹³², D. R. Davis⁵¹, B. Davis-Purcell³⁴, I. Dawson⁹³, K. De⁸, R. De Asmundis^{71a}, M. De Beurs¹¹³, S. De Castro^{23b,23a}, N. De Groot¹¹², P. de Jong¹¹³, H. De la Torre¹⁰⁶, A. De Maria^{14c}, A. De Salvo^{74a}, U. De Sanctis^{75a,75b}, A. De Santo¹⁴⁵, J. B. De Vivie De Regie⁶⁰, D. V. Dedovich³⁸, J. Degens¹¹³, A. M. Deiana⁴⁴, F. Del Corso^{23b,23a}, J. Del Peso⁹⁸, F. Del Rio^{63a}, F. Deliot¹³⁴, C. M. Delitzsch⁴⁹, M. Della Pietra^{71a,71b}, D. Della Volpe⁵⁶, A. Dell'Acqua³⁶, L. Dell'Asta^{70a,70b}, M. Delmastro⁴, P. A. Delsart⁶⁰, S. Demers¹⁷⁰, M. Demichev³⁸, S. P. Denisov³⁷, L. D'Eramo¹¹⁴, D. Derendarz⁸⁵, F. Derue¹²⁶, P. Dervan⁹¹, K. Desch²⁴, K. Dette¹⁵⁴, C. Deutsch²⁴, P. O. Deviveiros³⁶, F. A. Di Bello^{74a,74b}, A. Di Ciaccio^{75a,75b}, L. Di Ciaccio⁴, A. Di Domenico^{74a,74b}

- C. Di Donato^{71a,71b} A. Di Girolamo³⁶ G. Di Gregorio^{73a,73b} A. Di Luca^{77a,77b} B. Di Micco^{76a,76b}
 R. Di Nardo^{76a,76b} C. Diaconu¹⁰¹ F. A. Dias¹¹³ T. Dias Do Vale¹⁴¹ M. A. Diaz^{136a,136b} F. G. Diaz Capriles²⁴
 M. Didenko¹⁶¹ E. B. Diehl¹⁰⁵ L. Diehl⁵⁴ S. Díez Cornell⁴⁸ C. Díez Pardos¹⁴⁰ C. Dimitriadis^{24,159}
 A. Dimitrievska^{17a} W. Ding^{14b} J. Dingfelder²⁴ I-M. Dinu^{27b} S. J. Dittmeier^{63b} F. Dittus³⁶ F. Djama¹⁰¹
 T. Djobava^{148b} J. I. Djuvsland¹⁶ D. Dodsworth²⁶ C. Doglioni^{100,97} J. Dolejsi¹³² Z. Dolezal¹³² K. Domijan^{84a}
 M. Donadelli^{81c} B. Dong^{62c} J. Donini⁴⁰ A. D'Onofrio^{14c} M. D'Onofrio⁹¹ J. Dopke¹³³ A. Doria^{71a}
 M. T. Dova⁸⁹ A. T. Doyle⁵⁹ M. A. Draguet¹²⁵ E. Drechsler¹⁴¹ E. Dreyer¹⁶⁷ I. Drivas-koulouris¹⁰
 A. S. Drobac¹⁵⁷ D. Du^{62a} T. A. du Pree¹¹³ F. Dubinin³⁷ M. Dubovsky^{28a} E. Duchovni¹⁶⁷ G. Duckeck¹⁰⁸
 O. A. Ducu³⁶ D. Duda¹⁰⁹ A. Dudarev³⁶ M. D'uffizi¹⁰⁰ L. Duflot⁶⁶ M. Dührssen³⁶ C. Dülsen¹⁶⁹
 A. E. Dumitriu^{27b} M. Dunford^{63a} S. Dungs⁴⁹ K. Dunne^{47a,47b} A. Duperrin¹⁰¹ H. Duran Yildiz^{3a} M. Düren⁵⁸
 A. Durglishvili^{148b} B. L. Dwyer¹¹⁴ G. I. Dyckes^{17a} M. Dyndal^{84a} S. Dysch¹⁰⁰ B. S. Dziedzic⁸⁵
 Z. O. Earnshaw¹⁴⁵ B. Eckerova^{28a} M. G. Eggleston⁵¹ E. Egidio Purcino De Souza^{81b} L. F. Ehrke⁵⁶ G. Eigen¹⁶
 K. Einsweiler^{17a} T. Ekelof¹⁵⁹ P. A. Ekman⁹⁷ Y. El Ghazali^{35b} H. El Jarrari^{35e,147} A. El Moussaouy^{35a}
 V. Ellajosyula¹⁵⁹ M. Ellert¹⁵⁹ F. Ellinghaus¹⁶⁹ A. A. Elliot⁹³ N. Ellis³⁶ J. Elmsheuser²⁹ M. Elsing³⁶
 D. Emeliyanov¹³³ A. Emerman⁴¹ Y. Enari¹⁵² I. Ene^{17a} S. Epari¹³ J. Erdmann⁴⁹ A. Ereditato¹⁹
 P. A. Erland⁸⁵ M. Errenst¹⁶⁹ M. Escalier⁶⁶ C. Escobar¹⁶¹ E. Etzion¹⁵⁰ G. Evans^{129a} H. Evans⁶⁷
 M. O. Evans¹⁴⁵ A. Ezhilov³⁷ S. Ezzarqtouni^{35a} F. Fabbri⁵⁹ L. Fabbri^{23b,23a} G. Facini⁹⁵ V. Fadeyev¹³⁵
 R. M. Fakhrutdinov³⁷ S. Falciano^{74a} P. J. Falke²⁴ S. Falke³⁶ J. Faltova¹³² Y. Fan^{14a} Y. Fang^{14a,14d}
 G. Fanourakis⁴⁶ M. Fanti^{70a,70b} M. Faraj^{68a,68b} A. Farbin⁸ A. Farilla^{76a} T. Farooque¹⁰⁶ S. M. Farrington⁵²
 F. Fassi^{35e} D. Fassouliotis⁹ M. Fauci Giannelli^{75a,75b} W. J. Fawcett³² L. Fayard⁶⁶ O. L. Fedin^{37,1}
 G. Fedotov³⁷ M. Feickert¹⁶⁰ L. Feligioni¹⁰¹ A. Fell¹³⁸ D. E. Fellers¹²² C. Feng^{62b} M. Feng^{14b} Z. Feng¹¹³
 M. J. Fenton¹⁵⁸ A. B. Fenyuk³⁷ L. Ferencz⁴⁸ S. W. Ferguson⁴⁵ J. Pretel⁵⁴ J. Ferrando⁴⁸ A. Ferrari¹⁵⁹
 P. Ferrari¹¹³ R. Ferrari^{72a} D. Ferrere⁵⁶ C. Ferretti¹⁰⁵ F. Fiedler⁹⁹ A. Filipčič⁹² E. K. Filmer¹ F. Filthaut¹¹²
 M. C. N. Fiolhais^{129a,129c,f} L. Fiorini¹⁶¹ F. Fischer¹⁴⁰ W. C. Fisher¹⁰⁶ T. Fitschen^{20,66} I. Fleck¹⁴⁰
 P. Fleischmann¹⁰⁵ T. Flick¹⁶⁹ L. Flores¹²⁷ M. Flores^{33d} L. R. Flores Castillo^{64a} F. M. Follega^{77a,77b}
 N. Fomin¹⁶ J. H. Foo¹⁵⁴ B. C. Forland⁶⁷ A. Formica¹³⁴ A. C. Forti¹⁰⁰ E. Fortin¹⁰¹ A. W. Fortman⁶¹
 M. G. Foti^{17a} L. Fountas^{9,s} D. Fournier⁶⁶ H. Fox⁹⁰ P. Francavilla^{73a,73b} S. Francescato⁶¹ M. Franchini^{23b,23a}
 S. Franchino^{63a} D. Francis³⁶ L. Franco¹¹² L. Franconi¹⁹ M. Franklin⁶¹ G. Frattari²⁶ A. C. Freegard⁹³
 P. M. Freeman²⁰ W. S. Freund^{81b} N. Fritzsche⁵⁰ A. Froch⁵⁴ D. Froidevaux³⁶ J. A. Frost¹²⁵ Y. Fu^{62a}
 M. Fujimoto¹¹⁷ E. Fullana Torregrosa^{161,a} J. Fuster¹⁶¹ A. Gabrielli^{23b,23a} A. Gabrielli³⁶ P. Gadow⁴⁸
 G. Gagliardi^{57b,57a} L. G. Gagnon^{17a} G. E. Gallardo¹²⁵ E. J. Gallas¹²⁵ B. J. Gallop¹³³ R. Gamboa Goni⁹³
 K. K. Gan¹¹⁸ S. Ganguly¹⁵² J. Gao^{62a} Y. Gao⁵² F. M. Garay Walls^{136a,136b} B. Garcia^{29,h} C. García¹⁶¹
 J. E. García Navarro¹⁶¹ J. A. García Pascual^{14a} M. Garcia-Sciveres^{17a} R. W. Gardner³⁹ D. Garg⁷⁹
 R. B. Garg^{142,kk} S. Gargiulo⁵⁴ C. A. Garner¹⁵⁴ V. Garonne²⁹ S. J. Gasiorowski¹³⁷ P. Gaspar^{81b} G. Gaudio^{72a}
 V. Gautam¹³ P. Gauzzi^{74a,74b} I. L. Gavrilenko³⁷ A. Gavrilyuk³⁷ C. Gay¹⁶² G. Gaycken⁴⁸ E. N. Gazis¹⁰
 A. A. Geanta^{27b,27e} C. M. Gee¹³⁵ J. Geisen⁹⁷ M. Geisen⁹⁹ C. Gemme^{57b} M. H. Genest⁶⁰ S. Gentile^{74a,74b}
 S. George⁹⁴ W. F. George²⁰ T. Geralis⁴⁶ L. O. Gerlach⁵⁵ P. Gessinger-Befurt³⁶ M. Ghasemi Bostanabad¹⁶³
 M. Ghneimat¹⁴⁰ A. Ghosal¹⁴⁰ A. Ghosh¹⁵⁸ A. Ghosh⁷ B. Giacobbe^{23b} S. Giagu^{74a,74b} N. Giangiacomi¹⁵⁴
 P. Giannetti^{73a} A. Giannini^{62a} S. M. Gibson⁹⁴ M. Gignac¹³⁵ D. T. Gil^{84b} A. K. Gilbert^{84a} B. J. Gilbert⁴¹
 D. Gillberg³⁴ G. Gilles¹¹³ N. E. K. Gillwald⁴⁸ L. Ginabat¹²⁶ D. M. Gingrich^{2,e} M. P. Giordani^{68a,68c}
 P. F. Giraud¹³⁴ G. Giugliarelli^{68a,68c} D. Giugni^{70a} F. Giuli³⁶ I. Gkialas^{9,s} L. K. Gladilin³⁷ C. Glasman⁹⁸
 G. R. Gledhill¹²² M. Glisic¹²² I. Gnesi^{43b,i} Y. Go^{29,h} M. Goblirsch-Kolb²⁶ D. Godin¹⁰⁷ S. Goldfarb¹⁰⁴
 T. Golling⁵⁶ M. G. D. Gololo^{33g} D. Golubkov³⁷ J. P. Gombas¹⁰⁶ A. Gomes^{129a,129b} G. Gomes Da Silva¹⁴⁰
 A. J. Gomez Delegido¹⁶¹ R. Goncalves Gama⁵⁵ R. Gonçalves^{129a,129c} G. Gonella¹²² L. Gonella²⁰ A. Gongadze³⁸
 F. Gonnella²⁰ J. L. Gonski⁴¹ S. González de la Hoz¹⁶¹ S. Gonzalez Fernandez¹³ R. Gonzalez Lopez⁹¹
 C. Gonzalez Renteria^{17a} R. Gonzalez Suarez¹⁵⁹ S. Gonzalez-Sevilla⁵⁶ G. R. Gonzalvo Rodriguez¹⁶¹
 R. Y. González Andana⁵² L. Goossens³⁶ N. A. Gorasia²⁰ P. A. Gorbounov³⁷ B. Gorini³⁶ E. Gorini^{69a,69b}
 A. Gorišek⁹² A. T. Goshaw⁵¹ M. I. Gostkin³⁸ C. A. Gottardo³⁶ M. Goughri^{35b} V. Goumarre⁴⁸
 A. G. Goussiou¹³⁷ N. Govender^{33c} C. Goy⁴ I. Grabowska-Bold^{84a} K. Graham³⁴ E. Gramstad¹²⁴

