

Identification of cavities and inclusions in linear elasticity with a phase-field approach

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Abstract

In this paper we deal with the inverse problem of determining cavities and inclusions embedded in a linear elastic isotropic medium from boundary displacement's measurements. For, we consider a constrained minimization problem involving a boundary quadratic misfit functional with a regularization term that penalizes the perimeter of the cavity or inclusion to be identified. Then using a phase field approach we derive a robust algorithm for the reconstruction of elastic inclusions and of cavities modelled as inclusions with a very small elasticity tensor.

1 Introduction

The focus of this paper is the reconstruction of cavities and inclusions embedded in an elastic isotropic medium by means of boundary tractions and displacements. Identification of defects from boundary measurements plays an important role in non-destructive testing for damage assessment of mechanical specimens, which are possibly defective due to the presence of interior voids or cavities appearing during the manufacturing process, see, for instance, [33, 47, 55, 63] for possible applications to 3D-printing and additive manufacturing. This kind of inverse problems has application also in medical imaging and in particular in elastography, a modality mapping the elastic properties and stiffness of soft tissue, [6, 7, 8, 31, 59, 60, 64] (to cite a few), and in reflection seismology [20, 62], a non invasive technique used by the oil and gas industry to map petroleum deposits in the Earth's upper crust and based on seismic data from land acquisition, see for example [61]. We also mention some applications in volcanology, see for example [9, 10, 58] and references therein.

The underlying mathematical model is the following: Consider a bounded domain $\Omega \subset \mathbb{R}^d$, with $d = 2, 3$, representing the region occupied by an elastic isotropic medium and let $\partial\Omega = \Sigma_D \cup \Sigma_N$, with Σ_D closed. Let the displacement field u be solution to the following

AMS 2020 subject classifications: 35R30, 65N21, 74G75

Key words and phrases: inverse problems, cavity, phase-field, linear elasticity, primal dual active set method

mixed boundary value problem for the Lamé system of linearized elasticity:

$$\begin{cases} \operatorname{div}(\mathbb{C}_0 \widehat{\nabla} u) = 0 & \text{in } \Omega \setminus \overline{C}, \\ (\mathbb{C}_0 \widehat{\nabla} u)n = 0 & \text{on } \partial C, \\ (\mathbb{C}_0 \widehat{\nabla} u)\nu = g & \text{on } \Sigma_N, \\ u = 0 & \text{on } \Sigma_D, \end{cases} \quad (1.1)$$

where $C \Subset \Omega$ is a cavity with Lipschitz boundary, and $\widehat{\nabla} u$ is the strain tensor. \mathbb{C}_0 is a fourth-order isotropic elastic tensor, uniformly bounded and strongly convex, and n and ν are the outer unit normal vector to ∂C and Σ_N , respectively. The Neumann boundary datum g is assumed to be in $L^2(\Sigma_N)$.

The forward problem consists in finding the elastic displacement u in the elastic body occupying the region Ω induced by the tractions on Σ_N , given the cavity C . The inverse problem concerns the determination of the cavity C from partial observations of u on the boundary. More precisely, given measurements of the displacement, i.e. $u_{meas} \in L^2(\Sigma_N)$, find C contained in Ω , such that $u|_{\Sigma_N} = u_{meas}$, where $u \in H^1_{\Sigma_D}(\Omega \setminus C)$ is the solution to the forward problem.

It is well known that this problem is severely ill-posed and only a very weak logarithmic conditional stability holds, assuming a-priori $C^{1,\alpha}$ regularity of the unknown cavities [53]. A similar weak stability result holds also in the case of the determination of elastic inclusions, see for example [54]. Hence, in general, the reconstruction of cavities and inclusions turns out to be a challenging issue.

To solve the problem we follow a similar strategy as in [14, 30] and the one in [13] for the reconstruction of conductivity inclusions and cavities respectively. Specifically, we consider the problem of minimizing the functional

$$J(C) = \frac{1}{2} \int_{\Sigma_N} |u(C) - u_{meas}|^2 d\sigma(x) + \alpha \operatorname{Per}(C), \quad (1.2)$$

over a suitable set of cavities of finite perimeter and where $u(C)$ is the solution of (1.1) for a given cavity C , $\operatorname{Per}(C)$ indicates the perimeter of C , and α is a positive regularization parameter.

We first investigate the continuity of solutions to (1.1) with respect to perturbations of the cavity C in the Hausdorff distance topology and prove it using the Mosco convergence, see [21, 22, 37]. Similarly as in [13], continuity then allows us to prove existence of minima of the functional $J(C)$, stability with respect to noisy data and convergence of the minimizers as $\alpha \rightarrow 0$ to the solution of the inverse problem.

In the second part of the paper, we use a suitable phase-field relaxation of the functional J in order to overcome issues arising from non-convexity and non-differentiability. To be more precise, we employ an idea adopted by Bourdin and Chambolle, [18] in the context of topology optimization which consists in filling the cavity with a fictitious elastic material described by an elastic tensor $\mathbb{C}_1 := \delta \mathbb{C}_0$, where δ is a small positive parameter and \mathbb{C}_0 has been extended to the whole domain Ω . In this way, we transform the original inverse problem in the one of reconstructing an elastic inclusion. Then, since the identification of sharp interfaces is in general difficult to be treated numerically, we use a phase-field approach. Instead of binary (i.e., either 0 or 1) phase parameter v describing sharp interfaces between regions with two different materials we use a phase parameter v as a H^1 scalar field, taking values in the interval $[0, 1]$. Then, we approximate the functional J in (1.2) by means of a Ginzburg-Landau type functional (cf. [52])

$$J_{\delta,\varepsilon}(v) := \frac{1}{2} \int_{\Sigma_N} |u_\delta(v) - u_{meas}|^2 d\sigma(x) + \frac{4\alpha}{\pi} \int_{\Omega} \left(\varepsilon |\nabla v|^2 + \frac{1}{\varepsilon} v(1-v) \right) dx, \quad (1.3)$$

where ε is a small positive parameter, $\frac{4}{\pi}$ is a rescaled parameter in the Modica-Mortola relaxation of the perimeter, $u_\delta(v)$ denotes the solution of the modified boundary value

problem:

$$\begin{cases} \operatorname{div}(\mathbb{C}_\delta(v)\widehat{\nabla}u_\delta(v)) = 0 & \text{in } \Omega, \\ (\mathbb{C}_\delta(v)\widehat{\nabla}u_\delta(v))\nu = g & \text{on } \Sigma_N, \\ u_\delta(v) = 0 & \text{on } \Sigma_D, \end{cases} \quad (1.4)$$

where

$$\mathbb{C}_\delta(v) = \mathbb{C}_0 + (\mathbb{C}_1 - \mathbb{C}_0)v, \quad \text{with } \mathbb{C}_1 = \delta\mathbb{C}_0. \quad (1.5)$$

Here \mathbb{C}_0 and \mathbb{C}_1 are the elasticity tensors in $\Omega \setminus C$ and C , respectively. Ideally, the optimal phase variable v should be close to an ideal binary field. In fact, when ε is small the potential term $(\int_\Omega \frac{1}{\varepsilon} v(1-v) dx)$ prevails and the minimum is attained by a phase-field variable which takes mainly values close to 0 and 1 and the transition occurs in a thin layer of thickness of order ε .

The phase-field approach to structural optimization problems has been successfully used by different authors (cf., e.g., [12, 15, 25, 36]), the main advantage being the fact that it allows to handle topology changes as well as nucleation of new holes.

To implement our algorithm in Section 3.2 we provide first order necessary optimality conditions for the minimization problem associated to $J_{\delta,\varepsilon}$ whose discretized version is then employed in Section 4 in order to develop the reconstruction algorithm. Minima of the functional $J_{\delta,\varepsilon}$ exist and the numerical experiments of Section 5 indicate that they are accurate approximations of minima of J , for ε and δ sufficiently small. This fact could be rigorously justified proving that the Γ -convergence, as ε and δ tend to 0, to the functional J holds, but this is still an open issue and will be the subject of a future research. Some attempts along this direction have been done in the scalar case for example in [13, 56, 57].

The literature on reconstruction algorithms for identification of inclusions and cavities in elastostatic, viscoelastic and elastic waves systems is very rich and of big impact. In the case of small elastic inclusions or cavities, asymptotic expansions of the perturbed displacement have been used to detect position, size and shape from boundary measurements, see for example [45] and [8]. The method followed in [5] is based on a shape derivative approach, both for elastic and thermoelastic problems. A topological gradient method has been applied in [24], for the detection of an elastic scatterer, and in [50], for identification of a cavity in time-harmonic wave elastic systems. Ikehata and Itou use the so-called *enclosure method* for the reconstruction of polygonal cavities in an elastostatic setting [42] and of a general cavity in a homogeneous isotropic viscoelastic body [43]. More recently, Doubova and Fernández-Cara proposed an augmented Lagrangian method to identify rigid inclusions in a elastic waves system [31]. Eberle and Harrach applied the *monotonicity method* for the reconstruction of elastic inclusions using the monotonicity property of the Neumann-to-Dirichlet map [32], and in [46] the authors used the method of fundamental solutions for the reconstruction of elastic cavities. For other reconstruction approaches we refer to the review paper [17] and references therein. Identification of cavities and elastic inclusions could be interpreted as a special case of the determination of Lamé parameters from boundary measurements, see for example [7, 41] and [61].

The plan of the paper is the following. In Section 2 we investigate the continuity of the solution to the direct problem with respect to perturbations of the cavity in the Hausdorff topology and then derive the major properties of the misfit functional $J(C)$. In Section 3 we consider the approximation of the cavity with an inclusion of small elasticity tensor, the corresponding misfit functional and its properties. We then introduce its phase-field relaxation and analyze its differentiability and derive necessary optimality conditions related to the phase-field minimization problem. In Section 4 we propose an iterative reconstruction algorithm allowing for the numerical approximation of the solution and prove its convergence properties. Finally, in Section 5 we present some numerical results showing the efficiency and robustness of the proposed reconstruction algorithm.

Notation and geometrical setting

We introduce the principal notation used in the paper.

Notation. We denote scalar quantities, points, and vectors in italics, e.g. x, y and u, v , and fourth-order tensors in blackboard face, e.g. \mathbb{A}, \mathbb{B} .

The symmetric part of a second-order tensor A is denoted by $\widehat{A} := \frac{1}{2}(A + A^T)$, where A^T is the transpose matrix. In particular, $\widehat{\nabla}u$ represents the deformation tensor. We utilize standard notation for inner products, that is, $u \cdot v = \sum_i u_i v_i$, and $A : B = \sum_{i,j} a_{ij} b_{ij}$ (B is a second-order tensor). $|A|$ denotes the norm induced by the inner product on matrices:

$$|A| = \sqrt{A : A}.$$

Domains. To represent locally a boundary as a graph of function, we adopt the notation: $\forall x \in \mathbb{R}^d$, we set $x = (x', x_d)$, where $x' \in \mathbb{R}^{d-1}$, $x_d \in \mathbb{R}$, with $d = 2, 3$. Given $r > 0$, we denote by $B_r(x) \subset \mathbb{R}^d$ the set $B_r(x) := \{(x', x_d) / |x'|^2 + x_d^2 < r^2\}$ and by $B'_r(x') \subset \mathbb{R}^{d-1}$ the set $B'_r(x') := \{x' \in \mathbb{R}^{d-1} / |x'|^2 < r^2\}$.

Definition 1.1 ($C^{0,1}$ regularity).

Let Ω be a bounded domain in \mathbb{R}^d . We say that a portion Σ of $\partial\Omega$ is of Lipschitz class with constants r_0, L_0 , if for any $p \in \Sigma$ there exists a rigid transformation of coordinates under which we have that p is mapped to the origin and

$$\Omega \cap B_{r_0}(0) = \{x \in B_{r_0}(0) : x_d > \psi(x')\},$$

where ψ is a $C^{0,1}$ function on $B'_{r_0}(0) \subset \mathbb{R}^{d-1}$, such that

$$\begin{aligned} \psi(0) &= 0, \\ \|\psi\|_{C^{0,1}(B'_{r_0}(0))} &\leq L_0. \end{aligned}$$

The Hausdorff distance between two sets Ω_1 and Ω_2 is defined by

$$d_H(\Omega_1, \Omega_2) = \max\left\{ \sup_{x \in \Omega_1} \inf_{y \in \Omega_2} \text{dist}(x, y), \sup_{x \in \Omega_2} \inf_{y \in \Omega_1} \text{dist}(x, y) \right\}.$$

Functional setting: Let Ω be a bounded domain. We set

$$BV(\Omega) = \{v \in L^1(\Omega) : TV(v) < \infty\}, \quad (1.6)$$

where

$$TV(v) = \sup \left\{ \int_{\Omega} v \operatorname{div}(\varphi); \quad \varphi \in C_0^1(\Omega), \|\varphi\|_{L^\infty(\Omega)} \leq 1 \right\} \quad (1.7)$$

is the total variation of v . The BV space is endowed with the natural norm $\|v\|_{BV(\Omega)} = \|v\|_{L^1(\Omega)} + TV(v)$. We recall that the perimeter of Ω is defined as

$$\operatorname{Per}(\Omega) = TV(\chi_\Omega), \quad (1.8)$$

where χ_Ω is the characteristic function of the set Ω .

Setting $H_{\partial\Omega}^1(\Omega) := \{v \in H^1(\Omega) : v|_{\partial\Omega} = 0\}$, we recall the following inequalities.

Proposition 1.1. *Let Ω be a bounded Lipschitz domain. For every $v \in H_{\partial\Omega}^1(\Omega)$, there exists a positive constant $\bar{c} = \bar{c}(\Omega)$ such that*

$$\text{(Korn inequality)} \quad \|\nabla v\|_{L^2(\Omega)} \leq \bar{c} \|\widehat{\nabla}v\|_{L^2(\Omega)}. \quad (1.9)$$

$$\text{(Poincaré inequality)} \quad \|v\|_{H^1(\Omega)} \leq \bar{c} \|\nabla v\|_{L^2(\Omega)}. \quad (1.10)$$

Estimates (1.9) and (1.10) hold also in the case where v is zero, in the trace sense, only on a portion of $\partial\Omega$.

2 Elastic problem - detection of a cavity

The focus of this work is the reconstruction of a cavity in an elastic body from boundary measurements using a phase-field approach. We assume that Ω is a bounded domain and that $\partial\Omega := \Sigma_N \cup \Sigma_D$, with $|\Sigma_N| > 0$, $|\Sigma_D| > 0$, Σ_D closed, where $\partial\Omega$ is of Lipschitz class with constants r_0 and L_0 . Denoting by C the cavity, we consider the mixed boundary value problem

$$\begin{cases} \operatorname{div}(\mathbb{C}_0 \widehat{\nabla} u) = 0 & \text{in } \Omega \setminus \overline{C}, \\ (\mathbb{C}_0 \widehat{\nabla} u)n = 0 & \text{on } \partial C, \\ (\mathbb{C}_0 \widehat{\nabla} u)\nu = g & \text{on } \Sigma_N, \\ u = 0 & \text{on } \Sigma_D, \end{cases} \quad (2.1)$$

where n, ν are the outer unit normal vector to ∂C and $\partial\Omega$, respectively. We make the following assumptions.

Assumption 2.1. $\mathbb{C}_0 = \mathbb{C}_0(x)$ is a fourth-order tensor such that

$$(\mathbb{C}_0)_{ijkl}(x) = (\mathbb{C}_0)_{jikl}(x) = (\mathbb{C}_0)_{klij}(x), \quad \forall 1 \leq i, j, k, h \leq d, \text{ and } x \in \Omega.$$

Moreover, \mathbb{C}_0 is assumed to be uniformly bounded and uniformly strongly convex, that is, \mathbb{C}_0 defines a positive-definite quadratic form on symmetric matrices:

$$\mathbb{C}_0(x)\widehat{A} : \widehat{A} \geq \xi_0 |\widehat{A}|^2, \quad \text{a.e in } \Omega,$$

for $\xi_0 > 0$.

Remark 2.1. We require that \mathbb{C}_0 is defined in Ω , and not only in $\Omega \setminus C$, because we employ, in the second part of the paper, a reconstruction algorithm based on the strategy of filling the cavity with a fictitious elastic material.

Assumption 2.2.

$$g \in L^2(\Sigma_N). \quad (2.2)$$

We assume Lipschitz regularity of the cavity (see Definition 1.1), which is a typical requirement to prove uniqueness of the solution to the inverse problem, see [53]. More precisely, we make the following assumption.

Assumption 2.3. Let

$$C \in \mathcal{C} := \{D \subset \overline{\Omega} : \text{compact, simply connected } \partial D \in C^{0,1} \text{ with constants } r_0, L_0 \text{ and } \operatorname{dist}(D, \partial\Omega) \geq d_0 > 0\}.$$

We define

$$\Omega^{d_0/2} = \{x \in \Omega / \operatorname{dist}(x, \partial\Omega) \leq \frac{d_0}{2}\}. \quad (2.3)$$

For the class of admissible sets \mathcal{C} , the following result holds.

