



On an elastic model arising from volcanology: An analysis of the direct and inverse problem

A. Aspri ^{a,*}, E. Beretta ^b, E. Rosset ^c

^a *Johann Radon Institute for Computational and Applied Mathematics (RICAM), Altenbergerstraße 69, 4040 Linz, Austria*

^b *Dipartimento di Matematica, Politecnico di Milano, Via Edoardo Bonardi 9, 20133 Milano, Italy*

^c *Dipartimento di Matematica e Geoscienze, Università degli Studi di Trieste, via Valerio 12/1, 34127 Trieste, Italy*

Received 20 June 2017; revised 4 June 2018

Available online 23 July 2018

Abstract

In this paper we investigate a mathematical model arising from volcanology describing surface deformation effects generated by a magma chamber embedded into Earth's interior and exerting on it a uniform hydrostatic pressure. The modeling assumptions translate mathematically into a Neumann boundary value problem for the classical Lamé system in a half-space with an embedded pressurized cavity. We establish well-posedness of the problem in suitable weighted Sobolev spaces and analyse the inverse problem of determining the pressurized cavity from partial measurements of the displacement field proving uniqueness and stability estimates.

© 2018 Elsevier Inc. All rights reserved.

MSC: 35R30; 35J57; 74B05; 86A60

Keywords: Lamé system; Neumann problem; Half-space; Weighted Sobolev spaces; Inverse problem; Stability estimates

* Corresponding author.

E-mail addresses: andrea.aspri@ricam.oeaw.ac.at (A. Aspri), elena.beretta@polimi.it (E. Beretta), rossedi@units.it (E. Rosset).

1. Introduction

In this paper we investigate a linear elastic model describing surface deformations in a volcanic area induced by a magma chamber embedded in Earth’s crust. From the mathematical point of view we introduce a simplified version of the model assuming the crust to be a half-space in a homogeneous and isotropic medium and the magma chamber to be a cavity subjected to a constant pressure on its boundary (for more details see, for example, [7,8,12,17]). More precisely, given \mathbb{C} a fourth-order isotropic and homogeneous elastic tensor, with Lamé parameters λ and μ , denoting by \mathbf{u} the displacement vector and by \mathbb{R}^3_- the half-space, we end up with the following linear elastostatic boundary value problem

$$\begin{cases} \operatorname{div}(\mathbb{C}\widehat{\nabla}\mathbf{u}) = \mathbf{0} & \text{in } \mathbb{R}^3_- \setminus \overline{C} \\ (\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n} = p\mathbf{n} & \text{on } \partial C \\ (\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{e}_3 = \mathbf{0} & \text{on } \mathbb{R}^2 \end{cases} \tag{1}$$

where $\widehat{\nabla}\mathbf{u}$ is the strain tensor, C is the cavity, $p > 0$ represents the pressure acting on the boundary of the cavity, \mathbf{n} is the outer unit normal vector on ∂C and $\mathbf{e}_3 = (0, 0, 1)$.

The main purpose of this paper is to derive quantitative stability estimates for the inverse problem of identifying the pressurized cavity C from one measurement of the displacement provided on a portion of the boundary of the half-space.

In order to address this issue, we first analyse the well-posedness of (1) under the assumption that ∂C is Lipschitz. We highlight that for the well-posedness we can either impose explicitly some decay conditions at infinity for \mathbf{u} and $\nabla\mathbf{u}$ (see, for example, [7]) or, more suitably for our purposes, set the analysis in some weighted Sobolev spaces where the decay conditions are expressed by means of weights. In particular, we will show the well-posedness in this weighted Sobolev space

$$H_w^1(\mathbb{R}^3_- \setminus \overline{C}) = \left\{ \mathbf{u} \in \mathcal{D}'(\mathbb{R}^3_- \setminus \overline{C}), \frac{\mathbf{u}}{(1 + |\mathbf{x}|^2)^{1/2}} \in L^2(\mathbb{R}^3_- \setminus \overline{C}), \nabla\mathbf{u} \in L^2(\mathbb{R}^3_- \setminus \overline{C}) \right\}, \tag{2}$$

where $\mathcal{D}'(\mathbb{R}^3_- \setminus \overline{C})$ is the space of distributions in $\mathbb{R}^3_- \setminus \overline{C}$, with the norm given by

$$\|\mathbf{u}\|_{H_w^1(\mathbb{R}^3_- \setminus \overline{C})}^2 = \left(\|(1 + |\mathbf{x}|^2)^{-1/2}\mathbf{u}\|_{L^2(\mathbb{R}^3_- \setminus \overline{C})}^2 + \|\nabla\mathbf{u}\|_{L^2(\mathbb{R}^3_- \setminus \overline{C})}^2 \right). \tag{3}$$

Even if the analysis of the well-posedness of general elastic problems in the half-space via weighted Sobolev spaces is known, see [6], we would like to emphasize that in our framework we have two principal difficulties and novelties to treat: the first one concerns the fact that the problem is stated in an unbounded domain with unbounded boundary and with non homogeneous Neumann boundary conditions on the cavity. The second one is related to the derivation of quantitative stability estimates of the solution in $H_w^1(\mathbb{R}^3_- \setminus \overline{C})$. To this end, we need to derive a quantitative weighted Poincaré inequality and a Korn-type inequality in $\mathbb{R}^3_- \setminus \overline{C}$. As far as we know, these inequalities are known in a quantitative way only for bounded domains, see [4], and for conical domains, see [11].

Therefore, the first part of this paper is devoted to prove quantitative Poincaré and Korn inequalities in $\mathbb{R}^3_- \setminus \overline{C}$ using some a priori information on the cavity C . This allows us to derive the following estimate for the unique solution \mathbf{u} of problem (1)

$$\|\mathbf{u}\|_{H^1_{\omega}(\mathbb{R}^3 \setminus \bar{C})} \leq cP,$$

where the constant c depends on the Lamé parameters, on the Lipschitz character of ∂C and on the distance of C from the boundary of the half-space. This estimate is fundamental for the direct problem and it is also necessary to establish stability estimates for the cavity in terms of the measurements.

To prove the stability result for the inverse problem we need stronger regularity on the cavity. We follow and adapt when needed the results contained in [15,16]. From the point of view of the rate of convergence it is well known that, despite of smoothness assumptions on the cavity, only a weak rate of logarithmic type is expected. In our case we are able to prove a log–log type estimate and not the optimal logarithmic one proved for the scalar case (see [1]) due to the lack of a doubling inequality at the boundary for the solutions of the Lamé system.

The paper is organized as follows. In Section 2 we give the notation used in the rest of the paper and definition on the regularity of the domains. In Section 3, we set the analysis of the elastic problem in the weighted Sobolev space, proving first the constructive Poincaré and Korn inequalities and then giving the result of the well-posedness. Section 4 is devoted to the analysis of the inverse problem and the derivation of the uniqueness and stability results.

2. On some notation and useful definitions

In this section we set up notation and some definitions paying specific attention to the regularity of bounded domains.

In the sequel, we denote scalar quantities in italic type, e.g. λ, μ, ν , points and vectors in bold italic type, e.g. $\mathbf{x}, \mathbf{y}, \mathbf{z}$ and $\mathbf{u}, \mathbf{v}, \mathbf{w}$, matrices and second-order tensors in bold type, e.g. $\mathbf{A}, \mathbf{B}, \mathbf{C}$, and fourth-order tensors in blackboard bold type, e.g. $\mathbb{A}, \mathbb{B}, \mathbb{C}$.

The transpose of a second-order tensor \mathbf{A} is denoted by \mathbf{A}^T and its symmetric part by

$$\widehat{\mathbf{A}} = \frac{1}{2} (\mathbf{A} + \mathbf{A}^T).$$

To indicate the inner product between two vectors \mathbf{u} and \mathbf{v} we use $\mathbf{u} \cdot \mathbf{v} = \sum_i u_i v_i$ whereas for second-order tensors $\mathbf{A} : \mathbf{B} = \sum_{i,j} a_{ij} b_{ij}$. The tensor product of two vectors \mathbf{u} and \mathbf{v} is denoted by $\mathbf{u} \otimes \mathbf{v} = u_i v_j$. Similarly, $\mathbf{A} \otimes \mathbf{B} = A_{ij} B_{hk}$ represents the tensor product between matrices. With $|\mathbf{A}|$ we mean the norm induced by the inner product between second-order tensors, that is

$$|\mathbf{A}| = \sqrt{\mathbf{A} : \mathbf{A}}.$$

We denote the open half-space

$$\{\mathbf{x} = (x_1, x_2, x_3) \in \mathbb{R}^3 : x_3 < 0\} = \mathbb{R}^3_-$$

and we represent with \mathbb{R}^2 its boundary, that is the set $\{\mathbf{x} = (x_1, x_2, x_3) \in \mathbb{R}^3 : x_3 = 0\}$. The set $B_r^-(\mathbf{0})$ denotes the half ball of center $\mathbf{0}$ and radius r , that is

$$B_r^-(\mathbf{0}) = \{\mathbf{x} \in \mathbb{R}^3 : x_1^2 + x_2^2 + x_3^2 < r^2, x_3 < 0\}.$$

With $B'_r(\mathbf{0})$ we mean the circle of center $\mathbf{0}$ and radius r , namely

$$B'_r(\mathbf{0}) = \{\mathbf{x} \in \mathbb{R}^2 : x_1^2 + x_2^2 < r^2\}.$$

We denote with $d(A, B)$ the distance between the two sets A and B , that is

$$d(A, B) := \inf\{|\mathbf{x} - \mathbf{y}| : \mathbf{x} \in A, \mathbf{y} \in B\}$$

and with $d_{\mathcal{H}}(A, B)$ their Hausdorff distance, namely

$$d_{\mathcal{H}}(A, B) := \max\{\sup_{\mathbf{x} \in A} \inf_{\mathbf{y} \in B} |\mathbf{x} - \mathbf{y}|, \sup_{\mathbf{y} \in B} \inf_{\mathbf{x} \in A} |\mathbf{x} - \mathbf{y}|\}.$$

The unit outer normal vector at the boundary of a regular domain is represented by \mathbf{n} .