- S. Grancagnolo¹⁸, M. Grandi¹⁴⁵, V. Gratchev^{37,a}, P. M. Gravila^{27f}, F. G. Gravili^{69a,69b}, H. M. Gray^{17a},
 M. Greco^{69a,69b}, C. Grefe²⁴, I. M. Gregor⁴⁸, P. Grenier¹⁴², C. Grieco¹³, A. A. Grillo¹³⁵, K. Grimm^{31,u},
 S. Grinstein^{13,v}, J.-F. Grivaz⁶⁶, E. Gross¹⁶⁷, J. Grosse-Knetter⁵⁵, C. Grud¹⁰⁵, A. Grummer¹¹¹, J. C. Grundy¹²⁵,
 L. Guan¹⁰⁵, W. Guan¹⁶⁸, C. Gubbels¹⁶², J. G. R. Guerrero Rojas¹⁶¹, G. Guerrieri^{68a,68b}, F. Guescini¹⁰⁹,
 R. Gugel⁹⁹, J. A. M. Guhit¹⁰⁵, A. Guida⁴⁸, T. Guillemain⁴, E. Guilloton^{165,133}, S. Guindon³⁶, F. Guo^{14a,14d},
 J. Guo^{62c}, L. Guo⁶⁶, Y. Guo¹⁰⁵, R. Gupta⁴⁸, S. Gurbuz²⁴, S. S. Gurdasani⁵⁴, G. Gustavino³⁶, M. Guth⁵⁶,
 P. Gutierrez¹¹⁹, L. F. Gutierrez Zagazeta¹²⁷, C. Gutschow⁹⁵, C. Guyot¹³⁴, C. Gwenlan¹²⁵, C. B. Gwilliam⁹¹,
 E. S. Haaland¹²⁴, A. Haas¹¹⁶, M. Habedank⁴⁸, C. Haber^{17a}, H. K. Hadavand⁸, A. Hadeef⁹⁹, S. Hadzic¹⁰⁹,
 M. Haleem¹⁶⁴, J. Haley¹²⁰, J. J. Hall¹³⁸, G. D. Hallowell¹⁰¹, L. Halser¹⁹, K. Hamano¹⁶³, H. Hamdaoui^{35e},
 M. Hamer²⁴, G. N. Hamity⁵², J. Han^{62b}, K. Han^{62a}, L. Han^{14c}, L. Han^{62a}, S. Han^{17a}, Y. F. Han¹⁵⁴,
 K. Hanagaki⁸², M. Hance¹³⁵, D. A. Hangal^{41,d}, M. D. Hank³⁹, R. Hankache¹⁰⁰, J. B. Hansen⁴², J. D. Hansen⁴²,
 P. H. Hansen⁴², K. Hara¹⁵⁶, D. Harada⁵⁶, T. Harenberg¹⁶⁹, S. Harkusha³⁷, Y. T. Harris¹²⁵, N. M. Harrison¹¹⁸,
 P. F. Harrison¹⁶⁵, N. M. Hartman¹⁴², N. M. Hartmann¹⁰⁸, Y. Hasegawa¹³⁹, A. Hasib⁵², S. Haug¹⁹, R. Hauser¹⁰⁶,
 M. Havranek¹³¹, C. M. Hawkes²⁰, R. J. Hawkins³⁶, S. Hayashida¹¹⁰, D. Hayden¹⁰⁶, C. Hayes¹⁰⁵,
 R. L. Hayes¹⁶², C. P. Hays¹²⁵, J. M. Hays⁹³, H. S. Hayward⁹¹, F. He^{62a}, Y. He¹⁵³, Y. He¹²⁶, M. P. Heath⁵²,
 V. Hedberg⁹⁷, A. L. Heggelund¹²⁴, N. D. Hehir⁹³, C. Heidegger⁵⁴, K. K. Heidegger⁵⁴, W. D. Heidorn⁸⁰,
 J. Heilman³⁴, S. Heim⁴⁸, T. Heim^{17a}, J. G. Heinlein¹²⁷, J. J. Heinrich¹²², L. Heinrich^{109,jj}, J. Hejbal¹³⁰,
 L. Helary⁴⁸, A. Held¹⁶⁸, S. Hellesund¹²⁴, C. M. Helling¹⁶², S. Hellman^{47a,47b}, C. Helsens³⁶, R. C. W. Henderson⁹⁰,
 L. Henkelmann³², A. M. Henriques Correia³⁶, H. Herde¹⁴², L. Hermann⁵⁴, Y. Hernández Jiménez¹⁴⁴, H. Herr⁹⁹,
 M. G. Herrmann¹⁰⁸, T. Herrmann⁵⁰, G. Herten⁵⁴, R. Hertenberger¹⁰⁸, L. Hervas³⁶, N. P. Hessey^{155a}, H. Hibi⁸³,
 E. Higón-Rodríguez¹⁶¹, S. J. Hillier²⁰, I. Hinchliffe^{17a}, F. Hinterkeuser²⁴, M. Hirose¹²³, S. Hirose¹⁵⁶,
 D. Hirschbuehl¹⁶⁹, T. G. Hitchings¹⁰⁰, B. Hiti⁹², J. Hobbs¹⁴⁴, R. Hobincu^{27e}, N. Hod¹⁶⁷, M. C. Hodgkinson¹³⁸,
 B. H. Hodgkinson³², A. Hoecker³⁶, J. Hofer⁴⁸, D. Hohn⁵⁴, T. Holm²⁴, M. Holzbock¹⁰⁹, L. B. A. H. Hommels³²,
 B. P. Honan¹⁰⁰, J. Hong^{62c}, T. M. Hong¹²⁸, Y. Hong⁵⁵, J. C. Honig⁵⁴, A. Hönle¹⁰⁹, B. H. Hooberman¹⁶⁰,
 W. H. Hopkins⁶, Y. Horii¹¹⁰, S. Hou¹⁴⁷, A. S. Howard⁹², J. Howarth⁵⁹, J. Hoya⁸⁹, M. Hrabovsky¹²¹,
 A. Hrynevich³⁷, T. Hryn'ova⁴, P. J. Hsu⁶⁵, S.-C. Hsu¹³⁷, Q. Hu^{41,d}, Y. F. Hu^{14a,14d,w}, D. P. Huang⁹⁵,
 S. Huang^{64b}, X. Huang^{14c}, Y. Huang^{62a}, Y. Huang^{14a}, Z. Huang¹⁰⁰, Z. Hubacek¹³¹, M. Huebner²⁴,
 F. Huegging²⁴, T. B. Huffman¹²⁵, M. Huhtinen³⁶, S. K. Huiberts¹⁶, R. Hulsken¹⁰³, N. Huseynov¹²¹,
 J. Huston¹⁰⁶, J. Huth⁶¹, R. Hyneman¹⁴², S. Hyrych^{28a}, G. Iacobucci⁵⁶, G. Iakovidis²⁹, I. Ibragimov¹⁴⁰,
 L. Iconomidou-Fayard⁶⁶, P. Iengo^{71a,71b}, R. Iguchi¹⁵², T. Iizawa⁵⁶, Y. Ikegami⁸², A. Ilg¹⁹, N. Ilic¹⁵⁴,
 H. Imam^{35a}, T. Ingebretsen Carlson^{47a,47b}, G. Introzzi^{72a,72b}, M. Iodice^{76a}, V. Ippolito^{74a,74b}, M. Ishino¹⁵²,
 W. Islam¹⁶⁸, C. Issever^{18,48}, S. Istin^{21a,x}, H. Ito¹⁶⁶, J. M. Iturbe Ponce^{64a}, R. Iuppa^{77a,77b}, A. Ivina¹⁶⁷,
 J. M. Izen⁴⁵, V. Izzo^{71a}, P. Jacka^{130,131}, P. Jackson¹, R. M. Jacobs⁴⁸, B. P. Jaeger¹⁴¹, C. S. Jagfeld¹⁰⁸,
 G. Jäkel¹⁶⁹, K. Jakobs⁵⁴, T. Jakoubek¹⁶⁷, J. Jamieson⁵⁹, K. W. Janas^{84a}, G. Jarlskog⁹⁷, A. E. Jaspán⁹¹,
 T. Javůrek³⁶, M. Javurkova¹⁰², F. Jeanneau¹³⁴, L. Jeanty¹²², J. Jejelava^{148a,y}, P. Jenni^{54,z}, C. E. Jessiman³⁴,
 S. Jézéquel⁴, J. Jia¹⁴⁴, X. Jia⁶¹, X. Jia^{14a,14d}, Z. Jia^{14c}, Y. Jiang^{62a}, S. Jiggins⁵², J. Jimenez Pena¹⁰⁹, S. Jin^{14c},
 A. Jinaru^{27b}, O. Jinnouchi¹⁵³, H. Jivan^{33g}, P. Johansson¹³⁸, K. A. Johns⁷, C. A. Johnson⁶⁷, D. M. Jones³²,
 E. Jones¹⁶⁵, P. Jones³², R. W. L. Jones⁹⁰, T. J. Jones⁹¹, J. Jovicevic¹⁵, X. Ju^{17a}, J. J. Junggeburth³⁶,
 A. Juste Rozas^{13,v}, S. Kabana^{136e}, A. Kaczmarek⁸⁵, M. Kado^{74a,74b}, H. Kagan¹¹⁸, M. Kagan¹⁴², A. Kahn⁴¹,
 A. Kahn¹²⁷, C. Kahra⁹⁹, T. Kaji¹⁶⁶, E. Kajomovitz¹⁴⁹, N. Kakati¹⁶⁷, C. W. Kalderon²⁹, A. Kamenshchikov¹⁵⁴,
 N. J. Kang¹³⁵, Y. Kano¹¹⁰, D. Kar^{33g}, K. Karava¹²⁵, M. J. Kareem^{155b}, E. Karentzos⁵⁴, I. Karkanas¹⁵¹,
 S. N. Karpov³⁸, Z. M. Karpova³⁸, V. Kartvelishvili⁹⁰, A. N. Karyukhin³⁷, E. Kasimi¹⁵¹, C. Kato^{62d}, J. Katzy⁴⁸,
 S. Kaur³⁴, K. Kawade¹³⁹, K. Kawagoe⁸⁸, T. Kawaguchi¹¹⁰, T. Kawamoto¹³⁴, G. Kawamura⁵⁵, E. F. Kay¹⁶³,
 F. I. Kaya¹⁵⁷, S. Kazakos¹³, V. F. Kazanin³⁷, Y. Ke¹⁴⁴, J. M. Keaveney^{33a}, R. Keeler¹⁶³, G. V. Kehris⁶¹,
 J. S. Keller³⁴, A. S. Kelly⁹⁵, D. Kelsey¹⁴⁵, J. J. Kempster²⁰, J. Kendrick²⁰, K. E. Kennedy⁴¹, O. Kepka¹³⁰,
 B. P. Kerridge¹⁶⁵, S. Kersten¹⁶⁹, B. P. Kerševan⁹², L. Keszeghova^{28a}, S. Ketabchi Haghighat¹⁵⁴, M. Khandoga¹²⁶,
 A. Khanov¹²⁰, A. G. Kharlamov³⁷, T. Kharlamova³⁷, E. E. Khoda¹³⁷, T. J. Khoo¹⁸, G. Khorauli¹⁶⁴,
 J. Khubua^{148b}, Y. A. R. Khwaira⁶⁶, M. Kiehn³⁶, A. Kilgallon¹²², D. W. Kim^{47a,47b}, E. Kim¹⁵³, Y. K. Kim³⁹,
 N. Kimura⁹⁵, A. Kirchhoff⁵⁵, D. Kirchmeier⁵⁰, C. Kirfel²⁴, J. Kirk¹³³, A. E. Kiryunin¹⁰⁹, T. Kishimoto¹⁵²