Remark 2.2. \mathcal{C} is compact with respect to the Hausdorff topology [29, 51].

Remark 2.3. From now on, we will denote with c any constant possibly depending on Ω , r_0 , L_0 , d , ξ_0 , d_0 , \bar{c} , and on the uniform bounds of the elasticity tensor.

Well-posedness of (2.1) in $H_{\Sigma_D}^1(\Omega \setminus C)$ follows from an application of the Lax-Milgram theorem to the weak formulation of Problem (2.1):

Find $u \in H_{\Sigma_D}^1(\Omega \setminus C)$ solution to

$$\int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u : \widehat{\nabla} \varphi \, dx = \int_{\Sigma_N} g \cdot \varphi \, d\sigma(x), \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega \setminus C), \quad (2.4)$$

(see for example [28]). Moreover, it holds

$$\|u\|_{H^1(\Omega \setminus C)} \leq c \|g\|_{L^2(\Sigma_N)}. \quad (2.5)$$

Choosing $\varphi = u$ in (2.4), the last inequality follows from the strong convexity of the elasticity tensor \mathbb{C}_0 (see Assumption 2.1), from an application of the Korn and Poincaré inequality to the left-hand side of (2.4) (see Proposition 1.1), and from the use of a Cauchy-Schwarz inequality to the right-hand side. In fact,

$$\int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u : \widehat{\nabla} u \, dx \geq c \|\widehat{\nabla} u\|_{L^2(\Omega \setminus C)}^2 \geq c \|\nabla u\|_{L^2(\Omega \setminus C)}^2 \geq c \|u\|_{H^1(\Omega \setminus C)}^2, \quad (2.6)$$

and

$$\left| \int_{\Sigma_N} g \cdot u \, d\sigma(x) \right| \leq \|g\|_{L^2(\Sigma_N)} \|u\|_{L^2(\Sigma_N)} \leq c \|g\|_{L^2(\Sigma_N)} \|u\|_{H^1(\Omega \setminus C)}, \quad (2.7)$$

and so estimate (2.5) follows by (2.6) and (2.7).

Our aim is to tackle the following inverse problem:

Problem 2.1. *Under Assumptions 2.1, 2.2, and 2.3, given $u_{meas} \in L^2(\Sigma_N)$, find $C \in \mathcal{C}$ such that $u|_{\Sigma_N} = u_{meas}$, where $u \in H_{\Sigma_D}^1(\Omega \setminus C)$ solves (2.1).*

It has been proved in [53] (see also [11]) that Problem 2.1 has a unique solution when ∂C is of Lipschitz class. Logarithmic stability estimates have been proved under the assumption of $C^{1,\sigma}$ regularity, $0 < \sigma \leq 1$, on the cavity C , cf. [53].

For the reconstruction of the solution to the inverse problem we consider a standard approach based on the minimization of a quadratic misfit functional, with a Tikhonov regularization penalizing the perimeter of C . More precisely, let

$$\min_{C \in \mathcal{C}} J(C), \quad \text{where } J(C) = \frac{1}{2} \int_{\Sigma_N} |u(C) - u_{meas}|^2 \, d\sigma(x) + \alpha \text{Per}(C), \quad (2.8)$$

where $\alpha > 0$ represents a regularization parameter, $\text{Per}(C)$ the perimeter of the set C , see (1.8), and $u(C) \in H_{\Sigma_D}^1(\Omega \setminus C)$ the solution to (2.4).

2.1 Continuity property of solutions with respect to C

Adapting to our case some known results in literature, see for example [26, 23, 21, 37, 49] and references therein, in this section we will show the continuity of the boundary term in (2.8) with respect to perturbations of the cavity C in the Hausdorff distance.

To this purpose, we recall the definition of Mosco convergence and some of its properties (see [22, 21, 37, 51]). Let X be a reflexive Banach space, and G_n a sequence of closed subspaces of X . We define

$$G' := \{x \in X \mid x = w - \limsup y_{n_k}, \, y_{n_k} \in G_{n_k}, \, n_k \rightarrow +\infty\} \quad (2.9)$$

and

$$G'' := \{x \in X \mid x = s - \liminf y_n, \, y_n \in G_n \text{ for } n \text{ large}\}. \quad (2.10)$$

G', G'' are called the *weak-limsup* and the *strong-liminf* of the sequence G_n in the sense of Mosco.

Definition 2.1. *The sequence G_n converges in the sense of Mosco if $G' = G'' = G$. G is called the Mosco limit of G_n .*

In other words, G_n converges in the sense of Mosco to G when the following two conditions hold:

$$\text{If } u_{n_k} \in G_{n_k} \text{ is such that } u_{n_k} \rightharpoonup u \text{ in } X, \text{ then } u \in G; \quad (2.11)$$

$$\forall u \in G, \exists u_n \in G_n \text{ such that } u_n \rightarrow u \text{ in } X. \quad (2.12)$$

Given Ω and $\Omega \setminus C$, we can identify the Sobolev space $H_{\Sigma_D}^1(\Omega \setminus C)$ with a closed subspace of $L^2(\Omega, \mathbb{R}^{d+d^2})$ through the map

$$\begin{aligned} H_{\Sigma_D}^1(\Omega \setminus C) &\hookrightarrow L^2(\Omega, \mathbb{R}^{d+d^2}) \\ u &\rightarrow (u, \partial_j u_i), \quad \forall i, j = 1, \dots, d \end{aligned} \quad (2.13)$$

with the convention of extending u and ∇u to zero in C . The same identification holds for $\Omega \setminus C_n$, extending u_n and ∇u_n to zero in C_n .

Since we are considering the case of uniform Lipschitz domains, we have the following result, which is an adaptation of Theorem 7.2.7 in [21].

Theorem 2.4. *Let us assume that $C_n, C \subset \Omega$ belong to the class \mathcal{C} . If $C_n \rightarrow C$ in the Hausdorff metric, then $H_{\Sigma_D}^1(\Omega \setminus C_n)$ converges to $H_{\Sigma_D}^1(\Omega \setminus C)$ in the sense of Mosco.*

We can now prove the following continuity result.

Theorem 2.5. *Let $C_n \in \mathcal{C}$ be a sequence of sets converging to C in the Hausdorff metric (cf. Remark 2.2), and let $u(C_n) =: u_n \in H_{\Sigma_D}^1(\Omega \setminus C_n)$, $u(C) =: u \in H_{\Sigma_D}^1(\Omega \setminus C)$ be solutions of (2.4) in $\Omega \setminus C_n$, $\Omega \setminus C$, respectively. Then*

$$\lim_{n \rightarrow +\infty} \int_{\Sigma_N} |u_n - u|^2 d\sigma(x) = 0. \quad (2.14)$$

Proof. Thanks to the uniform Lipschitz regularity of $\partial(\Omega \setminus C_n)$ (and $\partial(\Omega \setminus C)$), we have that the Korn and Poincaré inequalities are uniform with respect to n in $H_{\Sigma_D}^1(\Omega \setminus C_n)$, since they depend only on the Lipschitz constants of the domain $\partial(\Omega \setminus C_n)$, see [2, 27]. Therefore, from (2.4) and (2.5), we have that

$$\|u_n\|_{H^1(\Omega \setminus C_n)} \leq c, \quad (2.15)$$

where c is independent of n .

Hence, from the identification (2.13), we get that $\|u_n\|_{L^2(\Omega, \mathbb{R}^{d+d^2})}$ is uniformly bounded. Up to subsequences, there exists $u^* \in L^2(\Omega, \mathbb{R}^{d+d^2})$ such that

$$u_n \rightharpoonup u^* \text{ in } L^2(\Omega, \mathbb{R}^{d+d^2}).$$

Thanks to Theorem 2.4 and from the first condition of the Mosco convergence applied to $G_n = H_{\Sigma_D}^1(\Omega \setminus C_n)$, $G = H_{\Sigma_D}^1(\Omega \setminus C)$, and $X = L^2(\Omega, \mathbb{R}^{d+d^2})$, see (2.11), we have that $u^* \in H_{\Sigma_D}^1(\Omega \setminus C)$.

Moreover, taking $\varphi \in H_{\Sigma_D}^1(\Omega \setminus C)$, there exists $\varphi_n \in H_{\Sigma_D}^1(\Omega \setminus C_n)$ by (2.12) such that

$$\varphi_n \rightarrow \varphi \text{ in } L^2(\Omega, \mathbb{R}^{d+d^2}). \quad (2.16)$$

Considering the weak formulation for u_n (see (2.4) specialized to the case with $C = C_n$ and $\varphi = \varphi_n$)

$$\int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi_n dx = \int_{\Sigma_N} g \cdot \varphi_n d\sigma(x), \quad (2.17)$$

and since $\varphi_n \in H_{\Sigma_D}^1(\Omega \setminus C_n)$ and $\varphi \in H_{\Sigma_D}^1(\Omega \setminus C)$, it holds

$$\int_{\Sigma_N} g \cdot \varphi_n d\sigma(x) = \int_{\Sigma_N} g \cdot (\varphi_n - \varphi) d\sigma(x) + \int_{\Sigma_N} g \cdot \varphi d\sigma(x).$$

Hence, thanks to Assumption 2.3 and (2.16), we have

$$\begin{aligned} \left| \int_{\Sigma_N} g \cdot (\varphi_n - \varphi) d\sigma(x) \right| &\leq c \|g\|_{L^2(\Sigma_N)} \|\varphi_n - \varphi\|_{L^2(\Sigma_N)} \\ &\leq c \|\varphi_n - \varphi\|_{H_{\Sigma_D}^1(\Omega^{d_0/2})} \rightarrow 0, \end{aligned}$$

as $n \rightarrow +\infty$, where $\Omega^{d_0/2}$ is defined as in (2.3). Therefore,

$$\int_{\Sigma_N} g \cdot \varphi_n d\sigma(x) \rightarrow \int_{\Sigma_N} g \cdot \varphi d\sigma(x), \quad \text{as } n \rightarrow +\infty. \quad (2.18)$$

The term on the left-hand side of (2.17) is equal to

$$\int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi_n dx = \int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} (\varphi_n - \varphi) dx + \int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi dx. \quad (2.19)$$

Then, by (2.15) and (2.16), it follows

$$\left| \int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} (\varphi_n - \varphi) dx \right| \leq c \|\widehat{\nabla} u_n\|_{L^2(\Omega \setminus C_n)} \|\widehat{\nabla} (\varphi_n - \varphi)\|_{L^2(\Omega \setminus C_n)} \rightarrow 0, \quad (2.20)$$

as $n \rightarrow +\infty$. Analogously, for the second integral on the right-hand side of (2.19), using the symmetries of the elasticity tensor, we get

$$\begin{aligned} \int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi dx &= \int_{\Omega \setminus C_n} \widehat{\nabla} u_n : \mathbb{C}_0 \widehat{\nabla} \varphi dx \\ &\rightarrow \int_{\Omega \setminus C} \widehat{\nabla} u^* : \mathbb{C}_0 \widehat{\nabla} \varphi dx = \int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u^* : \widehat{\nabla} \varphi dx, \end{aligned} \quad (2.21)$$

as $n \rightarrow +\infty$. Consequently, using (2.20) and (2.21) in (2.19), we get

$$\int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi_n dx \rightarrow \int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u^* : \widehat{\nabla} \varphi dx, \quad \text{as } n \rightarrow +\infty. \quad (2.22)$$

Therefore, we find that

$$\begin{aligned} \int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u^* : \widehat{\nabla} \varphi dx &= \int_{\Sigma_N} g \cdot \varphi d\sigma(x) \\ &= \int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u : \widehat{\nabla} \varphi dx, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega \setminus C), \end{aligned}$$

where the last equality comes from the weak formulation (2.4). Therefore,

$$\int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} (u^* - u) : \widehat{\nabla} \varphi dx = 0, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega \setminus C),$$

so that $u^* = u$. This conclusion comes from the choice $\varphi = u^* - u$, and the use of Assumption 2.1 and Korn and Poincaré inequalities (see Proposition 1.1).

Next, we prove that $u_n \rightarrow u$ in $L^2(\Sigma_N)$ by showing strong convergence of u_n to u in H^1 -norm in a neighborhood of the boundary of Ω . Consider the weak formulations

$$\int_{\Omega \setminus C_n} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} \varphi_1 dx = \int_{\Sigma_N} g \cdot \varphi_1 d\sigma(x), \quad \forall \varphi_1 \in H^1(\Omega \setminus C_n), \quad (2.23)$$

$$\int_{\Omega \setminus C} \mathbb{C}_0 \widehat{\nabla} u : \widehat{\nabla} \varphi_2 dx = \int_{\Sigma_N} g \cdot \varphi_2 d\sigma(x), \quad \forall \varphi_2 \in H^1(\Omega \setminus C). \quad (2.24)$$

Now, we define $\Phi = (u_n - u)\chi^2$, where χ is a smooth cut-off function, $\chi \in [0, 1]$ in Ω , such that

$$\chi = \begin{cases} 1 & \text{in } \overline{\Omega}^{d_0/4} \\ 0 & \text{in } \Omega \setminus \Omega^{d_0/2}. \end{cases}$$

Then, we choose $\varphi_1 = \varphi_2 = \Phi$ in (2.23) and (2.24), that is

$$\int_{\Omega^{d_0/2}} \mathbb{C}_0 \widehat{\nabla} u_n : \widehat{\nabla} ((u_n - u)\chi^2) dx = \int_{\Sigma_N} g \cdot (u_n - u) d\sigma(x),$$

$$\int_{\Omega^{d_0/2}} \mathbb{C}_0 \widehat{\nabla} u : \widehat{\nabla} ((u_n - u)\chi^2) dx = \int_{\Sigma_N} g \cdot (u_n - u) d\sigma(x).$$

Subtracting the last two equations, we find

$$\int_{\Omega^{d_0/2}} \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} ((u_n - u)\chi^2) dx = 0,$$

that is,

$$\begin{aligned} \int_{\Omega^{d_0/2}} \chi^2 \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) dx \\ + \int_{\Omega^{d_0/2}} 2\chi \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \overline{((u_n - u) \otimes \nabla \chi)} dx = 0. \end{aligned} \quad (2.25)$$

On the second integral, we apply the Young's inequality with a suitable parameter $\kappa > 0$, that is

$$\begin{aligned} \int_{\Omega^{d_0/2}} 2\chi \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \overline{((u_n - u) \otimes \nabla \chi)} dx \\ \leq 4\kappa \int_{\Omega^{d_0/2}} \chi^2 \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) \\ + \frac{1}{\kappa} \int_{\Omega^{d_0/2}} \mathbb{C}_0 \overline{((u_n - u) \otimes \nabla \chi)} : \overline{((u_n - u) \otimes \nabla \chi)}. \end{aligned}$$

Hence, using this last inequality in (2.25), we get

$$\begin{aligned} (1 - 4\kappa) \int_{\Omega^{d_0/2}} \chi^2 \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) dx \\ \leq \frac{1}{\kappa} \int_{\Omega^{d_0/2}} \mathbb{C}_0 \overline{((u_n - u) \otimes \nabla \chi)} : \overline{((u_n - u) \otimes \nabla \chi)}. \end{aligned}$$

The right-hand side integral goes to zero, noticing that

$$\begin{aligned} \int_{\Omega^{d_0/2}} \mathbb{C}_0 \overline{((u_n - u) \otimes \nabla \chi)} : \overline{((u_n - u) \otimes \nabla \chi)} \\ \leq c \int_{\Omega^{d_0/2}} \mathbb{C}_0 |u_n - u|^2 |\nabla \chi|^2 dx \leq c \int_{\Omega^{d_0/2}} |u_n - u|^2 dx \longrightarrow 0, \quad \text{as } n \rightarrow +\infty. \end{aligned} \quad (2.26)$$

The left-hand side can be estimated using the fact that

$$\int_{\Omega^{d_0/2}} \chi^2 \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) dx \geq \int_{\Omega^{d_0/4}} \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) dx$$

and, then, by means of the Korn inequality

$$\int_{\Omega^{d_0/4}} \mathbb{C}_0 \widehat{\nabla} (u_n - u) : \widehat{\nabla} (u_n - u) dx \geq c \|\nabla(u_n - u)\|_{L^2(\Omega^{d_0/4})}^2. \quad (2.27)$$

From (2.27) and (2.26), and recalling that u_n is converging strongly in L^2 -norm to u from the previous results, we find that

$$\|u_n - u\|_{H^1(\Omega^{d_0/4})} \rightarrow 0, \quad \text{as } n \rightarrow +\infty. \quad (2.28)$$

Finally, by the continuity of the trace theorem the proof is concluded. \square

Remark 2.6. *In the previous result, $u_n \rightarrow u$ in $L^2(\Sigma_N)$ can be also proved using the following arguments: note that the trace operator is a linear continuous operator from $H_{\Sigma_D}^1(\Omega \setminus C_n)$ to $H^{\frac{1}{2}}(\Sigma_N)$ (and, analogously, from $H_{\Sigma_D}^1(\Omega \setminus C)$ to $H^{\frac{1}{2}}(\Sigma_N)$), hence is also continuous in the weak topology, see [19]. Moreover, since $H^{\frac{1}{2}}(\Sigma_N) \hookrightarrow L^2(\Sigma_N)$ is compact, we find that $u_n \rightarrow u$ in $L^2(\Sigma_N)$.*

As a consequence of the continuity of the boundary functional, some properties of the functional $J(C)$ defined in (2.8) follow.