2.1. Domain regularity

In the following sections, the constants appearing in the inequalities will depend on some a priori information of the constitutive parameters of the linear elastic model and on the a priori geometric and regularity assumptions on the cavity. For this reason it is important to recall the definition of $C^{k,\alpha}$ regularity for a bounded domain.

Definition 2.1 ($C^{k,\alpha}$ regularity). Let Ω be a bounded domain in \mathbb{R}^3 . Given k, α , with $k \in \mathbb{N}$ and $0 < \alpha \leq 1$, we say that a portion S of $\partial\Omega$ is of class $C^{k,\alpha}$ with constant r_0, E_0 , if for any $\mathbf{P} \in S$, there exists a rigid transformation of coordinates under which we have $\mathbf{P} = \mathbf{0}$ and

$$\Omega \cap B_{r_0}(\mathbf{0}) = \{\mathbf{x} \in B_{r_0}(\mathbf{0}) : x_3 > \psi(\mathbf{x}')\},$$

where ψ is a $C^{k,\alpha}$ function on $B'_{r_0}(\mathbf{0}) \subset \mathbb{R}^2$ such that

$$\begin{aligned} \psi(\mathbf{0}) &= 0, \\ \nabla\psi(\mathbf{0}) &= \mathbf{0}, \quad \text{for } k \geq 1 \\ \|\psi\|_{C^{k,\alpha}(B_{r_0}(\mathbf{0}))} &\leq E_0. \end{aligned}$$

When $k = 0, \alpha = 1$, we also say that S is of Lipschitz class with constants r_0, E_0 .

3. The direct problem

In this section, we will analyse the well-posedness of the following linear elastostatic boundary value problem

$$\begin{cases} \operatorname{div}(\mathbb{C}\widehat{\nabla}\mathbf{u}) = \mathbf{0} & \text{in } \mathbb{R}^3_- \setminus \overline{C} \\ (\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n} = p\mathbf{n} & \text{on } \partial C \\ (\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{e}_3 = \mathbf{0} & \text{on } \mathbb{R}^2 \end{cases} \tag{4}$$

where $p > 0$ represents the pressure, C is the cavity, \mathbf{n} is the outer unit normal vector on ∂C and $\mathbf{e}_3 = (0, 0, 1)$. \mathbb{C} is the fourth-order isotropic and homogeneous elastic tensor given by

$$\mathbb{C} := \lambda \mathbf{I} \otimes \mathbf{I} + 2\mu \mathbb{I},$$

where λ and μ are the two constants Lamé parameters, \mathbf{I} is the identity matrix in \mathbb{R}^3 and \mathbb{I} is the fourth-order identity tensor such that $\mathbb{I}\mathbf{A} = \widehat{\mathbf{A}}$, for any second-order tensor \mathbf{A} .

We first provide some physical information on the cavity C and the elastic tensor \mathbb{C} .

3.1. Main assumptions and a priori information

For the study of the direct problem (4), we assume that the constant Lamé parameters satisfy the inequalities

$$3\lambda + 2\mu > 0 \quad \text{and} \quad \mu > 0, \tag{5}$$

that is the tensor \mathbb{C} is strongly convex:

$$\mathbb{C}\widehat{\mathbf{A}} : \widehat{\mathbf{A}} \geq \xi_0 |\widehat{\mathbf{A}}|^2, \tag{6}$$

with $\xi_0 = \min\{2\mu, 2\mu + 3\lambda\}$, see [9].

The cavity C is supposed to be a bounded domain with Lipschitz regularity, that is

$$\partial C \text{ is Lipschitz with constants } r_0 \text{ and } E_0. \tag{7}$$

Additionally, we impose some a priori information on the size of the cavity and its distance from the boundary of the half-space. In particular, denoting with $\text{diam}(A)$ the diameter of a set A , we require

$$B_{2D_0}^-(\mathbf{0}) \supset C, \tag{8}$$

$$d(C, \mathbb{R}^2) \geq D_0, \tag{9}$$

$$\text{diam}(C) < D_0, \tag{10}$$

where, without loss of generality, we can assume that the constant $D_0 > 1$.

Remark 3.1. From here on, for simplicity of reading, we omit the dependence of some constants on the Lamé coefficients λ and μ , on the parameters r_0, E_0, D_0 related to the a priori information on the cavity C and on s_0 which represents the radius of the circle where the measurements are collected. For its definition see Section 4.

We highlight that, since we are in the half-space, namely an unbounded domain with unbounded boundary, the study of the well-posedness of the direct problem (4) can be done either imposing some decay conditions at infinity for the function \mathbf{u} and $\nabla \mathbf{u}$ (see for example [7]) or setting the analysis in a suitable weighted Sobolev space. We choose this second strategy. So we recall the definition of the weighted Sobolev space for domains of type $\mathbb{R}^3 \setminus \overline{C}$. To do that, we denote the space of the indefinitely differentiable functions with compact support in $\mathbb{R}^3 \setminus \overline{C}$ by $\mathcal{D}(\mathbb{R}^3 \setminus \overline{C})$ and with $\mathcal{D}'(\mathbb{R}^3 \setminus \overline{C})$ its dual space, that is the space of distributions.

Definition 3.1 (*Weighted Sobolev space*). Given the function

$$\rho = (1 + |\mathbf{x}|^2)^{1/2}, \tag{11}$$

we define

$$H_w^1(\mathbb{R}_-^3 \setminus \overline{C}) = \left\{ \mathbf{u} \in \mathcal{D}'(\mathbb{R}_-^3 \setminus \overline{C}), \frac{\mathbf{u}}{\rho} \in L^2(\mathbb{R}_-^3 \setminus \overline{C}), \nabla \mathbf{u} \in L^2(\mathbb{R}_-^3 \setminus \overline{C}) \right\}. \tag{12}$$

This weighted Sobolev space is a reflexive Banach space, see for example [5] and references therein, equipped with its natural norm

$$\|\mathbf{u}\|_{H_w^1(\mathbb{R}_-^3 \setminus \overline{C})}^2 = \left(\|\rho^{-1}\mathbf{u}\|_{L^2(\mathbb{R}_-^3 \setminus \overline{C})}^2 + \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}_-^3 \setminus \overline{C})}^2 \right). \tag{13}$$

We recall that the weight is chosen so that the space $\mathcal{D}(\overline{\mathbb{R}_-^3} \setminus C)$ is dense in $H_w^1(\mathbb{R}_-^3 \setminus \overline{C})$, see [5,10]. When we deal with bounded domains $D \subset (\mathbb{R}_-^3 \setminus \overline{C})$ we emphasize that $H_w^1(D)$ reduces to $H^1(D)$ regularity, hence the usual trace theorems hold. For generalizations and more details on weighted Sobolev spaces see, for example, [5,6,10] and references therein.

The study of the well-posedness of (4) is therefore accomplished using Lax–Milgram theorem in $H_w^1(\mathbb{R}_-^3 \setminus \overline{C})$ space. We stress that the use of the weighted Sobolev space is the natural approach to obtain a quantitative $H_w^1(\mathbb{R}_-^3 \setminus \overline{C})$ estimate in terms of the boundary data. To this end we need first to have constructive Poincaré and Korn-type inequalities.

3.2. Weighted Poincaré inequality and Korn-type inequality

In this section we want to prove a weighted Poincaré inequality and a Korn-type inequality in $\mathbb{R}_-^3 \setminus \overline{C}$ using a suitable partition of unity.

Let us consider two half balls $B_r^-(\mathbf{0})$ and $B_R^-(\mathbf{0})$, with $r < R$, such that

$$C \subset B_r^-(\mathbf{0}) \subset B_R^-(\mathbf{0}).$$

Using the a priori information (9) and (10) on the cavity C , we fix

$$r = 3D_0 \quad \text{and} \quad R = 4D_0.$$

We consider a specific partition of unity of \mathbb{R}_-^3 . In particular, we take $\varphi_1, \varphi_2 \in C^\infty(\mathbb{R}_-^3)$ such that

$$0 \leq \varphi_1, \varphi_2 \leq 1 \quad \text{and} \quad \varphi_1 + \varphi_2 = 1 \quad \text{in } \mathbb{R}_-^3, \tag{14}$$

with

$$\varphi_2 = 0, \quad \varphi_1 = 1, \quad \text{in } B_r^-(\mathbf{0}), \tag{15}$$

$$\varphi_1 = 0, \quad \varphi_2 = 1, \quad \text{in } \{|\mathbf{x}| \geq R\} \cap \mathbb{R}_-^3, \tag{16}$$

$$|\nabla \varphi_1| \leq \frac{c}{\rho}, \quad |\nabla \varphi_2| \leq \frac{c}{\rho}, \quad \text{in } B_R^-(\mathbf{0}) \setminus B_r^-(\mathbf{0}), \tag{17}$$

where c is an absolute positive constant thanks to the choice made for r and R .