- D. P. Kisliuk,¹⁵⁴ C. Kitsaki¹⁰ O. Kivernyk²⁴ M. Klassen^{63a} C. Klein³⁴ L. Klein¹⁶⁴ M. H. Klein¹⁰⁵ M. Klein⁹¹
 U. Klein⁹¹ P. Klimek³⁶ A. Klimentov²⁹ F. Klimpel¹⁰⁹ T. Klingl²⁴ T. Klioutchnikova³⁶ F. F. Klitzner¹⁰⁸
 P. Kluit¹¹³ S. Kluth¹⁰⁹ E. Kneringer⁷⁸ T. M. Knight¹⁵⁴ A. Knue⁵⁴ D. Kobayashi⁸⁸ R. Kobayashi⁸⁶
 M. Kocian¹⁴² T. Kodama¹⁵² P. Kodyš¹³² D. M. Koeck¹⁴⁵ P. T. Koenig²⁴ T. Koffas³⁴ N. M. Köhler³⁶
 M. Kolb¹³⁴ I. Koletsou⁴ T. Komarek¹²¹ K. Köneke⁵⁴ A. X. Y. Kong¹ T. Kono¹¹⁷ N. Konstantinidis⁹⁵
 B. Konya⁹⁷ R. Kopeliānsky⁶⁷ S. Koperny^{84a} K. Korcyl⁸⁵ K. Kordas¹⁵¹ G. Koren¹⁵⁰ A. Korn⁹⁵ S. Korn⁵⁵
 I. Korolkov¹³ N. Korotkova³⁷ B. Kortman¹¹³ O. Kortner¹⁰⁹ S. Kortner¹⁰⁹ W. H. Kostecka¹¹⁴
 V. V. Kostyukhin¹⁴⁰ A. Kotskechagia⁶⁶ A. Kotwal⁵¹ A. Koulouris³⁶ A. Kourkouveli-Charalampidi^{72a,72b}
 C. Kourkouvelis⁹ E. Kourlitis⁶ O. Kovanda¹⁴⁵ R. Kowalewski¹⁶³ W. Kozanecki¹³⁴ A. S. Kozhin³⁷
 V. A. Kramarenko³⁷ G. Kramberger⁹² P. Kramer⁹⁹ M. W. Krasny¹²⁶ A. Krasznahorkay³⁶ J. A. Kremer⁹⁹
 T. Kresse⁵⁰ J. Kretschmar⁹¹ K. Kreul¹⁸ P. Krieger¹⁵⁴ F. Krieter¹⁰⁸ S. Krishnamurthy¹⁰² A. Krishnan^{63b}
 M. Krivos¹³² K. Krizka^{17a} K. Kroeninger⁴⁹ H. Kroha¹⁰⁹ J. Kroll¹³⁰ J. Kroll¹²⁷ K. S. Krowpman¹⁰⁶
 U. Kruchonak³⁸ H. Krüger²⁴ N. Krumnack⁸⁰ M. C. Kruse⁵¹ J. A. Krzysiak⁸⁵ A. Kubota¹⁵³ O. Kuchinskaja³⁷
 S. Kудay^{3a} D. Kuechler⁴⁸ J. T. Kuechler⁴⁸ S. Kuehn³⁶ T. Kuhl⁴⁸ V. Kukhtin³⁸ Y. Kulchitsky^{37,1}
 S. Kuleshov^{136d,136b} M. Kumar^{33g} N. Kumari¹⁰¹ M. Kuna⁶⁰ A. Kupco¹³⁰ T. Kupfer⁴⁹ A. Kupich³⁷
 O. Kuprash⁵⁴ H. Kurashige⁸³ L. L. Kurchaninov^{155a} Y. A. Kurochkin³⁷ A. Kurova³⁷ E. S. Kuwertz³⁶
 M. Kuze¹⁵³ A. K. Kvam¹⁰² J. Kvita¹²¹ T. Kwan¹⁰³ K. W. Kwok^{64a} N. G. Kyriacou¹⁰⁵ L. A. O. Laatu¹⁰¹
 C. Lacasta¹⁶¹ F. Lacava^{74a,74b} H. Lacker¹⁸ D. Lacour¹²⁶ N. N. Lad⁹⁵ E. Ladygin³⁸ B. Laforge¹²⁶
 T. Lagouri^{136e} S. Lai⁵⁵ I. K. Lakomic^{84a} N. Lalloue⁶⁰ J. E. Lambert¹¹⁹ S. Lammers⁶⁷ W. Lampl⁷
 C. Lampoudis¹⁵¹ A. N. Lancaster¹¹⁴ E. Lançon²⁹ U. Landgraf⁵⁴ M. P. J. Landon⁹³ V. S. Lang⁵⁴
 R. J. Langenberg¹⁰² A. J. Lankford¹⁵⁸ F. Lanni²⁹ K. Lantzsch²⁴ A. Lanza^{72a} A. Lapertosa^{57b,57a}
 J. F. Laporte¹³⁴ T. Lari^{70a} F. Lasagni Manghi^{23b} M. Lassnig³⁶ V. Latonova¹³⁰ T. S. Lau^{64a} A. Laudrain⁹⁹
 A. Laurier³⁴ S. D. Lawlor⁹⁴ Z. Lawrence¹⁰⁰ M. Lazzaroni^{70a,70b} B. Le¹⁰⁰ B. Leban⁹² A. Lebedev⁸⁰
 M. LeBlanc³⁶ T. LeCompte⁶ F. Ledroit-Guillon⁶⁰ A. C. A. Lee⁹⁵ G. R. Lee¹⁶ L. Lee⁶¹ S. C. Lee¹⁴⁷
 S. Lee^{47a,47b} T. F. Lee⁹¹ L. L. Leeuw^{33c} H. P. Lefebvre⁹⁴ M. Lefebvre¹⁶³ C. Leggett^{17a} K. Lehmann¹⁴¹
 G. Lehmann Miotto³⁶ W. A. Leight¹⁰² A. Leisos^{151,aa} M. A. L. Leite^{81c} C. E. Leitgeb⁴⁸ R. Leitner¹³²
 K. J. C. Leney⁴⁴ T. Lenz²⁴ S. Leone^{73a} C. Leonidopoulos⁵² A. Leopold¹⁴³ C. Leroy¹⁰⁷ R. Les¹⁰⁶
 C. G. Lester³² M. Levchenko³⁷ J. Levêque⁴ D. Levin¹⁰⁵ L. J. Levinson¹⁶⁷ M. P. Lewicki⁸⁵ D. J. Lewis²⁰
 B. Li^{14b} B. Li^{62b} C. Li^{62a} C-Q. Li^{62c,62d} H. Li^{62a} H. Li^{62b} H. Li^{14c} H. Li^{62b} J. Li^{62c} K. Li¹³⁷ L. Li^{62c}
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 M. Liberatore⁴⁸ B. Liberti^{75a} K. Lie^{64c} J. Lieber Marin^{81b} K. Lin¹⁰⁶ R. A. Linck⁶⁷ R. E. Lindley⁷
 J. H. Lindon² A. Linss⁴⁸ E. Lipeles¹²⁷ A. Lipniacka¹⁶ A. Lister¹⁶² J. D. Little⁴ B. Liu^{14a} B. X. Liu¹⁴¹
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 X. Liu^{62a} X. Liu^{62b} Y. Liu⁴⁸ Y. Liu^{14c,14d} Y. L. Liu¹⁰⁵ Y. W. Liu^{62a} M. Livan^{72a,72b} J. Llorente Merino¹⁴¹
 S. L. Lloyd⁹³ E. M. Lobodzinska⁴⁸ P. Loch⁷ S. Loffredo^{75a,75b} T. Lohse¹⁸ K. Lohwasser¹³⁸ M. Lokajicek¹³⁰
 J. D. Long¹⁶⁰ I. Longarini^{74a,74b} L. Longo^{69a,69b} R. Longo¹⁶⁰ I. Lopez Paz³⁶ A. Lopez Solis⁴⁸ J. Lorenz¹⁰⁸
 N. Lorenzo Martinez⁴ A. M. Lory¹⁰⁸ A. Lösle⁵⁴ X. Lou^{47a,47b} X. Lou^{14a,14d} A. Lounis⁶⁶ J. Love⁶
 P. A. Love⁹⁰ J. J. Lozano Bahilo¹⁶¹ G. Lu^{14a,14d} M. Lu⁷⁹ S. Lu¹²⁷ Y. J. Lu⁶⁵ H. J. Lubatti¹³⁷ C. Luci^{74a,74b}
 F. L. Lucio Alves^{14c} A. Lucotte⁶⁰ F. Luehring⁶⁷ I. Luise¹⁴⁴ O. Lukianchuk⁶⁶ O. Lundberg¹⁴³
 B. Lund-Jensen¹⁴³ N. A. Luongo¹²² M. S. Lutz¹⁵⁰ D. Lynn²⁹ H. Lyons⁹¹ R. Lysak¹³⁰ E. Lytken⁹⁷ F. Lyu^{14a}
 V. Lyubushkin³⁸ T. Lyubushkina³⁸ H. Ma²⁹ L. L. Ma^{62b} Y. Ma⁹⁵ D. M. Mac Donell¹⁶³ G. Maccarrone⁵³
 J. C. MacDonald¹³⁸ R. Madar⁴⁰ W. F. Mader⁵⁰ J. Maeda⁸³ T. Maeno²⁹ M. Maerker⁵⁰ V. Magerl⁵⁴
 J. Magro^{68a,68c} H. Maguire¹³⁸ D. J. Mahon⁴¹ C. Maidantchik^{81b} A. Maio^{129a,129b,129d} K. Maj^{84a}
 O. Majersky^{28a} S. Majewski¹²² N. Makovec⁶⁶ V. Maksimovic¹⁵ B. Malaescu¹²⁶ Pa. Malecki⁸⁵
 V. P. Maleev³⁷ F. Malek⁶⁰ D. Malito^{43b,43a} U. Mallik⁷⁹ C. Malone³² S. Maltezos¹⁰ S. Malyukov³⁸
 J. Mamuzic¹³ G. Mancini⁵³ G. Manco^{72a,72b} J. P. Mandalia⁹³ I. Mandić⁹² L. Manhaes de Andrade Filho^{81a}
 I. M. Maniatis¹⁵¹ M. Manisha¹³⁴ J. Manjarres Ramos⁵⁰ D. C. Mankad¹⁶⁷ K. H. Mankinen⁹⁷ A. Mann¹⁰⁸
 A. Manousos⁷⁸ B. Mansoulie¹³⁴ S. Manzoni³⁶ A. Marantis^{151,aa} G. Marchiori⁵ M. Marcisovsky¹³⁰
 L. Marcoccia^{75a,75b} C. Marcon^{70a,70b} M. Marinescu²⁰ M. Marjanovic¹¹⁹ Z. Marshall^{17a} S. Marti-Garcia¹⁶¹