Proposition 2.7. *For every $\alpha > 0$ there exists at least one solution of the minimization problem (2.8).*

Proof. Let $\{C_n\}_{n \geq 0} \in \mathcal{C}$ be a minimizing sequence. Then there exists a positive constant M such that

$$J(C_n) \leq M, \quad \forall n, \quad (2.29)$$

hence

$$\text{Per}(C_n) \leq M, \quad \forall n.$$

By compactness (see Theorem 3.39 in [4]), there exists a set of finite perimeter C_0 such that, possibly up to a subsequence,

$$|C_n \Delta C_0| \rightarrow 0, \quad n \rightarrow \infty,$$

where $C_n \Delta C_0$ is the symmetric difference of the two sets. Moreover, thanks to the compactness and equiboundedness of the sets C_n and the fact that $C_n \in \mathcal{C}$, there exists a further subsequence which converges in the Hausdorff metric to $C_0 \in \mathcal{C}$, thanks to [39, Theorem 2.4.10]. Moreover, by the lower semicontinuity of the perimeter functional (see Section 5.2.1, Theorem 1, in [34]) it follows that

$$\text{Per}(C_0) \leq \liminf_{n \rightarrow \infty} \text{Per}(C_n).$$

Using the continuity of the boundary functional, see (2.14), we also have

$$\int_{\Sigma_N} (u(C_n) - u_{meas})^2 d\sigma(x) \rightarrow \int_{\Sigma_N} (u(C_0) - u_{meas})^2 d\sigma(x), \quad \text{as } n \rightarrow \infty.$$

In conclusion, we find that

$$J(C_0) \leq \liminf_{n \rightarrow \infty} J(C_n) = \lim_{n \rightarrow \infty} J(C_n) = \inf_{C \in \mathcal{C}} J(C),$$

and the claim follows. \square

We also prove stability with respect to the measured data.

Proposition 2.8. *Solutions of (2.8) are stable with respect to perturbations of the data u_{meas} , i.e., if $u_n \rightarrow u_{meas}$ in $L^2(\Sigma_N)$ as $n \rightarrow \infty$ then the solutions C_n of (2.8) with datum u_n are such that, up to subsequences,*

$$d_H(C_n, \tilde{C}) \rightarrow 0, \quad \text{as } n \rightarrow \infty,$$

where $\tilde{C} \in \mathcal{C}$ is a solution of (2.8), with datum u_{meas} .

Proof. Using (2.8), we have that, for any n , C_n satisfies

$$\frac{1}{2} \int_{\Sigma_N} (u(C_n) - u_n)^2 d\sigma(x) + \alpha \text{Per}(C_n) \leq \frac{1}{2} \int_{\Sigma_N} (u(C) - u_n)^2 d\sigma(x) + \alpha \text{Per}(C),$$

for all $C \in \mathcal{C}$. Therefore, $\text{Per}(C_n) \leq M$ and hence, possibly up to subsequences,

$$d_H(C_n, \tilde{C}) \rightarrow 0, \quad n \rightarrow \infty,$$

for some $\tilde{C} \in \mathcal{C}$, and

$$\text{Per}(\tilde{C}) \leq \liminf_{n \rightarrow \infty} \text{Per}(C_n).$$

Moreover, by the continuity of the solution of (2.4) with respect to C , see Theorem 2.5, we get

$$\begin{aligned} J(\tilde{C}) &\leq \liminf_{n \rightarrow \infty} \frac{1}{2} \int_{\Sigma_N} (u(C_n) - u_n)^2 d\sigma(x) + \alpha \text{Per}(C_n) \\ &\leq \lim_{n \rightarrow \infty} \frac{1}{2} \int_{\Sigma_N} (u(C) - u_n)^2 d\sigma(x) + \alpha \text{Per}(C) \\ &= \frac{1}{2} \int_{\Sigma_N} (u(C) - u_{meas})^2 d\sigma(x) + \alpha \text{Per}(C), \end{aligned}$$

for all $C \in \mathcal{C}$. Summarizing, $\tilde{C} \in \mathcal{C}$ and it is a minimizer of the functional, hence the assertion follows. \square

Finally, we can prove that the solution of the minimization problem (2.8) converges to the unique solution of the inverse problem when the regularization parameter tends to zero.

Proposition 2.9. *Let us assume that there exists a solution $C^\sharp \in \mathcal{C}$ of the inverse problem corresponding to datum u_{meas} . Moreover, for any $\eta > 0$ let $(\alpha(\eta))_{\eta > 0}$ be such that $\alpha(\eta) = o(1)$ and $\frac{\eta^2}{\alpha(\eta)}$ is bounded as $\eta \rightarrow 0$.*

Furthermore, let C_η be a solution to the minimization problem (2.8) with $\alpha = \alpha(\eta)$ and datum $u_\eta \in L^2(\Sigma_N)$ satisfying $\|u_{meas} - u_\eta\|_{L^2(\Sigma_N)} \leq \eta$. Then

$$C_\eta \rightarrow C^\sharp$$

in the Hausdorff metric, as $\eta \rightarrow 0$.

Proof. From the definition of C_η , it immediately follows that

$$\begin{aligned} \frac{1}{2} \int_{\Sigma_N} (u(C_\eta) - u_\eta)^2 d\sigma(x) + \alpha \text{Per}(C_\eta) &\leq \frac{1}{2} \int_{\Sigma_N} (u(C^\sharp) - u_\eta)^2 d\sigma + \alpha \text{Per}(C^\sharp) \\ &= \frac{1}{2} \int_{\Sigma_N} (u_{meas} - u_\eta)^2 d\sigma + \alpha \text{Per}(C^\sharp) \\ &\leq \eta^2 + \alpha \text{Per}(C^\sharp). \end{aligned} \quad (2.30)$$

Straightforwardly, we find that

$$\text{Per}(C_\eta) \leq \frac{\eta^2}{\alpha} + \text{Per}(C^\sharp) \leq M. \quad (2.31)$$

Hence, up to subsequences, arguing as in Proposition 2.8, we get

$$d_H(C_\eta, C_0) \rightarrow 0, \quad \text{as } \eta \rightarrow 0,$$

for some $C_0 \in \mathcal{C}$. From (2.30) and (2.31), as $\eta \rightarrow 0$, we find

$$\int_{\Sigma_N} (u(C_\eta) - u_\eta)^2 d\sigma \rightarrow 0,$$

hence, also

$$\int_{\Sigma_N} (u(C_\eta) - u_{meas})^2 d\sigma(x) \leq \int_{\Sigma_N} (u(C_\eta) - u_\eta)^2 d\sigma(x) + \int_{\Sigma_N} (u_{meas} - u_\eta)^2 d\sigma(x) \rightarrow 0.$$

By the continuity result in Theorem 2.5 and using the last relation, we find that

$$u(C_0) = u_{meas}, \quad \text{on } \Sigma_N.$$

Therefore, thanks to the uniqueness result of the inverse problem in Lipschitz domains (cf. [53]) we get $C_0 = C^\sharp$. \square

3 Reconstruction of cavities - filling the void

From the numerical point of view, the minimization of the functional (2.8) is complicated due to its non-differentiability. A typical approach to overcome this issue is to consider a further regularization of the functional, where the perimeter is approximated by a Ginzburg-Landau type functional, see for example [18]. This approach is well-known in the literature and it has been applied in different contexts, see for example [3, 12, 14, 15, 16, 18, 25, 30, 35, 44, 48].

First, we note that Problem (2.8) is equivalent to the following formulation

$$\min_{v \in X_{0,1}} J(v), \text{ where } J(v) = \frac{1}{2} \int_{\Sigma_N} |u(v) - u_{meas}|^2 d\sigma(x) + \alpha TV(v), \quad (3.1)$$

where $X_{0,1} := \{v \in BV(\Omega) : v = \chi_C \text{ a.e. in } \Omega, C \in \mathcal{C}\}$, $TV(v)$ is defined in (1.7), and χ_C is the indicator function of C . Note that the space $X_{0,1}$ is endowed with the norm $\|v\|_{BV(\Omega)} = \|v\|_{L^1(\Omega)} + TV(v)$.

Remark 3.1. *By compactness properties of $BV(\Omega)$ (see, e.g., [4], Theorem 3.23), any uniformly bounded sequence in $X_{0,1}$ admits a subsequence converging in $L^1(\Omega)$ to an element in $X_{0,1}$. In fact, let v_n a sequence uniformly bounded in $X_{0,1}$, there exists, possibly up to a subsequence, $v \in BV(\Omega)$ such that*

$$v_n \rightarrow v \text{ in } L^1(\Omega) \Rightarrow v_n \rightarrow v \text{ a.e. in } \Omega.$$

Since v_n attains values 0 and 1 only, it follows that $v \in X_{0,1}$.

Following the approach proposed in [18], we fill the cavity with a fictitious material with elastic properties that are different from the background. Specifically, we take an elasticity tensor $\mathbb{C}_1 := \delta \mathbb{C}_0$, where $\delta > 0$ is sufficiently small. Therefore, the boundary value problem (2.1) is modified into

$$\begin{cases} \operatorname{div}(\mathbb{C}_\delta(v) \widehat{\nabla} u_\delta(v)) = 0 & \text{in } \Omega, \\ (\mathbb{C}_\delta(v) \widehat{\nabla} u_\delta(v)) \nu = g & \text{on } \Sigma_N, \\ u_\delta(v) = 0 & \text{on } \Sigma_D, \end{cases} \quad (3.2)$$

where

$$\mathbb{C}_\delta(v) = \mathbb{C}_0 + (\mathbb{C}_1 - \mathbb{C}_0)v, \quad \text{with } \mathbb{C}_1 = \delta \mathbb{C}_0. \quad (3.3)$$

Here \mathbb{C}_0 and \mathbb{C}_1 are the elasticity tensors in $\Omega \setminus C$ and C , respectively.

Remark 3.2. *Thanks to Assumption 2.1, the fact that $\delta > 0$, and by (3.3), the elasticity tensor $\mathbb{C}_\delta(v)$ is strongly convex.*

Remark 3.3. *The following analysis can be generalized to the case of a generic fourth-order elasticity tensor \mathbb{C}_1 which is strongly convex and uniformly bounded with the further hypothesis that*

$$\mathbb{C}_1 \widehat{A} : \widehat{A} \leq \mathbb{C}_0 \widehat{A} : \widehat{A} \quad \text{or} \quad \mathbb{C}_0 \widehat{A} : \widehat{A} \leq \mathbb{C}_1 \widehat{A} : \widehat{A}.$$

Remark 3.4. *When dealing with sequences, we will often use the simplified notation $u_n := u_\delta(v_n)$, $u := u_\delta(v)$, $\mathbb{C}_n := \mathbb{C}_\delta(v_n)$, $\mathbb{C} := \mathbb{C}_\delta(v)$.*

The elastic problem (3.2) has the following weak formulation:

Find $u_\delta(v) \in H_{\Sigma_D}^1(\Omega)$ solution to

$$\int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} u_\delta(v) : \widehat{\nabla} \varphi dx = \int_{\Sigma_N} g \cdot \varphi d\sigma(x), \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega). \quad (3.4)$$

Well-posedness of Problem (3.2) in $H_{\Sigma_D}^1(\Omega)$ follows in the same way as for Problem (2.1), and, in addition

$$\|u_\delta(v)\|_{H^1(\Omega)} \leq c \|g\|_{L_{\Sigma_N}^2}.$$

We now approximate Problem (3.1) with the following one

$$\min_{v \in X_{0,1}} J_\delta(v), \text{ where } J_\delta(v) = \frac{1}{2} \int_{\Sigma_N} |u_\delta(v) - u_{meas}|^2 d\sigma(x) + \alpha TV(v), \quad (3.5)$$

where $u_\delta(v) \in H_{\Sigma_D}^1(\Omega)$ is the solution of Problem (3.2).

We prove the existence of minima of $J_\delta(v)$ in $X_{0,1}$, on account of the ideas contained in [14]. The proof is a consequence of the following property.

Proposition 3.5. *Let $\{v_n\} \subset X_{0,1}$ be strongly convergent in $L^1(\Omega)$ to $v \in X_{0,1}$. Then $\{u_\delta(v_n)|_{\Sigma_N}\}$ strongly converges in $L^2(\Sigma_N)$ to $u_\delta(v)|_{\Sigma_N}$, i.e., the map $F : v \rightarrow u_\delta(v)|_{\Sigma_N}$ is continuous from $X_{0,1}$ to $L^2(\Sigma_N)$ in the L^1 topology.*

Proof. Consider the weak formulation (3.4) associated to v and v_n , respectively,

$$\begin{aligned} \int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} u_\delta(v) : \widehat{\nabla} \varphi &= \int_{\Sigma_N} g \cdot \varphi, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega), \\ \int_{\Omega} \mathbb{C}_\delta(v_n) \widehat{\nabla} u_\delta(v_n) : \widehat{\nabla} \varphi &= \int_{\Sigma_N} g \cdot \varphi, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega). \end{aligned}$$

Subtracting the two equations and setting $u_n := u_\delta(v_n)$, $u := u_\delta(v)$, $\mathbb{C}_n := \mathbb{C}_\delta(v_n)$, $\mathbb{C} := \mathbb{C}_\delta(v)$, we get

$$\int_{\Omega} \mathbb{C}_n \widehat{\nabla} (u_n - u) : \widehat{\nabla} \varphi + \int_{\Omega} (\mathbb{C}_n - \mathbb{C}) \widehat{\nabla} u : \widehat{\nabla} \varphi = 0, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega).$$

Thus, making the choice $\varphi = u_n - u$ and proceeding similarly as in (2.5) to get H^1 -estimates, we find

$$\|u_n - u\|_{H^1(\Omega)}^2 \leq c \|(\mathbb{C}_n - \mathbb{C}) \widehat{\nabla} u\|_{L^2(\Omega)} \|\widehat{\nabla} (u_n - u)\|_{L^2(\Omega)},$$

and then, by $\mathbb{C}_n - \mathbb{C} = (\mathbb{C}_1 - \mathbb{C}_0)(v_n - v)$ and the uniform bound on the elasticity tensor, see Assumption 2.1, we derive

$$\|u_n - u\|_{H^1(\Omega)} \leq c \|(\widehat{\nabla} u)(v_n - v)\|_{L^2(\Omega)}.$$

Observe now that $v_n - v \rightarrow 0$ in $L^1(\Omega)$ as $n \rightarrow +\infty$ so that, possibly up to a subsequence, $v_n - v \rightarrow 0$, a.e. in Ω . Moreover, recalling that v_n and v are bounded and $u \in H^1(\Omega)$, we deduce, by dominated convergence theorem, that

$$\|u_n - u\|_{H^1(\Omega)} \rightarrow 0, \quad \text{as } n \rightarrow +\infty.$$

Finally, the trace theorem implies

$$\|u_n - u\|_{L^2(\Sigma_N)} \rightarrow 0, \quad \text{as } n \rightarrow +\infty.$$

□

Proposition 3.6. $J_\delta(v)$ admits a minimum $v \in X_{0,1}$.