In the following it is useful to split $\partial B_R^-(\mathbf{0}) = \partial B_R^h(\mathbf{0}) \cup \partial B_R^b(\mathbf{0})$, where $\partial B_R^h(\mathbf{0})$ is the circle $B'_R(\mathbf{0})$ whereas $\partial B_R^b(\mathbf{0})$ is the spherical cap. We first prove Poincaré inequality

Theorem 3.2 (Weighted Poincaré inequality). *For any function $\mathbf{u} \in H_w^1(\mathbb{R}^3 \setminus \bar{C})$ there exists a positive constant c , with $c = c(r_0, E_0, D_0)$, such that*

$$\int_{\mathbb{R}^3 \setminus \bar{C}} \left| \frac{\mathbf{u}}{\rho} \right|^2 dx \leq c \int_{\mathbb{R}^3 \setminus \bar{C}} |\nabla \mathbf{u}|^2 dx, \tag{18}$$

where ρ is defined in (11).

Proof. From (14) we find that

$$\left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 \leq 2 \left(\left\| \varphi_1 \frac{\mathbf{u}}{\rho} \right\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 + \left\| \varphi_2 \frac{\mathbf{u}}{\rho} \right\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 \right) := 2(\mathcal{N}_1 + \mathcal{N}_2).$$

We study \mathcal{N}_1 and \mathcal{N}_2 .

From the property (16) and since $\rho^{-1} \leq 1$, we get

$$\mathcal{N}_1 = \left\| \varphi_1 \frac{\mathbf{u}}{\rho} \right\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \leq \|\varphi_1 \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2. \tag{19}$$

Therefore, since $\varphi_1 = 0$ on $\partial B_R^b(\mathbf{0})$, we use the quantitative Poincaré inequality for functions vanishing on a portion of the boundary of a bounded domain, see for instance [4] (Theorem 3.3, in particular Example 3.6), finding

$$\|\varphi_1 \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \leq c \|\nabla(\varphi_1 \mathbf{u})\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2, \tag{20}$$

where c is a positive constant such that $c = c(r_0, E_0, D_0)$. In this way, we obtain

$$\|\nabla(\varphi_1 \mathbf{u})\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \leq 2 \left(\|\mathbf{u} \otimes \nabla \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 + \|\varphi_1 \nabla \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \right). \tag{21}$$

Now, from the property (17), we have

$$\|\mathbf{u} \otimes \nabla \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 = \int_{B_R^-(\mathbf{0}) \setminus \bar{C}} |\mathbf{u}|^2 |\nabla \varphi_1|^2 dx \leq c \int_{B_R^-(\mathbf{0}) \setminus B_r^-(\mathbf{0})} \frac{|\mathbf{u}|^2}{\rho^2} dx. \tag{22}$$

By (20), (21), (22) and recalling (14), we have

$$\|\varphi_1 \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \leq c \left(\left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 + \|\nabla \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \right). \tag{23}$$

Applying Hardy’s inequality for the exterior of a half ball in the half-space (see [11], Lemma 3, p. 83) to the first term in the right-hand side of the inequality (23), we find

$$\left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)}^2 \leq c \|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)}^2. \tag{24}$$

Inserting (24) in (23) and then going back to (19), we have

$$\begin{aligned} \mathcal{N}_1 &= \left\| \varphi_1 \frac{\mathbf{u}}{\rho} \right\|_{L^2(B_R^-(\mathbf{0}) \setminus \overline{C})}^2 \leq c \left(\|\nabla \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \overline{C})}^2 + \|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)}^2 \right) \\ &\leq c \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2. \end{aligned} \tag{25}$$

Analogously, using the properties (14) and (15) and applying again Hardy’s inequality (see [11], Lemma 3, p. 83), we find

$$\begin{aligned} \mathcal{N}_2 &= \left\| \varphi_2 \frac{\mathbf{u}}{\rho} \right\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2 \leq \left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)}^2 \leq c \|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)}^2 \\ &\leq c \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2. \end{aligned} \tag{26}$$

Putting together the inequalities (25) and (26) we have the assertion. \square

Before proving a Korn-type inequality in the exterior domain of a half-space, we state, for the reader’s convenience, a slight modification of a lemma proved by Kondrat’ev and Oleinik in [11] (see Lemma 5. p. 85) for the case of a function $\mathbf{u} \in H_w^1(\mathbb{R}^3_+ \setminus \overline{C})$. This lemma will be useful in the proof of the Korn inequality.

Lemma 3.3. *Let $\mathbf{u} \in H_w^1(\mathbb{R}^3_+ \setminus \overline{C})$. For every $r' < r$ there exists a positive constant c such that*

$$\|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3_+)} \leq c \|\widehat{\nabla} \mathbf{u}\|_{L^2(\{|x|>r'\} \cap \mathbb{R}^3_+)},$$

where $c = c(r, r')$.

Now, we are ready to prove the following quantitative Korn inequality.

Theorem 3.4 (Korn-type inequality). *For any function $\mathbf{u} \in H_w^1(\mathbb{R}^3_+ \setminus \overline{C})$ there exists a positive constant c , with $c = c(r_0, E_0, D_0)$, such that*

$$\int_{\mathbb{R}^3_+ \setminus \overline{C}} |\nabla \mathbf{u}|^2 dx \leq c \int_{\mathbb{R}^3_+ \setminus \overline{C}} |\widehat{\nabla} \mathbf{u}|^2 dx. \tag{27}$$

Proof. From the definition of the functions φ_1, φ_2 , see (14), we have

$$\|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2 \leq 2 \left(\|\nabla(\varphi_1 \mathbf{u})\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2 + \|\nabla(\varphi_2 \mathbf{u})\|_{L^2(\mathbb{R}^3_+ \setminus \overline{C})}^2 \right) := 2(\mathcal{N}'_1 + \mathcal{N}'_2).$$

We study, separately, the two terms \mathcal{N}'_1 and \mathcal{N}'_2 .

By (16) we find

$$\mathcal{N}'_1 = \|\nabla(\varphi_1 \mathbf{u})\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 = \|\nabla(\varphi_1 \mathbf{u})\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2, \tag{28}$$

hence, since $\varphi_1 = 0$ on $\partial B_R^b(\mathbf{0})$, we apply the quantitative Korn inequality for functions vanishing on a portion of the boundary of a bounded domain, see for instance [4] (Theorem 5.7), getting for $c = c(r_0, E_0, D_0)$

$$\begin{aligned} \|\nabla(\varphi_1 \mathbf{u})\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 &\leq c \|\widehat{\nabla}(\varphi_1 \mathbf{u})\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \\ &\leq c \left(\|\mathbf{u} \otimes \nabla \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 + \|\varphi_1 \widehat{\nabla} \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \right), \end{aligned}$$

where in the right side of the previous inequality we have used

$$\|\widehat{\mathbf{u}} \otimes \widehat{\nabla} \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \leq \|\mathbf{u} \otimes \nabla \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2.$$

From (22), the properties (14) and Hardy’s inequality (24), we have

$$\begin{aligned} \|\mathbf{u} \otimes \nabla \varphi_1\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 + \|\varphi_1 \widehat{\nabla} \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \\ \leq c \left(\left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 + \|\widehat{\nabla} \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \right) \\ \leq c \left(\|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 + \|\widehat{\nabla} \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 \right). \end{aligned}$$

Applying Lemma 3.3 to the first term in the right side of the previous formula, we find

$$\|\nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 \leq c \|\widehat{\nabla} \mathbf{u}\|_{L^2(\{|x|>r'\} \cap \mathbb{R}^3)}^2,$$

where we choose $r' = 2D_0 < r$ so that $C \subset B_{r'}^-(\mathbf{0})$ (see the a priori information (8)). Putting together all these results and going back to (28), we have

$$\mathcal{N}'_1 \leq c \left(\|\widehat{\nabla} \mathbf{u}\|_{L^2(B_R^-(\mathbf{0}) \setminus \bar{C})}^2 + \|\widehat{\nabla} \mathbf{u}\|_{L^2(\{|x|>r'\} \cap \mathbb{R}^3)}^2 \right) \leq c \|\widehat{\nabla} \mathbf{u}\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2. \tag{29}$$

In a similar way, using the properties (14) and (15), we find

$$\begin{aligned} \mathcal{N}'_2 = \|\nabla(\varphi_2 \mathbf{u})\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 &= \|\nabla(\varphi_2 \mathbf{u})\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 \\ &\leq c \left(\|\mathbf{u} \otimes \nabla \varphi_2\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 + \|\varphi_2 \nabla \mathbf{u}\|_{L^2(\{|x|>r\} \cap \mathbb{R}^3)}^2 \right). \end{aligned}$$

From the properties (14) and (17) of φ_2 , we get

$$\begin{aligned} & \| \mathbf{u} \otimes \nabla \varphi_2 \|_{L^2(\{|x|>r\} \cap \mathbb{R}_-^3)}^2 + \| \varphi_2 \nabla \mathbf{u} \|_{L^2(\{|x|>r\} \cap \mathbb{R}_-^3)}^2 \\ & \leq c \left(\left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\{|x|>r\} \cap \mathbb{R}_-^3)}^2 + \| \nabla \mathbf{u} \|_{L^2(\{|x|>r\} \cap \mathbb{R}_-^3)}^2 \right). \end{aligned}$$

Using again Hardy’s inequality (24) and the result in Lemma 3.3 in the last two terms of the previous formula, we find

$$\mathcal{N}'_2 \leq c \| \widehat{\nabla} \mathbf{u} \|_{L^2(\{|x|>r'\} \cap \mathbb{R}_-^3)}^2 \leq c \| \widehat{\nabla} \mathbf{u} \|_{L^2(\mathbb{R}_-^3 \setminus \overline{C})}^2. \tag{30}$$

Finally, collecting the results in (29) and (30), we have the assertion. \square

3.3. Well-posedness

To study the well-posedness of problem (4) we use a variational approach. We suppose, for the moment, \mathbf{u} regular and the test functions \mathbf{v} in $\mathcal{D}(\overline{\mathbb{R}_-^3} \setminus C)$. Multiplying the equations in (4) for the functions \mathbf{v} and integrating in $\mathbb{R}_-^3 \setminus \overline{C}$, we obtain

$$\int_{\mathbb{R}_-^3 \setminus \overline{C}} \mathbb{C} \widehat{\nabla} \mathbf{u} : \widehat{\nabla} \mathbf{v} \, dx = -p \int_{\partial C} \mathbf{n} \cdot \mathbf{v} \, d\sigma(\mathbf{x}), \quad \forall \mathbf{v} \in \mathcal{D}(\overline{\mathbb{R}_-^3} \setminus C).$$

Now, from the density property of the functional space $\mathcal{D}(\overline{\mathbb{R}_-^3} \setminus C)$ into the weighted Sobolev space defined in (12), problem (4) becomes:

find $\mathbf{u} \in H_w^1(\mathbb{R}_-^3 \setminus \overline{C})$ such that

$$a(\mathbf{u}, \mathbf{v}) = f(\mathbf{v}), \quad \forall \mathbf{v} \in H_w^1(\mathbb{R}_-^3 \setminus \overline{C}), \tag{31}$$

where $a : H_w^1(\mathbb{R}_-^3 \setminus \overline{C}) \times H_w^1(\mathbb{R}_-^3 \setminus \overline{C}) \rightarrow \mathbb{R}$ is the bilinear form given by

$$a(\mathbf{u}, \mathbf{v}) = \int_{\mathbb{R}_-^3 \setminus \overline{C}} \mathbb{C} \widehat{\nabla} \mathbf{u} : \widehat{\nabla} \mathbf{v} \, dx, \tag{32}$$

and $f : H_w^1(\mathbb{R}_-^3 \setminus \overline{C}) \rightarrow \mathbb{R}$ is the linear functional given by

$$f(\mathbf{v}) = -p \int_{\partial C} \mathbf{n} \cdot \mathbf{v} \, d\sigma(\mathbf{x}). \tag{33}$$

Now, we can prove

Theorem 3.5. *Problem (4) admits a unique solution $\mathbf{u} \in H_w^1(\mathbb{R}_-^3 \setminus \overline{C})$ satisfying*

$$\| \mathbf{u} \|_{H_w^1(\mathbb{R}_-^3 \setminus \overline{C})} \leq c p, \tag{34}$$

where the constant $c = c(\lambda, \mu, r_0, E_0, D_0)$.

Proof. To prove the well-posedness of problem (4) we apply Lax–Milgram theorem to (31). Therefore, we need to prove the coercivity and the continuity property of the bilinear form (32) and the boundedness of the linear functional (33).

Continuity of (32).

From the Cauchy–Schwarz inequality we have

$$|a(\mathbf{u}, \mathbf{v})| = \left| \int_{\mathbb{R}^3 \setminus \bar{C}} \mathbb{C} \widehat{\nabla} \mathbf{u} : \widehat{\nabla} \mathbf{v} \, dx \right| \leq c \|\widehat{\nabla} \mathbf{u}\|_{L^2(\mathbb{R}^3 \setminus \bar{C})} \|\widehat{\nabla} \mathbf{v}\|_{L^2(\mathbb{R}^3 \setminus \bar{C})} \\ \leq c \|\mathbf{u}\|_{H_w^1(\mathbb{R}^3 \setminus \bar{C})} \|\mathbf{v}\|_{H_w^1(\mathbb{R}^3 \setminus \bar{C})},$$

where $c = c(\lambda, \mu)$.

Coercivity of (32).

We apply the constructive Poincaré and Korn inequalities proved in Theorem 3.2, Theorem 3.4 and the strong convexity condition of \mathbb{C} , see (6). In detail, we have

$$a(\mathbf{u}, \mathbf{u}) = \int_{\mathbb{R}^3 \setminus \bar{C}} \mathbb{C} \widehat{\nabla} \mathbf{u} : \widehat{\nabla} \mathbf{u} \, dx \geq c \|\widehat{\nabla} \mathbf{u}\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 \\ \geq c \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3 \setminus \bar{C})}^2 \geq c \|\mathbf{u}\|_{H_w^1(\mathbb{R}^3 \setminus \bar{C})}^2,$$

where the constant $c = c(\lambda, \mu, r_0, E_0, D_0)$.

Boundedness of (33).

Let us take $B_{2D_0}^-(\mathbf{0})$. Then applying the trace theorem for bounded domains, we find

$$\left| -p \int_{\partial C} \mathbf{n} \cdot \mathbf{v} \, d\sigma(\mathbf{x}) \right| \leq c p \|\mathbf{v}\|_{L^2(\partial C)} \leq c p \left(\left\| \frac{\mathbf{v}}{\rho} \right\|_{L^2((B_{2D_0}^-(\mathbf{0})) \setminus \bar{C})} + \|\nabla \mathbf{v}\|_{L^2((B_{2D_0}^-(\mathbf{0})) \setminus \bar{C})} \right) \\ \leq c p \|\mathbf{v}\|_{H_w^1(\mathbb{R}^3 \setminus \bar{C})}.$$

Applying the Lax–Milgram theorem we obtain the well-posedness of problem (4). Moreover, by means of the strong convexity condition of \mathbb{C} , see (6), and from the application of the Korn and Poincaré inequalities, see Theorem 3.2 and Theorem 3.4, we find that

$$\|\mathbf{u}\|_{H_w^1(\mathbb{R}^3 \setminus C)}^2 \leq \left| \int_{\mathbb{R}^3 \setminus \bar{C}} \mathbb{C} \widehat{\nabla} \mathbf{u} : \widehat{\nabla} \mathbf{u} \, dx \right| \leq \left| p \int_{\partial C} \mathbf{n} \cdot \mathbf{u} \, d\sigma(\mathbf{x}) \right| \leq c p \|\mathbf{u}\|_{H_w^1(\mathbb{R}^3 \setminus \bar{C})},$$

where the constant $c = c(\lambda, \mu, r_0, E_0, D_0)$, hence the assertion of the theorem follows. \square

4. The inverse problem: uniqueness and stability estimate

In this section we will investigate the following inverse problem: *given the displacement vector \mathbf{u} on a portion of the boundary of the half-space can we detect uniquely and in a stable way the cavity C ?*

We suppose to have the measurements on $B'_{s_0}(\mathbf{0}) = \{\mathbf{x} \in \mathbb{R}^2 : x_1^2 + x_2^2 < s_0^2\}$ contained in $\{x_3 = 0\}$, with $s_0 < D_0$. To prove a stability estimate for the inverse problem we need to require more regularity on C than (7). In particular, we suppose that:

$$\partial C \text{ is of class } C^3 \text{ with constant } r_0 \text{ and } E_0. \tag{35}$$

In addition, we recall that C satisfy the a priori information (8), (9) and (10). We also assume that

$$\mathbb{R}^3 \setminus \overline{C} \text{ is connected.}$$

Before proceeding, we highlight that the proof of the uniqueness and the stability result is based on the possibility to build the displacement field

$$\overline{\mathbf{u}} = \frac{p}{3\lambda + 2\mu} \mathbf{x} \tag{36}$$

so to reduce problem (4) to a problem with homogeneous Neumann boundary conditions on the boundary of the cavity. A straightforward calculation shows that $\overline{\mathbf{u}}$ satisfies the Lamé system and the boundary condition on C satisfied by \mathbf{u} .

In this way, the function

$$\mathbf{w} := \mathbf{u} - \overline{\mathbf{u}}, \tag{37}$$

satisfies the following boundary value problem

$$\begin{cases} \operatorname{div}(\mathbb{C}\widehat{\nabla}\mathbf{w}) = \mathbf{0} & \text{in } \mathbb{R}^3 \setminus \overline{C} \\ (\mathbb{C}\widehat{\nabla}\mathbf{w})\mathbf{n} = 0 & \text{on } \partial C \\ (\mathbb{C}\widehat{\nabla}\mathbf{w})\mathbf{e}_3 = -p\mathbf{e}_3 & \text{on } \mathbb{R}^2 \\ \mathbf{w} + \overline{\mathbf{u}} \in H_w^1(\mathbb{R}^3 \setminus \overline{C}), \end{cases} \tag{38}$$

where $\mathbf{e}_3 = (0, 0, 1)$. The inverse problem reduces therefore to determine the cavity C from a single pair of Cauchy data on $B'_{s_0}(\mathbf{0})$ of the solution to problem (38).

In the sequel, we denote $\mathbf{w}_i = \mathbf{u}_i - \overline{\mathbf{u}}$, for $i = 1, 2$, where \mathbf{w}_i and \mathbf{u}_i are respectively the solutions to (38) and (4) with $C = C_i$, for $i = 1, 2$. It immediately follows that

$$\mathbf{w}_1 - \mathbf{w}_2 = \mathbf{u}_1 - \mathbf{u}_2, \quad \text{in } \mathbb{R}^3 \setminus \overline{(C_1 \cup C_2)}. \tag{39}$$

Moreover, we indicate with

$$G \text{ the unbounded connected component of } \mathbb{R}^3 \setminus \overline{(C_1 \cup C_2)}. \tag{40}$$

Notice that $B'_{s_0}(\mathbf{0}) \subset \partial G$.

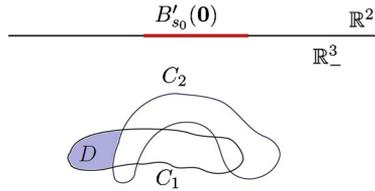


Fig. 1. The two domains C_1 and C_2 and the domain D .

4.1. Uniqueness

Although the procedure to get the uniqueness result for the inverse problem is known in the literature, for the reader’s convenience, we give a sketch of the proof.

Theorem 4.1. *Given the single pair of Cauchy data $\{\mathbf{w}, -pe_3\}$ on $B'_{s_0}(\mathbf{0})$ there exists at most one pair (C, \mathbf{w}) satisfying problem (38).*

Remark 4.2. Theorem 4.1 is true under weaker regularity assumptions on the cavity C . In fact, it is sufficient to have ∂C of Lipschitz class.

Proof of Theorem 4.1. Suppose by contradiction that there exist two cavities C_1 and C_2 , with $C_1 \neq C_2$, and the corresponding vector displacements $\mathbf{w}_1, \mathbf{w}_2$ such that

$$\mathbf{w}_1|_{B'_{s_0}(\mathbf{0})} = \mathbf{w}_2|_{B'_{s_0}(\mathbf{0})} = \mathbf{w}, \quad (\mathbb{C}\widehat{\nabla}\mathbf{w}_1)\mathbf{e}_3|_{B'_{s_0}(\mathbf{0})} = (\mathbb{C}\widehat{\nabla}\mathbf{w}_2)\mathbf{e}_3|_{B'_{s_0}(\mathbf{0})} = -pe_3.$$

From the unique continuation theorem for solution to the Lamé system, see [19], we have

$$\mathbf{w}_1 = \mathbf{w}_2, \quad \text{in } G,$$

where G is defined in (40). Next, we consider two different cases of intersection of the domains C_1 and C_2 . In fact, all the other possibilities can be reduced to these two configurations. For example, we take the domain D as in Fig. 1, where the function \mathbf{w}_2 is well-defined and satisfies the elastostatic equations, finding

$$\int_D (\mathbb{C}\widehat{\nabla}\mathbf{w}_2) : \widehat{\nabla}\mathbf{w}_2 \, dx = \int_{\partial D} (\mathbb{C}\widehat{\nabla}\mathbf{w}_2\mathbf{n}) \cdot \mathbf{w}_2 \, d\sigma(x).$$

Since $\mathbb{C}\widehat{\nabla}\mathbf{w}_2\mathbf{n} = 0$ on $\partial D \cap \partial C_2$ and by the unique continuation property $\mathbb{C}\widehat{\nabla}\mathbf{w}_2\mathbf{n} = \mathbb{C}\widehat{\nabla}\mathbf{w}_1\mathbf{n}$ on $\Gamma_1 = (\partial D \cap \partial C_1) \subset \partial G$, we get

$$\int_D (\mathbb{C}\widehat{\nabla}\mathbf{w}_2) : \widehat{\nabla}\mathbf{w}_2 \, dx = \int_{\Gamma_1} (\mathbb{C}\widehat{\nabla}\mathbf{w}_1\mathbf{n}) \cdot \mathbf{w}_2 \, d\sigma(x) = 0,$$

hence $\mathbf{w}_2 = \mathbf{A}\mathbf{x} + \mathbf{a}$ in D , where $\mathbf{A} \in \mathbb{R}^{3 \times 3}$ is a skew-symmetric matrix and $\mathbf{a} \in \mathbb{R}^3$. From the unique continuation principle applied to $\mathbf{w}_2 - \mathbf{A}\mathbf{x} - \mathbf{a}$ we obtain that $\mathbf{w}_2 = \mathbf{A}\mathbf{x} + \mathbf{a}$ in $\mathbb{R}^3 \setminus \overline{C_2}$, hence $(\mathbb{C}\widehat{\nabla}\mathbf{w}_2)\mathbf{n} = 0$ on $B'_{s_0}(\mathbf{0})$, that is a contradiction.

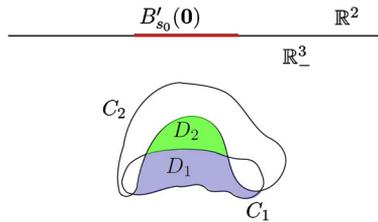


Fig. 2. The two domains C_1 and C_2 and the domain $D = D_1 \cup D_2$.

The second case we analyse is related to the setting in Fig. 2. To prove the contradiction, in this case we can consider, for example, the domain $D = D_1 \cup D_2$, where D_1 and D_2 are represented in Fig. 2. We emphasize that in this setting if we only consider the domain D_1 we will not be able to find a contradiction. In fact, from the unique continuation and the Green’s formula we would not be able to prove that the energy of the system related to the function w_2 is equal to zero in D_1 . Instead, considering the region D and taking w_2 , which satisfies $\text{div}(\mathbb{C}\widehat{\nabla}w_2) = \mathbf{0}$ in D , from the Green’s formula we find, as before,

$$\int_D (\mathbb{C}\widehat{\nabla}w_2) : \widehat{\nabla}w_2 \, dx = 0,$$

hence $w_2 = \mathbf{A}x + \mathbf{a}$ in D , with $\mathbf{A} \in \mathbb{R}^{3 \times 3}$ a skew-symmetric matrix and $\mathbf{a} \in \mathbb{R}^3$. Applying the unique continuation principle to $w_2 - \mathbf{A}x - \mathbf{a}$ we have that $w_2 = \mathbf{A}x + \mathbf{a}$ in $\mathbb{R}^3 \setminus \overline{C_2}$, hence $(\mathbb{C}\widehat{\nabla}w_2)\mathbf{n} = 0$ on $B'_{s_0}(\mathbf{0})$, that is a contradiction. \square

4.2. Stability estimate

In this section we state and prove the stability estimates for our inverse problem by adapting the arguments contained in [15] and [16] to the case of a pressurized cavity in an unbounded domain. In order to keep the proof of the main result as readable as possible and since the strategy to get the stability theorem is similar of the one obtained in [15] we will not repeat all the details of the proofs of the auxiliary results we need. The main idea behind the stability result (Theorem 4.8) is to give a quantitative version of the uniqueness argument. More precisely, we combine two steps:

1. the propagation of the smallness of the Cauchy data up to the boundary of the cavities, leading to an integral estimate of the solutions (see Propositions 4.6 and 4.7);
2. an estimate of continuation from the interior (Proposition 4.4).

The basic tool for both steps is the three spheres inequality stated in Lemma 4.5.

In the sequel, for any $\varrho > 0$, we denote by

$$\Omega_\varrho = \{x \in \Omega : \text{dist}(x, \partial\Omega) > \varrho\}.$$

We first prove a regularity result on the solution of problem (4). To this end, we consider a bounded domain $Q \subset \mathbb{R}^3_-$ such that

$$\partial Q \in C^3 \text{ with constants } r_0, E_0, \tag{41}$$

$$\overline{B_{\alpha D_0}^-(\mathbf{0})} \subset\subset \overline{Q} \subset\subset \overline{B_{\beta D_0}^-(\mathbf{0})}, \tag{42}$$

where $\alpha > 2$ and $\beta \geq 3$, with $\alpha < \beta$. Now, we have the following regularity estimate.

Proposition 4.3. *Under the assumptions (35) for C and (41), (42) for Q , the solution of problem (4), satisfies*

$$\|\mathbf{u}\|_{C^{1,1/2}(\overline{Q} \setminus C)} \leq c p, \tag{43}$$

where the constant $c = c(\lambda, \mu, \alpha, \beta, r_0, E_0, D_0)$.

Before proving this theorem, we briefly recall the integral representation formula for the solution to problem (4) derived in [7]. In particular, we define first the Neumann function, solution to

$$\begin{cases} \operatorname{div}(\mathbb{C}\widehat{\nabla}\mathbf{N}(\cdot, \mathbf{y})) = \delta_{\mathbf{y}}\mathbf{I} & \text{in } \mathbb{R}_-^3 \\ (\mathbb{C}\widehat{\nabla}\mathbf{N}(\cdot, \mathbf{y}))\mathbf{n} = 0 & \text{on } \mathbb{R}^2 \\ \mathbf{N} = O(|\mathbf{x}|^{-1}), \quad |\nabla\mathbf{N}| = O(|\mathbf{x}|^{-2}) & |\mathbf{x}| \rightarrow \infty, \end{cases} \tag{44}$$

where $\delta_{\mathbf{y}}$ is the delta function centered at $\mathbf{y} \in \mathbb{R}_-^3$ and \mathbf{I} is the identity matrix (see Theorem 4.9 in Appendix for its explicit expression). Then, for any $\mathbf{y} \in \overline{\mathbb{R}_-^3} \setminus \overline{C}$ we have

$$\mathbf{u}(\mathbf{y}) = p \int_{\partial C} \mathbf{N}(\mathbf{x}, \mathbf{y})\mathbf{n}(\mathbf{x}) d\sigma(\mathbf{x}) - \int_{\partial C} [(\mathbb{C}\widehat{\nabla}_{\mathbf{x}}\mathbf{N}(\mathbf{x}, \mathbf{y}))\mathbf{n}(\mathbf{x})]^T \mathbf{f}(\mathbf{x}) d\sigma(\mathbf{x}), \tag{45}$$

where \mathbf{f} is the trace of \mathbf{u} on ∂C and \mathbf{n} is the outer unit normal vector on ∂C (for more details see [7]).

Proof of Theorem 4.3. Since the kernels of the integral operators in (45) are regular for $\mathbf{y} \in \partial Q \setminus \partial C$, we can estimate $D^k\mathbf{u}(\mathbf{y})$, for $k = 0, 1, 2, 3$, getting easily

$$|D^k\mathbf{u}(\mathbf{y})| \leq p|\partial C| \sup_{\substack{\mathbf{x} \in \partial C \\ \mathbf{y} \in \partial Q \setminus \partial C}} |D_{\mathbf{y}}^k\mathbf{N}(\mathbf{x}, \mathbf{y})| + |\partial C|^{1/2} \|\mathbf{f}\|_{L^2(\partial C)} \sup_{\substack{\mathbf{x} \in \partial C \\ \mathbf{y} \in \partial Q \setminus \partial C}} |D_{\mathbf{y}}^k(\mathbb{C}\widehat{\nabla}_{\mathbf{x}}\mathbf{N}(\mathbf{x}, \mathbf{y}))|.$$

From the regularity properties of \mathbf{N} and Theorem 4.9 in Appendix, we have

$$\sup_{\substack{\mathbf{x} \in \partial C \\ \mathbf{y} \in \partial Q \setminus \partial C}} |D_{\mathbf{y}}^k\mathbf{N}(\mathbf{x}, \mathbf{y})| \leq \frac{c}{D_0^{k+1}}, \quad \sup_{\substack{\mathbf{x} \in \partial C \\ \mathbf{y} \in \partial Q \setminus \partial C}} |D_{\mathbf{y}}^k(\mathbb{C}\widehat{\nabla}_{\mathbf{x}}\mathbf{N}(\mathbf{x}, \mathbf{y}))| \leq \frac{c}{D_0^{k+2}}, \tag{46}$$

where the constant $c = c(\lambda, \mu, \alpha)$. From the trace estimate applied in Q , we have

$$\|\mathbf{f}\|_{L^2(\partial C)} \leq c\|\mathbf{u}\|_{H_w^1(\mathbb{R}_-^3 \setminus \overline{C})} \leq c p, \tag{47}$$

hence, from (46) and (47), we get

$$|D^k \mathbf{u}(\mathbf{y})| \leq cp,$$

where the constant $c = c(\lambda, \mu, \alpha, r_0, E_0, D_0)$. Therefore

$$\|D^k \mathbf{u}\|_{L^\infty(\partial Q \setminus \partial C)} \leq cp, \quad k = 0, 1, 2, 3. \tag{48}$$

Next, we apply the following global regularity estimate for the elastostatic system with Neumann boundary conditions (see [18], Theorem 6.6, p. 79) for \mathbf{u} in $Q \setminus \bar{C}$, that is

$$\|\mathbf{u}\|_{H^3(Q \setminus \bar{C})} \leq c \left(\|\mathbf{u}\|_{L^2(Q \setminus \bar{C})} + \|(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n}\|_{W^{3/2,2}(\partial(Q \setminus \bar{C}))} \right), \tag{49}$$

with $c = c(\lambda, \mu, \alpha, \beta, r_0, E_0, D_0)$. Now, observe that by (34)

$$\|\mathbf{u}\|_{L^2(Q \setminus \bar{C})} \leq c \left\| \frac{\mathbf{u}}{\rho} \right\|_{L^2(\mathbb{R}^3 \setminus \bar{C})} \leq cp, \tag{50}$$

where $c = c(\lambda, \mu, \alpha, \beta, r_0, E_0, D_0)$, while, for the term $\|(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n}\|_{W^{3/2,2}(\partial(Q \setminus \bar{C}))}$ we have

$$(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n} = \mathbf{0}, \quad \text{on } \partial Q \cap \{x_3 = 0\}, \quad (\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n} = p\mathbf{n}, \quad \text{on } \partial C,$$

and since ∂C is of class C^3 (see (35)) it follows $(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n} \in C^2(\partial C)$, hence

$$\|(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n}\|_{W^{3/2,2}(\partial C)} \leq cp. \tag{51}$$

Analogously, from the estimate (48) and the regularity of the boundary of Q , we find

$$\|(\mathbb{C}\widehat{\nabla}\mathbf{u})\mathbf{n}\|_{W^{3/2,2}(\partial Q \setminus \partial C)} \leq cp. \tag{52}$$

Therefore, collecting (50), (51) and (52), the estimate (49) gives

$$\|\mathbf{u}\|_{H^3(Q \setminus \bar{C})} \leq cp.$$

Finally, applying the Sobolev embedding theorem (43) follows. \square

Proposition 4.4 (*Lipschitz propagation of smallness*). *Under the assumptions (5), (8), (9), (10) and (35), let \mathbf{w} be the solution to (38). There exist $R \geq 3D_0$, $R = R(\lambda, \mu, r_0, E_0, D_0)$, and $s > 1$, $s = s(\lambda, \mu, E_0)$, such that for every $\varrho > 0$ and every $\bar{\mathbf{x}} \in (B_R^-(\mathbf{0}) \setminus \bar{C})_{s\varrho}$, we have*

$$\int_{B_\varrho(\bar{\mathbf{x}})} |\widehat{\nabla}\mathbf{w}|^2 dx \geq \frac{c}{e^{a\varrho^{-b}}} \int_{B_R^-(\mathbf{0}) \setminus \bar{C}} |\widehat{\nabla}\mathbf{w}|^2 dx, \tag{53}$$

where $a, b, c > 0$ depend on λ, μ, r_0, E_0 and D_0 .

The proof of this proposition is based on the application of the three-spheres inequality. For the reader’s convenience we recall here only the statement of the theorem in the case of the linear elasticity in a homogeneous and isotropic medium; for a more general case and its proof one can refer to [2,3].

Lemma 4.5 (*Three spheres inequality*). *Let Ω be a bounded domain in \mathbb{R}^3 . Let $\mathbf{w} \in H^1(\Omega)$ be a solution to the Lamé system. There exists ϑ^* , $0 < \vartheta^* \leq 1$, only depending on λ and μ such that for every r_1, r_2, r_3, \bar{r} , $0 < r_1 < r_2 < r_3 \leq \vartheta^* \bar{r}$, and for every $\mathbf{x} \in \Omega_{\bar{r}}$ we have*

$$\int_{B_{r_2}(\mathbf{x})} |\widehat{\nabla} \mathbf{w}|^2 \leq c \left(\int_{B_{r_1}(\mathbf{x})} |\widehat{\nabla} \mathbf{w}|^2 \right)^\delta \left(\int_{B_{r_3}(\mathbf{x})} |\widehat{\nabla} \mathbf{w}|^2 \right)^{1-\delta}, \tag{54}$$

where $c > 0$ and δ , $0 < \delta < 1$, only depend on λ , μ , $\frac{r_2}{r_3}$ and are monotone increasing functions of $\frac{r_1}{r_3}$.

Now, the Lipschitz propagation of smallness inequality can be proved.

Proof of Proposition 4.4. Let us denote $\Omega = B_R^-(\mathbf{0}) \setminus \bar{C}$, with $R \geq 3D_0$ to be chosen later. Since the hemisphere $B_1^-(\mathbf{0})$ has Lipschitz boundary with absolute constants r^* , E^* , $B_R^-(\mathbf{0})$ has Lipschitz boundary with constants r^*R , E^* . Possibly worsening the regularity parameters of C , see (35), we can assume $E_0 \geq E^*$ and $r_0 \leq r^*R$, so that the boundary of Ω is of Lipschitz class with constants r_0 and E_0 .

Following similar arguments as those in [15] with the simplification of maintaining as integrand function $|\widehat{\nabla} \mathbf{w}|^2$, we find that there exist $\varrho_0 = \varrho_0(\lambda, \mu, r_0, E_0, R)$, with $0 < \varrho_0 < 1$, and $s = s(\lambda, \mu, E_0)$, $s > 1$, such that for all $0 < \varrho \leq \varrho_0$ and for all $\bar{\mathbf{x}} \in \Omega_{s\varrho}$, it holds

$$\int_{B_\varrho(\bar{\mathbf{x}})} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x} \geq c_1 \int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x} \left(\frac{c\varrho^3 \int_{\Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x}}{\int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x}} \right)^{\sigma - A_1 - B_1 \log(1/\varrho)}, \tag{55}$$

where $c_1 > 0$ only depend on λ, μ ; $\sigma \in (0, 1)$, $s > 1$ depend on λ, μ, E_0 and $c, A_1, B_1 > 0$ depend on $\lambda, \mu, E_0, r_0, R$. The main goal is to give a lower estimate of the ratio

$$\frac{\int_{\Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x}}{\int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 d\mathbf{x}}$$

to get the assertion of the theorem. To this end, we first notice that

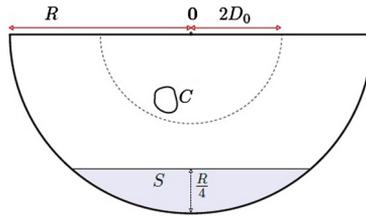


Fig. 3. The region S .

$$\frac{\int_{\Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 dx}{\int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 dx} = 1 - \frac{\int_{\Omega \setminus \Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 dx}{\int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 dx} := 1 - \frac{\mathcal{I}_1}{\mathcal{I}_2}. \tag{56}$$

Let us consider the integral \mathcal{I}_2 . Since $\mathbf{w} = \mathbf{u} - \bar{\mathbf{u}}$, see (37), we use the integral representation formula (45) for the function \mathbf{u} and the explicit expression of $\bar{\mathbf{u}}$ in (36). In detail, we consider

$$S = \left\{ \mathbf{x} \in B_R^-(\mathbf{0}) : x_3 \leq -\frac{3}{4}R \right\}. \tag{57}$$

see Fig. 3. By a simple calculation we have $|S| = (7/128)\pi R^3$.

If $\mathbf{x} \in \partial C$, $\mathbf{y} \in S$, then by (45) and Theorem 4.9, it is easy to see that

$$|\widehat{\nabla} \mathbf{u}(\mathbf{y})| \leq |\nabla \mathbf{u}(\mathbf{y})| \leq \frac{c p}{R^2}, \quad \forall \mathbf{y} \in S,$$

where $c = c(\lambda, \mu, r_0, E_0, D_0)$, hence

$$|\widehat{\nabla} \mathbf{w}(\mathbf{y})| \geq |\widehat{\nabla} \bar{\mathbf{u}}(\mathbf{y})| - |\widehat{\nabla} \mathbf{u}(\mathbf{y})| \geq \frac{p}{3\lambda + 2\mu} - \frac{c p}{R^2} \geq \frac{p}{2(3\lambda + 2\mu)},$$

where the last inequality holds choosing $R = \max\{3D_0, (2c^{-1}(3\lambda + 2\mu))^{1/2}\}$. In this way, we have that

$$\mathcal{I}_2 = \int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 dx \geq \int_S |\widehat{\nabla} \mathbf{w}|^2 dx \geq c p^2 R^3, \tag{58}$$

where $c = c(\lambda, \mu)$.

Now, we estimate the integral \mathcal{I}_1 using the regularity result of the Proposition 4.3. First, we split the integral domain as

$$\Omega \setminus \Omega_{(s+1)p} = F_1 \cup F_2,$$

where

$$F_1 = \{\mathbf{x} \in \Omega : d(\mathbf{x}, C) \leq (s + 1)\varrho\}, \quad F_2 = \{\mathbf{x} \in \Omega : d(\mathbf{x}, \partial B_R^-(\mathbf{0})) \leq (s + 1)\varrho\},$$

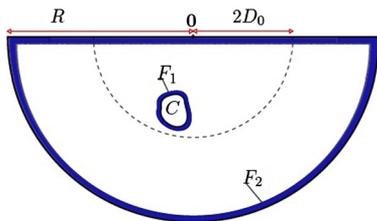


Fig. 4. The region $\Omega \setminus \Omega_{(s+1)\varrho}$.

see Fig. 4.

From (35), we notice that

$$|F_1 \cup F_2| \leq c(r_0, E_0, D_0)\varrho R^2. \tag{59}$$

Choosing in (42) $\alpha = R/D_0$, where $R = \max\{3D_0, (2c^{-1}(3\lambda + 2\mu))^{1/2}\}$, and $\beta = 2\alpha$, then

$$\overline{B_R^-(\mathbf{0})} \subset\subset \overline{Q} \subset\subset \overline{B_{2R}^-(\mathbf{0})},$$

hence we apply the regularity estimate (43) for the two regions F_1 and F_2 . Now, we have that

$$\|\widehat{\nabla} \mathbf{w}\|_{L^\infty(\Omega \setminus \Omega_{(s+1)\varrho})} \leq (\|\widehat{\nabla} \mathbf{u}\|_{L^\infty(\Omega \setminus \Omega_{(s+1)\varrho})} + \|\widehat{\nabla} \overline{\mathbf{u}}\|_{L^\infty(\Omega \setminus \Omega_{(s+1)\varrho})}) \leq cp,$$

where $c = c(\lambda, \mu, r_0, E_0, D_0)$. Therefore, from (59), we find

$$\mathcal{I}_1 = \int_{\Omega \setminus \Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 dx \leq cp^2 |\Omega \setminus \Omega_{(s+1)\varrho}| \leq cp^2 \varrho R^2. \tag{60}$$

Putting together inequalities (58) and (60), there exists $\varrho^* = \varrho^*(\lambda, \mu, r_0, E_0, D_0) > 0$, such that for any $\varrho \leq \varrho^*$, we have

$$\frac{\mathcal{I}_1}{\mathcal{I}_2} = \frac{\int_{\Omega \setminus \Omega_{(s+1)\varrho}} |\widehat{\nabla} \mathbf{w}|^2 dx}{\int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 dx} \leq \frac{1}{2}.$$

Going back to (55) and (56) we have

$$\int_{B_\varrho(\overline{\mathbf{x}})} |\widehat{\nabla} \mathbf{w}|^2 dx \geq (c\varrho^3)^{\sigma^{-A_1 - B_1 \log(1/\varrho)}} \int_{\Omega} |\widehat{\nabla} \mathbf{w}|^2 dx,$$

where c, A_1, B_1 depend on $\lambda, \mu, E_0, r_0, D_0$, for all $\varrho \leq \varrho^*$. To conclude, we take $\varrho \leq c$, hence, for every $\varrho \leq \min(c, \varrho^*)$ and noticing that $\log \varrho \geq -1/\varrho$ for $0 < \varrho < 1$, we get the assertion choosing

$$a = 6e^{A_1|\log \sigma|} \quad \text{and} \quad b = B_1|\log \sigma| + 1. \quad \square$$

We omit the proof of the following two propositions, since they can be obtained using the same strategy adopted in proving Propositions 3.5 and 3.6 in [15].

Proposition 4.6 (Stability estimates of continuation from Cauchy data). *Under the assumption (5) let C_1 and C_2 be two domains satisfying (8), (9), (10) and (35). Moreover, let w_i , for $i = 1, 2$, be the solution to (38) with $C = C_i$. Then, for $\varepsilon < e^{-1}p$, we have*

$$\begin{aligned} \int_{C_2 \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx &\leq cp^2 \left(\log \left| \log \frac{\varepsilon}{p} \right| \right)^{-1/6}, \\ \int_{C_1 \setminus \bar{C}_2} |\widehat{\nabla} w_2|^2 dx &\leq cp^2 \left(\log \left| \log \frac{\varepsilon}{p} \right| \right)^{-1/6}, \end{aligned} \tag{61}$$

where the constant $c = c(\lambda, \mu, r_0, E_0, D_0, s_0)$.

The stability estimates in (61) can be improved when ∂G is of Lipschitz class, where G is defined by (40), as stated in the proposition below.

Proposition 4.7 (Improved stability estimates of continuation from Cauchy data). *Under the assumption (5) let C_1 and C_2 be two domains satisfying (8), (9), (10) and (35). In addition, let us assume that there exist $L > 0$ and \tilde{r}_0 , with $0 < \tilde{r}_0 \leq r_0$, such that ∂G is of Lipschitz class with constants \tilde{r}_0, L . Then, we have*

$$\begin{aligned} \int_{C_2 \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx &\leq cp^2 \left| \log \frac{\varepsilon}{p} \right|^{-\gamma}, \\ \int_{C_1 \setminus \bar{C}_2} |\widehat{\nabla} w_2|^2 dx &\leq cp^2 \left| \log \frac{\varepsilon}{p} \right|^{-\gamma}, \end{aligned} \tag{62}$$

where $c, \gamma > 0$ depend on $\lambda, \mu, r_0, E_0, D_0, s_0, L, \tilde{r}_0$.

Now, we have all the preliminary results to prove the stability theorem.

Theorem 4.8 (Stability estimate). *Under the assumption (5) let C_1 and C_2 be two domains satisfying (8), (9), (10) and (35). Moreover, let u_i , for $i = 1, 2$, be the solution to (4) with $C = C_i$. If, given $\varepsilon > 0$, we have*

$$\|u_1 - u_2\|_{L^2(B'_{s_0}(\mathbf{0}))} \leq \varepsilon, \tag{63}$$

then it holds

$$d_{\mathcal{H}}(\partial C_1, \partial C_2) \leq c \left(\log \left| \log \frac{\varepsilon}{p} \right| \right)^{-\eta}, \tag{64}$$

for every $\varepsilon < e^{-1} p$, where the constants c and η , with $0 < \eta \leq 1$, depend on $\lambda, \mu, r_0, E_0, D_0$ and s_0 .

Proof. Since (39) holds, we prove the assertion using the function w_i , for $i = 1, 2$. In this way we can apply the same proof strategy contained in [16]. In the sequel we simply denote with $d_{\mathcal{H}}$ the Hausdorff distance $d_{\mathcal{H}}(\partial C_1, \partial C_2)$.

Let us prove that if $\eta > 0$ is such that

$$\int_{C_2 \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx \leq \eta, \quad \int_{C_1 \setminus \bar{C}_2} |\widehat{\nabla} w_2|^2 dx \leq \eta, \tag{65}$$

then we have

$$d_{\mathcal{H}} \leq c \left(\log \frac{cp^2}{\eta} \right)^{-1/b}, \tag{66}$$

where b, c depend on $\lambda, \mu, r_0, E_0, D_0$.

We may assume, with no loss of generality, that there exists $x_0 \in \partial C_1$ such that

$$\text{dist}(x_0, \partial C_2) = d_{\mathcal{H}}.$$

In this setting, we have to distinguish two cases:

1. $B_{d_{\mathcal{H}}}(x_0) \subset C_2$;
2. $B_{d_{\mathcal{H}}}(x_0) \cap C_2 = \emptyset$.

Let us consider case (1). By the regularity assumption made on ∂C_1 , see (35), there exists $x_1 \in C_2 \setminus \bar{C}_1$ such that

$$B_{td_{\mathcal{H}}}(x_1) \subset (C_2 \setminus \bar{C}_1), \text{ with } t = \frac{1}{1 + \sqrt{1 + E_0^2}}.$$

By the first inequality in (65), taking $\varrho = td_{\mathcal{H}}/s$ in Proposition 4.4, we have

$$\eta \geq \int_{C_2 \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx \geq \int_{B_{\varrho}(x_1)} |\widehat{\nabla} w_1|^2 dx \geq \frac{c}{e^{a\varrho^{-b}}} \int_{B_R^- \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx, \tag{67}$$

where we recall that $R = R(\lambda, \mu, r_0, E_0, D_0)$. By (58), we find that

$$\int_{B_R^- \setminus \bar{C}_1} |\widehat{\nabla} w_1|^2 dx \geq cp^2,$$

so that, going back to (67), we have

$$\eta \geq \frac{cp^2}{e^{ae^{-b}}} = \frac{cp^2}{e^{a(td_{\mathcal{H}}/s)^{-b}}}. \tag{68}$$

From this inequality it is straightforward to find (66).

Case (2) can be proved in a similar way by substituting w_1 with w_2 in the previous calculations and employing the second inequality in (65).

Now, applying (61), that is taking

$$\eta = cp^2 \left(\log \left| \log \frac{\varepsilon}{p} \right| \right)^{-1/6},$$

we obtain from (66) that

$$d_{\mathcal{H}} \leq c \left(\log \log \left| \log \frac{\varepsilon}{p} \right| \right)^{-1/b}, \tag{69}$$

where we require $\varepsilon < e^{-e} p$ to have a positive quantity in right side of the previous inequality; the positive constants b, c depend on $\lambda, \mu, r_0, E_0, s_0$ and D_0 .

Next, to improve the modulus of continuity of this estimate we recall a geometrical result, first introduced and proved in [1], ensuring that there exists $d_0 > 0, d_0 = d_0(r_0, E_0)$ such that if $d_{\mathcal{H}}(\partial C_1, \partial C_2) \leq d_0$, then the boundary of G is of Lipschitz class with constants \tilde{r}_0, L , only depending on r_0 and E_0 . By (69), there exists $\varepsilon_0 > 0$ only depending on $\lambda, \mu, r_0, E_0, s_0$ and D_0 such that if $\varepsilon \leq \varepsilon_0$ then $d_{\mathcal{H}} \leq d_0$. In this way G satisfies the hypotheses of Proposition 4.7 hence the assertion follows. \square

Acknowledgments

The authors thank Prof. Cherif Amrouche for his kindness in suggesting and providing some useful papers and for his enlightening advice. Andrea Aspri and Elena Beretta thank the New York University in Abu Dhabi (EAU) for its kind hospitality that permitted a further development of the present research. Andrea Aspri thanks ÖAW (Austrian Academy of Sciences) and RICAM for giving him the possibility to finish this paper. Edi Rosset is supported by FRA2016 “Problemi inversi, dalla stabilità alla ricostruzione”, Università degli Studi di Trieste and by Progetto GNAMPA 2017 “Analisi di problemi inversi: stabilità e ricostruzione”, Istituto Nazionale di Alta Matematica (INdAM).

Appendix A. Neumann function for the Lamé operator in the half-space

This appendix is devoted to the explicit expression of the Neumann function for the Lamé operator in the half-space presented in [13,14]. Before doing that we recall the fundamental solution Γ of the Lamé operator, that is the so called Kelvin–Somigliana matrix. Γ is the solution to the equation

$$\operatorname{div}(\mathbb{C}\widehat{\nabla}\Gamma) = \delta_0\mathbf{I}, \quad \mathbf{x} \in \mathbb{R}^3 \setminus \{\mathbf{0}\},$$

where δ_0 is the Dirac function centered at $\mathbf{0}$ and \mathbf{I} is the identity matrix. Setting $C_{\mu,\nu} := 1/\{16\pi\mu(1-\nu)\}$, where ν is the Poisson ratio $\nu = \lambda/(2(\lambda + \mu))$, the explicit expression of $\Gamma = (\Gamma_{ij})$ is

$$\Gamma_{ij}(\mathbf{x}) = -C_{\mu,\nu} \left\{ \frac{(3-4\nu)\delta_{ij}}{|\mathbf{x}|} + \frac{x_i x_j}{|\mathbf{x}|^3} \right\}, \quad i, j = 1, 2, 3, \quad (\text{A.1})$$

where δ_{ij} is the Kronecker symbol.

Given $\mathbf{y} = (y_1, y_2, y_3)$, we set $\tilde{\mathbf{y}} = (y_1, y_2, -y_3)$. Now, we have

Theorem 4.9 ([7]). *The Neumann function \mathbf{N} of problem (44) can be decomposed as*

$$\mathbf{N}(\mathbf{x}, \mathbf{y}) = \mathbf{\Gamma}(\mathbf{x} - \mathbf{y}) + \mathbf{R}^1(\mathbf{x} - \tilde{\mathbf{y}}) + y_3 \mathbf{R}^2(\mathbf{x} - \tilde{\mathbf{y}}) + y_3^2 \mathbf{R}^3(\mathbf{x} - \tilde{\mathbf{y}}),$$

where $\mathbf{\Gamma}$ is the Kelvin matrix, see (A.1), and \mathbf{R}^k , $k = 1, 2, 3$, have components R_{ij}^k given by

$$\begin{aligned} R_{ij}^1(\boldsymbol{\eta}) &:= C_{\mu,\nu} \left\{ -(\tilde{f} + c_\nu \tilde{g})\delta_{ij} - (3-4\nu)\eta_i \eta_j \tilde{f}^3 \right. \\ &\quad \left. + c_\nu [\delta_{i3}\eta_j - \delta_{j3}(1-\delta_{i3})\eta_i] \tilde{f} \tilde{g} + c_\nu (1-\delta_{i3})(1-\delta_{j3})\eta_i \eta_j \tilde{f} \tilde{g}^2 \right\} \\ R_{ij}^2(\boldsymbol{\eta}) &:= 2C_{\mu,\nu} \left\{ (3-4\nu) [\delta_{i3}(1-\delta_{j3})\eta_j + \delta_{j3}(1-\delta_{i3})\eta_i] \tilde{f}^3 - (1-2\delta_{3j})\delta_{ij}\eta_3 \tilde{f}^3 \right. \\ &\quad \left. + 3(1-2\delta_{3j})\eta_i \eta_j \eta_3 \tilde{f}^5 \right\} \\ R_{ij}^3(\boldsymbol{\eta}) &:= 2C_{\mu,\nu} (1-2\delta_{j3}) \left\{ \delta_{ij} \tilde{f}^3 - 3\eta_i \eta_j \tilde{f}^5 \right\}. \end{aligned}$$

for $i, j = 1, 2, 3$, where $c_\nu := 4(1-\nu)(1-2\nu)$ and

$$\tilde{f}(\boldsymbol{\eta}) := \frac{1}{|\boldsymbol{\eta}|}, \quad \tilde{g}(\boldsymbol{\eta}) := \frac{1}{|\boldsymbol{\eta}| - \eta_3}.$$

References

- [1] G. Alessandrini, E. Beretta, E. Rosset, S. Vessella, Optimal stability for inverse elliptic boundary value problems with unknown boundaries, *Ann. Sc. Norm. Super. Pisa Cl. Sci.* 4 (2000) 755–806.
- [2] G. Alessandrini, A. Morassi, Strong unique continuation for Lamé system of elasticity, *Comm. Partial Differential Equations* 26 (2001) 1787–1810.
- [3] G. Alessandrini, A. Morassi, E. Rosset, Detecting an inclusion in an elastic body by boundary measurements, *SIAM J. Math. Anal.* 33 (2002) 1247–1268.
- [4] G. Alessandrini, A. Morassi, E. Rosset, The linear constraints in Poincaré and Korn type inequalities, *Forum Math.* 20 (2008) 557–569.
- [5] C. Amrouche, F. Bonzom, Exterior problems in the half-space for the Laplace operator in weighted Sobolev spaces, *J. Differential Equations* 246 (2009) 1894–1920.
- [6] C. Amrouche, M. Dambrine, Y. Raudin, An L^p theory of linear elasticity in the half-space, *J. Differential Equations* 253 (2012) 906–932.
- [7] A. Aspri, E. Beretta, C. Mascia, Analysis of a Mogi-type model describing surface deformations induced by a magma chamber embedded in an elastic half-space, *J. Éc. Polytech., Math.* 4 (2017) 223–255.
- [8] M. Battaglia, D.P. Hill, Analytical modeling of gravity changes and crustal deformation at volcanoes: the Long Valley caldera, California, case study, *Tectonophysics* 471 (2009) 45–57.
- [9] M.E. Gurtin, *The Linear Theory of Elasticity*, *Encyclopedia of Physics*, vol. VI a/2, Springer-Verlag, 1972.
- [10] B. Hanouzet, *Espaces de Sobolev avec poids. Application au problème de Dirichlet dans un demi-espace*, *Rend. Semin. Mat. Univ. Padova* 46 (1971) 227–272.
- [11] V.A. Kondrat'ev, O.A. Oleinik, Boundary-value problems for the system of elasticity theory in unbounded domains. Korn's inequalities, *Russian Math. Surveys* 43 (1988) 65–119.
- [12] M. Lisowski, Analytical volcano deformation source models, in: D. Dzurisin (Ed.), *Volcano Deformation. Geodetic Monitoring Techniques*, Springer Praxis Books, 2006, pp. 279–304.
- [13] R.D. Mindlin, Force at a point in the interior of a semiinfinite solid, *J. Appl. Phys.* 7 (1936) 195–202.

- [14] R.D. Mindlin, Force at a point in the interior of a semi-infinite solid, in: Proceedings of The First Midwestern Conference on Solid Mechanics, University of Illinois, Urbana, IL, April, 1954.
- [15] A. Morassi, E. Rosset, Stable determination of cavities in elastic bodies, *Inverse Probl.* 20 (2004) 453–480.
- [16] A. Morassi, E. Rosset, Uniqueness and stability in determining a rigid inclusion in an elastic body, *Mem. Amer. Math. Soc.* 938 (2009) 453–480.
- [17] P. Segall, *Earthquake and Volcano Deformation*, Princeton University Press, 2010.
- [18] T. Valent, *Boundary Value Problems of Finite Elasticity, Local Theorems on Existence, Uniqueness and Analytic Dependence on Data*, Springer-Verlag, New York–Berlin–Heidelberg, 1988.
- [19] N. Weck, Außenraumaufgaben in der Theorie stationärer Schwingungen inhomogener elastischer Körper, *Math. Z.* 111 (1969) 387–398.