- T. A. Martin¹⁶⁵ V. J. Martin⁵² B. Martin dit Latour¹⁶ L. Martinelli^{74a,74b} M. Martinez^{13,v} P. Martinez Agullo¹⁶¹
V. I. Martinez Outschoorn¹⁰² P. Martinez Suarez¹³ S. Martin-Haugh¹³³ V. S. Martoiu^{27b} A. C. Martyniuk⁹⁵
A. Marzin³⁶ S. R. Maschek¹⁰⁹ L. Masetti⁹⁹ T. Mashimo¹⁵² J. Masik¹⁰⁰ A. L. Maslennikov³⁷ L. Massa^{23b}
P. Massarotti^{71a,71b} P. Mastrandrea^{73a,73b} A. Mastroberardino^{43b,43a} T. Masubuchi¹⁵² T. Mathisen¹⁵⁹
A. Matic¹⁰⁸ N. Matsuzawa¹⁵² J. Maurer^{27b} B. Maček⁹² D. A. Maximov³⁷ R. Mazini¹⁴⁷ I. Maznas¹⁵¹
M. Mazza¹⁰⁶ S. M. Mazza¹³⁵ C. Mc Ginn^{29,h} J. P. Mc Gowan¹⁰³ S. P. Mc Kee¹⁰⁵ T. G. McCarthy¹⁰⁹
W. P. McCormack^{17a} E. F. McDonald¹⁰⁴ A. E. McDougall¹¹³ J. A. Mcfayden¹⁴⁵ G. Mchedlidze^{148b}
R. P. McKenzie^{33g} T. C. McLachlan⁴⁸ D. J. McLaughlin⁹⁵ K. D. McLean¹⁶³ S. J. McMahon¹³³
P. C. McNamara¹⁰⁴ R. A. McPherson^{163,p} J. E. Mdhluli^{33g} S. Meehan³⁶ T. Megy⁴⁰ S. Mehlhase¹⁰⁸
A. Mehta⁹¹ B. Meirose⁴⁵ D. Melini¹⁴⁹ B. R. Mellado Garcia^{33g} A. H. Melo⁵⁵ F. Meloni⁴⁸
E. D. Mendes Gouveia^{129a} A. M. Mendes Jacques Da Costa²⁰ H. Y. Meng¹⁵⁴ L. Meng⁹⁰ S. Menke¹⁰⁹
M. Mentink³⁶ E. Meoni^{43b,43a} C. Merlassino¹²⁵ L. Merola^{71a,71b} C. Meroni^{70a} G. Merz¹⁰⁵ O. Meshkov³⁷
J. K. R. Meshreki¹⁴⁰ J. Metcalfe⁶ A. S. Mete⁶ C. Meyer⁶⁷ J-P. Meyer¹³⁴ M. Michetti¹⁸ R. P. Middleton¹³³
L. Mijović⁵² G. Mikenberg¹⁶⁷ M. Mikesikova¹³⁰ M. Mikuž⁹² H. Mildner¹³⁸ A. Milic¹⁵⁴ C. D. Milke⁴⁴
D. W. Miller³⁹ L. S. Miller³⁴ A. Milov¹⁶⁷ D. A. Milstead^{47a,47b} T. Min^{14c} A. A. Minaenko³⁷ I. A. Minashvili^{148b}
L. Mince⁵⁹ A. I. Mincer¹¹⁶ B. Mindur^{84a} M. Mineev³⁸ Y. Minegishi¹⁵² Y. Mino⁸⁶ L. M. Mir¹³
M. Miralles Lopez¹⁶¹ M. Mironova¹²⁵ T. Mitani¹⁶⁶ A. Mitra¹⁶⁵ V. A. Mitsou¹⁶¹ O. Miu¹⁵⁴ P. S. Miyagawa⁹³
Y. Miyazaki⁸⁸ A. Mizukami⁸² J. U. Mjörnmark⁹⁷ T. Mkrtchyan^{63a} M. Mlynarikova¹¹⁴ T. Moa^{47a,47b}
S. Mobius⁵⁵ K. Mochizuki¹⁰⁷ P. Moder⁴⁸ P. Mogg¹⁰⁸ A. F. Mohammed^{14a,14d} S. Mohapatra⁴¹
G. Mokgatitswane^{33g} B. Mondal¹⁴⁰ S. Mondal¹³¹ K. Mönig⁴⁸ E. Monnier¹⁰¹ L. Monsonis Romero¹⁶¹
J. Montejo Berlingen³⁶ M. Montella¹¹⁸ F. Monticelli⁸⁹ N. Morange⁶⁶ A. L. Moreira De Carvalho^{129a}
M. Moreno Llácer¹⁶¹ C. Moreno Martinez¹³ P. Morettini^{57b} S. Morgenstern¹⁶⁵ M. Morii⁶¹ M. Morinaga¹⁵²
V. Morisbak¹²⁴ A. K. Morley³⁶ F. Morodei^{74a,74b} L. Morvaj³⁶ P. Moschovakos³⁶ B. Moser³⁶ M. Mosidze^{148b}
T. Moskalets⁵⁴ P. Moskvitina¹¹² J. Moss^{31,cc} E. J. W. Moyse¹⁰² S. Muanza¹⁰¹ J. Mueller¹²⁸
D. Muenstermann⁹⁰ R. Müller¹⁹ G. A. Mullier⁹⁷ J. J. Mullin¹²⁷ D. P. Mungo^{70a,70b} J. L. Munoz Martinez¹³
D. Munoz Perez¹⁶¹ F. J. Munoz Sanchez¹⁰⁰ M. Murin¹⁰⁰ W. J. Murray^{165,133} A. Murrone^{70a,70b} J. M. Muse¹¹⁹
M. Muškinja^{17a} C. Mwewa²⁹ A. G. Myagkov^{37,1} A. J. Myers⁸ A. A. Myers¹²⁸ G. Myers⁶⁷ M. Myska¹³¹
B. P. Nachman^{17a} O. Nackenhorst⁴⁹ A. Nag⁵⁰ K. Nagai¹²⁵ K. Nagano⁸² J. L. Nagle^{29,h} E. Nagy¹⁰¹
A. M. Nairz³⁶ Y. Nakahama⁸² K. Nakamura⁸² H. Nanjo¹²³ R. Narayan⁴⁴ E. A. Narayanan¹¹¹ I. Naryshkin³⁷
M. Naseri³⁴ C. Nass²⁴ G. Navarro^{22a} J. Navarro-Gonzalez¹⁶¹ R. Nayak¹⁵⁰ P. Y. Nechaeva³⁷ F. Nechansky⁴⁸
T. J. Neep²⁰ A. Negri^{72a,72b} M. Negrini^{23b} C. Nellist¹¹² C. Nelson¹⁰³ K. Nelson¹⁰⁵ S. Nemecek¹³⁰
M. Nessi^{36,dd} M. S. Neubauer¹⁶⁰ F. Neuhaus⁹⁹ J. Neundorff⁴⁸ R. Newhouse¹⁶² P. R. Newman²⁰ C. W. Ng¹²⁸
Y. S. Ng¹⁸ Y. W. Y. Ng¹⁵⁸ B. Ngair^{35e} H. D. N. Nguyen¹⁰⁷ R. B. Nickerson¹²⁵ R. Nicolaidou¹³⁴ J. Nielsen¹³⁵
M. Niemeyer⁵⁵ N. Nikiforou³⁶ V. Nikolaenko^{37,1} I. Nikolic-Audit¹²⁶ K. Nikolopoulos²⁰ P. Nilsson²⁹
H. R. Nindhito⁵⁶ A. Nisati^{74a} N. Nishu² R. Nisius¹⁰⁹ J-E. Nitschke⁵⁰ E. K. Nkadimeng^{33g}
S. J. Noacco Rosende⁸⁹ T. Nobe¹⁵² D. L. Noel³² Y. Noguchi⁸⁶ T. Nommensen¹⁴⁶ M. A. Nomura²⁹
M. B. Norfolk¹³⁸ R. R. B. Norisam⁹⁵ B. J. Norman³⁴ J. Novak⁹² T. Novak⁴⁸ O. Novgorodova⁵⁰
L. Novotny¹³¹ R. Novotny¹¹¹ L. Nozka¹²¹ K. Ntekas¹⁵⁸ E. Nurse⁹⁵ F. G. Oakham^{34,e} J. Ocariz¹²⁶ A. Ochi⁸³
I. Ochoa^{129a} S. Oerdek¹⁵⁹ A. Ogrodnik^{84a} A. Oh¹⁰⁰ C. C. Ohm¹⁴³ H. Oide¹⁵³ R. Oishi¹⁵² M. L. Ojeda⁴⁸
Y. Okazaki⁸⁶ M. W. O'Keefe⁹¹ Y. Okumura¹⁵² A. Olariu^{27b} L. F. Oleiro Seabra^{129a} S. A. Olivares Pino^{136e}
D. Oliveira Damazio²⁹ D. Oliveira Goncalves^{81a} J. L. Oliver¹⁵⁸ M. J. R. Olsson¹⁵⁸ A. Olszewski⁸⁵
J. Olszowska^{85,a} Ö. Ö. Öncel⁵⁴ D. C. O'Neil¹⁴¹ A. P. O'Neill¹⁹ A. Onofre^{129a,129e} P. U. E. Onyisi¹¹
M. J. Oreglia³⁹ G. E. Orellana⁸⁹ D. Orestano^{76a,76b} N. Orlando¹³ R. S. Orr¹⁵⁴ V. O'Shea⁵⁹ R. Ospanov^{62a}
G. Otero y Garzon³⁰ H. Otono⁸⁸ P. S. Ott^{63a} G. J. Ottino^{17a} M. Ouchrif^{35d} J. Ouellette^{29,h} F. Ould-Saada¹²⁴
M. Owen⁵⁹ R. E. Owen¹³³ K. Y. Oyulmaz^{21a} V. E. Ozcan^{21a} N. Ozturk⁸ S. Ozturk^{21d} J. Pacalt¹²¹
H. A. Pacey³² K. Pachal⁵¹ A. Pacheco Pages¹³ C. Padilla Aranda¹³ G. Padovano^{74a,74b} S. Pagan Griso^{17a}
G. Palacino⁶⁷ A. Palazzo^{69a,69b} S. Palazzo⁵² S. Palestini³⁶ M. Palka^{84b} J. Pan¹⁷⁰ T. Pan^{64a} D. K. Panchal¹¹