Proof. Observe that $J_\delta(v)$ is bounded from below, by definition. Moreover, $J_\delta(v) \neq +\infty$, for $v \in X_{0,1}$. So, let $\{v_n\} \subset X_{0,1}$ be a minimizing sequence of $J_\delta(v)$, that is

$$J_\delta(v_n) \rightarrow \inf_{v \in X_{0,1}} J_\delta(v) = M, \quad \text{as } n \rightarrow +\infty.$$

Then

$$0 \leq J_\delta(v_n) \leq 2M \quad \text{and} \quad 0 \leq \alpha TV(v_n) \leq 2M.$$

Hence, there exists a positive constant c , independent on n , such that

$$\|v_n\|_{BV(\Omega)} = \|v_n\|_{L^1(\Omega)} + TV(v_n) \leq c. \quad (3.6)$$

This implies that v_n is uniformly bounded in $X_{0,1}$. Therefore, thanks to Remark 3.1, there exists $v \in X_{0,1}$ such that $v_n \rightarrow v$ in $L^1(\Omega)$. Due to the lower semicontinuity of $TV(v)$ with respect to the L^1 -convergence, we have

$$TV(v) \leq \liminf_{n \rightarrow +\infty} TV(v_n)$$

and, using Proposition 3.5, we get

$$\begin{aligned} J_\delta(v) &= \frac{1}{2} \int_{\Sigma_N} |u_\delta(v) - u_{meas}|^2 + \alpha TV(v) \\ &\leq \liminf_{n \rightarrow +\infty} \left(\frac{1}{2} \int_{\Sigma_N} |u_\delta(v_n) - u_{meas}|^2 + \alpha TV(v_n) \right) \\ &= \lim_{n \rightarrow +\infty} J_\delta(v_n) = M. \end{aligned} \tag{3.7}$$

□

3.1 Phase-field relaxation

Proceeding as in [30, 14], we now consider a phase-field relaxation of the optimization problem (3.5). More precisely, we define a minimization problem for a differentiable cost functional defined on a convex subspace of $H^1(\Omega)$, namely on the set

$$\mathcal{K} = \{v \in H^1(\Omega) : 0 \leq v(x) \leq 1 \text{ a.e. in } \Omega, v(x) = 0 \text{ a.e. in } \Omega^{d_0/2}\},$$

where $\Omega^{d_0/2}$ has been defined in (2.3), and, for every $\varepsilon > 0$, we replace the total variation term with the following Modica-Mortola functional.

Problem 3.1. *Given $u_{meas} \in L^2(\Sigma_N)$, and $\varepsilon, \delta > 0$, find*

$$\min_{v \in \mathcal{K}} J_{\delta,\varepsilon}(v), \quad J_{\delta,\varepsilon}(v) := \frac{1}{2} \|u_\delta(v) - u_{meas}\|_{L^2(\Sigma_N)}^2 + \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v|^2 + \frac{1}{\varepsilon} v(1-v) \right), \tag{3.8}$$

$u_\delta(v) \in H_{\Sigma_D}^1(\Omega)$ being the solution to (3.2), for $v \in \mathcal{K}$, and $\tilde{\alpha} = \frac{4}{\pi} \alpha$, where $4/\pi = (2 \int_0^1 \sqrt{v(1-v)} dv)^{-1}$ is a rescaling parameter, see [1].

Remark 3.7. *We expect Γ -convergence of the functional $J_{\delta,\varepsilon}$ to J , given in (3.1). However, this analysis is involved in the elastic context and is still an open issue that needs a specific accurate study.*

The following result holds

Proposition 3.8. *For any $\delta, \varepsilon > 0$, Problem (3.8) admits a solution $v = v_{\delta,\varepsilon} \in \mathcal{K}$.*

Proof. Let us fix $\delta, \varepsilon > 0$ and consider a minimizing sequence $\{v_n\} \subset \mathcal{K}$ for $J_{\delta,\varepsilon}(v)$ (we omit the dependence of v_n on δ and ε). We have

$$J_{\delta,\varepsilon}(v_n) \rightarrow \inf_{v \in \mathcal{K}} J_{\delta,\varepsilon}(v) = M.$$

Hence, by definition of minimizing sequence, $0 \leq J_{\delta,\varepsilon}(v_n) \leq 2M$ independently of n , which implies that also $\|\nabla v_n\|_{L^2(\Omega)}^2$ is bounded. Moreover, recalling that $v_n \in \mathcal{K}$ and $0 \leq v_n(x) \leq 1$ a.e. in Ω , we deduce that $\|v_n\|_{L^2(\Omega)} \leq M_1$, with M_1 independent of n and hence $\|v_n\|_{H^1(\Omega)} \leq M_2$, with M_2 independent of n . Due to the weak compactness of $H^1(\Omega)$, there exists $v \in H^1(\Omega)$ such that, possibly up to a subsequence, $v_n \rightharpoonup v$ in $H^1(\Omega)$. Hence, $v_n \rightarrow v$ strongly in $L^2(\Omega)$ and $v_n \rightarrow v$ a.e. in Ω . Since $v_n(1-v_n) \leq 1/4$, by means of the Lebesgue's dominated convergence theorem, we get

$$\int_{\Omega} v_n(1-v_n) \rightarrow \int_{\Omega} v(1-v).$$

Moreover, by the lower semicontinuity of the $H^1(\Omega)$ norm with respect to the weak convergence, we obtain

$$\begin{aligned} \|v\|_{H^1(\Omega)}^2 &\leq \liminf_{n \rightarrow +\infty} \|v_n\|_{H^1(\Omega)}^2, \\ \|v\|_{L^2(\Omega)}^2 + \|\nabla v\|_{L^2(\Omega)}^2 &\leq \lim_{n \rightarrow +\infty} \|v_n\|_{L^2(\Omega)}^2 + \liminf_{n \rightarrow +\infty} \|\nabla v_n\|_{L^2(\Omega)}^2 \\ &= \|v\|_{L^2(\Omega)}^2 + \liminf_{n \rightarrow +\infty} \|\nabla v_n\|_{L^2(\Omega)}^2, \\ \|\nabla v\|_{L^2(\Omega)}^2 &\leq \liminf_{n \rightarrow +\infty} \|\nabla v_n\|_{L^2(\Omega)}^2. \end{aligned}$$

By the last inequality and the convergence of v_n to v a.e., by the use of Proposition 3.5 and the fact that v_n is a minimizing sequence, we have

$$\begin{aligned} J_{\delta,\varepsilon}(v) &= \frac{1}{2} \|u_\delta(v) - u_{meas}\|_{L^2(\Sigma_N)}^2 + \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v|^2 + \frac{1}{\varepsilon} v(1-v) \right) \\ &\leq \liminf_{n \rightarrow +\infty} \left(\frac{1}{2} \|u_\delta(v_n) - u_{meas}\|_{L^2(\Sigma_N)}^2 + \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v_n|^2 + \frac{1}{\varepsilon} v_n(1-v_n) \right) \right) \\ &= \lim_{n \rightarrow +\infty} J_{\delta,\varepsilon}(v_n) = M. \end{aligned}$$

Finally, by pointwise convergence, we know that $0 \leq v \leq 1$ a.e. in Ω and $v = 1$ a.e. in $\Omega^{d_0/2}$. Hence, v is a minimum of $J_{\delta,\varepsilon}$ in \mathcal{K} . \square

3.2 Necessary optimality conditions

In this section we provide an expression for the first order necessary optimality condition associated with the minimization problem (3.8), formulated as a variational inequality involving the Fréchet derivative of $J_{\delta,\varepsilon}$.

Proposition 3.9. *Define the map $F : \mathcal{K} \rightarrow H^1(\Omega)$, $F(v) = u_\delta(v)$, $u_\delta(v)$ solution to (3.2). Then the operators F and $J_{\delta,\varepsilon}$ (for every $\delta, \varepsilon > 0$) are Fréchet-differentiable on $\mathcal{K} \subset L^\infty(\Omega) \cap H^1(\Omega)$.*

Moreover, any minimizer $v_{\delta,\varepsilon}$ of $J_{\delta,\varepsilon}$ satisfies the variational inequality

$$J'_{\delta,\varepsilon}(v_{\delta,\varepsilon})[\omega - v_{\delta,\varepsilon}] \geq 0, \quad \forall \omega \in \mathcal{K}, \quad (3.9)$$

where

$$J'_{\delta,\varepsilon}(v)[\vartheta] = \int_{\Omega} \vartheta (\mathbb{C}_0 - \mathbb{C}_1) \widehat{\nabla} u_\delta(v) : \widehat{\nabla} p_\delta(v) + 2\tilde{\alpha}\varepsilon \int_{\Omega} \widehat{\nabla} v : \widehat{\nabla} \vartheta + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1-2v)\vartheta. \quad (3.10)$$

Here $\vartheta \in \mathcal{K} - v = \{z \text{ s.t. } z + v \in \mathcal{K}\}$ and $p_\delta \in H_{\Sigma_D}^1(\Omega)$ is the solution to the adjoint problem

$$\int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} p_\delta : \widehat{\nabla} \psi = \int_{\Sigma_N} (u_\delta(v) - u_{meas}) \psi, \quad \forall \psi \in H_{\Sigma_D}^1(\Omega). \quad (3.11)$$

Proof. First we prove that F is Fréchet differentiable in $L^\infty(\Omega)$. More precisely,

$$F'(v)[\vartheta] = u^\sharp(v), \text{ for } \vartheta \in L^\infty(\Omega) \cap (\mathcal{K} - v),$$

where $u^\sharp(v)$ is the solution in $H_{\Sigma_D}^1(\Omega)$ of

$$\int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} u^\sharp(v) : \widehat{\nabla} \varphi = \int_{\Omega} \vartheta (\mathbb{C}_0 - \mathbb{C}_1) \widehat{\nabla} u_\delta(v) : \widehat{\nabla} \varphi, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega), \quad (3.12)$$

namely,

$$\|F(v + \vartheta) - F(v) - u^\sharp(v)\|_{H^1(\Omega)} = o(\|\vartheta\|_{L^\infty(\Omega)}). \quad (3.13)$$

To this aim, we first show that

$$\|u_\delta(v + \vartheta) - u_\delta(v)\|_{H^1(\Omega)} \leq c\|\vartheta\|_{L^\infty(\Omega)}, \text{ for } \vartheta \in L^\infty(\Omega) \cap (\mathcal{K} - v).$$

Indeed, the difference $u_\delta(v + \vartheta) - u_\delta(v)$ satisfies

$$\begin{aligned} & \int_{\Omega} \mathbb{C}_\delta(v + \vartheta) \widehat{\nabla}(u_\delta(v + \vartheta) - u_\delta(v)) : \widehat{\nabla} \varphi \\ & + \int_{\Omega} (\mathbb{C}_\delta(v + \vartheta) - \mathbb{C}_\delta(v)) \widehat{\nabla} u_\delta(v) : \widehat{\nabla} \varphi = 0, \quad \forall \varphi \in H_{\Sigma_D}^1(\Omega). \end{aligned} \quad (3.14)$$

Taking $\varphi = u_\delta(v + \vartheta) - u_\delta(v)$ and recalling that $\mathbb{C}_\delta(v + \vartheta) - \mathbb{C}_\delta(v) = (\mathbb{C}_1 - \mathbb{C}_0)\vartheta$, we obtain

$$\begin{aligned} & \int_{\Omega} \mathbb{C}_\delta(v + \vartheta) \widehat{\nabla}(u_\delta(v + \vartheta) - u_\delta(v)) : \widehat{\nabla}(u_\delta(v + \vartheta) - u_\delta(v)) \\ & = - \int_{\Omega} \vartheta(\mathbb{C}_1 - \mathbb{C}_0) \widehat{\nabla} u_\delta(v) : \widehat{\nabla}(u_\delta(v + \vartheta) - u_\delta(v)). \end{aligned} \quad (3.15)$$

Hence, by using the assumptions on the elasticity tensors, Korn and Poincaré inequalities, and the fact that $v + \vartheta \in \mathcal{K}$, we obtain

$$\begin{aligned} \|u_\delta(v + \vartheta) - u_\delta(v)\|_{H^1(\Omega)} & \leq c \|\vartheta\|_{L^\infty(\Omega)} \|\widehat{\nabla} u_\delta(v)\|_{L^2(\Omega)} \\ & \leq c \|\vartheta\|_{L^\infty(\Omega)} \|u_\delta(v)\|_{H^1(\Omega)} \leq c \|\vartheta\|_{L^\infty(\Omega)} \|g\|_{L^2(\Sigma_N)} \leq c \|\vartheta\|_{L^\infty(\Omega)}. \end{aligned} \quad (3.16)$$

We now estimate $u_\delta(v + \vartheta) - u_\delta(v) - u^\sharp(v)$. Subtracting (3.12) from (3.14) and setting $\omega = u_\delta(v + \vartheta) - u_\delta(v)$, then it holds

$$\int_{\Omega} \mathbb{C}_\delta(v + \vartheta) \widehat{\nabla} \omega : \widehat{\nabla} \varphi - \int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} u^\sharp(v) : \widehat{\nabla} \varphi = 0, \quad (3.17)$$

from which

$$\begin{aligned} \int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla}(\omega - u^\sharp(v)) : \widehat{\nabla} \varphi & = - \int_{\Omega} (\mathbb{C}_\delta(v + \vartheta) - \mathbb{C}_\delta(v)) \widehat{\nabla} \omega : \widehat{\nabla} \varphi \\ & = \int_{\Omega} \vartheta(\mathbb{C}_0 - \mathbb{C}_1) \widehat{\nabla} \omega : \widehat{\nabla} \varphi. \end{aligned} \quad (3.18)$$

Choosing now $\varphi = \omega - u^\sharp(v)$, we get

$$\int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla}(\omega - u^\sharp(v)) : \widehat{\nabla}(\omega - u^\sharp(v)) = \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1) \vartheta \widehat{\nabla} \omega : \widehat{\nabla}(\omega - u^\sharp(v)), \quad (3.19)$$

and again by the boundedness of the elasticity tensors and the use of Korn and Poincaré inequalities it follows

$$\|\omega - u^\sharp(v)\|_{H^1(\Omega)} = \|u(v + \vartheta) - u_\delta(v) - u^\sharp(v)\|_{H^1(\Omega)} \leq c \|\vartheta\|_{L^\infty(\Omega)}^2, \quad (3.20)$$

so that $F'(v)[\vartheta] = u^\sharp(v)$.

We now prove that $J_{\delta,\varepsilon}$ is Fréchet differentiable. By means of the chain rule and the Fréchet differentiability of F , we compute the expression of $J'_{\delta,\varepsilon}(v)$, i.e.,

$$J'_{\delta,\varepsilon}(v)[\vartheta] = \int_{\Sigma_N} (F(v) - u_{meas}) F'(v)[\vartheta] + \tilde{\alpha} \int_{\Omega} \left(2\varepsilon \nabla v : \nabla \vartheta + \frac{1}{\varepsilon} (1 - 2v) \vartheta \right), \quad (3.21)$$

where, with abuse of notation, $F(v)$ and $F'(v)[\vartheta]$ denote the trace of $F(v)$ and $F'(v)[\vartheta]$ on Σ_N , respectively. By the definition of the adjoint problem and of $u^\sharp(v)$, we get

$$\begin{aligned} \int_{\Sigma_N} (F(v) - u_{meas}) F'(v)[\vartheta] & = \int_{\Sigma_N} (F(v) - u_{meas}) u^\sharp(v) = \\ & = \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1) \vartheta \widehat{\nabla} F(v) : \widehat{\nabla} p_\delta \end{aligned} \quad (3.22)$$

and hence

$$J'_{\delta,\varepsilon}(v)[\vartheta] = \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1) \vartheta \widehat{\nabla} F(v) : \widehat{\nabla} p_\delta + \tilde{\alpha} \int_{\Omega} \left(2\varepsilon \nabla v \cdot \nabla \vartheta + \frac{1}{\varepsilon} (1 - 2v) \vartheta \right).$$

Finally, by standard arguments, since $J_{\delta,\varepsilon}$ is a continuous and Fréchet differentiable functional on a convex subset \mathcal{K} of the Banach space $H^1(\Omega)$, the optimality conditions for the optimization problem (3.8) are expressed in terms of the variational inequality (3.9). \square

4 Discretization and reconstruction algorithm

4.1 Convergence analysis

Here, we assume that Ω is a polygonal ($d = 2$) or polyhedral ($d = 3$) domain. Again, for simplifying the notation, we denote by $u := u_\delta$ and $p := p_\delta$.

Let $(\mathcal{T}_h)_{0 < h \leq h_0}$ be a regular triangulation of Ω and define

$$\mathcal{V}_h := \{v_h \in C(\bar{\Omega}) : v_h|_{\mathcal{T}} \in \mathcal{P}_1(\mathcal{T}), \forall \mathcal{T} \in \mathcal{T}_h\}, \quad (4.1)$$

where $\mathcal{P}_1(\mathcal{T})$ is the set of polynomials of first degree on \mathcal{T} , and

$$\mathcal{K}_h := \mathcal{V}_h \cap \mathcal{K}, \quad \mathcal{V}_{h, \Sigma_D} := \mathcal{V}_h \cap H_{\Sigma_D}^1(\Omega). \quad (4.2)$$

For every $h > 0$, we set $u_h := (u_\delta)_h : \mathcal{K} \rightarrow \mathcal{V}_{h, \Sigma_D}$ where u_h is solution to

$$\int_{\Omega} \mathbb{C}_\delta(v) \widehat{\nabla} u_h(v) : \widehat{\nabla} \varphi_h = \int_{\Sigma_N} g_h \cdot \varphi_h, \quad \forall \varphi_h \in \mathcal{V}_{h, \Sigma_D}. \quad (4.3)$$

Here g_h is a piecewise linear, continuous approximation of g such that $g_h \rightarrow g$ in $L^2(\Sigma_N)$ as $h \rightarrow 0$.