- C. E. Pandini¹¹³ J. G. Panduro Vazquez⁹⁴ H. Pang^{14b} P. Pani⁴⁸ G. Panizzo^{68a,68c} L. Paolozzi⁵⁶
 C. Papadatos¹⁰⁷ S. Parajuli⁴⁴ A. Paramonov⁶ C. Paraskevopoulos¹⁰ D. Paredes Hernandez^{64b} T. H. Park¹⁵⁴
 M. A. Parker³² F. Parodi^{57b,57a} E. W. Parrish¹¹⁴ V. A. Parrish⁵² J. A. Parsons⁴¹ U. Parzefall⁵⁴
 B. Pascual Dias¹⁰⁷ L. Pascual Dominguez¹⁵⁰ V. R. Pascuzzi^{17a} F. Pasquali¹¹³ E. Pasqualucci^{74a} S. Passaggio^{57b}
 F. Pastore⁹⁴ P. Pasuwan^{47a,47b} J. R. Pater¹⁰⁰ J. Patton⁹¹ T. Pauly³⁶ J. Pearkes¹⁴² M. Pedersen¹²⁴ R. Pedro^{129a}
 S. V. Peleganchuk³⁷ O. Penc³⁶ C. Peng^{64b} H. Peng^{62a} K. E. Penski¹⁰⁸ M. Penzin³⁷ B. S. Peralva^{81d}
 A. P. Pereira Peixoto⁶⁰ L. Pereira Sanchez^{47a,47b} D. V. Perepelitsa^{29,h} E. Perez Codina^{155a} M. Perganti¹⁰
 L. Perini^{70a,70b,a} H. Pernegger³⁶ S. Perrella³⁶ A. Perrevoort¹¹² O. Perrin⁴⁰ K. Peters⁴⁸ R. F. Y. Peters¹⁰⁰
 B. A. Petersen³⁶ T. C. Petersen⁴² E. Petit¹⁰¹ V. Petousis¹³¹ C. Petridou¹⁵¹ A. Petrukhin¹⁴⁰ M. Pettee^{17a}
 N. E. Pettersson³⁶ A. Petukhov³⁷ K. Petukhova¹³² A. Peyaud¹³⁴ R. Pezoa^{136f} L. Pezzotti³⁶ G. Pezzullo¹⁷⁰
 T. Pham¹⁰⁴ P. W. Phillips¹³³ M. W. Phipps¹⁶⁰ G. Piacquadio¹⁴⁴ E. Pianori^{17a} F. Piazza^{70a,70b} R. Piegai³⁰
 D. Pietreanu^{27b} A. D. Pilkington¹⁰⁰ M. Pinamonti^{68a,68c} J. L. Pinfold² B. C. Pinheiro Pereira^{129a}
 C. Pitman Donaldson⁹⁵ D. A. Pizzi³⁴ L. Pizzimento^{75a,75b} A. Pizzini¹¹³ M.-A. Pleier²⁹ V. Plesanovs⁵⁴
 V. Pleskot¹³² E. Plotnikova³⁸ G. Poddar⁴ R. Poettgen⁹⁷ L. Poggioli¹²⁶ I. Pogrebnyak¹⁰⁶ D. Pohl²⁴
 I. Pokharel⁵⁵ S. Polacek¹³² G. Polesello^{72a} A. Poley^{141,155a} R. Polifka¹³¹ A. Polini^{23b} C. S. Pollard¹²⁵
 Z. B. Pollock¹¹⁸ V. Polychronakos²⁹ D. Ponomarenko³⁷ L. Pontecorvo³⁶ S. Popa^{27a} G. A. Popeneciu^{27d}
 D. M. Portillo Quintero^{155a} S. Pospisil¹³¹ P. Postolache^{27c} K. Potamianos¹²⁵ I. N. Potrap³⁸ C. J. Potter³²
 H. Potti¹ T. Poulsen⁴⁸ J. Poveda¹⁶¹ G. Pownall⁴⁸ M. E. Pozo Astigarraga³⁶ A. Prades Ibanez¹⁶¹
 M. M. Prapa⁴⁶ D. Price¹⁰⁰ M. Primavera^{69a} M. A. Principe Martin⁹⁸ M. L. Proffitt¹³⁷ N. Proklova³⁷
 K. Prokofiev^{64c} G. Proto^{75a,75b} S. Protopopescu²⁹ J. Proudfoot⁶ M. Przybycien^{84a} J. E. Puddefoot¹³⁸
 D. Pudzha³⁷ P. Puzo⁶⁶ D. Pyatiizbyantseva³⁷ J. Qian¹⁰⁵ Y. Qin¹⁰⁰ T. Qiu⁹³ A. Quadt⁵⁵
 M. Queitsch-Maitland¹⁰⁰ G. Rabanal Bolanos⁶¹ D. Rafanoharana⁵⁴ F. Ragusa^{70a,70b} J. L. Rainbolt³⁹
 J. A. Raine⁵⁶ S. Rajagopalan²⁹ E. Ramakoti³⁷ K. Ran^{48,14d} V. Raskina¹²⁶ D. F. Rassloff^{63a} S. Rave⁹⁹
 B. Ravina⁵⁹ I. Ravinovich¹⁶⁷ M. Raymond³⁶ A. L. Read¹²⁴ N. P. Readioff¹³⁸ D. M. Rebuffi^{72a,72b}
 G. Redlinger²⁹ K. Reeves⁴⁵ J. A. Reidelsturz¹⁶⁹ D. Reikher¹⁵⁰ A. Reiss⁹⁹ A. Rej¹⁴⁰ C. Rembser³⁶
 A. Renardi⁴⁸ M. Renda^{27b} M. B. Rendel¹⁰⁹ A. G. Rennie⁵⁹ S. Resconi^{70a} M. Ressegotti^{57b,57a}
 E. D. Resseguie^{17a} S. Rettie⁹⁵ B. Reynolds¹¹⁸ E. Reynolds^{17a} M. Rezaei Estabragh¹⁶⁹ O. L. Rezanova³⁷
 P. Reznicek¹³² E. Ricci^{77a,77b} R. Richter¹⁰⁹ S. Richter^{47a,47b} E. Richter-Was^{84b} M. Ridel¹²⁶ P. Rieck¹¹⁶
 P. Riedler³⁶ M. Rijssenbeek¹⁴⁴ A. Rimoldi^{72a,72b} M. Rimoldi⁴⁸ L. Rinaldi^{23b,23a} T. T. Rinn²⁹
 M. P. Rinnagel¹⁰⁸ G. Ripellino¹⁴³ I. Riu¹³ P. Rivadeneira⁴⁸ J. C. Rivera Vergara¹⁶³ F. Rizatdinova¹²⁰
 E. Rizvi⁹³ C. Rizzi⁵⁶ B. A. Roberts¹⁶⁵ B. R. Roberts^{17a} S. H. Robertson^{103,p} M. Robin⁴⁸ D. Robinson³²
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 J. T. Roemer¹⁵⁸ A. R. Roepe-Gier¹¹⁹ J. Roggel¹⁶⁹ O. Röhne¹²⁴ R. A. Rojas¹⁶³ B. Roland⁵⁴ C. P. A. Roland⁶⁷
 J. Roloff²⁹ A. Romaniouk³⁷ E. Romano^{72a,72b} M. Romano^{23b} A. C. Romero Hernandez¹⁶⁰ N. Rompotis⁹¹
 L. Roos¹²⁶ S. Rosati^{74a} B. J. Rosser³⁹ E. Rossi⁴ E. Rossi^{71a,71b} L. P. Rossi^{57b} L. Rossini⁴⁸ R. Rosten¹¹⁸
 M. Rotaru^{27b} B. Rottler⁵⁴ D. Rousseau⁶⁶ D. Rousso³² G. Rovelli^{72a,72b} A. Roy¹⁶⁰ A. Rozanov¹⁰¹
 Y. Rozen¹⁴⁹ X. Ruan^{33g} A. Rubio Jimenez¹⁶¹ A. J. Ruby⁹¹ V. H. Ruelas Rivera¹⁸ T. A. Ruggeri¹ F. Rühr⁵⁴
 A. Ruiz-Martinez¹⁶¹ A. Rummler³⁶ Z. Rurikova⁵⁴ N. A. Rusakovich³⁸ H. L. Russell¹⁶³ J. P. Rutherford⁷
 E. M. Rüttinger¹³⁸ K. Rybacki⁹⁰ M. Rybar¹³² E. B. Rye¹²⁴ A. Ryzhov³⁷ J. A. Sabater Iglesias⁵⁶ P. Sabatini¹⁶¹
 L. Sabetta^{74a,74b} H. F.-W. Sadrozinski¹³⁵ F. Safai Tehrani^{74a} B. Safarzadeh Samani¹⁴⁵ M. Safdari¹⁴² S. Saha¹⁰³
 M. Sahinsoy¹⁰⁹ M. Saimpert¹³⁴ M. Saito¹⁵² T. Saito¹⁵² D. Salamani³⁶ G. Salamanna^{76a,76b} A. Salnikov¹⁴²
 J. Salt¹⁶¹ A. Salvador Salas¹³ D. Salvatore^{43b,43a} F. Salvatore¹⁴⁵ A. Salzburger³⁶ D. Sammel⁵⁴
 D. Sampsonidis¹⁵¹ D. Sampsonidou^{62d,62c} J. Sánchez¹⁶¹ A. Sanchez Pineda⁴ V. Sanchez Sebastian¹⁶¹
 H. Sandaker¹²⁴ C. O. Sander⁴⁸ J. A. Sandesara¹⁰² M. Sandhoff¹⁶⁹ C. Sandoval^{22b} D. P. C. Sankey¹³³
 A. Sansoni⁵³ L. Santi^{74a,74b} C. Santoni⁴⁰ H. Santos^{129a,129b} S. N. Santpur^{17a} A. Santra¹⁶⁷ K. A. Saoucha¹³⁸
 J. G. Saraiva^{129a,129d} J. Sardain⁷ O. Sasaki⁸² K. Sato¹⁵⁶ C. Sauer^{63b} F. Sauerburger⁵⁴ E. Sauvan⁴
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 T. Scanlon⁹⁵ J. Schaarschmidt¹³⁷ P. Schacht¹⁰⁹ D. Schaefer³⁹ U. Schäfer⁹⁹ A. C. Schaffer⁶⁶ D. Schaile¹⁰⁸

R. D. Schamberger¹⁴⁴ E. Schanet¹⁰⁸ C. Scharf¹⁸ V. A. Schegelsky³⁷ D. Scheirich¹³² F. Schenck¹⁸
 M. Schernau¹⁵⁸ C. Scheulen⁵⁵ C. Schiavi^{57b,57a} Z. M. Schillaci²⁶ E. J. Schioppa^{69a,69b} M. Schioppa^{43b,43a}
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 A. Schoening^{63b} P. G. Scholer⁵⁴ E. Schopf¹²⁵ M. Schott⁹⁹ J. Schovancova³⁶ S. Schramm⁵⁶ F. Schroeder¹⁶⁹
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 D. Sengupta⁵⁶ V. Senthilkumar¹⁶¹ L. Serin⁶⁶ L. Serkin^{68a,68b} M. Sessa^{76a,76b} H. Severini¹¹⁹ S. Sevova¹⁴²
 F. Sforza^{57b,57a} A. Sfyrla⁵⁶ E. Shabalina⁵⁵ R. Shaheen¹⁴³ J. D. Shahinian¹²⁷ N. W. Shaikh^{47a,47b}
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 P. B. Shatalov³⁷ K. Shaw¹⁴⁵ S. M. Shaw¹⁰⁰ Q. Shen^{62c} P. Sherwood⁹⁵ L. Shi⁹⁵ C. O. Shimmin¹⁷⁰
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 N. D. Simpson⁹⁷ S. Simsek^{21d} S. Sindhu⁵⁵ P. Sinervo¹⁵⁴ V. Sinetckii³⁷ S. Singh¹⁴¹ S. Singh¹⁵⁴ S. Sinha⁴⁸
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 P. Skubic¹¹⁹ M. Slawinska⁸⁵ V. Smakhtin¹⁶⁷ B. H. Smart¹³³ J. Smiesko¹³² S. Yu. Smirnov³⁷ Y. Smirnov³⁷
 L. N. Smirnova^{37,1} O. Smirnova⁹⁷ A. C. Smith⁴¹ E. A. Smith³⁹ H. A. Smith¹²⁵ J. L. Smith⁹¹ R. Smith¹⁴²
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 M. Stankaityte¹²⁵ B. Stapf⁴⁸ E. A. Starchenko³⁷ G. H. Stark¹³⁵ J. Stark^{101,ii} D. M. Starko^{155b} P. Staroba¹³⁰
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 A. L. Steinhebel¹²² B. Stelzer^{141,155a} H. J. Stelzer¹²⁸ O. Stelzer-Chilton^{155a} H. Stenzel⁵⁸ T. J. Stevenson¹⁴⁵
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 J. Strandberg¹⁴³ S. Strandberg^{47a,47b} M. Strauss¹¹⁹ T. Strebler¹⁰¹ P. Strizenec^{28b} R. Ströhmer¹⁶⁴
 D. M. Strom¹²² L. R. Strom⁴⁸ R. Stroynowski⁴⁴ A. Strubig^{47a,47b} S. A. Stucci²⁹ B. Stugu¹⁶ J. Stupak¹¹⁹
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 O. Sunneborn Gudnadottir¹⁵⁹ M. R. Sutton¹⁴⁵ M. Svatos¹³⁰ M. Swiatlowski^{155a} T. Swirski¹⁶⁴ I. Sykora^{28a}
 M. Sykora¹³² T. Sykora¹³² D. Ta⁹⁹ K. Tackmann^{48,ff} A. Taffard¹⁵⁸ R. Tafirout^{155a} J. S. Tafoya Vargas⁶⁶
 R. H. M. Taibah¹²⁶ R. Takashima⁸⁷ K. Takeda⁸³ E. P. Takeva⁵² Y. Takubo⁸² M. Talby¹⁰¹ A. A. Talyshv³⁷
 K. C. Tam^{64b} N. M. Tamir¹⁵⁰ A. Tanaka¹⁵² J. Tanaka¹⁵² R. Tanaka⁶⁶ M. Tanasini^{57b,57a} J. Tang^{62c} Z. Tao¹⁶²
 S. Tapia Araya⁸⁰ S. Tapprogge⁹⁹ A. Tarek Abouelfadl Mohamed¹⁰⁶ S. Tarem¹⁴⁹ K. Tariq^{62b} G. Tarna^{27b}
 G. F. Tartarelli^{70a} P. Tas¹³² M. Tasevsky¹³⁰ E. Tassi^{43b,43a} A. C. Tate¹⁶⁰ G. Tateno¹⁵² Y. Tayalati^{35e}
 G. N. Taylor¹⁰⁴ W. Taylor^{155b} H. Teagle⁹¹ A. S. Tee¹⁶⁸ R. Teixeira De Lima¹⁴² P. Teixeira-Dias⁹⁴ J. J. Teoh¹⁵⁴
 K. Terashi¹⁵² J. Terron⁹⁸ S. Terzo¹³ M. Testa⁵³ R. J. Teuscher^{154,p} A. Thaler⁷⁸ O. Theiner⁵⁶
 N. Themistokleous⁵² T. Thevenaux-Pelzer¹⁸ O. Thielmann¹⁶⁹ D. W. Thomas⁹⁴ J. P. Thomas²⁰
 E. A. Thompson⁴⁸ P. D. Thompson²⁰ E. Thomson¹²⁷ E. J. Thorpe⁹³ Y. Tian⁵⁵ V. Tikhomirov^{37,1}
 Yu. A. Tikhonov³⁷ S. Timoshenko³⁷ E. X. L. Ting¹ P. Tipton¹⁷⁰ S. Tisserant¹⁰¹ S. H. Tlou^{33g} A. Tnourji⁴⁰
 K. Todome^{23b,23a} S. Todorova-Nova¹³² S. Todt⁵⁰ M. Togawa⁸² J. Tojo⁸⁸ S. Tokár^{28a} K. Tokushuku⁸²
 R. Tombs³² M. Tomoto^{82,110} L. Tompkins^{142,kk} K. W. Topolnicki^{84b} P. Tornambe¹⁰² E. Torrence¹²²

- H. Torres⁵⁰ E. Torró Pastor¹⁶¹ M. Toscani³⁰ C. Toscirì³⁹ D. R. Tovey¹³⁸ A. Traeet¹⁶ I. S. Trandafir^{27b}
 T. Trefzger¹⁶⁴ A. Tricoli²⁹ I. M. Trigger^{155a} S. Trincaz-Duvold¹²⁶ D. A. Trischuk¹⁶² B. Trocmé⁶⁰
 A. Trofymov⁶⁶ C. Troncon^{70a} L. Truong^{33c} M. Trzebinski⁸⁵ A. Trzupek⁸⁵ F. Tsai¹⁴⁴ M. Tsai¹⁰⁵
 A. Tsiamis¹⁵¹ P. V. Tsiarehska³⁷ S. Tsigaridas^{155a} A. Tsirigotis^{151,aa} V. Tsiskaridze¹⁴⁴ E. G. Tskhadadze^{148a}
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 Y. Tu^{64b} A. Tudorache^{27b} V. Tudorache^{27b} A. N. Tuna³⁶ S. Turchikhin³⁸ I. Turk Cakir^{3a} R. Turra^{70a}
 T. Turtuvshin³⁸ P. M. Tuts⁴¹ S. Tzamarias¹⁵¹ P. Tzanis¹⁰ E. Tzovara⁹⁹ K. Uchida¹⁵² F. Ukegawa¹⁵⁶
 P. A. Ulloa Poblete^{136c} G. Unal³⁶ M. Unal¹¹ A. Undrus²⁹ G. Unel¹⁵⁸ K. Uno¹⁵² J. Urban^{28b} P. Urquijo¹⁰⁴
 G. Usai⁸ R. Ushioda¹⁵³ M. Usman¹⁰⁷ Z. Uysal^{21b} V. Vacek¹³¹ B. Vachon¹⁰³ K. O. H. Vadla¹²⁴
 T. Vafeiadis³⁶ C. Valderanis¹⁰⁸ E. Valdes Santurio^{47a,47b} M. Valente^{155a} S. Valentini^{23b,23a} A. Valero¹⁶¹
 A. Vallier^{101,ii} J. A. Valls Ferrer¹⁶¹ T. R. Van Daalen¹³⁷ P. Van Gemmeren⁶ M. Van Rijnbach^{124,36}
 S. Van Stroud⁹⁵ I. Van Vulpen¹¹³ M. Vanadia^{75a,75b} W. Vandelli³⁶ M. Vandenbroucke¹³⁴ E. R. Vandewall¹²⁰
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 L. Varriale¹⁶¹ K. E. Varvell¹⁴⁶ M. E. Vasile^{27b} L. Vaslin⁴⁰ G. A. Vasquez¹⁶³ F. Vazeille⁴⁰
 T. Vazquez Schroeder³⁶ J. Veatch³¹ V. Vecchio¹⁰⁰ M. J. Veen¹¹³ I. Veliscek¹²⁵ L. M. Veloce¹⁵⁴
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 M. Verissimo De Araujo^{81b} W. Verkerke¹¹³ J. C. Vermeulen¹¹³ C. Vernieri¹⁴² P. J. Verschuuren⁹⁴
 M. Vessella¹⁰² M. L. Vesterbacka¹¹⁶ M. C. Vetterli^{141,e} A. Vgenopoulos¹⁵¹ N. Viaux Maira^{136f} T. Vickey¹³⁸
 O. E. Vickey Boeriu¹³⁸ G. H. A. Viehhauser¹²⁵ L. Vigani^{63b} M. Villa^{23b,23a} M. Villaplana Perez¹⁶¹
 E. M. Villhauer⁵² E. Vilucchi⁵³ M. G. Vinciter³⁴ G. S. Virdee²⁰ A. Vishwakarma⁵² C. Vittori^{23b,23a}
 I. Vivarelli¹⁴⁵ V. Vladimirov¹⁶⁵ E. Voevodina¹⁰⁹ F. Vogel¹⁰⁸ P. Vokac¹³¹ J. Von Ahnen⁴⁸ E. Von Toerne²⁴
 B. Vormwald³⁶ V. Vorobel¹³² K. Vorobev³⁷ M. Vos¹⁶¹ J. H. Vosseveld⁹¹ M. Vozak¹¹³ L. Vozdecky⁹³
 N. Vranjes¹⁵ M. Vranjes Milosavljevic¹⁵ M. Vreeswijk¹¹³ R. Vuillermet³⁶ O. Vujinovic⁹⁹ I. Vukotic³⁹
 S. Wada¹⁵⁶ C. Wagner¹⁰² W. Wagner¹⁶⁹ S. Wahdan¹⁶⁹ H. Wahlberg⁸⁹ R. Wakasa¹⁵⁶ M. Wakida¹¹⁰
 V. M. Walbrecht¹⁰⁹ J. Walder¹³³ R. Walker¹⁰⁸ W. Walkowiak¹⁴⁰ A. M. Wang⁶¹ A. Z. Wang¹⁶⁸ C. Wang^{62a}
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 A. T. Watson²⁰ M. F. Watson²⁰ G. Watts¹³⁷ B. M. Waugh⁹⁵ A. F. Webb¹¹ C. Weber²⁹ M. S. Weber¹⁹
 S. A. Weber³⁴ S. M. Weber^{63a} C. Wei^{62a} Y. Wei¹²⁵ A. R. Weidberg¹²⁵ J. Weingarten⁴⁹ M. Weirich⁹⁹
 C. Weiser⁵⁴ C. J. Wells⁴⁸ T. Wenaus²⁹ B. Wendland⁴⁹ T. Wengler³⁶ N. S. Wenke¹⁰⁹ N. Wermes²⁴
 M. Wessels^{63a} K. Whalen¹²² A. M. Wharton⁹⁰ A. S. White⁶¹ A. White⁸ M. J. White¹ D. Whiteson¹⁵⁸
 L. Wickremasinghe¹²³ W. Wiedenmann¹⁶⁸ C. Wiel⁵⁰ M. Wielers¹³³ N. Wieseotte⁹⁹ C. Wiglesworth⁴²
 L. A. M. Wiik-Fuchs⁵⁴ D. J. Wilbern¹¹⁹ H. G. Wilkens³⁶ D. M. Williams⁴¹ H. H. Williams¹²⁷ S. Williams³²
 S. Willocq¹⁰² P. J. Windischhofer¹²⁵ F. Winklmeier¹²² B. T. Winter⁵⁴ M. Wittgen¹⁴² M. Wobisch⁹⁶ A. Wolf⁹⁹
 R. Wölker¹²⁵ J. Wollrath¹⁵⁸ M. W. Wolter⁸⁵ H. Wolters^{129a,129c} V. W. S. Wong¹⁶² A. F. Wongel⁴⁸
 S. D. Worm⁴⁸ B. K. Wosiek⁸⁵ K. W. Woźniak⁸⁵ K. Wraight⁵⁹ J. Wu^{14a,14d} M. Wu^{64a} S. L. Wu¹⁶⁸ X. Wu⁵⁶
 Y. Wu^{62a} Z. Wu^{134,62a} J. Wuerzinger¹²⁵ T. R. Wyatt¹⁰⁰ B. M. Wynne⁵² S. Xella⁴² L. Xia^{14c} M. Xia^{14b}
 J. Xiang^{64c} X. Xiao¹⁰⁵ M. Xie^{62a} X. Xie^{62a} J. Xiong^{17a} I. Xiotidis¹⁴⁵ D. Xu^{14a} H. Xu^{62a} L. Xu^{62a}
 R. Xu¹²⁷ T. Xu¹⁰⁵ W. Xu¹⁰⁵ Y. Xu^{14b} Z. Xu^{62b} Z. Xu¹⁴² B. Yabsley¹⁴⁶ S. Yacoob^{33a} N. Yamaguchi⁸⁸
 Y. Yamaguchi¹⁵³ H. Yamauchi¹⁵⁶ T. Yamazaki^{17a} Y. Yamazaki⁸³ J. Yan^{62c} S. Yan¹²⁵ Z. Yan²⁵
 H. J. Yang^{62c,62d} H. T. Yang^{17a} S. Yang^{62a} T. Yang^{64c} X. Yang^{62a} X. Yang^{14a} Y. Yang⁴⁴ Z. Yang^{62a,105}
 W.-M. Yao^{17a} Y. C. Yap⁴⁸ H. Ye^{14c} J. Ye⁴⁴ S. Ye²⁹ X. Ye^{62a} Y. Yeh⁹⁵ I. Yeletsikh³⁸ M. R. Yexley⁹⁰
 P. Yin⁴¹ K. Yorita¹⁶⁶ C. J. S. Young⁵⁴ C. Young¹⁴² M. Yuan¹⁰⁵ R. Yuan^{62b,gg} L. Yue⁹⁵ X. Yue^{63a}
 M. Zaazoua^{35e} B. Zabinski⁸⁵ E. Zaid⁵² T. Zakareishvili^{148b} N. Zakharchuk³⁴ S. Zambito⁵⁶
 J. A. Zamora Saa^{136d} J. Zang¹⁵² D. Zanzi⁵⁴ O. Zaplatilek¹³¹ S. V. Zeißner⁴⁹ C. Zeitnitz¹⁶⁹ J. C. Zeng¹⁶⁰
 D. T. Zenger Jr.²⁶ O. Zenin³⁷ T. Ženiš^{28a} S. Zenz⁹³ S. Zerradi^{35a} D. Zerwas⁶⁶ B. Zhang^{14c} D. F. Zhang¹³⁸
 G. Zhang^{14b} J. Zhang⁶ K. Zhang^{14a,14d} L. Zhang^{14c} P. Zhang^{14a,14d} R. Zhang¹⁶⁸ S. Zhang¹⁰⁵ T. Zhang¹⁵²