As in [30], one can show that for every $v \in \mathcal{K}$ there exists a sequence $v_h \in \mathcal{K}_h$ such that $v_h \rightarrow v$ in $H^1(\Omega)$. Most of the following results are an adaptation of those presented in [30] for a scalar equation to the case of the elasticity system, hence we do not provide the proofs for some of them.

The following lemma is a consequence of the continuity and coercivity of the bilinear form on the left-hand side of (4.3) and Céa's Lemma (see, e.g., [14]).

Lemma 4.1. *Let $g \in L^2(\Sigma_N)$. Then, $\forall v \in \mathcal{K}$, $u_h(v) \rightarrow u(v)$ strongly in $H^1(\Omega)$ as $h \rightarrow 0$.*

Next we state a result concerning the continuity of u_h in the space $\mathcal{V}_{h, \Sigma_D}$.

Proposition 4.2. *Let h_k, v_{h_k} be two sequences such that $\lim_{k \rightarrow +\infty} h_k = 0$ and $v_{h_k} \in \mathcal{K}_{h_k}$ with $v_{h_k} \rightarrow v$ in $L^1(\Omega)$. Then $u_{h_k}(v_{h_k}) \rightarrow u(v)$ in $H_{\Sigma_D}^1(\Omega)$ for $k \rightarrow +\infty$.*

Proof. The proof can be obtained reasoning similarly as in Lemma 3.1 of [30]. \square

Let $J_{\delta, \varepsilon, h} : \mathcal{K}_h \rightarrow \mathbb{R}$ be the approximation to $J_{\delta, \varepsilon}$ defined as follows

$$J_{\delta, \varepsilon, h} := \frac{1}{2} \|u_h(v_h) - u_{meas, h}\|_{L^2(\Sigma_N)}^2 + \tilde{\alpha} \int_{\Omega} \varepsilon |\nabla v_h|^2 + \frac{1}{\varepsilon} v_h(1 - v_h), \quad (4.4)$$

where we assume that $u_{meas, h} \rightarrow u_{meas}$, as $h \rightarrow 0$. Similarly as in Theorem 3.2 of [30], we can show the following result.

Theorem 4.3. *There exists $v_h \in \mathcal{K}_h$ such that $J_{\delta, \varepsilon, h}(v_h) = \min_{\eta_h \in \mathcal{K}_h} J_{\delta, \varepsilon, h}(\eta_h)$. Moreover, let h_k be such that $\lim_{k \rightarrow +\infty} h_k = 0$. Then every sequence v_{h_k} has a subsequence converging strongly in $H^1(\Omega)$ and a.e. in Ω to a minimum of $J_{\delta, \varepsilon}$.*

In our numerical algorithm we approximately solve (3.8) and so we look for an admissible point $v_h \in \mathcal{K}_h$ that satisfies the first order necessary condition

$$J'_{\delta, \varepsilon, h}(v_h)[\omega_h - v_h] \geq 0, \quad \forall \omega_h \in \mathcal{K}_h, \quad (4.5)$$

rather than trying to locate a global minimum of $J_{\delta, \varepsilon, h}$. To this aim, we consider the discrete adjoint problem: find $p_h := (p_\delta)_h \in \mathcal{V}_{h, \Sigma_D}$ such that

$$\int_{\Omega} \mathbb{C}_\delta(v_h) \widehat{\nabla} p_h : \widehat{\nabla} \psi_h = \int_{\Sigma_N} (u_h(v_h) - u_{meas, h}) \psi_h, \quad \forall \psi_h \in \mathcal{V}_{h, \Sigma_D}. \quad (4.6)$$

Then using $v_h \in \mathcal{K}_h$, we can prove the discrete version of Proposition 3.9, where the discrete variational inequality reads as:

$$\begin{aligned} & \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega_h - v_h) \widehat{\nabla} u_h(v_h) : \widehat{\nabla} p_h + 2\tilde{\alpha}\varepsilon \int_{\Omega} \nabla v_h \cdot \nabla(\omega_h - v_h) \\ & + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v_h)(\omega_h - v_h) \geq 0, \quad \forall \omega_h \in \mathcal{K}_h. \end{aligned} \quad (4.7)$$

Then, we can prove the following theorem:

Theorem 4.4. *Let h_k be such that $\lim_{k \rightarrow +\infty} h_k = 0$ and v_{h_k} be a sequence satisfying (4.5). Then there exists a subsequence of v_{h_k} that converges strongly in $H^1(\Omega)$ and a.e. in Ω to a solution v of (3.9).*

Proof. We set $v_k := v_{h_k}$, $u_k := u_{h_k}(v_{h_k})$ and $p_k := p_{h_k}(v_{h_k})$. Testing (4.6) with $\psi_h = p_k$ we get

$$\int_{\Omega} \mathbb{C}_{\delta}(v_k) \widehat{\nabla} p_k : \widehat{\nabla} p_k = \int_{\Sigma_N} (u_k - u_{meas,k}) \cdot p_k,$$

which yields, arguing as in (2.5) to get H^1 -estimates,

$$c \|p_k\|_{H^1(\Omega)}^2 \leq \|u_k - u_{meas,k}\|_{L^2(\Sigma_N)} \|p_k\|_{L^2(\Sigma_N)}.$$

As the problem for u_k is well-posed with $u_k \in H_{\Sigma_D}^1(\Omega)$ and $u_{meas,k} \rightarrow u_{meas}$ (implying that $\|u_{meas,k}\|_{L^2(\Sigma_N)}$ is uniformly bounded with respect to k), we get

$$\|p_k\|_{H^1(\Omega)} \leq c.$$

A similar result holds for $\|u_k\|_{H^1(\Omega)}$. Therefore

$$\|p_k\|_{H^1(\Omega)} + \|u_k\|_{H^1(\Omega)} \leq c, \quad \text{uniformly in } k. \quad (4.8)$$

From (4.7), employing $(1 - 2v_k)(w_k - v_k) \leq w_k + 2v_k^2$ and testing with $w_k = 0 \in \mathcal{K}_h$, we get

$$2\tilde{\alpha}\varepsilon \int_{\Omega} |\nabla v_k|^2 \leq c \|\widehat{\nabla} u_k\|_{L^2(\Omega)} \|\widehat{\nabla} p_k\|_{L^2(\Omega)} + \frac{2\tilde{\alpha}}{\varepsilon} |\Omega| \leq c_{\varepsilon}, \quad (4.9)$$

where we used (4.8). Therefore, v_k is bounded in $H^1(\Omega)$, hence there exists a subsequence (still denoted by v_k) and $v \in \mathcal{K}$ such that

$$\begin{aligned} v_k & \rightharpoonup v \quad \text{in } H^1(\Omega), & v_k & \rightarrow v \quad \text{in } L^2(\Omega) \quad (\text{and in } L^1(\Omega)), \\ v_k & \rightarrow v \quad \text{a.e. in } \Omega. \end{aligned}$$

Thanks to Proposition 4.2 we have

$$u_k \rightarrow u \quad \text{in } H_{\Sigma_D}^1(\Omega). \quad (4.10)$$

Now, let $p \in H_{\Sigma_D}^1(\Omega)$ be the solution of the continuous adjoint problem and let $\hat{p}_k \in \mathcal{V}_{h_k, \Sigma_D}$ be such that $\hat{p}_k \rightarrow p$ in $H_{\Sigma_D}^1(\Omega)$. Taking the difference of the problems solved by p and p_k , after some standard manipulation we get

$$\begin{aligned} & \int_{\Omega} \mathbb{C}_{\delta}(v_k) \widehat{\nabla} (p_k - \hat{p}_k) : \widehat{\nabla} \psi \\ & = \int_{\Omega} \mathbb{C}_{\delta}(v_k) \widehat{\nabla} (p - \hat{p}_k) : \widehat{\nabla} \psi + \int_{\Omega} (\mathbb{C}_{\delta}(v) - \mathbb{C}_{\delta}(v_k)) \widehat{\nabla} p : \widehat{\nabla} \psi \\ & + \int_{\Sigma_N} (u_k - u) \cdot \psi + \int_{\Sigma_N} (u_{meas} - u_{meas,k}) \cdot \psi, \end{aligned}$$

for all $\psi \in \mathcal{V}_{h_k, \Sigma_D}$. Taking $\psi = p_k - \hat{p}_k$, we get

$$\begin{aligned} \|p_k - \hat{p}_k\|_{H^1(\Omega)} & \leq c \left(\|\widehat{\nabla} (p - \hat{p}_k)\|_{L^2(\Omega)} + \int_{\Omega} |v - v_k|^2 |\widehat{\nabla} p|^2 \right. \\ & \left. + \|u_k - u\|_{L^2(\Sigma_N)} + \|u_{meas} - u_{meas,k}\|_{L^2(\Sigma_N)} \right). \end{aligned} \quad (4.11)$$

By hypothesis, we have $\|u_{meas} - u_{meas,k}\|_{L^2(\Sigma_N)} \rightarrow 0$ and $\|p - \hat{p}_k\|_{H^1(\Omega)} \rightarrow 0$ for $k \rightarrow +\infty$. Hence, invoking Proposition 4.2 and observing that $\int_{\Omega} |v - v_k|^2 |\widehat{\nabla} p|^2 \rightarrow 0$ for $k \rightarrow +\infty$, we deduce $p_k \rightarrow p$ in $H^1(\Omega)$.

Next, we have to show that v satisfies the variational inequality (3.9). Given $\omega \in \mathcal{K}$, there exists a sequence $\hat{\omega}_k \in \mathcal{K}_{h_k}$ such that $\hat{\omega}_k \rightarrow \omega$ in $H^1(\Omega)$ and a.e. in Ω . Then, from the discrete variational inequality (4.7) we have for v_k that

$$\begin{aligned} & \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\hat{\omega}_k - v_k) \widehat{\nabla} u_k : \widehat{\nabla} p_k + 2\tilde{\alpha}\varepsilon \int_{\Omega} \nabla v_k \cdot \nabla(\hat{\omega}_k - v_k) \\ & + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v_k)(\hat{\omega}_k - v_k) \geq 0. \end{aligned} \quad (4.12)$$

Now, observe that

$$\begin{aligned} & \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\hat{\omega}_k - v_k) \widehat{\nabla} u_k : \widehat{\nabla} p_k - \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega - v) \widehat{\nabla} u : \widehat{\nabla} p \\ & = \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\hat{\omega}_k - v_k) [\widehat{\nabla}(u_k - u) : \widehat{\nabla} p_k + \widehat{\nabla} u : \widehat{\nabla}(p_k - p)] \\ & + \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)[(\hat{\omega}_k - \omega) - (v_k - v)] \widehat{\nabla} u : \widehat{\nabla} p. \end{aligned} \quad (4.13)$$

The first integral on the right hand side converges to zero by (4.10) and $p_k \rightarrow p$ in $H^1(\Omega)$. To show that also the second integral converges to zero, we invoke the dominated convergence theorem. Hence, from (4.13), we obtain

$$\int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\hat{\omega}_k - v_k) \widehat{\nabla} u_k : \widehat{\nabla} p_k - \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega - v) \widehat{\nabla} u : \widehat{\nabla} p \rightarrow 0, \quad (4.14)$$

as $k \rightarrow +\infty$. Then, utilizing (4.14) into (4.12), together with the fact that $v_k \rightarrow v$ in $H^1(\Omega)$, and the lower semicontinuity of the norm, we find

$$\|\nabla v\|_{L^2(\Omega)}^2 \leq \liminf_{k \rightarrow +\infty} \|\nabla v_k\|_{L^2(\Omega)}^2.$$

Moreover, noticing that $\int_{\Omega} v_k \hat{\omega}_k \rightarrow \int_{\Omega} v \omega$ for $k \rightarrow +\infty$, we get

$$\begin{aligned} & \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega - v) \widehat{\nabla} u : \widehat{\nabla} p + 2\tilde{\alpha}\varepsilon \int_{\Omega} \nabla v \cdot \nabla(\omega - v) \\ & + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v)(\omega - v) \\ & \geq \liminf_{k \rightarrow +\infty} \left\{ \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\hat{\omega}_k - v_k) \widehat{\nabla} u_k : \widehat{\nabla} p_k \right. \\ & \left. + 2\tilde{\alpha}\varepsilon \int_{\Omega} \nabla v_k \cdot \nabla(\hat{\omega}_k - v_k) + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v_k)(\hat{\omega}_k - v_k) \right\} \geq 0. \end{aligned} \quad (4.15)$$

Finally, it remains to show that $v_k \rightarrow v$ strongly in $H^1(\Omega)$. We choose a sequence $\hat{v}_k \in \mathcal{K}_{h_k}$ such that $\hat{v}_k \rightarrow v$ in $H^1(\Omega)$ and using the discrete variational inequality (4.7) with $\omega_{h_k} = \hat{v}_k$, we easily get $\nabla v_k \rightarrow \nabla v$ in $L^2(\Omega)$, implying the result. \square

4.2 Reconstruction Algorithm

In order to solve the discrete optimization problem we follow the method used in [14] and [30]. The method is based on solving the following parabolic obstacle problem. For $\delta, \varepsilon > 0$ fixed, let v be the solution to

$$\begin{aligned} & \int_{\Omega} \partial_t v(\omega - v) + J'_{\delta, \varepsilon}(v)[\omega - v] \geq 0, \quad \forall \omega \in \mathcal{K}, t \in (0 + \infty), \\ & v(\cdot, 0) = v_0 \in \mathcal{K}. \end{aligned}$$

An easy computation shows that the value of the objective functional decreases in time. Hence, we expect that if the limit as $t \rightarrow +\infty$ of its solution $v(\cdot, t)$ exists and it is equal to the asymptotic state v_∞ , then this should satisfy the continuous optimality conditions (3.9). We now discretize the above problem by using a semi-implicit time discretization scheme. We denote by $\{v_h^n\}_{n \in \mathbb{N}} \subset \mathcal{K}_h$ the sequence of approximations $v_h^n \simeq v(\cdot, t^n)$ obtained as follows:

$$\begin{aligned}
v_h^0 &= v_0 \in \mathcal{K}_h \\
v_h^{n+1} &\in \mathcal{K}_h : \frac{1}{\tau_n} \int_{\Omega} (v_h^{n+1} - v_h^n)(\omega_h - v_h^{n+1}) \\
&+ \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega_h - v_h^{n+1}) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + 2\tilde{\alpha}\varepsilon \int_{\Omega} \nabla v_h^{n+1} \cdot \nabla(\omega_h - v_h^{n+1}) \\
&+ \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v_h^n)(\omega_h - v_h^{n+1}) \geq 0, \quad \forall \omega_h \in \mathcal{K}_h, \quad n \geq 0,
\end{aligned} \tag{4.16}$$

where τ_n is the time step, and $u_h^n, p_h^n \in \mathcal{V}_{h, \Sigma_D}$ are the discrete solutions of the forward problem (4.3) and adjoint problem (4.6), respectively, for $v_h = v_h^n$. We now prove a monotonicity property of the method.

Lemma 4.5. *For each $n \in \mathbb{N}$, there exists a constant $c^n > 0$ such that, if $\tau_n \leq (1 + c^n)^{-1}$, then*

$$\|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + J_{\delta, \varepsilon, h}(v_h^{n+1}) \leq J_{\delta, \varepsilon, h}(v_h^n), \tag{4.17}$$

where $c^n = c^n(\Omega, \delta, \xi_0, h, \|\mathbb{C}_0 - \mathbb{C}_1\|_{L^\infty(\Omega)}, \|p_h^n\|_{W^{1, \infty}(\Omega)}, \|u_h^n\|_{W^{1, \infty}(\Omega)})$.

Proof. Choosing $\omega_h = v_h^n$ in (4.16), after some simple manipulations we obtain

$$\begin{aligned}
&\frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha}\varepsilon \|\nabla(v_h^{n+1} - v_h^n)\|_{L^2(\Omega)}^2 \\
&+ \frac{\tilde{\alpha}}{\varepsilon} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v_h^{n+1}|^2 - \frac{1}{\varepsilon} v_h^{n+1} (1 - v_h^{n+1}) \right) \\
&- \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v_h^n|^2 - \frac{1}{\varepsilon} v_h^n (1 + v_h^n) \right) \\
&\leq \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(v_h^n - v_h^{n+1}) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n.
\end{aligned}$$

Adding and subtracting $\frac{1}{2} \|u_h^{n+1} - u_{meas, h}\|_{L^2(\Sigma_N)}^2$ and $\frac{1}{2} \|u_h^n - u_{meas, h}\|_{L^2(\Sigma_N)}^2$, we get

$$\begin{aligned}
&\frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha}\varepsilon \|\nabla(v_h^{n+1} - v_h^n)\|_{L^2(\Omega)}^2 + \frac{\tilde{\alpha}}{\varepsilon} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 \\
&+ J_{\delta, \varepsilon, h}(v_h^{n+1}) - \frac{1}{2} \|u_h^{n+1} - u_{meas, h}\|_{L^2(\Sigma_N)}^2 - J_{\delta, \varepsilon, h}(v_h^n) \\
&+ \frac{1}{2} \|u_h^n - u_{meas, h}\|_{L^2(\Sigma_N)}^2 \leq \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(v_h^n - v_h^{n+1}) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n,
\end{aligned}$$

which implies

$$\begin{aligned}
& \frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha}\varepsilon \|\nabla(v_h^{n+1} - v_h^n)\|_{L^2(\Omega)}^2 + \frac{\tilde{\alpha}}{\varepsilon} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 \\
& \quad + J_{\delta,\varepsilon,h}(v_h^{n+1}) - J_{\delta,\varepsilon,h}(v_h^n) \\
& \leq \int_{\Omega} (v_h^n - v_h^{n+1})(\mathbb{C}_0 - \mathbb{C}_1) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \quad + \int_{\Sigma_N} (u_h^{n+1} - u_h^n) \cdot (u_h^n - u_{meas,h}) \\
& = \int_{\Omega} (v_h^n - v_h^{n+1})(\mathbb{C}_0 - \mathbb{C}_1) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \quad + \int_{\Omega} \mathbb{C}_{\delta}(v_h^n) \widehat{\nabla} p_h^n : \widehat{\nabla} (u_h^{n+1} - u_h^n) \\
& = \int_{\Omega} (\mathbb{C}_{\delta}(v_h^{n+1}) - \mathbb{C}_{\delta}(v_h^n)) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \quad + \int_{\Omega} \mathbb{C}_{\delta}(v_h^n) \widehat{\nabla} p_h^n : \widehat{\nabla} (u_h^{n+1} - u_h^n), \tag{4.18}
\end{aligned}$$

where in the last step we employed

$$\mathbb{C}_{\delta}(v_h^{n+1}) - \mathbb{C}_{\delta}(v_h^n) = (\mathbb{C}_0 - \mathbb{C}_1)(v_h^n - v_h^{n+1}).$$

It is easy to verify that it holds

$$\begin{aligned}
& \int_{\Omega} (\mathbb{C}_{\delta}(v_h^{n+1}) - \mathbb{C}_{\delta}(v_h^n)) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \quad + \int_{\Omega} \mathbb{C}_{\delta}(v_h^n) \widehat{\nabla} p_h^n : \widehat{\nabla} (u_h^{n+1} - u_h^n) \\
& = \int_{\Omega} (\mathbb{C}_{\delta}(v_h^n) - \mathbb{C}_{\delta}(v_h^{n+1})) \widehat{\nabla} (u_h^{n+1} - u_h^n) : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \quad + \int_{\Omega} \mathbb{C}_{\delta}(v_h^{n+1}) \widehat{\nabla} u_h^{n+1} : \widehat{\nabla} p_h^n - \int_{\Omega} \mathbb{C}_{\delta}(v_h^n) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n \\
& = \int_{\Omega} (\mathbb{C}_{\delta}(v_h^n) - \mathbb{C}_{\delta}(v_h^{n+1})) \widehat{\nabla} (u_h^{n+1} - u_h^n) : \widehat{\nabla} p_h^n + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& =: I_1,
\end{aligned}$$

where the last step follows from the definition of the discrete adjoint problem.

Then, using the Cauchy-Schwarz inequality, the trace theorem and the fact that in finite dimensional spaces all norms are equivalent, we have

$$\begin{aligned}
|I_1| & \leq c_0^n \|\mathbb{C}_1 - \mathbb{C}_0\|_{L^\infty(\Omega)} \|\widehat{\nabla} p_h^n\|_{L^\infty(\Omega)} \|v_h^n - v_h^{n+1}\|_{L^2(\Omega)} \|\widehat{\nabla} (u_h^{n+1} - u_h^n)\|_{L^2(\Omega)} \\
& \quad + \frac{1}{2} \|u_h^{n+1} - u_h^n\|_{L^2(\Sigma_N)}^2 \\
& \leq c_1^n \|v_h^n - v_h^{n+1}\|_{L^2(\Omega)} \|u_h^{n+1} - u_h^n\|_{H^1(\Omega)} + \frac{c_2^n}{2} \|u_h^{n+1} - u_h^n\|_{H^1(\Omega)}^2 \tag{4.19}
\end{aligned}$$

where $c_0^n = c_0^n(\Omega, h)$, $c_1^n = c_1^n(\|\mathbb{C}_1 - \mathbb{C}_0\|_{L^\infty(\Omega)}, \|\widehat{\nabla} p_h^n\|_{L^\infty(\Omega)}, \Omega, h)$ and c_2^n is the constant in the trace theorem.

In the sequel we bound $\|u_h^{n+1} - u_h^n\|_{H^1(\Omega)}$ by means of the term $\|v_h^n - v_h^{n+1}\|_{L^2(\Omega)}$. To this aim, we subtract the equations for u_h^{n+1} and u_h^n (cf. (4.3)) and employ $\varphi = u_h^{n+1} - u_h^n$ as a test function. A standard manipulation yields

$$\|u_h^{n+1} - u_h^n\|_{H^1(\Omega)} \leq c_3^n \|v_h^n - v_h^{n+1}\|_{L^2(\Omega)}, \tag{4.20}$$

with $c_3^n = c_3^n(\Omega, \delta, \xi_0, h, \|\mathbb{C}_1 - \mathbb{C}_0\|_{L^\infty(\Omega)}, \|\widehat{\nabla} p_h^n\|_{L^\infty(\Omega)})$. Employing (4.20) into (4.19), we obtain

$$|I_1| \leq c_4^n \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2, \tag{4.21}$$

where $c_4^n = c_4^n(\Omega, \delta, \xi_0, h, \|(\mathbb{C}_1 - \mathbb{C}_0)\|_{L^\infty(\Omega)}, \|\widehat{\nabla} p_h^n\|_{L^\infty(\Omega)}, c_2^n)$.

Using (4.21) into (4.18), we deduce

$$\begin{aligned} \frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha}\varepsilon \|\nabla(v_h^{n+1} - v_h^n)\|_{L^2(\Omega)}^2 + \frac{\tilde{\alpha}}{\varepsilon} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 \\ + J_{\delta,\varepsilon,h}(v_h^{n+1}) \leq J_{\delta,\varepsilon,h}(v_h^n) + c_4^n \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2. \end{aligned} \quad (4.22)$$

Now, since

$$\begin{aligned} \frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + \tilde{\alpha}\varepsilon \|\nabla(v_h^{n+1} - v_h^n)\|_{L^2(\Omega)}^2 + \frac{\tilde{\alpha}}{\varepsilon} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 \\ \geq \frac{1}{\tau_n} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2, \end{aligned} \quad (4.23)$$

we get

$$\left(\frac{1}{\tau_n} - c_4^n\right) \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)}^2 + J_{\delta,\varepsilon,h}(v_h^{n+1}) \leq J_{\delta,\varepsilon,h}(v_h^n). \quad (4.24)$$

Finally, choosing $\tau_n \leq \frac{1}{1+c_4^n}$, the assertion of the lemma follows, just setting $c^n := c_4^n$. \square

We are now ready to state a convergence result for our numerical scheme.

Theorem 4.6. *Let $v_h^0 \in \mathcal{K}_h$ be an initial guess. Then there exists a collection of timesteps τ_n such that $0 < \gamma \leq \tau_n \leq (1 + c^n)^{-1}$, $\forall n > 0$, where c^n is the constant appearing in Lemma 4.5, and γ depends on the data and possibly on h . The corresponding sequence v_h^n generated by (4.16) has a convergence subsequence (still denoted by v_h^n) in $W^{1,\infty}$ such that*

$$v_h^n \rightarrow v_h, \quad n \rightarrow +\infty,$$

where $v_h \in \mathcal{K}_h$ satisfies the discrete optimality condition

$$J'_{\delta,\varepsilon,h}(v_h)[\omega_h - v_h] \geq 0, \quad \forall \omega_h \in \mathcal{K}_h.$$

Proof. Consider a collection of timesteps bounded by $(1 + c^n)^{-1}$, for all $n > 0$. Employing Lemma 4.5, we have

$$\sum_{n=0}^{+\infty} \|v_h^n - v_h^{n+1}\|_{L^2(\Omega)}^2 \leq J_{\delta,\varepsilon,h}(v_h^0), \quad (4.25)$$

$$\sup_{n \in \mathbb{N}} J_{\delta,\varepsilon,h}(v_h^n) \leq J_{\delta,\varepsilon,h}(v_h^0). \quad (4.26)$$

Hence, the sequence v_h^n is bounded in $H_0^1(\Omega)$ and it holds

$$\lim_{n \rightarrow +\infty} \|v_h^n - v_h^{n+1}\|_{L^2(\Omega)}^2 = 0. \quad (4.27)$$

From the weak formulation of the forward and adjoint problems, the previous relations give that u_h^n and p_h^n are bounded in $H^1(\Omega)$, hence in $W^{1,\infty}(\Omega)$ as we are in finite dimensional spaces. Therefore, thanks to the definition of the constant c^n , reported in the last part of the proof of Lemma 4.5, this gives that there exists a constant $M > 0$ such that $c^n < M$, and equivalently there exists a positive constant $\gamma > 0$, independent of n , such that $\gamma \leq (1 + c^n)^{-1}$. Hence, there exists a subsequence of (v_h^n, u_h^n, p_h^n) (still denoted by the same symbol) such that

$$(v_h^n, u_h^n, p_h^n) \rightarrow (v_h, u_h, p_h) \quad \text{in } W^{1,\infty}(\Omega),$$

and in particular

$$u_h^n \rightarrow u_h \text{ a.e. in } \Omega, \quad p_h^n \rightarrow p_h \text{ a.e. in } \Omega.$$

Hence, u_h is the solution of the discrete forward problem and p_h is the solution of the discrete adjoint problem. Finally, from (4.16) and $\tau_n \geq \gamma$ we get

$$\begin{aligned} \int_{\Omega} (\mathbb{C}_0 - \mathbb{C}_1)(\omega_h - v_h^{n+1}) \widehat{\nabla} u_h^n : \widehat{\nabla} p_h^n + 2\tilde{\alpha}\varepsilon \int_{\Omega} \widehat{\nabla} v_h^{n+1} \cdot \widehat{\nabla}(\omega_h - v_h^{n+1}) \\ + \frac{\tilde{\alpha}}{\varepsilon} \int_{\Omega} (1 - 2v_h^n)(\omega_h - v_h^{n+1}) \geq -\frac{C}{\gamma} \|v_h^{n+1} - v_h^n\|_{L^2(\Omega)} \|\omega_h - v_h^{n+1}\|_{L^2(\Omega)}. \end{aligned}$$

From (4.27) and recalling that $v_h^n \rightarrow v_h$, we deduce that v_h satisfies the discrete optimality condition (4.5). \square

5 Numerical Examples

In this section we show the numerical results which are obtained from an application of the Primal Dual Active Set Method (PDASM) to the variational inequality (4.16). This method has been presented in [40] and later applied for the detection of conductivity inclusions in [30] and [14] for a linear and a semilinear elliptic equation, respectively. Primal dual active set methods represent a very good choice in engineering applications due to their effectiveness and robustness (cf., e.g., [38]). Here, we show that choosing the parameter δ sufficiently small we are able to reconstruct elastic cavities of different shapes. Given a tolerance $\text{tol} > 0$, the reconstruction algorithm is based on the following steps.

Algorithm 1 Discrete Parabolic Obstacle Problem

```

Set  $n = 0$  and  $v_h^0 = v_0$ , the initial guess for the inclusion
while  $\|v_h^n - v_h^{n-1}\| > \text{tol}$  do
    find solution of the forward problem (4.3) with  $v = v_h^n$ 
    find solution of the adjoint problem (4.6) with  $v = v_h^n$ 
    find  $v_h^{n+1}$  solving (4.16) via PDASM algorithm
    update  $n = n + 1$ ;
end while

```

In the implementation of Algorithm 1, the numerical experiments are performed for $d = 2$ in the domain $\Omega = (-1, 1)^2$, using a triangular tessellation \mathcal{T}_h of Ω . As boundary measurements, we use synthetic data. They are generated by solving via the Finite Element method the forward problem (2.1), with boundary conditions prescribed as in Figure 2a on the square, with one or more cavities of given geometries. We use a tessellation \mathcal{T}_h^{ref} which is more refined than \mathcal{T}_h on the common part outside the cavities (see Figure 1 for an example of the two tessellations) in order not to commit inverse crime. Once extracting the values of the solution of the forward problem on the boundary of the domain Ω obtained by the mesh \mathcal{T}_h^{ref} , we interpolate these values on the mesh \mathcal{T}_h . Therefore, by u_{meas} we denote

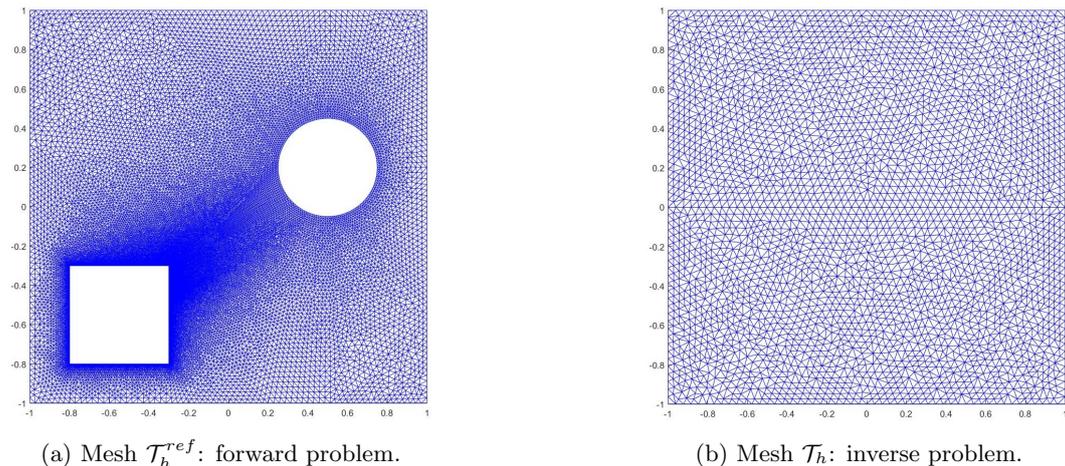
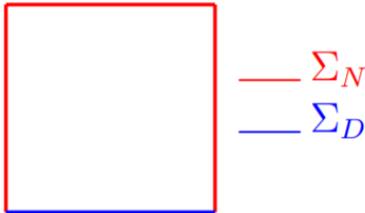
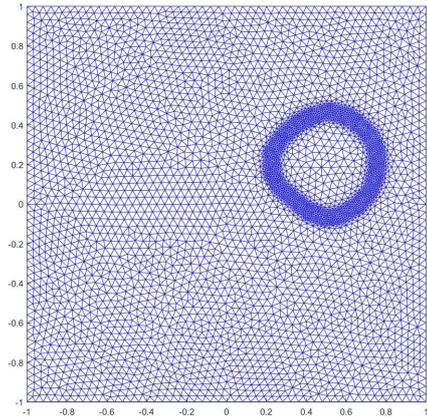


Figure 1: Example of the meshes adopted.

the resulting boundary datum on the mesh \mathcal{T}_h . We also mention that the triangular mesh is adaptively refined during the reconstruction procedure using the values of ∇v_h after an a-priori fixed number of iterations which depend on the specific numerical example. See, as example, Figure 2b related to the reconstruction of a circular cavity.



(a) Boundary condition in numerical experiments: Neumann boundary conditions are assigned on the red part. Homogeneous Dirichlet conditions are assigned on the blue part.



(b) Refinement of the mesh around the reconstructed domain. This is the mesh at the final iteration of the experiment in Figure 4b.

Figure 2: Geometrical setting and refinement of the mesh.

In the reconstruction procedure, i.e. for the implementation of the Algorithm 1, we assume to know two different boundary measurements. In fact, in the context of inverse boundary value problems of this kind, it is reasonable to use $N_g > 1$ different boundary measurements u_{meas}^i , for $i = 1, \dots, N_g$ which clearly improve the numerical reconstruction results. Thus, we consider a slight modification of the original optimization problem (3.8), assuming the knowledge of N_g different Neumann boundary data g^i , for $i = 1, \dots, N_g$ and hence considering

$$\begin{aligned} & \min_{v \in \mathcal{K}} J_{\delta, \varepsilon}^{sum}(v), \\ J_{\delta, \varepsilon}^{sum}(v) & := \frac{1}{N_g} \sum_{i=1}^{N_g} \left(\frac{1}{2} \|u_{\delta}^i(v) - u_{meas}^i\|_{L^2(\Sigma_N)}^2 \right) + \tilde{\alpha} \int_{\Omega} \left(\varepsilon |\nabla v|^2 + \frac{1}{\varepsilon} v(1-v) \right), \end{aligned} \quad (5.1)$$

where $u_{\delta}^i(v) \in H_{\Sigma_D}^1(\Omega)$ is the solution to (3.2) with $g = g^i$ and for $v \in \mathcal{K}$. The necessary optimality condition related to (5.1) can be equivalently obtained reasoning similarly as we did to derive (3.10).

In Table 1, we collect some of the parameters utilized in most numerical tests. Possible changes in these values are highlighted in the text related to each specific experiment.

Finally, all the numerical experiments are performed choosing, as initial guess, the phase-field variable $v_0 = 0$.

tol	$\tilde{\alpha}$	τ_n	ε	δ
10^{-5}	10^{-2}	10^{-3}	$\frac{1}{16\pi}$ or $\frac{1}{8\pi}$	10^{-2}

Table 1: Values of some parameters utilized in Algorithm 1.

5.1 Numerical experiments with $N_g = 2$ and without noise in the measurements.

Test 1: reconstruction of a circular cavity. The elastic medium is described by the Lamé parameters $\mu = 0.2$ and $\lambda = 1$. The Neumann boundary conditions are $g^1(x, y) = (0, \frac{1}{10} - \frac{3}{10}y)$ and $g^2(x, y) = (-\frac{1}{2}x^2, y^2)$. We set the parameter $\varepsilon = \frac{1}{16\pi}$. The mesh is refined with respect to the gradient of the phase-field variable every 1000 iterations. The algorithm

stops after $n = 3544$ iterations. In Figure 3 we show the numerical results at three different time steps.

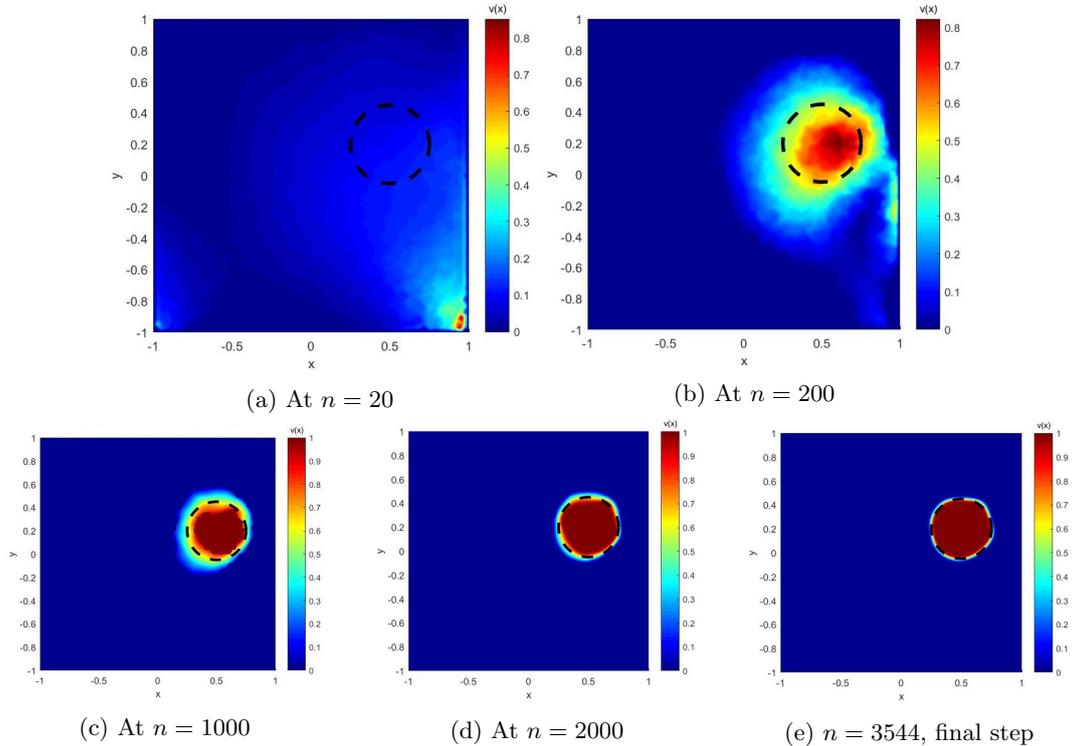


Figure 3: Test 1. Reconstruction of a circular cavity. Dotted line represents the target cavity.

Test 2: reconstruction of a circular cavity - changing boundary conditions and Lamé parameters. We propose the same numerical experiments of Test 1, showing how the results change using different Neumann boundary conditions and Lamé parameters. We report in the captions of Figure 4 the selected parameters, data, and also the number of time steps needed for reaching the tolerance. Note that the three experiments consider different values for the Poisson coefficient $\nu := \frac{\lambda}{2(\lambda + \mu)}$, that is $\nu = \frac{1}{4}$, $\nu = \frac{1}{3}$, and $\nu = -\frac{1}{18}$, respectively. In the three numerical examples of Figure 4, the refinement of the mesh happens every 1500, 1000, 2000 iterations, respectively.

Test 3: reconstruction of a Lipschitz domain. This experiment aims at reconstructing a square-shaped cavity. We show several numerical tests, choosing different values for ε , different boundary conditions and different values of the number of iterations for the refinement of the mesh. We have already shown results based on different choices for the values of the Lamé parameters in the previous numerical tests, so we fix the values of Lamé coefficients to be $\mu = 0.5$ and $\lambda = 1$. In fact, recalling that the range of the Poisson coefficient is $-1 < \nu < \frac{1}{2}$ ($\nu = \frac{1}{2}$ represents the incompressible case), we have considered four relevant cases for the Poisson coefficient: one test on an elastic material close to incompressible case ($\nu = \frac{5}{12}$ in Figure 3e), two tests on elastic coefficients of common materials ($\nu = \frac{1}{4}$ and $\nu = \frac{1}{3}$ in Figures 4a and 4b, respectively), and one test on auxetic materials, that is materials with negative Poisson ratio ($\nu = -\frac{1}{18}$ in Figure 4c). In the results of Figure 5, the refinement of the mesh happens every 6000 for the first two experiments and every 3000 iterations for the last one. The second numerical result, see Figure 5b, has the same parameters of the numerical example of Figure 5a except $\tilde{\alpha}$ which is chosen $\tilde{\alpha} = 5 \times 10^{-2}$.

Test 4: reconstruction of two cavities. This test provides results when the two cavities to be reconstructed are a square and a circle. Neumann boundary conditions are given by $g^1(x, y) = (x, y)$ and $g^2(x, y) = (-y, -x)$. We propose two numerical reconstruction procedures, see Figure 6. In Figure 6a, we report the results obtained by the standard

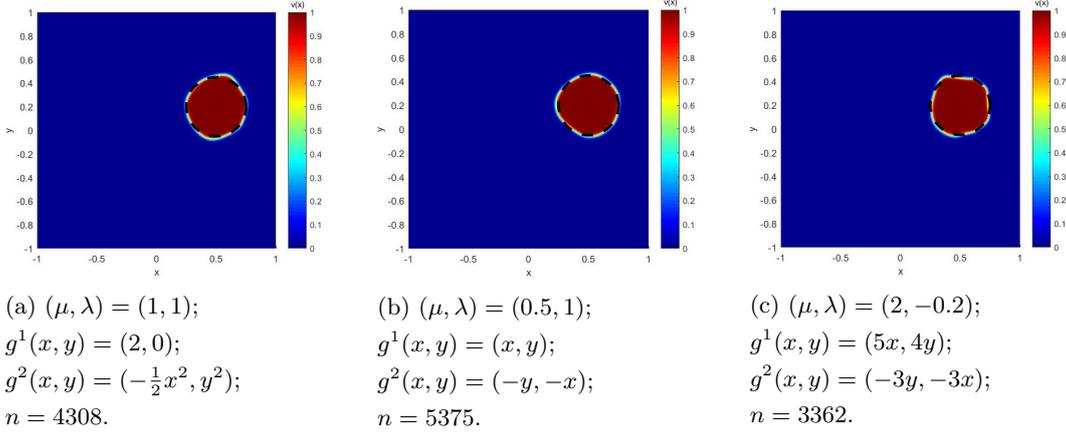


Figure 4: Test 2. Reconstruction of a circular cavity using several parameters and data. For each experiment, we report the configuration at the final step n . Dotted line represents the target cavity.

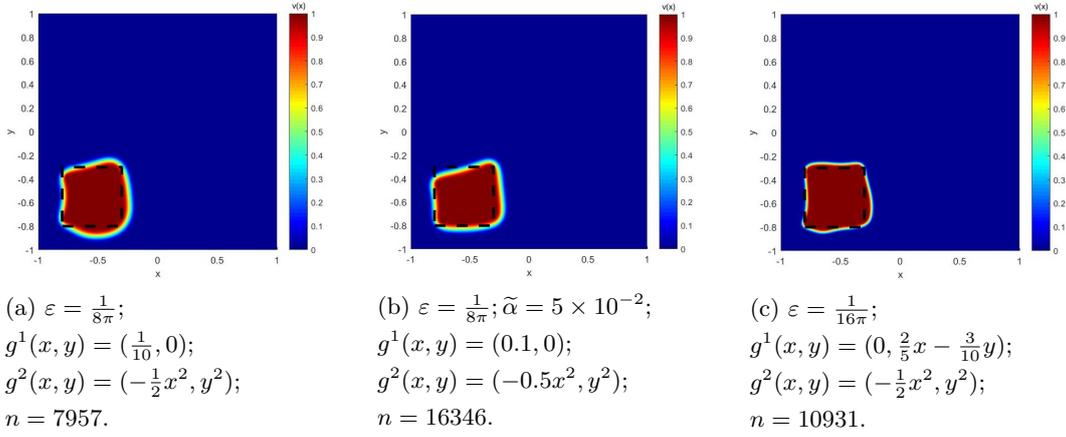


Figure 5: Test 3. Reconstruction of a square-shaped cavity. Dotted line represents the target cavity.

algorithm, while in Figure 6b we use a variant of the Algorithm 1 where the parameter ε is initially set $\varepsilon = \frac{1}{4\pi}$ but after a fixed and a-priori chosen number of iterations (8000 iterations) is updated and set $\varepsilon = \varepsilon/4$. In both cases the mesh is refined after 5000 iterations. It is worth noting that the variant of Algorithm 1 does not produce the visible oscillations of the test in Figure 6a.

Note that we also change a little bit the value of δ . We have observed that δ cannot be chosen too small otherwise numerical instability can appear. Numerically we have seen that, in order to overcome this issue, τ_n has to be chosen always smaller than δ . However, choosing τ too small increases the number of necessary iterations to satisfy the stopping criterium.

Test 5: reconstruction of a non-convex domain. We finally propose the reconstruction of a cavity which is not convex, see Figure 7. We use $g^1(x, y) = (x, y)$ and $g^2(x, y) = (-y, -x)$ as Neumann boundary conditions and $\mu = 0.5$ and $\lambda = 1$. Parameters have the following values: $\varepsilon = \frac{1}{16\pi}$, and $\tau_n = 5 \times 10^{-4}$. Mesh is refined every 5000 iterations. The stopping criterium is satisfied after $n = 6825$ iterations.

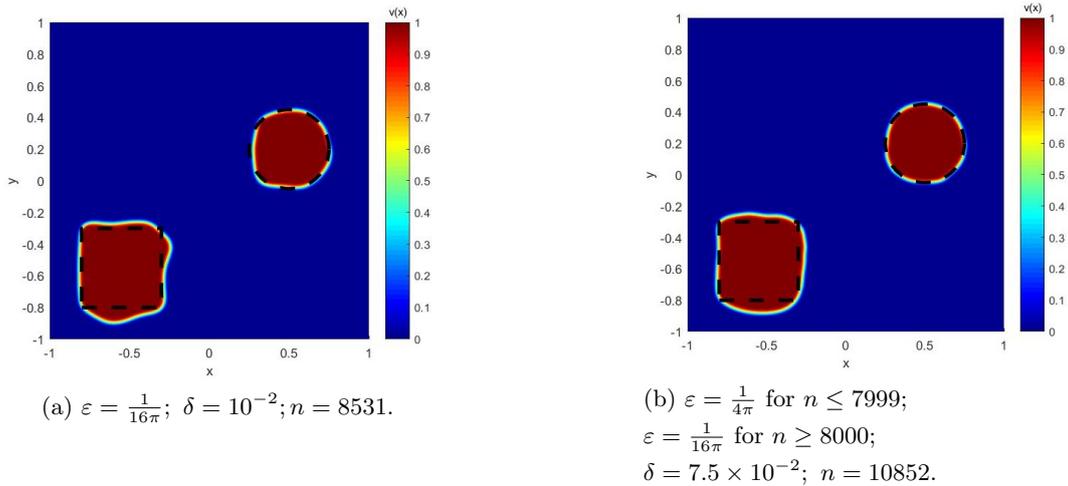


Figure 6: Test 4. Reconstruction of two cavities. Dotted line represents the target cavity.

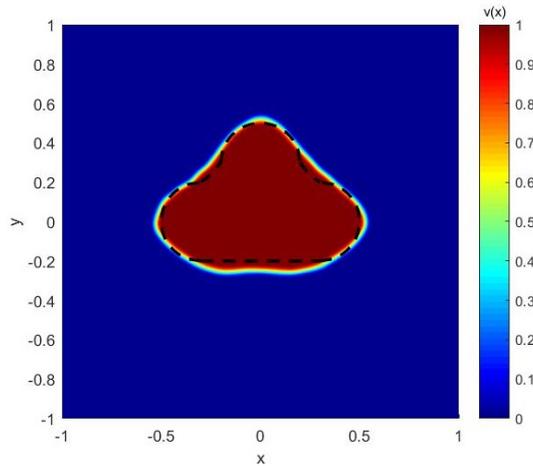
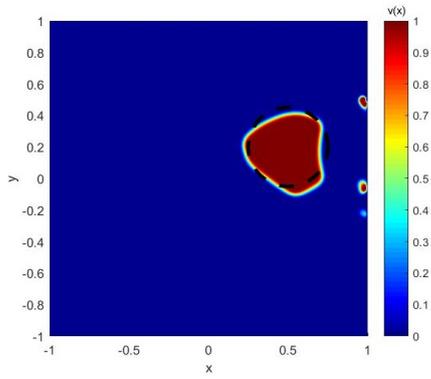


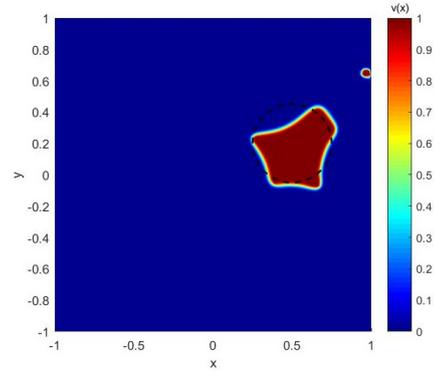
Figure 7: Test 5. Reconstruction of a non-convex domain. It seems that the algorithm tends to reconstruct a convex domain. Dotted line represents the target cavity.

5.2 Numerical experiments with $N_g = 2$ and noise in the measurements.

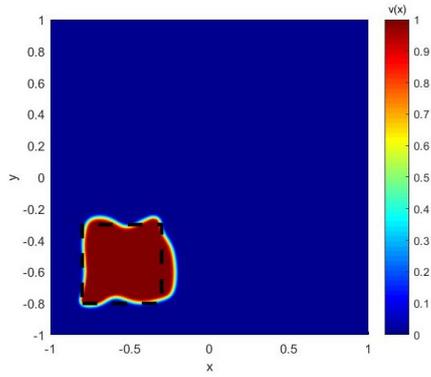
Test 6: reconstruction of cavities of different shapes using noisy measurements. Here we run some of the numerical tests showed in the previous section, adding to the boundary measurements a normal distributed noise with zero mean and variance equal to one. We choose two different noise levels: 2% and 5%. The results are reported in Figure 8. For the the test in Figure 8a and Figure 8b, we use values of parameters as in Test 1 and refine the mesh every 2000 and 2500 iterations, respectively. The reconstruction of a square-shaped cavity, that is Figure 8c and Figure 8d, are obtained by means of parameters of Test 3 - Figure 5c, refining the mesh every 3000 and 10000 iterations. Lastly, to get the results in Figure 8e and Figure 8f we use the same parameters of Test 4 - Figure 6b. The mesh is refined every 5000 and 8000 iterations, while the value of the parameter ε is adapted after 8000 and 10000 iterations, respectively. In the captions of the single figures, we specify the values that are changed with respect to the ones proposed in the Tests 1, 3, and 4.



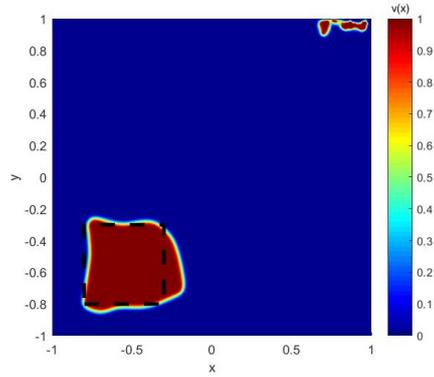
(a) $(\mu, \lambda) = (0.2, 1)$, $\tau_n = 5 \times 10^{-4}$.
 $n = 11071$. Noise level 2%.



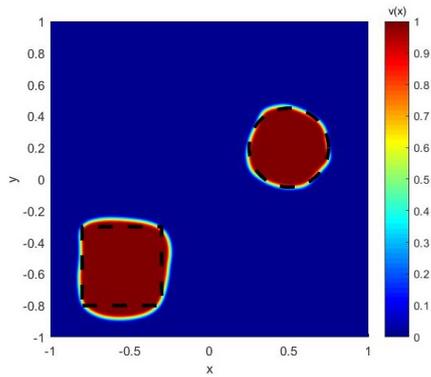
(b) $(\mu, \lambda) = (0.2, 1)$, $\tau_n = 5 \times 10^{-4}$.
 $\tilde{\alpha} = 5 \times 10^{-2}$, $n = 14361$.
 Noise level 5%.



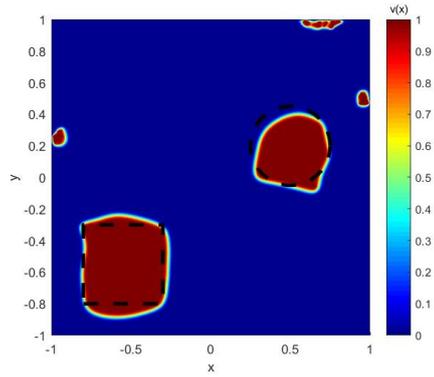
(c) $(\mu, \lambda) = (0.5, 1)$, $\tau_n = 5 \times 10^{-4}$.
 $n = 23652$. Noise level 2%.



(d) $(\mu, \lambda) = (0.5, 1)$, $\tau_n = 5 \times 10^{-4}$.
 $n = 15854$. Noise level 5%.



(e) $(\mu, \lambda) = (0.5, 1)$.
 $n = 11776$. Noise level 2%.



(f) $(\mu, \lambda) = (0.5, 1)$, $\tau_n = 5 \times 10^{-4}$.
 $n = 25480$. Noise level 5%.

Figure 8: Test 6. Reconstruction of cavities by means of noisy measurements. Dotted line represents the target cavity.

Acknowledgments

The authors deeply thank Dorin Bucur and Alessandro Giacomini for suggesting relevant literature and for useful discussions that led us to improve some of the results in this work.

This research has been partially performed in the framework of the MIUR-PRIN Grant 2020F3NCPX “Mathematics for industry 4.0 (Math4I4)”.

Andrea Aspri, Cecilia Cavaterra and Elisabetta Rocca are members of GNAMPA (Gruppo Nazionale per l’Analisi Matematica, la Probabilità e le loro Applicazioni) of INdAM (Istituto Nazionale di Alta Matematica).

Marco Verani has been partially funded by MIUR PRIN research grants n. 201744KLJL and n. 20204LN5N5. Marco Verani is a member of GNCS (Gruppo Nazionale per il Calcolo Scientifico) of INdAM.

References

- [1] G. Alberti. Variational models for phase transitions, an approach via Γ -convergence. In *Calculus of variations and partial differential equations (Pisa, 1996)*, pages 95–114. Springer, Berlin, 2000.
- [2] G. Alessandrini, A. Morassi, and E. Rosset. The linear constraints in Poincaré and Korn type inequalities. *Forum Math.*, 20(3):557–569, 2008.
- [3] S. Almi and U. Stefanelli. Topology optimization for incremental elastoplasticity: a phase-field approach. *SIAM J. Control Optim.*, 59(1):339–364, 2021.
- [4] L. Ambrosio, N. Fusco, and D. Pallara. *Functions of bounded variation and free discontinuity problems*. Oxford Mathematical Monographs. The Clarendon Press, Oxford University Press, New York, 2000.
- [5] H.B. Ameer, M. Burger, and B. Hackl. Cavity identification in linear elasticity and thermoelasticity. *Math. Methods Appl. Sci.*, 30(6):625–647, 2007.
- [6] H. Ammari. *An introduction to mathematics of emerging biomedical imaging*, volume 62 of *Mathématiques & Applications*. Springer, Berlin, 2008.
- [7] H. Ammari, E. Bretin, J. Garnier, H. Kang, H. Lee, and A. Wahab. *Mathematical methods in elasticity imaging*. Princeton Series in Applied Mathematics. Princeton University Press, Princeton, NJ, 2015.
- [8] H. Ammari, H. Kang, G. Nakamura, and K. Tanuma. Complete asymptotic expansions of solutions of the system of elastostatics in the presence of an inclusion of small diameter and detection of an inclusion. *J. Elasticity*, 67(2):97–129 (2003), 2002.
- [9] A. Aspri. *An elastic model for volcanology*. Lecture Notes in Geosystems Mathematics and Computing. Birkhäuser/Springer, Cham, 2019.
- [10] A. Aspri, E. Beretta, and C. Mascia. Analysis of a Mogi-type model describing surface deformations induced by a magma chamber embedded in an elastic half-space. *J. Éc. polytech. Math.*, 4:223–255, 2017.
- [11] A. Aspri, E. Beretta, and E. Rosset. On an elastic model arising from volcanology: an analysis of the direct and inverse problem. *J. Differential Equations*, 265(12):6400–6423, 2018.
- [12] F. Auricchio, E. Bonetti, M. Carraturo, D. Hömberg, A. Reali, and E. Rocca. A phase-field-based graded-material topology optimization with stress constraint. *Math. Models Methods Appl. Sci.*, 30(8):1461–1483, 2020.
- [13] E. Beretta, M.C. Cerutti, and D. Pierotti. Detection of cavities in a nonlinear model arising from cardiac electrophysiology via γ -convergence. *arXiv 2106.04213*, 2021.

- [14] E. Beretta, L. Ratti, and M. Verani. Detection of conductivity inclusions in a semi-linear elliptic problem arising from cardiac electrophysiology. *Commun. Math. Sci.*, 16(7):1975–2002, 2018.
- [15] L. Blank, H. Garcke, M.H. Farshbaf-Shaker, and V. Styles. Relating phase field and sharp interface approaches to structural topology optimization. *ESAIM Control Optim. Calc. Var.*, 20(4):1025–1058, 2014.
- [16] L. Blank, H. Garcke, C. Hecht, and C. Rupprecht. Sharp interface limit for a phase field model in structural optimization. *SIAM J. Control Optim.*, 54(3):1558–1584, 2016.
- [17] M. Bonnet and A. Constantinescu. Inverse problems in elasticity. *Inverse Problems*, 21(2):R1–R50, 2005.
- [18] B. Bourdin and A. Chambolle. Design-dependent loads in topology optimization. *ESAIM Control Optim. Calc. Var.*, 9:19–48, 2003.
- [19] H. Brezis. *Analyse fonctionnelle*. Collection Mathématiques Appliquées pour la Maîtrise. [Collection of Applied Mathematics for the Master’s Degree]. Masson, Paris, 1983. Théorie et applications. [Theory and applications].
- [20] B. M. Brown, M. Jais, and I. W. Knowles. A variational approach to an elastic inverse problem. *Inverse Problems*, 21(6):1953–1973, 2005.
- [21] D. Bucur and G. Buttazzo. *Variational methods in shape optimization problems*, volume 65 of *Progress in Nonlinear Differential Equations and their Applications*. Birkhäuser Boston, Inc., Boston, MA, 2005.
- [22] D. Bucur, A. Henrot, J. Sokółowski, and A. Żochowski. Continuity of the elasticity system solutions with respect to the geometrical domain variations. *Adv. Math. Sci. Appl.*, 11(1):57–73, 2001.
- [23] D. Bucur and N. Varchon. Stabilité de la solution d’un problème de Neumann pour des variations de frontière. *C. R. Acad. Sci. Paris Sér. I Math.*, 331(5):371–374, 2000.
- [24] A. Carpio and M.L. Rapún. Topological derivatives for shape reconstruction. In *Inverse problems and imaging*, volume 1943 of *Lecture Notes in Math.*, pages 85–133. Springer, Berlin, 2008.
- [25] M. Carraturo, E. Rocca, E. Bonetti, D. Hömberg, A. Reali, and F. Auricchio. Graded-material design based on phase-field and topology optimization. *Comput. Mech.*, 64(6):1589–1600, 2019.
- [26] A. Chambolle and F. Doveri. Continuity of Neumann linear elliptic problems on varying two-dimensional bounded open sets. *Comm. Partial Differential Equations*, 22(5-6):811–840, 1997.
- [27] D. Chenaïs. On the existence of a solution in a domain identification problem. *J. Math. Anal. Appl.*, 52(2):189–219, 1975.
- [28] P.G. Ciarlet. *Mathematical elasticity. Vol. I*, volume 20 of *Studies in Mathematics and its Applications*. North-Holland Publishing Co., Amsterdam, 1988. Three-dimensional elasticity.
- [29] G. Dal Maso. *An introduction to Γ -convergence*, volume 8 of *Progress in Nonlinear Differential Equations and their Applications*. Birkhäuser Boston, Inc., Boston, MA, 1993.
- [30] K. Deckelnick, C.M. Elliott, and V. Styles. Double obstacle phase field approach to an inverse problem for a discontinuous diffusion coefficient. *Inverse Problems*, 32(4):045008, 26, 2016.

- [31] A. Doubova and E. Fernández-Cara. Some geometric inverse problems for the Lamé system with applications in elastography. *Appl. Math. Optim.*, 82(1):1–21, 2020.
- [32] S. Eberle and B. Harrach. Shape reconstruction in linear elasticity: standard and linearized monotonicity method. *Inverse Problems*, 37(4):045006, 27, 2021.
- [33] H. Eiliat and J. Urbanic. Visualizing, analyzing, and managing voids in the material extrusion process. *Int J Adv Manuf Technol*, 96:4095–4109, 2018.
- [34] L.C. Evans and R.F. Gariepy. *Measure theory and fine properties of functions*. Textbooks in Mathematics. CRC Press, Boca Raton, FL, revised edition, 2015.
- [35] H. Garcke, C. Hecht, M. Hinze, and C. Kahle. Numerical approximation of phase field based shape and topology optimization for fluids. *SIAM J. Sci. Comput.*, 37(4):A1846–A1871, 2015.
- [36] H. Garcke, K. Lam Fong, R. Nürnberg, and A. Signori. Overhang penalization in additive manufacturing via phase field structural topology optimization with anisotropic energies. <https://arxiv.org/pdf/2111.14070>, 2021.
- [37] A. Giacomini. A stability result for Neumann problems in dimension $N \geq 3$. *J. Convex Anal.*, 11(1):41–58, 2004.
- [38] Xiahui He and Peng Yang. The primal-dual active set method for a class of nonlinear problems with T -monotone operators. *Math. Probl. Eng.*, pages Art. ID 2912301, 8, 2019.
- [39] A. Henrot and M. Pierre. *Shape variation and optimization*, volume 28 of *EMS Tracts in Mathematics*. European Mathematical Society (EMS), Zürich, 2018. A geometrical analysis, English version of the French publication [MR2512810] with additions and updates.
- [40] M. Hintermüller, K. Ito, and K. Kunisch. The primal-dual active set strategy as a semismooth Newton method. *SIAM J. Optim.*, 13(3):865–888 (2003), 2002.
- [41] S. Hubmer, E. Sherina, A. Neubauer, and O. Scherzer. Lamé parameter estimation from static displacement field measurements in the framework of nonlinear inverse problems. *SIAM J. Imaging Sci.*, 11(2):1268–1293, 2018.
- [42] M. Ikehata and H. Itou. On reconstruction of an unknown polygonal cavity in a linearized elasticity with one measurement. *Journal of Physics: Conference Series*, 290:012005, apr 2011.
- [43] M. Ikehata and H. Itou. On reconstruction of a cavity in a linearized viscoelastic body from infinitely many transient boundary data. *Inverse Problems*, 28(12):125003, nov 2012.
- [44] B. Jin and J. Zou. Numerical estimation of the Robin coefficient in a stationary diffusion equation. *IMA J. Numer. Anal.*, 30(3):677–701, 2010.
- [45] H. Kang, E. Kim, and J.-Y. Lee. Identification of elastic inclusions and elastic moment tensors by boundary measurements. *Inverse Problems*, 19(3):703–724, 2003.
- [46] A. Karageorghis, D. Lesnic, and L. Marin. The method of fundamental solutions for the detection of rigid inclusions and cavities in plane linear elastic bodies. *Computers & Structures*, 106, 2012.
- [47] T. Kurahashi, K. Maruoka, and T. Iyama. Numerical shape identification of cavity in three dimensions based on thermal non-destructive testing data. *Engineering Optimization*, 49(3):434–448, 2017.

- [48] K.F. Lam and I. Yousept. Consistency of a phase field regularisation for an inverse problem governed by a quasilinear Maxwell system. *Inverse Problems*, 36(4):045011, 33, 2020.
- [49] H. Liu, L. Rondi, and J. Xiao. Mosco convergence for $H(\text{curl})$ spaces, higher integrability for Maxwell’s equations, and stability in direct and inverse EM scattering problems. *J. Eur. Math. Soc. (JEMS)*, 21(10):2945–2993, 2019.
- [50] A.E. Martínez-Castro, I.H. Faris, and R. Gallego. Identification of cavities in a three-dimensional layer by minimization of an optimal cost functional expansion. *Computer Modeling in Engineering & Sciences*, 87(3):177–206, 2012.
- [51] G. Menegatti and L. Rondi. Stability for the acoustic scattering problem for sound-hard scatterers. *Inverse Probl. Imaging*, 7(4):1307–1329, 2013.
- [52] L. Modica. The gradient theory of phase transitions and the minimal interface criterion. *Arch. Rational Mech. Anal.*, 98(2):123–142, 1987.
- [53] A. Morassi and E. Rosset. Stable determination of cavities in elastic bodies. *Inverse Problems*, 20(2):453–480, 2004.
- [54] A. Morassi and E. Rosset. Stable determination of an inclusion in an inhomogeneous elastic body by boundary measurements. *Rend. Istit. Mat. Univ. Trieste*, 48:101–120, 2016.
- [55] T.D. Ngo, A. Kashani, G. Imbalzano, K.T.Q. Nguyen, and D. Hui. Additive manufacturing (3d printing): A review of materials, methods, applications and challenges. *Composites Part B: Engineering*, 143:172–196, 2018.
- [56] W. Ring and L. Rondi. Reconstruction of cracks and material losses by perimeter-like penalizations and phase-field methods: numerical results. *Interfaces Free Bound.*, 13(3):353–371, 2011.
- [57] L. Rondi. Reconstruction of material losses by perimeter penalization and phase-field methods. *J. Differential Equations*, 251(1):150–175, 2011.
- [58] P. Segall. *Earthquake and volcano deformation*. Princeton University Press, Princeton, NJ, 2010.
- [59] J. Shao, G. Shi, Z. Qi, J. Zheng, and S. Chen. Advancements in the application of ultrasound elastography in the cervix. *Ultrasound in Medicine & Biology*, 47(8):2048–2063, 2021.
- [60] E. Sherina, L. Krainz, S. Hubmer, W. Drexler, and O. Scherzer. Displacement field estimation from OCT images utilizing speckle information with applications in quantitative elastography. *Inverse Problems*, 36(12):124003, 27, 2020.
- [61] J. Shi, E. Beretta, M.V. de Hoop, E. Francini, and S. Vessella. A numerical study of multi-parameter full waveform inversion with iterative regularization using multi-frequency vibroseis data. *Comput. Geosci.*, 24(1):89–107, 2020.
- [62] W. W. Symes. The seismic reflection inverse problem. *Inverse Problems*, 25(12):123008, 39, 2009.
- [63] S.A. Tronvoll, T. Welo, and C.W. Elverum. The effects of voids on structural properties of fused deposition modelled parts: a probabilistic approach. *The International Journal of Advanced Manufacturing Technology*, 97(9):3607–3618, Aug 2018.
- [64] T. Widlak and O. Scherzer. Stability in the linearized problem of quantitative elastography. *Inverse Problems*, 31(3):035005, 27, 2015.