X. Zhang^{62c}, X. Zhang^{62b}, Z. Zhang^{17a}, Z. Zhang⁶⁶, H. Zhao¹³⁷, P. Zhao⁵¹, T. Zhao^{62b}, Y. Zhao¹³⁵,
Z. Zhao^{62a}, A. Zhemchugov³⁸, Z. Zheng¹⁴², D. Zhong¹⁶⁰, B. Zhou¹⁰⁵, C. Zhou¹⁶⁸, H. Zhou⁷, N. Zhou^{62c},
Y. Zhou⁷, C. G. Zhu^{62b}, C. Zhu^{14a,14d}, H. L. Zhu^{62a}, H. Zhu^{14a}, J. Zhu¹⁰⁵, Y. Zhu^{62c}, Y. Zhu^{62a}, X. Zhuang^{14a},
K. Zhukov³⁷, V. Zhulanov³⁷, N. I. Zimine³⁸, J. Zinsser^{63b}, M. Ziolkowski¹⁴⁰, L. Živković¹⁵, A. Zoccoli^{23b,23a},
K. Zoch⁵⁶, T. G. Zorbas¹³⁸, O. Zormpa⁴⁶, W. Zou⁴¹ and L. Zwalinski³⁶

(ATLAS Collaboration)

¹*Department of Physics, University of Adelaide, Adelaide, Australia*

²*Department of Physics, University of Alberta, Edmonton AB, Canada*

^{3a}*Department of Physics, Ankara University, Ankara, Türkiye*

^{3b}*Division of Physics, TOBB University of Economics and Technology, Ankara, Türkiye*

⁴*LAPP, Univ. Savoie Mont Blanc, CNRS/IN2P3, Annecy, France*

⁵*APC, Université Paris Cité, CNRS/IN2P3, Paris, France*

⁶*High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA*

⁷*Department of Physics, University of Arizona, Tucson, Arizona, USA*

⁸*Department of Physics, University of Texas at Arlington, Arlington, Texas, USA*

⁹*Physics Department, National and Kapodistrian University of Athens, Athens, Greece*

¹⁰*Physics Department, National Technical University of Athens, Zografou, Greece*

¹¹*Department of Physics, University of Texas at Austin, Austin, Texas, USA*

¹²*Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*

¹³*Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain*

^{14a}*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*

^{14b}*Physics Department, Tsinghua University, Beijing, China*

^{14c}*Department of Physics, Nanjing University, Nanjing, China*

^{14d}*University of Chinese Academy of Science (UCAS), Beijing, China*

¹⁵*Institute of Physics, University of Belgrade, Belgrade, Serbia*

¹⁶*Department for Physics and Technology, University of Bergen, Bergen, Norway*

^{17a}*Physics Division, Lawrence Berkeley National Laboratory, Berkeley, California, USA*

^{17b}*University of California, Berkeley, California, USA*

¹⁸*Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany*

¹⁹*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*

²⁰*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*

^{21a}*Department of Physics, Bogazici University, Istanbul, Türkiye*

^{21b}*Department of Physics Engineering, Gaziantep University, Gaziantep, Türkiye*

^{21c}*Department of Physics, Istanbul University, Istanbul, Türkiye*

^{21d}*Istinye University, Sariyer, Istanbul, Türkiye*

^{22a}*Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá, Colombia*

^{22b}*Departamento de Física, Universidad Nacional de Colombia, Bogotá, Colombia*

^{23a}*Dipartimento di Fisica e Astronomia A. Righi, Università di Bologna, Bologna, Italy*

^{23b}*INFN Sezione di Bologna, Italy*

²⁴*Physikalisches Institut, Universität Bonn, Bonn, Germany*

²⁵*Department of Physics, Boston University, Boston, Massachusetts, USA*

²⁶*Department of Physics, Brandeis University, Waltham, Massachusetts, USA*

^{27a}*Transilvania University of Brasov, Brasov, Romania*

^{27b}*Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania*

^{27c}*Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania*

^{27d}*National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca, Romania*

^{27e}*University Politehnica Bucharest, Bucharest, Romania*

^{27f}*West University in Timisoara, Timisoara, Romania*

^{27g}*Faculty of Physics, University of Bucharest, Bucharest, Romania*

^{28a}*Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic*

^{28b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*

²⁹*Physics Department, Brookhaven National Laboratory, Upton, New York, USA*

³⁰*Universidad de Buenos Aires, Facultad de Ciencias Exactas y Naturales, Departamento de Física, y CONICET, Instituto de Física de Buenos Aires (IFIBA), Buenos Aires, Argentina*

- ³¹California State University, California, USA
- ³²Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- ^{33a}Department of Physics, University of Cape Town, Cape Town, South Africa
- ^{33b}iThemba Labs, Western Cape, South Africa
- ^{33c}Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa
- ^{33d}National Institute of Physics, University of the Philippines Diliman (Philippines), Philippines
- ^{33e}University of South Africa, Department of Physics, Pretoria, South Africa
- ^{33f}University of Zululand, KwaDlangezwa, South Africa
- ^{33g}School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³⁴Department of Physics, Carleton University, Ottawa ON, Canada
- ^{35a}Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies—Université Hassan II, Casablanca, Morocco
- ^{35b}Faculté des Sciences, Université Ibn-Tofail, Kénitra, Morocco
- ^{35c}Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco
- ^{35d}LPMR, Faculté des Sciences, Université Mohamed Premier, Oujda, Morocco
- ^{35e}Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ^{35f}Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir, Morocco
- ³⁶CERN, Geneva, Switzerland
- ³⁷Affiliated with an institute covered by a cooperation agreement with CERN
- ³⁸Affiliated with an international laboratory covered by a cooperation agreement with CERN
- ³⁹Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
- ⁴⁰LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ⁴¹Nevis Laboratory, Columbia University, Irvington New York, USA
- ⁴²Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ^{43a}Dipartimento di Fisica, Università della Calabria, Rende, Italy
- ^{43b}INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴⁴Physics Department, Southern Methodist University, Dallas, Texas, USA
- ⁴⁵Physics Department, University of Texas at Dallas, Richardson, Texas, USA
- ⁴⁶National Centre for Scientific Research “Demokritos”, Agia Paraskevi, Greece
- ^{47a}Department of Physics, Stockholm University, Sweden
- ^{47b}Oskar Klein Centre, Stockholm, Sweden
- ⁴⁸Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁹Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany
- ⁵⁰Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁵¹Department of Physics, Duke University, Durham, North Carolina, USA
- ⁵²SUPA—School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁵³INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵⁴Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ⁵⁵II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁵⁶Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ^{57a}Dipartimento di Fisica, Università di Genova, Genova, Italy
- ^{57b}INFN Sezione di Genova, Italy
- ⁵⁸II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁹SUPA—School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁶⁰LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁶¹Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA
- ^{62a}Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei, China
- ^{62b}Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao, China
- ^{62c}School of Physics and Astronomy, Shanghai Jiao Tong University, Key Laboratory for Particle Astrophysics and Cosmology (MOE), SKLPPC, Shanghai, China
- ^{62d}Tsung-Dao Lee Institute, Shanghai, China
- ^{63a}Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ^{63b}Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ^{64a}Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China
- ^{64b}Department of Physics, University of Hong Kong, Hong Kong, China

- ^{64c}*Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
- ⁶⁵*Department of Physics, National Tsing Hua University, Hsinchu, Taiwan*
- ⁶⁶*IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France*
- ⁶⁷*Department of Physics, Indiana University, Bloomington, Indiana, USA*
- ^{68a}*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
- ^{68b}*ICTP, Trieste, Italy*
- ^{68c}*Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine, Italy*
- ^{69a}*INFN Sezione di Lecce, Italy*
- ^{69b}*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
- ^{70a}*INFN Sezione di Milano, Italy*
- ^{70b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
- ^{71a}*INFN Sezione di Napoli, Italy*
- ^{71b}*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
- ^{72a}*INFN Sezione di Pavia, Italy*
- ^{72b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- ^{73a}*INFN Sezione di Pisa, Italy*
- ^{73b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ^{74a}*INFN Sezione di Roma, Italy*
- ^{74b}*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
- ^{75a}*INFN Sezione di Roma Tor Vergata, Italy*
- ^{75b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{76a}*INFN Sezione di Roma Tre, Italy*
- ^{76b}*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
- ^{77a}*INFN-TIFPA, Italy*
- ^{77b}*Università degli Studi di Trento, Trento, Italy*
- ⁷⁸*Universität Innsbruck, Department of Astro and Particle Physics, Innsbruck, Austria*
- ⁷⁹*University of Iowa, Iowa City, Iowa, USA*
- ⁸⁰*Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA*
- ^{81a}*Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- ^{81b}*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- ^{81c}*Instituto de Física, Universidade de São Paulo, São Paulo, Brazil*
- ^{81d}*Rio de Janeiro State University, Rio de Janeiro, Brazil*
- ⁸²*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
- ⁸³*Graduate School of Science, Kobe University, Kobe, Japan*
- ^{84a}*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- ^{84b}*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- ⁸⁵*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- ⁸⁶*Faculty of Science, Kyoto University, Kyoto, Japan*
- ⁸⁷*Kyoto University of Education, Kyoto, Japan*
- ⁸⁸*Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan*
- ⁸⁹*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
- ⁹⁰*Physics Department, Lancaster University, Lancaster, United Kingdom*
- ⁹¹*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
- ⁹²*Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia*
- ⁹³*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
- ⁹⁴*Department of Physics, Royal Holloway University of London, Egham, United Kingdom*
- ⁹⁵*Department of Physics and Astronomy, University College London, London, United Kingdom*
- ⁹⁶*Louisiana Tech University, Ruston, Louisiana, USA*
- ⁹⁷*Fysiska institutionen, Lunds universitet, Lund, Sweden*
- ⁹⁸*Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain*
- ⁹⁹*Institut für Physik, Universität Mainz, Mainz, Germany*
- ¹⁰⁰*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
- ¹⁰¹*CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France*
- ¹⁰²*Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA*
- ¹⁰³*Department of Physics, McGill University, Montreal QC, Canada*
- ¹⁰⁴*School of Physics, University of Melbourne, Victoria, Australia*

- ¹⁰⁵*Department of Physics, University of Michigan, Ann Arbor, Michigan, USA*
- ¹⁰⁶*Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA*
- ¹⁰⁷*Group of Particle Physics, University of Montreal, Montreal QC, Canada*
- ¹⁰⁸*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
- ¹⁰⁹*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
- ¹¹⁰*Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan*
- ¹¹¹*Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA*
- ¹¹²*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen, Netherlands*
- ¹¹³*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
- ¹¹⁴*Department of Physics, Northern Illinois University, DeKalb, Illinois, USA*
- ^{115a}*New York University Abu Dhabi, Abu Dhabi, United Arab Emirates*
- ^{115b}*United Arab Emirates University, Al Ain, United Arab Emirates*
- ^{115c}*University of Sharjah, Sharjah, United Arab Emirates*
- ¹¹⁶*Department of Physics, New York University, New York, New York, USA*
- ¹¹⁷*Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan*
- ¹¹⁸*The Ohio State University, Columbus, Ohio, USA*
- ¹¹⁹*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA*
- ¹²⁰*Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA*
- ¹²¹*Palacký University, Joint Laboratory of Optics, Olomouc, Czech Republic*
- ¹²²*Institute for Fundamental Science, University of Oregon, Eugene, Oregon, USA*
- ¹²³*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹²⁴*Department of Physics, University of Oslo, Oslo, Norway*
- ¹²⁵*Department of Physics, Oxford University, Oxford, United Kingdom*
- ¹²⁶*LPNHE, Sorbonne Université, Université Paris Cité, CNRS/IN2P3, Paris, France*
- ¹²⁷*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA*
- ¹²⁸*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA*
- ^{129a}*Laboratório de Instrumentação e Física Experimental de Partículas—LIP, Lisboa, Portugal*
- ^{129b}*Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal*
- ^{129c}*Departamento de Física, Universidade de Coimbra, Coimbra, Portugal*
- ^{129d}*Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal*
- ^{129e}*Departamento de Física, Universidade do Minho, Braga, Portugal*
- ^{129f}*Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain), Spain*
- ^{129g}*Instituto Superior Técnico, Universidade de Lisboa, Lisboa, Portugal*
- ¹³⁰*Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic*
- ¹³¹*Czech Technical University in Prague, Prague, Czech Republic*
- ¹³²*Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic*
- ¹³³*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ¹³⁴*IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France*
- ¹³⁵*Santa Cruz Institute for Particle Physics, University of California, Santa Cruz, California, USA*
- ^{136a}*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- ^{136b}*Millennium Institute for Subatomic physics at high energy frontier (SAPHIR), Santiago, Chile*
- ^{136c}*Instituto de Investigación Multidisciplinario en Ciencia y Tecnología, y Departamento de Física, Universidad de La Serena, Chile*
- ^{136d}*Universidad Andres Bello, Department of Physics, Santiago, Chile*
- ^{136e}*Instituto de Alta Investigación, Universidad de Tarapacá, Arica, Chile*
- ^{136f}*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ¹³⁷*Department of Physics, University of Washington, Seattle, Washington, USA*
- ¹³⁸*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹³⁹*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁴⁰*Department Physik, Universität Siegen, Siegen, Germany*
- ¹⁴¹*Department of Physics, Simon Fraser University, Burnaby BC, Canada*
- ¹⁴²*SLAC National Accelerator Laboratory, Stanford, California, USA*
- ¹⁴³*Department of Physics, Royal Institute of Technology, Stockholm, Sweden*
- ¹⁴⁴*Departments of Physics and Astronomy, Stony Brook University, Stony Brook, New York, USA*
- ¹⁴⁵*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
- ¹⁴⁶*School of Physics, University of Sydney, Sydney, Australia*
- ¹⁴⁷*Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ^{148a}*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
- ^{148b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
- ^{148c}*University of Georgia, Tbilisi, Georgia*
- ¹⁴⁹*Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel*

- ¹⁵⁰*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
- ¹⁵¹*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
- ¹⁵²*International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan*
- ¹⁵³*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
- ¹⁵⁴*Department of Physics, University of Toronto, Toronto ON, Canada*
- ^{155a}*TRIUMF, Vancouver BC, Canada*
- ^{155b}*Department of Physics and Astronomy, York University, Toronto ON, Canada*
- ¹⁵⁶*Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
- ¹⁵⁷*Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA*
- ¹⁵⁸*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
- ¹⁵⁹*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
- ¹⁶⁰*Department of Physics, University of Illinois, Urbana, Illinois, USA*
- ¹⁶¹*Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia—CSIC, Valencia, Spain*
- ¹⁶²*Department of Physics, University of British Columbia, Vancouver BC, Canada*
- ¹⁶³*Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada*
- ¹⁶⁴*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany*
- ¹⁶⁵*Department of Physics, University of Warwick, Coventry, United Kingdom*
- ¹⁶⁶*Waseda University, Tokyo, Japan*
- ¹⁶⁷*Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot, Israel*
- ¹⁶⁸*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
- ¹⁶⁹*Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
- ¹⁷⁰*Department of Physics, Yale University, New Haven, Connecticut, USA*

^aDeceased.

^bAlso at Department of Physics, King's College London, London, United Kingdom.

^cAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^dAlso at Lawrence Livermore National Laboratory, Livermore, United States of America.

^eAlso at TRIUMF, Vancouver BC, Canada.

^fAlso at Department of Physics, University of Thessaly, Greece.

^gAlso at Department of Physics, University of Fribourg, Fribourg, Switzerland.

^hAlso at University of Colorado Boulder, Department of Physics, Colorado, United States of America.

ⁱAlso at Department of Physics and Astronomy, University of Louisville, Louisville, KY, United States of America.

^jAlso at Department of Physics, Westmont College, Santa Barbara, United States of America.

^kAlso at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain.

^lAlso at Affiliated with an institute covered by a cooperation agreement with CERN.

^mAlso at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China.

ⁿAlso at Department of Physics, Ben Gurion University of the Negev, Beer Sheva, Israel.

^oAlso at Università di Napoli Parthenope, Napoli, Italy.

^pAlso at Institute of Particle Physics (IPP), Canada.

^qAlso at Bruno Kessler Foundation, Trento, Italy.

^rAlso at Borough of Manhattan Community College, City University of New York, New York NY, United States of America.

^sAlso at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

^tAlso at Centro Studi e Ricerche Enrico Fermi, Italy.

^uAlso at Department of Physics, California State University, East Bay, United States of America.

^vAlso at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

^wAlso at University of Chinese Academy of Sciences (UCAS), Beijing, China.

^xAlso at Yeditepe University, Physics Department, Istanbul, Türkiye.

^yAlso at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.

^zAlso at CERN, Geneva, Switzerland.

^{aa}Also at Hellenic Open University, Patras, Greece.

^{bb}Also at Center for High Energy Physics, Peking University, China.

^{cc}Also at Department of Physics, California State University, Sacramento, United States of America.

^{dd}Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland.

^{ee}Also at Washington College, Maryland, United States of America.

^{ff}Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

^{gg}Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America.

^{hh}Also at Physics Department, An-Najah National University, Nablus, Palestine.

ⁱⁱ Also at L2IT, Université de Toulouse, CNRS/IN2P3, UPS, Toulouse, France.

^{jj} Also at Technical University of Munich, Munich, Germany.

^{kk} Also at Department of Physics, Stanford University, Stanford CA, United States of America.

^{ll} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria.