

# 1 Consistent determination of the heating rate of light-absorbing 2 aerosol using wavelength- and time-dependent Aethalometer 3 multiple-scattering correction

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22 **Abstract.** Accurate and temporally consistent measurements of light absorbing aerosol (LAA) heating  
23 rate (HR) and of its source apportionment (fossil-fuel, FF; biomass-burning, BB) and speciation (black  
24 and brown Carbon; BC, BrC) are needed to evaluate LAA short-term climate forcing. For this purpose,

wavelength- and time-dependent accurate LAA absorption coefficients are required. HR was experimentally determined and apportioned (sources/species) in the EMEP/ACTRIS/COLOSSAL-2018 winter campaign in Milan (urban-background site). Two Aethalometers (AE31/AE33) were installed together with a MAAP, CPC, OPC, a low volume sampler (PM<sub>2.5</sub>) and radiation instruments. AE31/AE33 multiple-scattering correction factors (*C*) were determined using two reference systems for the absorption coefficient: 1) 5-wavelength PP\_UniMI with low time resolution (12 h, applied to PM<sub>2.5</sub> samples); 2) timely-resolved MAAP data at a single wavelength. Using wavelength- and time-independent *C* values for the AE31 and AE33 obtained with the same reference device, the total HR showed a consistency (i.e. reproducibility) with average values comparable at 95% probability. However, if different reference devices/approaches are used, i.e. MAAP is chosen as reference instead of a PP\_UniMI, the HR can be overestimated by 23-30% factor (by both AE31/AE33). This became more evident focusing on HR apportionment: AE33 data (corrected by a wavelength- and time-independent *C*) showed higher HR<sub>FF</sub> (+24±1%) and higher HR<sub>BC</sub> (+10±1%) than that of AE31. Conversely, HR<sub>BB</sub> and HR<sub>BrC</sub> were -28±1% and -29±1% lower for AE33 compared to AE31. These inconsistencies were overcome by introducing a wavelength-dependent *C<sub>λ</sub>* for both AE31 and AE33, or using multi-wavelength apportionment methods, highlighting the need for further studies on the influence of wavelength corrections for HR determination. Finally, the temporally-resolved determination of *C* resulted in a diurnal cycle of the HR not statistically different whatever the source- speciation- apportionment used.

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**KEYWORDS:** Heating rate; Black Carbon; Brown Carbon; Fossil Fuel; Biomass Burning; COLOSSAL.

## 1. Introduction

Light absorbing aerosols (LAA: e.g. black carbon, BC; brown carbon, BrC; and dust) are important due to their optical and radiative forcing properties (Laskin et al., 2015; Petzold et al., 2013; Bond et al., 2013). Accordingly to IPCC (2013) and several authors (Bond et al., 2013; Ramanathan and Carmichael, 2008), the BC is the second most important positive anthropogenic climate-forcing agent after CO<sub>2</sub> (top-of-

atmosphere direct forcing of  $\sim 0.71 \text{ W m}^{-2}$ ; range  $0.08 - 1.27 \text{ W m}^{-2}$ ). In addition to BC, BrC may contribute  $\sim 10\text{-}30\%$  to the total solar absorption induced by LAA (Chung et al., 2012; Shamjad et al., 2015; Laskin et al., 2015; Ferrero et al., 2018). Finally, dust aerosol could promote an atmospheric forcing in the range  $-0.1 \pm 0.2 \text{ W m}^{-2}$  (Myhre et al., 2013), rising up regionally to  $\sim 80\text{-}100 \text{ W m}^{-2}$  (Huang et al., 2009; Mallet et al., 2009). Unlike  $\text{CO}_2$ , which is distributed quite homogeneously in the atmosphere, LAA are short-lived climate forcers ( $\sim 1$  to several weeks of residence time) (Samset et al., 2013; Cape et al., 2012); thus, their effect could be counteracted by short-term climate strategies (Ødemark et al., 2012; Shindell et al., 2012; Jacobson, 2010; Quinn et al., 2008). This would lead enough time to  $\text{CO}_2$  related sustainability strategies and renewable energy to become fruitful (Qerimi et al., 2020; Kibaara et al., 2020; Burciaga, 2020) to avoid the dramatic environmental effects of climate change (e.g. Javadinejad et al., 2020; Oo et al., 2020). However, to adopt the right mitigation strategies on LAA, the identification of the relative importance of different sources and species of LAA on the total HR is required with a clear need to reduce the uncertainties related to their determination (Ferrero et al., 2018; Myhre et al., 2013) as well as to properly quantify their atmospheric feedbacks (Su et al., 2020; Tian et al., 2019; Diemoz et al., 2019; Ferrero et al., 2011a, 2014, 2016, 2019a,b; Ran et al., 2016; Costabile et al., 2013; Lemaître et al., 2010). For example, LAA can affect: the atmospheric stability (Su et al., 2020), the pattern of synoptic winds such as the monsoon (Bond et al., 2013; Ramanathan and Feng, 2009; Koch et al., 2009; Ramanathan and Carmichael, 2008; Koren et al. 2008; Koren et al., 2004; Kaufman et al., 2002), the cloud distribution (Matus et al., 2015; Koch and Del Genio, 2010) thus modulating the amount and the spectrum of the short-wave radiation that reaches the ground (Ferrero et al., 2021; Črnivec and Mayer, 2019). Finally, the high UV-VIS absorption of BrC can influence the photolysis rate in the atmosphere affecting the concentration of other atmospheric pollutants, such as the ozone (Laskin et al., 2015 and references therein). Thus, in order to deal with the aforementioned feedbacks, a proper determination of the HR induced by LAA is required.

To achieve this goal, accurate measurements of the wavelength-dependent absorption coefficient ( $b_{abs}(\lambda)$ ) in the UV-VIS-IR region are required. In fact, when the HR is directly experimentally determined (Ferrero

et al., 2018, 2021), or when radiative transfer models are used to compute the HR (Tian et al., 2019; Gao et al., 2008),  $b_{abs}(\lambda)$  is integrated over the whole shortwave spectrum (in synergy with the radiation; see section 2.1); thus, even a small bias on any  $b_{abs}(\lambda)$  can be magnified by the integral operation.

For measuring the LAA  $b_{abs}(\lambda)$  in the whole UV-VIS-NIR region, Aethalometers (e.g. AE31 and AE33) are the most common and widely used devices. They feature a wide range of spectral channels (7- $\lambda$ : 370, 470, 520, 590, 660, 880 and 950 nm) not available in other instruments (e.g. MAAP, PSAP, photoacoustic spectrometer) (Petzold et al., 2005; Virkkula et al., 2010; Arnott et al., 1999). This spectral range is needed for the experimental HR determination (and radiative transfer models) as it includes the portion of the solar spectrum mainly responsible for incoming energy (Ferrero et al., 2018). Moreover, Aethalometers take also the advantage of global long-term data series (Ferrero et al., 2016; Eleftheriadis et al., 2009; Collaud-Coen et al., 2010; Junker et al., 2006) that could be used in the near future to retrieve historical series of the HR.

Despite the aforementioned advantages, the disadvantage (common to any filter-based absorption technique) is the need to compensate for three kinds of artefacts (Liousse et al., 1993; Petzold et al., 1997; Bond et al., 1999): 1) the multiple scattering (enhanced light attenuation and optical path by multiple scattering by the filter fibres); 2) the aerosol backscattering (light attenuation is further enhanced due to scattering of aerosols embedded in the filter); 3) the loading effect (the non-linear optical path reduction induced by absorbing particles accumulating in the filter). To account for the loading effects, Aethalometers data are traditionally corrected applying the procedures reported in Weingartner et al. (2003) and Collaud-Coen et al. (2010) for AE31 or in Drinovec et al. (2015) for AE33. All these procedures, require an accurate determination of the filter multiple scattering correction parameter  $C$  (Schmid et al., 2006; Arnott et al., 2005, Weingartner et al., 2003). Several estimates of  $C$  values are reported in the literature, ranging from 2.14 up to 5 (Weingartner et al., 2003; Arnott et al., 2005; Schmid et al., 2006; Sandradewi et al., 2008a, 2008b; Collaud Coen et al., 2010; GAW, 2016). As recently demonstrated,  $C$  values can vary with the Aethalometer type, Aethalometer filter material, time (diurnal/seasonal pattern) and wavelength (Bernardoni et al., 2021; Yus-Díez et al., 2021). Any arbitrary

102 assumption on the  $C$  values (avoiding to control the influence of the aforementioned parameters) will  
 103 influence the determination of the LAA optical properties and their apportionment in terms of source and  
 104 species (Sandradewi et al., 2008a; Massabò et al., 2015; Bernardoni et al., 2017; Tian et al., 2019); this  
 105 will reflect on the HR source-species quantification too. Thus, accurate measurements of the wavelength-  
 106 dependent absorption coefficient ( $b_{abs}(\lambda)$ ) in the UV-VIS-IR region are required.

107 Moreover, Andreae and Ramanathan (2013) report that radiative model results can be characterized by  
 108 large discrepancies among each other, amplifying the uncertainties related to implications for the global  
 109 and regional climate; among the whole ensemble of uncertainties, LAA optical properties play a major  
 110 role (Nordmann et al., 2014; Koch et al., 2009). As a conclusion, Andreae and Ramanathan (2013)  
 111 declared the need for better observational constraints at all scales, while Chung et al. (2012) and Samset  
 112 et al. (2018) called for model independent, observationally based estimate of the LAA properties.

113 Recently, Ferrero et al. (2018, 2021) developed a novel methodology to experimentally quantify the HR  
 114 induced by LAA. This method enables to provide experimental LAA-HR data at high-time resolution  
 115 with an accessible technique applicable in any established aerosol monitoring network (e.g. EMEP). The  
 116 method also allows a source-related (i.e. fossil-fuel, FF; biomass-burning, BB) or a species-related (e.g.  
 117 BC, BrC) apportionment of the HR.

118 Thus, the current paper applies the aforementioned experimental method for HR determination in function  
 119 of different wavelength and temporal parametrizations of the multiple scattering parameter  $C$  with the  
 120 aims to:

- 121 - quantify the  $C$  influence on the assessment of the experimental total, source (FF, BB), and species (BC,
- 122 BrC) HR;
- 123 - determine the best practice procedures for Aethalometer multiple-scattering correction for radiative
- 124 forcing assessment purposes.

125 A special attention is given both to the temporal consistency – important in any measurement time series  
 126 where AE31 Aethalometers are supplanted by the newest AE33 version – and to the accuracy of  
 127 source/species apportionment. The obtained results were determined within the framework of the winter

128 2018 European EMEP/ACTRIS/COLOSSAL (COST Action CA16109; Chemical On-Line cOmpoSition  
129 and Source Apportionment of fine aerosol) WG3 campaign.

130 The manuscript is organized as follows: in section 2 we will describe the experimental methodology to  
131 determine the LAA HR (section 2.1) and the experimental set-up used for determining both the HR and  
132 the multiple scattering parameter  $C$  (sections 2.2 and 2.3); results and discussion will follow in section 3,  
133 describing first the environmental context of atmospheric pollution during the campaign (section 3.1) and  
134 then (sections 3.2, 3.3 and 3.4) the HR values (total and from different species and sources of LAA) as a  
135 function of different  $C$  parametrizations. Conclusions follows with a best practice procedure for  
136 Aethalometer correction for radiative forcing assessment purposes.

137

## 138 **2. Methodology**

139 The LAA heating rate measurements were performed in the city of Milan from 18 January 2018 to 15  
140 February 2018. Milan is situated in the Po Valley which is a European air pollution hotspot, especially  
141 during wintertime, when stable atmospheric conditions occur, promoting the accumulation of aerosol  
142 within the mixing layer (Crova et al., 2021; Diémoz et al., 2019; Ferrero et al., 2011b, 2012; Vecchi et  
143 al., 2004, 2019) which is well visible even from satellites (Ferrero et al., 2019a; Di Nicolantonio et al.,  
144 2007, 2009; Barnaba and Gobbi 2004).

145 The sampling site is located on the rooftop (10 m a.g.l.) of the U9-bulding of the University of Milano-  
146 Bicocca (45°30'38" N, 127 9°12'42" E, Italy). The site is characterized by a full hemispherical sky view;  
147 moreover, its roof is characterized by a Lambertian concrete surface: due to its flat and homogeneous  
148 characteristics, it well represents the average spectral reflectance of the Milano urban area (Ferrero et al,  
149 2018). The U9 is equipped with instrumentation for the measurement of the absorption coefficient and  
150 radiation needed to determine the HR (section 2.2) but before introducing them, the HR determination  
151 method is detailed in the next section (2.1).

152

### 153 **2.1 Heating Rate determination**

154 The HR ( $\text{K day}^{-1}$ ) induced by LAA is determined using the experimental methodology reported in Ferrero  
 155 et al. (2018, 2020). Here we briefly summarize the method and the reader is referred to the aforementioned  
 156 publications for the physical demonstration of the approach.

157 The instantaneous aerosol total HR is determined by combining the LAA multi-wavelength absorption  
 158 coefficient ( $b_{abs}(\lambda)$ ) and radiation measurements (direct, diffuse or reflected) integrating them over the  
 159 whole ensemble of shortwave wavelengths and incident angles within the  $2\pi$  hemisphere from sky and  
 160 those reflected from the ground as follows:

$$161 \quad HR = \frac{1}{\rho C_p} \cdot \sum_{n=1}^3 \int_{\theta} \int_{\lambda} \frac{F_n(\lambda, \theta)}{\cos \theta} b_{abs}(\lambda) d\lambda d\theta \quad (1)$$

162 where  $\rho$  represents the air density ( $\text{kg m}^{-3}$ ),  $C_p$  ( $1005 \text{ J kg}^{-1} \text{ K}^{-1}$ ) is the isobaric specific heat of dry air,  $n$   
 163 is the index indicating the  $n^{\text{th}}$  type of radiation ( $n$  stands for direct, diffuse or reflected) impinging the  
 164 absorbing aerosol,  $\lambda$  and  $\theta$  represent the wavelength and zenith angle of the radiation,  $F_n(\lambda, \theta)$  is the  $n^{\text{th}}$  type  
 165 (direct or diffuse or reflected) of monochromatic irradiance of wavelength  $\lambda$  that strikes with an angle  $\theta$   
 166 the aerosol layer.

167 Under the isotropic (for diffuse radiation from sky and clouds) and Lambertian (for land) assumptions (as  
 168 used in Ferrero et al., 2018) equation 1 can be solved (see Ferrero et al., 2021) becoming:

$$169 \quad HR = \frac{1}{\rho C_p} \cdot \left[ \frac{1}{\cos(\theta_z)} \int_{\lambda} F_{dir}(\lambda) b_{abs}(\lambda) d\lambda + 2 \int_{\lambda} F_{dif}(\lambda) b_{abs}(\lambda) d\lambda + 2 \int_{\lambda} F_{ref}(\lambda) b_{abs}(\lambda) d\lambda \right] \quad (2)$$

170 where  $\theta_z$  refers to the solar zenith angle, while  $F_{dir}(\lambda)$ ,  $F_{dif}(\lambda)$  and  $F_{ref}(\lambda)$  are the spectral direct, diffuse  
 171 and reflected irradiances.

172 As already pointed out in Ferrero et al. (2018), it is worth recalling that in the present method (equations  
 173 1-2), the HR is independent from the thickness of the investigated atmospheric aerosol layer and thus  
 174 referred to the altitude at which it is determined. In the present work they were applied to the near-surface  
 175 atmospheric layer. BC and HR vertical profiles data previously collected at the same site revealed that  
 176 they were constant inside the mixing layer (Ferrero et al., 2014). Despite this limit, the HR measurement  
 177 has significant advantages: 1) no radiative transfer assumptions are needed (i.e. clear sky situation), 2)

178 measurements are carried out at high time resolution allowing to follow the HR dynamic, 3) measurements  
 179 are carried out in any sky conditions allowing to manage a continuous time series.  
 180 In addition to these advantages, the method allows a source-related (i.e. fossil-fuel, biomass-burning) or  
 181 a species-related (e.g. BC, BrC) apportionment of the HR. In fact, the spectral dependence of the aerosol  
 182 absorption coefficient ( $b_{abs}(\lambda)$ ) is generally described by the power-law relationship  $b_{abs} \propto \lambda^{-AAE}$  where  
 183 AAE is the Ångström absorption exponent (Moosmüller et al., 2009, 2011). The AAE is a parameter  
 184 useful for the apportionment of carbonaceous aerosol into different species (BC and BrC) and sources (FF  
 185 and BB) (Tian et al., 2019; Massabò et al., 2015, Sandradewi et al., 2008a). Thus, any  $b_{abs}(\lambda)$   
 186 apportionment method also allows to apportion the contribution to the HR of the different LAA species  
 187 ( $HR_{BC}$  and  $HR_{BrC}$ ) and sources ( $HR_{FF}$  and  $HR_{BB}$ ) by applying eq. 1-2. Source specific AAE values ( $AAE_{FF}$   
 188 and  $AAE_{BB}$ ) in Milan were set as 1 and 2 on the basis of previous investigations in the same area  
 189 (Bernardoni et al., 2017, 2021). BrC is characterized by  $AAE_{BrC}$  values which can vary in literature from  
 190 3 to 10 (Ferrero et al., 2018; Shamjad et al., 2015, Massabò et al., 2015; Srinivas et al., 2013; Yang et al.,  
 191 2009; Kirchstetter et al., 2004).

192 From a general point of view, any traditional apportionment of the species and sources is based on four  
 193 equations:

$$194 \quad \frac{b_{abs}(\lambda_1)_A}{b_{abs}(\lambda_2)_A} = \left(\frac{\lambda_1}{\lambda_2}\right)^{-AAE_A} \quad (3)$$

$$195 \quad \frac{b_{abs}(\lambda_1)_B}{b_{abs}(\lambda_2)_B} = \left(\frac{\lambda_1}{\lambda_2}\right)^{-AAE_B} \quad (4)$$

$$196 \quad b_{abs}(\lambda_{1,2}) = b_{abs}(\lambda_{1,2})_A + b_{abs}(\lambda_{1,2})_B \quad (5,6)$$

197 where A and B two different LAA species (BC, BrC) or sources (FF, BB),  $\lambda_1$  and  $\lambda_2$  are two different  
 198 wavelengths. If  $AAE_A$  and  $AAE_B$  are chosen a-priori as well as  $\lambda_1$  and  $\lambda_2$ , we are in the case of the so-  
 199 called Aethalometer model for FF and BB apportionment (Sandradewi et al., 2008a): in this paper  
 200  $AAE_{FF}=1$ ,  $AAE_{BB}=2$ ,  $\lambda_1=470$  nm,  $\lambda_2=950$  nm. If  $AAE_A$  is chosen a-priori and only  $\lambda_2$  is used, as in Tian  
 201 et al., (2019), the method allows a BC/BrC apportionment, based on the calculation of  $b_{abs}(\lambda)_{BC}$  assuming



202 that  $b_{abs}(950)_{BrC}=0$  at high wavelengths and  $AAE_{BC}=1$ ;  $b_{abs(\lambda)}_{BrC}$  is obtained as difference from the total  
203 and  $AAE_{BrC}$  is estimated on its basis.

204 If finally,  $\lambda_1$  and  $\lambda_2$  are not chosen a-priori, but all possible Aethalometer wavelengths are used, two  
205 methods can be applied by fitting  $b_{abs}(\lambda)$  over the whole spectrum:

206 1) the Multi- $\lambda$  fit approach reported in Bernardoni et al., 2021 from which it follows that

$$207 \quad b_{abs}(\lambda) = A' \cdot \lambda^{-AAE_{FF}} + B' \cdot \lambda^{-AAE_{BB}} \quad (7)$$

208 where  $A'$  and  $B'$  in equation (7) are obtained by multi- $\lambda$  fit of the input data, assuming a-priori  $AAE_{FF}$   
209 and  $AAE_{BB}$  values (1 and 2).

210 2) the MWAA (Multi-Wavelength Absorption Analyzer) model described in previous works  
211 (Massabò et al., 2015; Bernardoni et al., 2017) enables to assess the contributions of BC and BrC  
212 and to provide information on the  $AAE_{BrC}$  by

$$213 \quad b_{abs}(\lambda) = A'' \cdot \lambda^{-AAE_{BC}} + B'' \cdot \lambda^{-AAE_{BrC}} \quad (8)$$

214 where  $A''$ ,  $B''$  and  $AAE_{BrC}$  are obtained by multi- $\lambda$  fit of the input data, provided that a value for  $AAE_{BC}$   
215 is assumed a-priori. It is noteworthy that in the multi- $\lambda$  fit for the aethalometer model source  
216 apportionment (eq. 7),  $AAE$  for both sources ( $AAE_{FF}$  and  $AAE_{BB}$ ) must be assumed a priori. Opposite, the  
217 MWAA model for LAA species apportionment also provides as output the information on  $AAE$  for one  
218 of the species ( $AAE_{BrC}$ ). In this case,  $AAE_{BC}=AAE_{FF}=1$  was chosen as done in previous applications  
219 (Bernardoni et al., 2017, Massabò et al., 2015).

220 Given all the aforementioned methodologies it is clear that the total HR accuracy (eq. 1-2), and thus its  
221 apportionment ( $HR=HR_{BC}+HR_{BrC}$  or  $HR=HR_{FF}+HR_{BB}$ ), is strictly dependent on the reliability of  $b_{abs}(\lambda)$ .  
222 Thus sections 2.2 and 2.3 are dedicated to the HR instrumentation, Aethalometer correction schemes and  
223 the determination of the multiple scattering correction parameter  $C$ , respectively. The flow chart of the  
224 methodological investigation is finally resumed in Figure 1 to improve the readability of the present work.

225

## 226 2.2 HR instrumentation

227 The methodology for the determination of HR (section 2.1) requires the measurements of  $b_{abs}(\lambda)$  in the  
 228 UV-VIS-NIR region at high time resolution. Therefore, as explained in the introduction (section 1), the  
 229 only instrument able to fully cover this range is the Aethalometer that performs measurements at 7  
 230 wavelengths (370, 470, 520, 590, 660, 880 and 950 nm). During the COLOSSAL campaign two models  
 231 of Aethalometers (AE) were used in parallel, the AE31 and AE33 (Magee Scientific, Aerosol) with two  
 232 different filter tapes: a quartz-fibre filter (Pallflex Q250F) for the model AE31, and a tetrafluoroethylene  
 233 (TFE)-coated glass filter (Pallflex “Fiberfilm” T60A20) for the model AE33 (Drinovec et al., 2015). AE31  
 234 and AE33 provide as output equivalent BC (eBC) (Petzold et al., 2013) concentration, and the  
 235 corresponding optical parameters, at 5- and 1-minute time resolution, respectively.

236 The AE measure the ATN of light at 7 wavelengths transmitted through a particle laden filter spot and  
 237 through a sample free portion of the tape acting as a reference. ATN is calculated as:

$$238 \quad ATN = \ln(I_0/I) \quad (9)$$

239 where  $I_0$  and  $I$  are the intensity of light transmitted through the reference and aerosol blank spot of the  
 240 filter, respectively. In the Aethalometer the aerosol sample is continuously deposited on one/two spots  
 241 (AE31/AE33) of a filter tape which moves once an attenuation (ATN) threshold (120) is reached at 370  
 242 nm.

243 The attenuation coefficient of the aerosol particles collected on the filter tape,  $b_{ATN(\lambda)}$ , is then defined as  
 244 follows (Weingartner, et al., 2003):

$$245 \quad b_{ATN(\lambda)} = \frac{A}{Q} \frac{\Delta ATN(\lambda)}{\Delta t} \quad (10)$$

246 where A is the filter spot area, Q the volumetric flow rate and  $\Delta ATN$  is the change in attenuation during  
 247 the time interval  $\Delta t$ .

248 It is noteworthy that  $b_{ATN}$  differs from the aerosol absorption coefficient of airborne particles because it is  
 249 determined from attenuation of light passing through a particle-laden filter; the resulting artifacts can be  
 250 corrected with different procedures (Weingartner, et al. 2003, Arnott, et al., 2005, Schmid, et al. 2006,

251 Collaud Coen, et al. 2010; Drinovec et al., 2015) in function of the Aethalometer model and filter types  
252 (section 1).

253 Aethalometer AE31 artifacts were corrected applying the Weingartner et al. (2003) procedure that  
254 introduces the parameters  $C$  and  $R(ATN, \lambda)$ , used to convert  $b_{ATN}(\lambda)$  to  $b_{abs}(\lambda)$  as follows:

$$255 \quad b_{abs}(\lambda) = \frac{b_{ATN}(\lambda)}{C \cdot R(ATN, \lambda)} \quad (11)$$

256 where  $C$  and  $R(ATN, \lambda)$  are the filter multiple scattering correction parameter and the  $\lambda$ -dependent loading  
257 effect correction parameter, respectively. The parameter  $R(ATN, \lambda)$  compensates for the loading effect due  
258 to the reduction in the optical path due to an increase of the sample collected on the filter over time  
259 (Weingartner et al., 2003).

260  $R(ATN, \lambda)$  was dynamically determined following the Sandradewi et al. (2008b) algorithm. This approach  
261 was recognised to be one of the best approaches as correction does not affect data in terms of the  
262 absorption Ångström exponent (AAE) (Collaud Coen et al., 2010) and it was previously applied to data  
263 collected at the investigated site since 2015 (Ferrero et al., 2018): for the experimental assessment of HR  
264 any artificial perturbation of the AAE must be avoided.

265 The parameter  $C$  compensates for the enhanced optical path through the filter caused by multiple  
266 scattering of the light induced by the filter fibers and by the particles embedded in it. Its determination, in  
267 function of time and wavelength, using reference absorption instruments (and application to HR) is  
268 discussed in the next section (2.2). Here we recall (see also section 1) that several  $C$  values are reported  
269 in literature, ranging from 2.14 up to 5 (Weingartner et al., 2003; Arnott et al., 2005; Schmid et al., 2006;  
270 Sandradewi et al., 2008a, 2008a; Collaud Coen et al., 2010) with a suggested  $C=3.5$  from GAW (2016)  
271 constant for all wavelengths of AE31.

272 In the Aethalometer AE33, the loading artifact is measured and corrected internally due to the DualSpot  
273 technology as reported in Drinovec et al. (2015). Thus, in the AE33,  $b_{abs}(\lambda)$  is retrieved as follows:

$$274 \quad b_{abs}(\lambda) = \frac{b_{ATN\_LC}(\lambda)}{C} \quad (12)$$

275 where  $b_{ATN\_LC(\lambda)}$  is the attenuation coefficient already corrected for the loading effect (Drinovec et al.,  
 276 2015), and  $C$  is set at 1.57 for tetrafluoroethylene (TFE)-coated glass filter (Pallflex “Fiberfilm” T60A20).  
 277 The mass absorption cross-section at 880 nm is assumed to be  $7.77 \text{ m}^2 \text{ g}^{-1}$ .  
 278 As the main aim of the present work is the determination of the influence of  $C$  on the HR for both AE31  
 279 and AE33, the methodology to investigate this topic is discussed in the next section (2.3). Here below,  
 280 other needed measurements for HR determination are described.  
 281 In order to determine the HR, irradiance measurements were obtained with a Multiplexer-Radiometer-  
 282 Irradiometer (MRI) and a meteorological station (LSI-Lastem) equipped with broadband radiometers.  
 283 The MRI detects downward and upward irradiance in the UV, visible and in the near-infrared spectrum  
 284 (UVNIR), from 350 to 1000 nm in 3648 spectral bands. The instrument uses an optical switch to  
 285 sequentially select among many different fiber optical inputs fixed to the up-looking and the down-  
 286 looking entrance foreoptics. The fiber optic multiplexer MPM-2000-2x8-VIS (Ocean Optics Inc., USA)  
 287 is used as an optical switch, connecting up to 8 different input channels into two different spectrometers.  
 288 In addition, a rotating shadow-band is implemented to allow measurements of direct, diffuse, and reflected  
 289 irradiance. Further details of the MRI are reported in Cogliati et al. (2015).  
 290 The LSI-Lastem meteorological station is equipped with a global pyranometer (DPA154 model) for  
 291 measuring global irradiance and a net radiometer (C201R model) for net radiation measurements. The  
 292 combination of a shadow band with another DPA154 pyranometer allows the measurement of diffuse  
 293 irradiance (either from sky and clouds) and the retrieval of direct irradiance after subtraction from the  
 294 global one. Before the subtraction the shadow band effect must be corrected (Ferrero et al., 2018) to  
 295 determine the true amount of diffuse irradiance.  
 296 The LSI-Lastem is also equipped with temperature (DMA580 sensor with data range from -30 to +70 °C  
 297 with an accuracy of  $\pm 0.1 \text{ °C}$  and with a sensitivity of  $0.025 \text{ °C}$ ), humidity (DMA570 sensor with data  
 298 range includes values from 10% to 98% with an accuracy of  $\pm 2.5\%$  and sensibility of 0.2%), pressure

299 (CX110P barometer with data range from 800hPa to 1100 hPa with an accuracy of 1 hPa) and wind (speed  
300 and direction, CombiSD anemometer with data range from 0 to 60 m/s and from 0 to 360°) sensors.  
301 Finally, the U9 sampling site is equipped with a Condensation Particle Counter (CPC, 3775 TSI) which  
302 provides particle number concentration ( $N$ ;  $d_{50}=4$  nm) in a concentration range from 0 to  $10^7$  particles/cm<sup>3</sup>;  
303 in addition, an optical particle counter (OPC, Grimm 1.107) which measured the particle number  
304 concentration and size distribution in the range 0.25 - 32  $\mu$ m and estimated PM<sub>10</sub>, PM<sub>2.5</sub> and PM<sub>1</sub> mass  
305 concentrations.

306

### 307 **2.3 Determination of the multiple-scattering correction parameter $C$**

308 The multiple scattering correction  $C$  is determined by comparing the attenuation coefficient, already  
309 corrected for the loading effect ( $b_{ATN\_LC}(\lambda)$ ; either from AE31 and AE33), with the absorption coefficient  
310 measured simultaneously at the same wavelength with a reference instruments ( $b_{abs\_ref}(\lambda)$ ) as follows:

$$311 \quad C(\lambda) = \frac{b_{ATN\_LC}(\lambda)}{b_{abs\_ref}(\lambda)} \quad (13)$$

312 Recalling equations 1-2, the HR is directly proportional to the integral (over the whole shortwave  
313 spectrum) of the product of  $b_{abs}(\lambda)$  and the radiation  $F_n(\lambda, \theta)$ . Thus, any inappropriate assumption on the  $C$   
314 value, e.g. it does not vary over wavelengths and time, will impact the HR and its apportionment to sources  
315 (FF, BB) and species (BC, BrC). The experimental determination of these impacts and best practice for  
316 climate data harmonization is the aim of the present work.

317 For this reason, two reference instruments for  $b_{abs\_ref}(\lambda)$  were used in order to give complementary  
318 information:

- 319 1- Highly time-resolved measurements at a single, fixed wavelength, using a Multi-Angle  
320 Absorption Photometer (MAAP, Thermo Scientific Model 5012)
- 321 2- Wavelength-resolved measurements with low time resolution (12 h), using the polar photometer  
322 (PP\_UniMI, set up at the Department of Physics, University of Milan)

323 Here below, the two reference instruments are briefly described.

324 The first reference, the Multi-Angle Absorption Photometer (MAAP, Thermo Scientific Model 5012)  
325 measures the aerosol absorption coefficient at a wavelength of 637 nm (Müller et al., 2011) and provides  
326 the atmospheric eBC at high time resolution (5 minutes). It simultaneously measures the transmission and  
327 scattering of light in the forward and back hemisphere (at 3 different angles) of a blank and a particle  
328 laden filter tape. The absorption coefficient at 637 nm is thereafter obtained from a radiative transfer  
329 scheme (Petzold and Schönlinner, 2004; Petzold et al., 2005). The eBC is obtained from absorption  
330 coefficient considering the deposit area and sampling air flow and setting at  $6.6 \text{ m}^2 \text{ g}^{-1}$  the eBC mass-  
331 specific absorption coefficient. Compared to AE, instrumental artifacts are reduced in the MAAP so that  
332 the absorption coefficient measured should be closer to the true one and therefore it is traditionally used  
333 as a reference among filter-based instruments in literature works (Collaud-Coen et al. 2010). With respect  
334 to *C* determination with AE-MAAP comparison we underlined that the nominal AE31 and AE33 660 nm  
335 channel is provided by a Kingbright led (APT 1608SRC PRV 1.6 x 0.8 mm SMD Chip LED Lamp;  
336 <https://www.kingbrightusa.com/images/catalog/SPEC/APT1608SRCPRV.pdf>) which is characterized by  
337 a wavelength peak emission at 655 nm with a 20 nm spectral full bandwidth at half maximum under 20mA  
338 of supplied current (information from manufacturer). This is in agreement with the absorption  
339 photometers intercomparison reported by Muller et al. (2011) in which the nominal AE red channel was  
340 determined to be at 654 nm with a 23 nm spectral full bandwidth at half maximum; they also reported the  
341 same data for the red MAAP channel which is characterized by a wavelength peak emission at 637 nm  
342 with an 18 nm spectral full bandwidth at half maximum. Thus, for practical purpose the AE and MAAP  
343 work in the red channel approximately at the same wavelength allowing their direct comparison for *C*  
344 determination avoiding to introduce computation noise by reporting the nominal 660 nm channel to 637  
345 nm in function of a delicate AAE determination as demonstrated in the results and discussion sections.

346 The second reference is the polar photometer (PP\_UniMI, set up at the Department of Physics, University  
347 of Milan) where different laser sources are mounted on a sliding motor (at the time of this work, 4 laser  
348 sources at 405, 532, 635 and 780 nm were used). The chosen laser beam impinges perpendicularly on the  
349 aerosol sample and a photodetector is mounted on a rotating arm; it performs measurements of the

transmitted and of the scattered radiation at scattering angles between 0° and 173° (with 0.4° resolution). This allows to retrieve the total light scattered in the front- and backward hemispheres with no further assumptions. The radiative transfer model used to obtain the aerosol absorption optical depth of the sample is the same as the one used in the MAAP (Petzold and Schölinner, 2004); the aerosol absorption coefficient is then calculated considering the area of the deposit and the sampled volume (Bernardoni et al. 2017; Vecchi et al. 2014).

However, the radiative transfer model by Petzold and Schölinner (2004) can be applied to the PP\_UniMi data both considering the same assumptions performed by the MAAP (in which fixed values for the fraction of the back-scattered radiation by filter matrix and the asymmetry factor are imposed a-priori) or exploiting the whole available information given by the measurement of the whole phase function from the PP\_UniMi (Valentini et al., 2021). Within the framework of EMEP/ACTRIS/COLOSSAL 2018 winter campaign, the PP\_UniMi  $b_{abs\_ref}(\lambda)$  was determined using two different approaches indicated as PP (Polar Photometer) and PaM (Photometer as MAAP). The PP approach fully exploits highly angular-resolved measurements, while PaM calculation introduces the same approximations as the ones used in the MAAP - i.e. the reconstruction (by analytical functions) from measurements at the same 3 angles and with the same fixed value between backward and total diffused radiation for blank filter ( $B_M=0.7$ ) used in MAAP (Valentini et al., 2021; Bernardoni et al., 2021). It noteworthy that, recently, Valentini et al. (2021) demonstrated that by fixing  $B_M=0.7$  affected  $b_{abs}$  obtained from MAAP with an average positive artifact of ~15%. Indeed, blank spots measured by PP\_UniMi (fully exploiting highly angular-resolved measurements) showed  $B_M=0.88$ . As a consequence, Bernardoni et al. (2021) found a linear relationship between PP and MAAP, which was confirmed by Valentini et al. (2021) independently from site and season, showing an overestimation of  $b_{abs}(637)$  obtained by MAAP. As reported in Bernardoni et al. (2021), at the U9 sampling site it was as:

$$b_{abs}(PP\ 635)=0.928* b_{abs}(MAAP\ 637)-2.07 \quad (14)$$

this equation will be used in section 3.4 to determine the time-dependent  $C$  values. Whereas an absolute and fully consistent reference method for  $b_{abs}(\lambda)$  and eBC is not actually present as

reported by Petzold et al. (2013), the aforementioned considerations enable to consider PP\_UniMI (PP approach) as the most promising reference for  $b_{abs\_ref}(\lambda)$  and  $C$  determination in the present work.

Thus, as the study aims at determining the response of the HR values as function of all the possible variations of the  $C$  values (in wavelength and in time), the following nomenclature will be applied to  $C$ :

- 1-  $C_{fix}$  indicates fixed values of  $C$  determined at 660 nm and applied to all AE wavelengths without any time-dependency, obtained using either MAAP (comparisons at high time resolution) or PP\_UniMI (PP and PaM approaches) as reference systems
- 2-  $C_\lambda$  indicates wavelength-dependent  $C$  values obtained from 370 to 950 nm (not varying in time) by a comparison with the multi- $\lambda$  PP\_UniMI  $b_{abs}(\lambda)$  reference system (section 2.3; PP approach).
- 3-  $C_t$  indicates time-dependent  $C$  values obtained at 660 nm by comparison with MAAP corrected data (eq. 13)
- 4-  $C_{\lambda\_t}$  finally indicates wavelength- and time-dependent  $C$  values by applying the wavelength-dependence of  $C_\lambda$  (found at point 2) to the time-dependent  $C_t$  values (point 3).

Among the aforementioned  $C$ , the wavelength-dependent  $C_\lambda$  determination needs to be clarified (details are reported in Bernardoni et al., 2021). In fact, PP\_UniMI, measures  $b_{abs}(\lambda)$  at four wavelengths (405, 532, 635 and 780 nm) while both AE31 and AE33 works at seven wavelengths (370, 470, 520, 590, 660, 880 and 950 nm) not coincident with those of PP\_UniMI. Thus, for the aim of the present paper PP\_UniMI data were reported to Aethalometer wavelengths finding the best experimental absorption Ångström exponent for each sample (multi- $\lambda$  fit of the four  $b_{abs}(\lambda)$  as  $b_{abs}(\lambda) = A\lambda^{-AAE}$  for both PP and PaM approaches) and applying it to derive  $b_{abs}(\lambda)$  at all Aethalometer wavelengths with the following generic relationships:

$$b_{abs}(\lambda_{AE}) = b_{abs}(\lambda_{PP\_UniMI}) \left( \frac{\lambda_{AE}}{\lambda_{PP\_UniMI}} \right)^{-AAE} \quad (15)$$

where  $\lambda_{AE}$  and  $\lambda_{PP\_UniMI}$  are each pair of the nearest couple of Aethalometer and PP\_UniMI wavelengths (e.g. 470 and 405 nm, 660 and 635 nm, etc.).



### 3. Results and discussions

Results are presented describing first the environmental context of atmospheric pollution during the campaign (section 3.1); next, sections 3.2, 3.3 and 3.4 present the HR values (total and apportioned to different species and sources of LAA) as a function of  $C_{fix}$ ,  $C_\lambda$ ,  $C_b$ , and  $C_{\lambda_t}$ . Conclusions follow with a best practice procedure for Aethalometer correction for radiative forcing assessment purposes. All the data are reported hereinafter as mean $\pm$ 95% confidence interval. All times are reported in local solar time (LST).

#### 3.1 Environmental context

Highly-time resolved eBC (880 nm) loading corrected data recorded during the campaign (by both AE31 and AE33) are reported in Figure 2 together with PM<sub>2.5</sub> concentrations, N, wind speed, temperature and relative humidity. At the beginning of the campaign, a Föen event (4 m s<sup>-1</sup> wind, 30°N) cleared the atmosphere bringing the concentrations to negligible eBC and PM<sub>2.5</sub> values; after this event stable atmospheric conditions occurred favoring a pollution accumulation within the mixing layer; in fact, eBC concentrations ranged from close to zero ( $\sim$ 20 ng m<sup>-3</sup>) up to 16.27  $\mu$ g m<sup>-3</sup>; similarly, PM<sub>2.5</sub> concentrations ranged from 1.6 up to 88.3  $\mu$ g m<sup>-3</sup> with an average value of 30.9 $\pm$ 0.4  $\mu$ g m<sup>-3</sup>. The large pollution range poses the EMEP2018 COLOSSAL campaign in a good position both for  $C$  determination in different conditions (Bernardoni et al., 2021) and for the quantification of the LAA (and related  $b_{abs}(\lambda)$ ) impact on the HR, with broader implications for any radiative transfer computation.

Figure 2 showed that eBC data, by both AE31 and AE33, were well correlated, leading to the same environmental eBC trend along time, thus appearing in good agreement (considering the different principle of operations, filter tape and constant attenuation/absorption cross sections of AE31 and AE33) for environmental monitoring purposes (and long term data series continuity). More in detail, results showed just slightly higher eBC readings ( $\sim$ 10%) obtained from AE33 raw data; the average eBC concentrations were 3.47 $\pm$ 0.05  $\mu$ g m<sup>-3</sup> (AE31) and 3.81 $\pm$ 0.05  $\mu$ g m<sup>-3</sup> (AE33) being statistically different at 95% confidence level. This feature can be evidenced both with the boxplot reported in Figure 3a, and

with the linear regression reported in Figure 3b: it shows high correlation ( $R^2 = 0.977$ ) with a slope slightly higher than one ( $1.049 \pm 0.003$ ) and a negligible intercept ( $0.073 \pm 0.015 \mu\text{g m}^{-3}$ ).

The diurnal pattern of eBC (averaged between AE31 and AE33, given the above results) is reported in Figure 4a-b, together with that of  $\text{PM}_{2.5}$ , N, wind speed, temperature and relative humidity. The eBC, N and  $\text{PM}_{2.5}$  experienced a morning peak (rush hour,  $4.61 \pm 0.12 \mu\text{g m}^{-3}$ ,  $26 \cdot 10^3 \pm 666 \text{ cm}^{-3}$  and  $32.2 \pm 0.9 \mu\text{g m}^{-3}$ , respectively) followed by a decrease of 45, 44 and 17% during the mid-afternoon ( $2.56 \pm 0.08 \mu\text{g m}^{-3}$ ,  $14.6 \cdot 10^3 \pm 284 \text{ cm}^{-3}$  and  $26.8 \pm 1.1 \mu\text{g m}^{-3}$ , respectively) (Figure 4a); conversely the wind speed and temperature increased, with double values compared to morning time (Figure 4b) thus increasing atmospheric dispersion potential as typically reported in previous studies (Ferrero et al., 2011, 2019a; Carbone et al., 2010, Diemoz et al., 2019, Vecchi et al., 2019).

Within the aforementioned environmental context, the following sections investigate the impact of  $C_{fix}$ ,  $C_\lambda$ ,  $C_b$ , and  $C_{\lambda-t}$  on the HR determination in function of species and sources of LAA.

439

## 440 **3.2 $C_{fix}$ and resulting heating rate**

### 441 **3.2.1 $C_{fix}$ values**

The multiple scattering parameters  $C_{fix}$  were obtained using different instruments as reference. They were determined by comparing the AE31 and AE33 loading-corrected attenuation coefficients ( $b_{ATN\_LC}$ ) at 660 nm with the absorption coefficient measured: 1) by MAAP and 2) PP\_UniMI, in both cases with no further wavelength adjustment (section 2.3).

The parameters  $C_{fix}$  obtained by a direct 5-min time-resolved comparison with MAAP were  $3.43 \pm 0.01$  and  $2.64 \pm 0.01$  for AE31 and AE33 respectively. The AE31 value is comparable with  $3.50 \pm 0.25$  suggested by GAW (2016) and with the expected range in Milan during wintertime recently reported in Ferrero et al. (2021):  $3.20 \pm 0.35$ . The similarity between the obtained C value ( $3.43 \pm 0.01$ ) and the GAW one confirms that Milan (in the middle of the Po Valley) is characterized by continental type aerosols (e.g.

451 Carbone et al., 2010) and that the Milan value is consistent with the global average as already verified in  
 452 Ferrero et al. (2021).

453 The parameters  $C_{fix}$  was also determined by a comparison with PP\_UniMI  $b_{abs\_ref}(\lambda)$  using both the PP  
 454 and PaM approaches (PP: full angular resolution; PaM: three angles as MAAP; section 2.3; Bernardoni  
 455 et al., 2021) considering both all the available data ("all") and for daytime ("day") data only. Isolation of  
 456 daytime data was performed as HR is determined considering aerosol-radiation interaction, which occurs  
 457 only during daytime, due to radiation absence during night-time. Table 1 resumes all the results also  
 458 including the traditional values for the instruments and tape in use:  $C_{AE31}=2.14$  (Weingartner et al., 2003)  
 459 and  $C_{AE33}= 1.57$  (Drinovec et al., 2015), for further comparison on the HR data (see below). It is  
 460 noteworthy that the  $C_{fix}$  obtained by a comparison with PP\_UniMI under the PP approach for AE31  
 461 ( $4.30\pm0.07$  and  $4.44\pm0.11$ , all day and daytime, respectively) and AE33 ( $3.37\pm0.05$  and  $3.43\pm0.08$ , all  
 462 day and daytime, respectively) were higher than the corresponding values obtained by a comparison with  
 463 both PP\_UniMI under the PaM approach and the ones obtained by a direct 5-min time-resolved  
 464 comparison with MAAP. Thus,  $C_{fix\_PaM}$  were comparable with  $C_{fix\_MAAP}$  while  $C_{fix\_PP}$  were 21% higher  
 465 than  $C_{fix\_PaM}$ . This is in agreement with results reported by Valentini et al. (2021) and recalled in section  
 466 2.3; assumptions made in the MAAP lead to higher  $b_{abs\_ref}(\lambda)$  and thus lower  $C_{fix}$  (eq. 13) compared to PP  
 467 approach.

468 Daytime values were higher than all day ones (and, consequently, they are higher than night-time results)  
 469 for both aethalometers. This may be due to two different causes: a different aerosol size distributions and  
 470 single scattering albedo (SSA). We begin with the role of different size distribution. A higher fraction of  
 471 bigger particles can change the scattering properties of the aerosol deposited on the filter matrix. In this  
 472 respect, the daytime (06:00-18:00 local solar time) and night-time (18:00-06:00 local solar time) average  
 473 number size distribution measured by optical particle counter (OPC Grimm 1.107) showed a difference  
 474 for particle size  $> 5.0 \mu\text{m}$  optical equivalent diameter (Figure S1). Ferrero et al. (2014) showed that  $5.0 \mu\text{m}$   
 475 optical equivalent diameter (measured by an OPC Grimm 1.107) corresponds to an ambient geometric  
 476 size of  $5.9 \mu\text{m}$ , in Milan during wintertime. Thus, as the experimental set-up for the COLOSSAL winter

2018 campaign was equipped with PM<sub>2.5</sub> cutoff inlet (Bernardoni et al., 2021), the size distribution differences observed in the coarse size range between daytime and night-time is expected to have a negligible effect. The second cause, a different SSA, is instead more reliable. Thus, this aspect was investigated given the findings reported in Collaud Coen et al. (2010) which defined the  $C$  values in function of the aerosol single scattering albedo (SSA) as follows:

$$C = C_{ref} + \alpha \frac{SSA}{1-SSA} \quad (16)$$

where  $C_{ref}$  is the reference value of  $C$  for the AE31 tape (determined in the pristine atmosphere of Jungfraujoch and Hohenpeissenberg sites where aerosol has a single scattering albedo of  $\sim 1$ ;  $C_{ref}$  was equal to  $2.81 \pm 0.11$ ),  $\alpha$  is the parameter for the Arnott (2005) scattering correction ( $0.0713$  at  $660$  nm).

From the aforementioned statement, it is clear that the analysis can be carried out on the AE31  $C$  data only; nonetheless, being the SSA and intrinsic property of the aerosol, the results also allow the interpretation of AE33  $C$  data. Moreover, as Collaud Coen et al. (2010) derived  $C_{ref}$  from a comparison with MAAP data, eq. 16 was applied to PaM (section 2.3)  $C$  daytime and night-time values to derive the corresponding SSA.

By inverting eq. 16, the retrieval of SSA from the experimental  $C$  and  $C_{ref}$  led to a SSA value of  $0.92 \pm 0.01$  during daytime,  $0.02 \pm 0.01$  higher than the  $0.90 \pm 0.01$  SSA during night-time. The same day to night SSA difference ( $0.02 \pm 0.01$ ) can also be achieved by using eq. 1 on high-time resolution  $C_t$  data (section 3.4) obtained by a direct AE31-MAAP comparison. As the SSA is related to change in the aerosol chemical composition and especially in the aerosol fraction due to the eBC, the daytime and night-time eBC fraction in PM<sub>2.5</sub> was investigated. Results showed a percentage of eBC in PM<sub>2.5</sub> of  $11.9 \pm 0.3\%$  and  $12.7 \pm 0.3\%$  (being statistically different at 95% of confidence) during daytime and night-time respectively; they are in keeping with the aforementioned SSA behaviour. Finally, the total number concentration data ( $N$ ) were also used to evaluate the  $N/eBC$  ratio (Rodríguez and Cuevas, 2007; Reche et al., 2011; Dall'Osto et al., 2011, 2013; Ferrero et al., 2016). The highest values of the  $N/eBC$  ratio are expected to occur during secondary aerosol formation in the atmosphere (Reche et al., 2011) which enhance the aerosol scattering,

thus lowering the SSA (Ferrero et al., 2011a; 2014). N/eBC were  $5.58 \pm 0.04$  and  $4.86 \pm 0.04 \text{ cm}^{-3} \text{ ng}^{-1} \text{ m}^3$  during daytime and night-time, respectively. Finally, due to Aethalometer type and filter differences all the  $C_{fix}$  values found for AE31 were 30% higher than the corresponding for AE33.

506

### 3.2.2 $C_{fix}$ and total HR assessment

The different  $C_{fix}$  values will impact the determination of the HR, as described by eq. 1 (section 2.1). To this purpose, in the  $C_{fix}$  case only, eq. 1 can be rewritten as follows:

$$HR = \frac{1}{C_{fix}} \cdot \frac{1}{\rho C_p} \cdot \sum_{n=1}^3 \int_{\theta} \int_{\lambda} \frac{F_{n(\lambda, \theta)}}{\cos \theta} b_{ATN\_LC(\lambda)} d\lambda d\theta \quad (17)$$

thus showing that, once the HR with a  $C_{fix}$  value has been determined, it is possible to rescale the data using a different  $C_{fix}$  and easily to determine the HR. This represents a simple application for long-term data series harmonization of radiative forcing and HR. Thus, all HR data with all possible  $C_{fix}$  estimates were obtained by applying eq. 17; they are resumed in Table 1 where HR values range from  $0.58 \pm 0.02$  up to  $1.21 \pm 0.05 \text{ K day}^{-1}$  for AE31 and from  $0.61 \pm 0.02$  up to  $1.34 \pm 0.05 \text{ K day}^{-1}$  for AE33 (relative variability above 100%). These data are reported in Figure 5a which excludes the extremes  $C_{fix}$  values of 2.14 (Weingartner et al., 2003) and 1.57 (Drinovec et al., 2015) for AE31 and AE33, respectively, which should not be used for  $b_{abs}(\lambda)$  determination from AE31 and AE33 (they are anyway reported in Table 1 and included in Figure S2). When these extremes  $C_{fix}$  values (2.14 and 1.57) are excluded, by applying all the  $C_{fix}$  obtained with the other approaches, all the HR determination are constrained within a  $\Delta HR$  of  $0.17 \pm 0.04$  and  $0.18 \pm 0.04 \text{ K day}^{-1}$  (for AE31 and AE33) which turns into an ambient climatic error within  $\sim 30\%$  at maximum.

If the reference  $C_{fix}$  values determined from PP at 660 nm are used, the best estimates for the AE31 and AE33 HRs become:  $0.58 \pm 0.02 \text{ K day}^{-1}$  and  $0.61 \pm 0.02 \text{ K day}^{-1}$ , respectively. This HR data obtained using the PP  $C_{fix}$  were the lowest obtained among all the data reported in Table 1 and Figure 5a. In fact,

the HR are overestimated by a  $23\pm6\%$  and  $31\pm6\%$  factor (for AE31 and AE33) if the  $C_{fix}$  values obtained by the PaM approach and high-resolution MAAP are considered as reference, respectively.

Thus, two consequences can be highlighted from these results for the total HR:

- 1- Consistency: the total HR determined for each of the  $C_{fix}$  pair obtained for the two aethalometers with the same reference device/approach presented average values not statistically different at 95% probability.
- 2- Over-estimation: if MAAP/PaM is chosen as reference instead of a PP approach, the HR can be overestimated by 23-30% factor.

Both results can be considered site-independent and discussed as follows:

- 1- Consistency: this property is important for long-term data series in which e.g. the AE31 is supplanted by the AE33. Due to the absence of standard reference material for eBC and to the lack of gold standard instrumentation for the measurement of ambient aerosol absorption coefficient, at the current state of the art, the consistency refers to the reproducibility between the total HR, obtained using the AE31 and AE33, not to the accuracy. To obtain the consistency (i.e. reproducibility), it is necessary that, for both AE31 and AE33, the  $C_{fix}$  values are obtained with the same reference device (e.g. a MAAP, a photoacoustic, polar photometer) at the same time; it is important that the temporal length of overlapping in the calibration with the reference device covers all the possible aerosol scenarios (i.e. chemical composition) in the different seasons (over the investigated site) to obtain a temporal consistency. This makes the consistency property site-independent.
- 2- Over-estimation: this happens if MAAP/PaM is chosen as reference instead of a PP approach. This observation completes the one related to the consistency as it refers to the accuracy of the total HR in function of the reference device used to derive the C factor applied to the Aethalometers (AE31 and AE33). It is noteworthy that the MAAP is a widely accepted reference device in the scientific community. Nevertheless, Valentini et al. (2021) derived a relationship between the MAAP absorption coefficient (that is obtained with limitation of imposing: 3 fixed angles, the

backscattered fraction from blank filter at 0.7 and the asymmetry factor at 0.75) and the one obtained from the PP\_UniMI (that allows to retrieve the total light scattered in the front- and backward hemispheres with no further assumptions – thus in principle improving measurement accuracy). This relationship was derived considering aerosol sampled at different urban and regional background sites in southern Europe, by a comparison between MAAP data and PP\_UniMI analysis of the MAAP spots collected during the MAAP measurements. Thus, we expect that the relationship found by Valentini et al. (2021) has wide applicability (it generalizes eq. 14, section 2.3, which instead is valid only for the Milan case). Such relationship is:

$$b_{abs}(PP)=0.86(\pm 0.01)*b_{abs}(MAAP)-0.19(\pm 0.03) \quad (18)$$

where the absorption coefficient  $b_{abs}$  is expressed in  $Mm^{-1}$ . Thus, once MAAP data are available in each site, they can be converted in the corresponding PP values enabling to improve the accuracy of  $b_{abs}$  determination and thus that of the total HR.

The possibility to solve the over-estimation of the total HR is important as the HR quantification will reflect on the assessment of the atmospheric feedbacks related to the HR; for example, the influence on the liquid water content (Jacobson et al., 2002), planetary boundary layer dynamics (Su et al., 2020; Ferrero et al., 2014), regional circulation systems (Ramanathan and Carmichael, 2008; Ramanathan and Feng, 2009) and cloud dynamic and evolution itself (Koren et al., 2008; Bond et al., 2013). Thus, any arbitrary, not verified, use of  $C_{fix}$  will also be reflected in the modelled HR-triggered feedbacks with the aforementioned uncertainty range of ~30% (at minimum, if climatic amplification phenomena in the triggered feedbacks are not present).

Finally, the all vs. day  $C_{fix}$  impact on the HR reported in Table 1 highlighted that the use of a day-time  $C_{fix}$  turns into a total HR that was 3.4% lower with respect to that obtained using the whole day (“all”)  $C_{fix}$ . A detailed discussion on this topic is included in section 3.4.

Here below, the impact of using a wavelength independent  $C_{fix}$  on the HR apportionment between FF and BB and between BC and BrC contribution is discussed.

### 578 3.2.3 $C_{fix}$ and HR apportionment

579 The previous results reported for the total HR highlighted the possibility to obtain a temporal consistency  
580 in long-term data series. Conversely, the data analysis showed a different scenario when the LAA sources  
581 (FF and BB) and speciation (BC and BrC) are considered by using the “Aethalometer model” (Sandradewi  
582 et al., 2008a) or the method proposed by Tian et al. (2019) (section 2.1). Their values are resumed in  
583 Figure 5b (for PP  $C_{fix}$  daytime data only) and in Table 1 for all the  $C_{fix}$  used above.

584 Focusing on the single  $C_{fix}$  calculation approach for data correction,  $HR_{FF}$  obtained by the AE33 data  
585 showed higher values ( $+24\pm1\%$ ;  $+0.10\pm0.02$  K day<sup>-1</sup>) than that of AE31; conversely,  $HR_{BB}$  for AE33 was  
586  $-28\pm1\%$  ( $-0.08\pm0.01$  K day<sup>-1</sup>) lower than for AE31. When the LAA species (BC and BrC) were  
587 apportioned, AE33  $HR_{BC}$  values were higher ( $+0.06\pm0.01$  K day<sup>-1</sup>;  $+10\pm1\%$ ) than for AE31, and  $HR_{BrC}$   
588 for AE33 was  $-29\pm1\%$  ( $-0.03\pm0.01$  K day<sup>-1</sup>) lower than for AE31. These variations reflect on the  
589 corresponding contribution of FF/BB/BC/BrC to the total HR. In this respect,  $HR_{FF}$  and  $HR_{BB}$  accounted  
590 for  $60\pm3\%$  and  $40\pm2\%$  in the AE31 data series, while for AE33 their contribution changed to  $72\pm4\%$  and  
591  $28\pm2\%$ , respectively. A similar behavior was observed for the LAA speciation, with the  $HR_{BC}$  and  $HR_{BrC}$   
592 contribution of  $85\pm5\%$  and  $15\pm3\%$  in the AE31 data and of  $91\pm5\%$  and  $9\pm1\%$  in the AE33 data,  
593 respectively.

594 The reasons behind these behaviours lie in the  $b_{abs}(\lambda)$  spectra of AE31 and AE33 that, once  $C_{fix}$  is applied,  
595 differ at the extreme wavelengths: 370 and 950 nm (Figure 6a). These differences are respectively of 4%  
596 and 6% ( $1.5$  and  $0.7$  Mm<sup>-1</sup>), with higher values measured by the AE33. Between the two wavelengths, the  
597 highest one (950 nm) affects the classical apportionment methods (the “Aethalometer model” and the BrC  
598 determination by Tian et al., 2019) while the lowest one (370 nm) affects the increase of the AAE of the  
599 source/species. In this respect, Figure 6b-e report the average  $b_{abs}(\lambda)$  spectra (from both AE31 and AE33)  
600 for FF, BB, BC and BrC obtained using the classical apportionment methods: they are in agreement with  
601 the HR apportionment differences described above. To explain the phenomenon, it is necessary to recall  
602 that, when the HR is determined (eq. 1-2),  $b_{abs}(\lambda)$  is integrated over the whole aethalometer spectrum. In  
603 this respect, the FF and BC  $b_{abs}$  integrals were  $15\pm4$  and  $6\pm4\%$  higher for AE33 when compared to AE31.



604 Conversely, the BB and BrC integrals were  $-23 \pm 4$  and  $-26 \pm 4\%$  lower for AE33 when compared to AE31.  
605 As the HR (eq. 1-2) is determined as an integral over the whole absorption spectrum, even a small bias  
606 on  $b_{abs}(950)$  can be magnified (after the application of the “Aethalometer model” and the BrC  
607 determination by Tian et al., 2019) by the integral operator on the obtained source and species  $b_{abs}(\lambda)$   
608 which are reported in Figure 6b-e.

609 From the aforementioned results, two important consequences can be derived for the HR apportionment  
610 (sources and species):

- 611 1- Inconsistency in long temporal series when a  $C_{fix}$  value is used and older AE31 are replaced by  
612 the new AE33. This inconsistency generates when the traditional source (FF and BB) and  
613 speciation (BC and BrC) apportionment models (“Aethalometer model” and the Tian et al. (2019)  
614 method) are considered. The  $HR_{FF}$ ,  $HR_{BB}$  and  $HR_{BrC}$  presented average values statistically  
615 different at 95% probability when determined from the two aethalometers with  $C_{fix}$  obtained with  
616 the same reference device. The only exception is given by  $HR_{BC}$  probably because BC accounts  
617 for the large majority of LAA in the atmosphere.
- 618 2- Over-estimation for the AE33 concerning  $HR_{FF}$  and  $HR_{BC}$  (and vice versa for  $HR_{BB}$  and  $HR_{BrC}$ )  
619 relative to the AE31. This poses an important issue concerning the long-term quantification of the  
620 HR improvement due to the policies implemented in the context of climate change mitigation.

621 Therefore, consistent HR source apportionment/speciation approaches are discussed in the next section.

622

### 623 **3.3 Overcoming the inconsistencies in HR source and speciation apportionment**

624 As mentioned in the previous section, the inconsistency (between AE31 and AE33) of the HR source and  
625 speciation apportionment represents an important issue, especially when long-term data series are  
626 considered in the context of climatic change.

627 Two approaches can be used to overreach this problem:

- 628 1- by introducing a wavelength dependent  $C_\lambda$  parameter for each aethalometer type using PP\_UniMI  
629 (PP approach) as a reference instrument,

2- by smoothing out the role of extreme wavelengths using all the  $b_{abs}(\lambda)$  information (based on  $C_{fix}$ )

and applying multi- $\lambda$  regression fitting as “multi- $\lambda$  fit” and MWAA described in section 2.1.

As shown in the following, these approaches are alternative and not counteracting, and they feature both advantages and disadvantages.

### 3.3.1 $C_\lambda$ and HR apportionment

$C_\lambda$  data, obtained using the methodology reported in section 2.3, are reported in Table 2. When the wavelength dependent parameters  $C_\lambda$  are applied, by definition the total  $b_{abs}(\lambda)$  spectra agrees for the AE31 and AE33 (Figure 6f), as all wavelengths of the same instrument are reported to the same value. As a consequence, the application of both the “Aethalometer model” and Tian et al. (2019) method becomes consistent between AE31 to AE33 (Figure 6g-l), with negligible differences even when  $b_{abs}(\lambda)$  is integrated over the aethalometer spectrum (1/10 of the differences compared to the case in which a fixed correction  $C_{fix}$  is applied). The obtained results (in terms of HR) are resumed in Table 1 and Figure 7.

First, using  $C_\lambda$  with traditional source and speciation apportionment approaches, the total HR value lies at  $0.59 \pm 0.02$  K day<sup>-1</sup> for both AE31 and AE33, resulting in a consistent temporal evolution in a climate change monitoring scenario. With respect to the HR source apportionment, HR<sub>FF</sub> and HR<sub>BB</sub> became  $0.40 \pm 0.02$  and  $0.42 \pm 0.02$  K day<sup>-1</sup>, and  $0.19 \pm 0.01$  and  $0.17 \pm 0.01$  K day<sup>-1</sup> (AE31 and AE33). The speciated HR<sub>BC</sub> became  $0.53 \pm 0.02$  and  $0.54 \pm 0.02$  K day<sup>-1</sup>, and HR<sub>B<sub>r</sub>C</sub>  $0.06 \pm 0.01$  and  $0.06 \pm 0.01$  K day<sup>-1</sup> (AE31 and AE33), respectively. These results demonstrate the need for using accurate wavelength dependent multiple scattering parameters  $C_\lambda$ ; this is important especially when the source apportionment “Aethalometer model” and speciation (Tian et al., 2019) models are used. Obviously, there is the need to determine proper  $C_\lambda$  values using appropriate reference methods.

### 3.3.2 Application of multi-wavelength models

The second approach to overcome the inconsistency shown in section 3.2.3 is to use advanced models: the “multi- $\lambda$  fit” and MWAA apportionment models recalled in section 2.1 (Bernardoni et al., 2021;

Massabò et al., 2015). In this case, their application (Table 1 and Figure 7), based on the appropriate PP daytime  $C_{fix}$ , resulted in a  $HR_{FF}$  of  $0.39 \pm 0.02$  K day<sup>-1</sup> (both AE31 and AE33) and in a  $HR_{BB}$  of  $0.19 \pm 0.01$  and  $0.21 \pm 0.01$  K day<sup>-1</sup> (AE31 and AE33); in the light of HR speciation, the MWAA resulted in  $HR_{BC}$  of  $0.46 \pm 0.02$  and  $0.53 \pm 0.02$  K day<sup>-1</sup>, and  $HR_{BrC}$  of  $0.14 \pm 0.01$  and  $0.09 \pm 0.01$  K day<sup>-1</sup> (AE31 and AE33). Finally, the total HR did not vary with respect to previous section, being again  $0.58 \pm 0.02$  and  $0.61 \pm 0.02$  K day<sup>-1</sup> for AE31 and AE33, respectively. The aforementioned data demonstrate the consistency for FF/BB apportionment, while the MWAA application from AE31 to AE33 by using a simple  $C_{fix}$  generates a discrepancy up to 58% when considering the BrC. However the AE33  $HR_{BrC}$  ( $0.09 \pm 0.01$  K day<sup>-1</sup>) was in keeping with the values obtained in section 3.3.2 using the Tian et al. (2019) method. It is noteworthy, however, that while the wavelength dependent  $C_\lambda$  values obtained for each aethalometer type are a site-dependent result, that need to be investigated in other sites and seasons around the world (to this aim, multi-wavelength reference information is needed), the use of multi- $\lambda$  regression fitting and MWAA (Massabò et al., 2015), with care and limitations to BrC apportionment (as above mentioned) can be applied worldwide with no need of further instrumentation and check.

670

### 671 3.3.3 Comparison with literature data

The obtained results (sections 3.3.1 and 3.3.2 and Table 2) can be compared with similar studies concerning both the C values and the related effect on radiative transfer. The C values of the Aethalometer AE31 (resumed in Table 1) obtained using the MAAP/PaM reference are close to that reported in Weingartner et al. (2003) ( $3.66 \pm 1.53$  and  $3.90 \pm 1.13$ ) referring to its experiment in smog chamber when soot particles were internally mixed with organic matter that in Milan accounts for more than 35% in winter (Perrone et al., 2012; Ferrero et al., 2019a) due to the aged nature of the Po Valley aerosol. Arnott et al. (2005) arrived to a similar conclusion addressing a  $C=3.69$  as the most suitable for ambient measurements in urban sites. Schmid et al. (2006) found AE31 C values up to 4.0-4.4 for internally mixed aerosol (e.g., for aged ambient aerosol), very close to the  $C_\lambda$  PP range (4.27-4.63) reported in the present

work and detailed in Bernardoni et al. (2021). For what concerns the Aethalometer AE33  $C$  values, a recent work of Valentini et al. (2021) reported a  $C$  value (at 660 nm) of 2.66 (on M8060 filter material) using the PaM reference, while Yus-Díez et al. (2021) investigated the  $C_\lambda$  dependency, finding it statistically significant for aged aerosol in the range of 3.47-4.03.

Finally, Rajesh and Ramachandran (2018) recently investigated the aerosol absorption coefficients obtained from AE31 and AE33 using  $C=2.14$  and  $C=1.57$  (Weingartner et al., 2003 and Drinovec et al., 2015), the aerosol SSA, and their role on the radiative forcing. They showed an average difference of 28% between AE31 and AE33 absorption coefficients. These differences, when not corrected (as instead done in the present work) can lead to an aerosol radiative forcing perturbation changing the estimation of the atmospheric warming up to 25%.

As a final conclusion, these results and the comparison with literature ones point towards the direction of both 1) using of multi- $\lambda$  regression fitting and MWAA approaches (limited to BC) to improve the HR apportionment and 2) exploring on a global scale the wavelength dependent multiple scattering correction factors ( $C_\lambda$ ) and their effect on climate forcing apportionment determination.

### 3.4 The temporal variation of $C$ and the heating rate

#### 3.4.1 $C_t$ features

As highlighted in section 3.2, there was a  $\sim 4\%$  difference between HRs determined using  $C_{fix}$  values obtained on the whole dataset and on daytime data only. To further investigate this difference, temporally-resolved multiple scattering coefficients  $C_t$  were obtained for each hour of day by comparing the AE31 and AE33 data ( $b_{ATN\_LC}$ ) at 660 nm with the absorption coefficient ( $b_{abs}$ ) measured by the MAAP (rescaled to the PP reference as detailed in section 2.3, eq. 14). This step was performed for consistency with the  $C_{fix}$  values chosen as reference in section 3.2.

Figure 8 shows the daily behavior of the multiple scattering parameters  $C_t$  for AE31 and AE33. Daily trends of  $C_t$  for both AE31 and AE33 were characterized by two peaks, one in the morning (6:00-8:00

706 LST) and one in the afternoon (17:00-19:00 LST) roughly in correspondence of the rush hours.  $C_t$  of  
 707 AE31 presented a maximum variation of 9% (range: 4.21 - 4.59), while for the  $C_t$  of AE33 is of 6%  
 708 constrained in a narrower range: 3.31- 3.51. The reason behind the presence of the two peaks (e.g. change  
 709 in chemical composition) is beyond the aim of the present paper; the possible artifacts related to negligible  
 710 relative humidity variations were avoided due to the presence of dryers which turned into an absence of  
 711 relationship with the daily variation of ambient RH (Figure S3). An insight into the possible explanation  
 712 of the  $C_t$  daily trend comes from the observation of the trend of N/eBC ratio. Results demonstrated that  
 713 also the N/eBC ratio experienced the same two peaks, the one in the morning (6:00-8:00 LST) and the  
 714 one in the afternoon (17:00-19:00 LST):  $5.78 \pm 0.14$  and  $6.11 \pm 0.12 \text{ cm}^{-3} \text{ ng}^{-1} \text{ m}^3$ , respectively. This  
 715 suggests that a higher fraction of scattering aerosol (primary or secondary) in the atmosphere can be  
 716 responsible of the observed correction of the multiple scattering parameter at the same time-slots.

717

### 718 **3.4.2 Total HR and its apportionment considering C temporal variability**

719 In order to obtain a proper time-dependent investigation of the total HR and its apportionment (FF, BB,  
 720 BC and BrC) the time-independent  $C_\lambda$  values were used to generate a  $C_{\lambda\_t}$  correction by following the  
 721 relative trend of  $C_t$  in Figure 8 (as in section 3.3 it has been demonstrated the need for a wavelength  
 722 dependent multiple scattering correction factors in order to obtain a consistent source/species  
 723 apportionment of LAA using the Aethalometer model and the Tian et al., 2019 method). Moreover, the  
 724  $C_t$  values were applied for using the Multi- $\lambda$  fit and the MWAA methods.

725 Figures 9a-l resumes the obtained HR and its daily average source and speciation apportionment  
 726 considering different C calculation approaches for both AE31 and AE33; Figure 10a-d reports the linear  
 727 correlation (AE33 vs. AE31) of the HR daily apportionment (FF, BB, BC and BrC) by using the  
 728 Aethalometer model and the Tian et al. (2019) method on the basis of both time-dependent  $C_t$  and  $C_{\lambda\_t}$   
 729 correction and by using the Multi- $\lambda$  fit and the MWAA methods on the basis of the time-dependent  $C_t$ .

730 Figures 9a and 9f show that for both AE31 and AE33 total HR, the interaction between eBC concentration  
 731 and the radiation intensity (Figures 4a-b) resulted in an asymmetric temporal behavior (with respect to

732 midday) characterized by a HR peak between the rush hour and midday ( $\sim 10$  LST;  $1.04 \pm 0.02$  and  
 733  $1.01 \pm 0.02$  K day<sup>-1</sup> peaks for AE31 and AE33, respectively) followed by a constant decrease until sunset.  
 734 Such behavior appears very important for understanding atmospheric feedbacks induced by the HR, e.g.  
 735 the influence on the planetary boundary layer evolution during the morning and on the aerosol liquid  
 736 water content (when morning fog events may dissipate). This asymmetric temporal behavior of total HR  
 737 is due to the synergy between the eBC concentration featuring a morning peak (Figure 4a) and  $F_{glo}$   
 738 showing larger values at midday (Figure 4b). Indeed, the amount of available radiation increased in the  
 739 late morning and early afternoon while the eBC concentrations decreased, leading to a compensatory  
 740 effect. It noteworthy that the  $F_{glo}$  radiation curve peaked when  $C_t$  reached minimum values. Thus, the  
 741 impact of  $C_t$  variability (9% and 6% for AE31 and AE33, respectively) on total HRs and their  
 742 apportionment from both instruments is here below investigated.

743 In this respect, Figures 9a-l show that the daily cycle of total HR average values (together with HR<sub>FF</sub>,  
 744 HR<sub>BB</sub>, HR<sub>BC</sub> and HR<sub>BrC</sub>) did not statistically differ at 95% probability when the AE31/AE33 instrument  
 745 and the apportionment method are fixed. In other words, the daily trend is unaffected by the use of fixed  
 746 or time-varying C values. In addition to this, the linear correlations reported in Figure 10a-b demonstrate,  
 747 even along the diurnal cycle, that the HR<sub>FF</sub> and HR<sub>BB</sub> reached the maximum agreement between the two  
 748 Aethalometers (slopes close to 1,  $R^2 > 0.98$ ) only when a wavelength- and time-dependent multiple  
 749 scattering parameters  $C_{\lambda,t}$  was used for the application of the “Aethalometer model” or when  $C_t$  is applied  
 750 together with the “multi- $\lambda$  fit” apportionment method, in keeping with Figure 7 and Table 1. Similarly,  
 751 Figure 10c-d demonstrate, even along the diurnal cycle, that the HR<sub>BC</sub> and HR<sub>BrC</sub> reached the maximum  
 752 agreement (slopes close to 1,  $R^2 > 0.98$ ) only when wavelength- and time-dependent multiple scattering  
 753 parameters  $C_{\lambda,t}$  are used with the Tian et al., 2019 model, in keeping with Figure 7 and Table 1. Instead,  
 754 the MWAA model application (Figure 10c-d) generates a discrepancy up to 17% considering the BC and  
 755 of about 50% for BrC, as already discussed in section 3.4. As a final remark, it noteworthy that the  $C_t$   
 756 impact on the HR should be investigated in the future in different locations, especially far from

757 midlatitudes (e.g. northern Europe, Arctic and Antarctica), where the diurnal cycle of sunshine and  
758 anthropogenic/natural LAA source can be very different.

759

## 760 **Conclusions**

761 Experimental heating rate (HR) values were obtained in the Po Valley (Milan), in the framework of the  
762 EMEP/ACTRIS/COLOSSAL 2018 winter campaign, using the Aethalometers AE31 and AE33. Source  
763 (fossil fuel, FF and biomass burning, BB) and speciation (black and brown carbon, BC and BrC)  
764 apportionment models were applied to derive the corresponding heating rates:  $HR_{FF}$ ,  $HR_{BB}$ ,  $HR_{BC}$  and  
765  $HR_{BrC}$ .

766 The effect of different wavelength- and time-dependent multiple scattering parameters ( $C$ ) on the HRs  
767 was investigated using different filter photometers (PP\_UniMI and MAAP) as reference.

768 As a conclusion, the best practices for accurate and temporally consistent HR data (from Aethalometers  
769 AE31/AE33) are resumed here below:

770 1) The use of single wavelength- and time-independent value of  $C$  (i.e.  $C_{fix}$  at 660 nm) results in temporally  
771 consistent total HR data between AE31 and AE33 version. However, this correction approach also results  
772 in statistically significant biases in the source (FF, BB) and speciation (BC and BrC) apportionment of  
773 the HR between the AE31 and the AE33, when traditional apportionment methods (“Aethalometer model”  
774 and Tian et al., 2019 model) are used. If not accounted for, these biases could give a false impression of  
775 ambient changes in terms of climate forcing due to different sources or nature of LAA when long-term  
776 data series in which AE33 supplanted AE31 are analysed. To overcome this problem, whenever further  
777 investigation on the  $C$  wavelength-dependence is not possible, advanced source and species  
778 apportionment methods (Multi- $\lambda$  fit or MWAA) have to be applied to obtain more consistent results, with  
779 limitation with respect to the  $HR_{BrC}$  from MWAA.

780 2) The use of a wavelength dependent multiple scattering correction parameters  $C_\lambda$  with traditional  
781 apportionments methods overcome all the aforementioned problems, generating a fully consistent dataset  
782 of HR (together with  $HR_{FF}$ ,  $HR_{BB}$ ,  $HR_{BC}$  and  $HR_{BrC}$ ). Nevertheless,  $C_\lambda$  determination requires the

783 availability multi-wavelength absorption measurements by a reference instrument in the season and  
 784 sampling site of interest.

785 3) The use of a time dependent multiple scattering correction parameter  $C_t$ , at the latitude of Milan, has a  
 786 negligible impact on the total HR, and its source/species apportionment.

787

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## 799 **Appendix A: Nomenclature**

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### **Nomenclature**

#### *Aerosol Acronyms*

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AAE	Absorption Angstrom Exponent
AAE <sub>BC</sub>	Absorption Angstrom Exponent of Black Carbon
AAE <sub>BrC</sub>	Absorption Angstrom Exponent of Brown Carbon
ATN	Attenuation
$b_{abs}(\lambda)$	wavelength dependent aerosol absorption coefficient ( $Mm^{-1}$ )
$b_{abs\_ref}(\lambda)$	wavelength dependent aerosol absorption coefficient from reference system ( $Mm^{-1}$ )
$b_{ATN}(\lambda)$	wavelength dependent aerosol attenuation coefficient ( $Mm^{-1}$ )
$b_{ATN\_LC}(\lambda)$	wavelength dependent aerosol attenuation coefficient corrected for loading ( $Mm^{-1}$ )
$C_{fix}$	multiple scattering correction parameter at 660 nm, time independent
$C_\lambda$	wavelength dependent multiple scattering correction parameter, time independent
$C_t$	time dependent multiple scattering correction parameter at 660 nm



$C_{\lambda,t}$	wavelength and time dependent multiple scattering correction parameter
BB	Biomass Burning
BC	Black Carbon
BrC	Brown Carbon
eBC	equivalent Black Carbon concentration ( $\mu\text{g m}^{-3}$ ) (Petzold et al., 2013)
FF	Fossil Fuel
LAA	Light Absorbing Aerosol
HR	Heating Rate ( $\text{K day}^{-1}$ )
$\text{HR}_{\text{BB}}$	Heating Rate of Biomass Burning ( $\text{K day}^{-1}$ )
$\text{HR}_{\text{BC}}$	Heating Rate of Black Carbon ( $\text{K day}^{-1}$ )
$\text{HR}_{\text{BrC}}$	Heating Rate of Brown Carbon ( $\text{K day}^{-1}$ )
$\text{HR}_{\text{FF}}$	Heating Rate of Fossil Fuel ( $\text{K day}^{-1}$ )
N	Total particle number concentration $d > 4 \text{ nm}$ ( $\text{cm}^{-3}$ )
PaM	Photometer as MAAP
PP	Polar Photometer
R(ATN)	Loading correction parameter
SSA	Single Scattering Albedo

### Other

### Symbols/Acronyms

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$\alpha$	Arnott (2005) scattering parameter
$\phi$	Azimuth angle (rad)
$\lambda$	Wavelength (nm)
$\rho$	Air Density ( $\text{kg m}^{-3}$ )
$\theta$	Zenith angle (rad)
$\theta_z$	Solar zenith angle (rad)
$C_p$	Isobaric specific heat of dry air ( $1005 \text{ J kg}^{-1} \text{ K}^{-1}$ )
dif	diffuse
dir	direct
$F_{\text{glo}}$	Global broadband irradiance; $F_{\text{glo}} = F_{\text{dir}} + F_{\text{dif}}$ ( $\text{W m}^{-2}$ )
$F_{\text{dif}}$	Diffuse broadband irradiance ( $\text{W m}^{-2}$ )
$F_{\text{dir}}$	Direct broadband irradiance ( $\text{W m}^{-2}$ )
$F_{\text{ref}}$	Reflected broadband irradiance ( $\text{W m}^{-2}$ )
$F_{\text{dir,dif,ref}}(\lambda)$	Spectral irradiance in function of $\lambda$ ( $\text{W m}^{-2} \text{ nm}^{-1}$ )
ref	reflected

800

801

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1110

1111 **Figure captions**

1112 Figure 1. Flow chart of the research methodology.

1113

1114 Figure 2. High-time resolution time series of data for the whole EMEP/ACTRIS/COLOSSAL 2018  
1115 campaign for equivalent black carbon (eBC, from AE31 and AE33), together with wind speed,  
1116 temperature, relative humidity, PM<sub>2.5</sub> and total particle concentrations (N).

1117

1118 Figure 3. AE33 vs. AE33 eBC concentrations in box-plot (a) and linear correlation (b).

1119

1120 Figure 4. Daily trend of eBC, total particle number (N) and PM<sub>2.5</sub> concentrations (a); daily trend of wind  
1121 speed (WS), temperature (T) and global radiation (F<sub>glo</sub>) values (b). The shaded area correspond to the  
1122 confidence interval at 95%.

1123

1124 Figure 5. a) HR determined by using different  $C_{fix}$  values obtained from different reference (MAAP, PP  
1125 all, PP day, PaM all, PaM day) for both AE31 and AE33; b) HR apportioned between Fossil Fuel (FF),  
1126 Biomass Burning (BB), BC and BrC by applying the “Aethalometer model” and the Tian et al. (2019)  
1127 method under the PP day  $C_{fix}$  (4.44 for AE31 and 3.43 for AE33) application to the whole  $b_{abs}(\lambda)$  spectrum.  
1128 The error bars represents the confidence interval at 95%.

1129

1130 Figure 6. Absorption coefficients ( $b_{abs}(\lambda)$ ) for both AE31 and AE33 for the whole spectrum and its  
1131 apportionment (FF, BB, BC and BrC) by using the Aethalometer model and the Tian et al. (2019) method  
1132 on the basis of a single  $C_{fix}$  correction (panels from a to e); absorption coefficients ( $b_{abs}(\lambda)$ ) for both AE31

1133 and AE33 for the whole spectrum and its apportionment (FF, BB, BC and BrC) by using the Aethalometer  
1134 model and the Tian et al. (2019) method on the basis of multi-wavelength  $C_\lambda$  correction (panels from f to  
1135 l). The shaded area correspond to the confidence interval at 95%.

1136

1137 Figure 7. Total and apportioned HR obtained by applying different apportionment methods using the most  
1138 suitable multiple scattering correction factor (C) for each method. The error bars represents the confidence  
1139 interval at 95%.

1140

1141 Figure 8.  $C_t$  values at 660 nm obtained from MAAP reference once reported at equivalent PP ( $b_{abs(PP}$   
1142  $_{660})=0.928*b_{abs(MAAP,660)}-2.07$ ; eq. 14, section 2.3) for both AE31 and AE33. The shaded area correspond  
1143 to the confidence interval at 95%.

1144

1145 Figure 9. Daily trend of heating rate (HR) data for both AE31 (panels a-e, left side) and AE33 (panels f-  
1146 l, right side) and its apportionment (FF, BB, BC and BrC) by using the Aethalometer model and the Tian  
1147 et al. (2019) method on the basis of a time-independent  $C_\lambda$  correction or a temporal  $C_{\lambda\_t}$  correction; the  
1148 same apportionment (FF, BB, BC and BrC) is reported by using the Multi- $\lambda$  fit and the MWAA methods  
1149 on the basis of a time-independent PP daytime  $C_{fix}$  (4.44 for AE31 and 3.43 for AE33) correction or a  
1150 temporal  $C_t$ . The shaded area correspond to the confidence interval at 95%.

1151

1152 Figure 10. Linear correlation (AE33 vs. AE31) of the heating rate (HR) daily trend data, apportionment  
1153 (FF, BB, BC and BrC) by using the Aethalometer model and the Tian et al. (2019) method on the basis  
1154 of a time-dependent  $C_t$  and  $C_{\lambda\_t}$  correction; the same apportionment (FF, BB, BC and BrC) is reported by  
1155 using the Multi- $\lambda$  fit and the MWAA methods on the basis of a time-dependent PP  $C_t$ .

1156

1157 **Table captions**

1158 Table 1. Multiple scattering correction factors  $C$  ( $C_{fix}$  and  $C_{\lambda}$ ) with the corresponding HR values ( $K \text{ day}^{-1}$ )  $\pm$  the confidence interval at 95% apportioned with respect to sources (FF and BB) and species (BC and BrC) of LAA by using different models (Aethalometer model, Multi- $\lambda$  fit, Tian et al., (2019), MWAA).

1159

1160

1161 The values that ensure a full consistency in both the total HR and its apportionment between AE31 and AE33 are highlighted in the table in bold.

1162

1163

1164 Table 2. Multiple scattering  $C_{\lambda}$  values for all the AE31 and AE33 wavelengths (from 370 to 950 nm) during diurnal time for radiative forcing applications.

1165

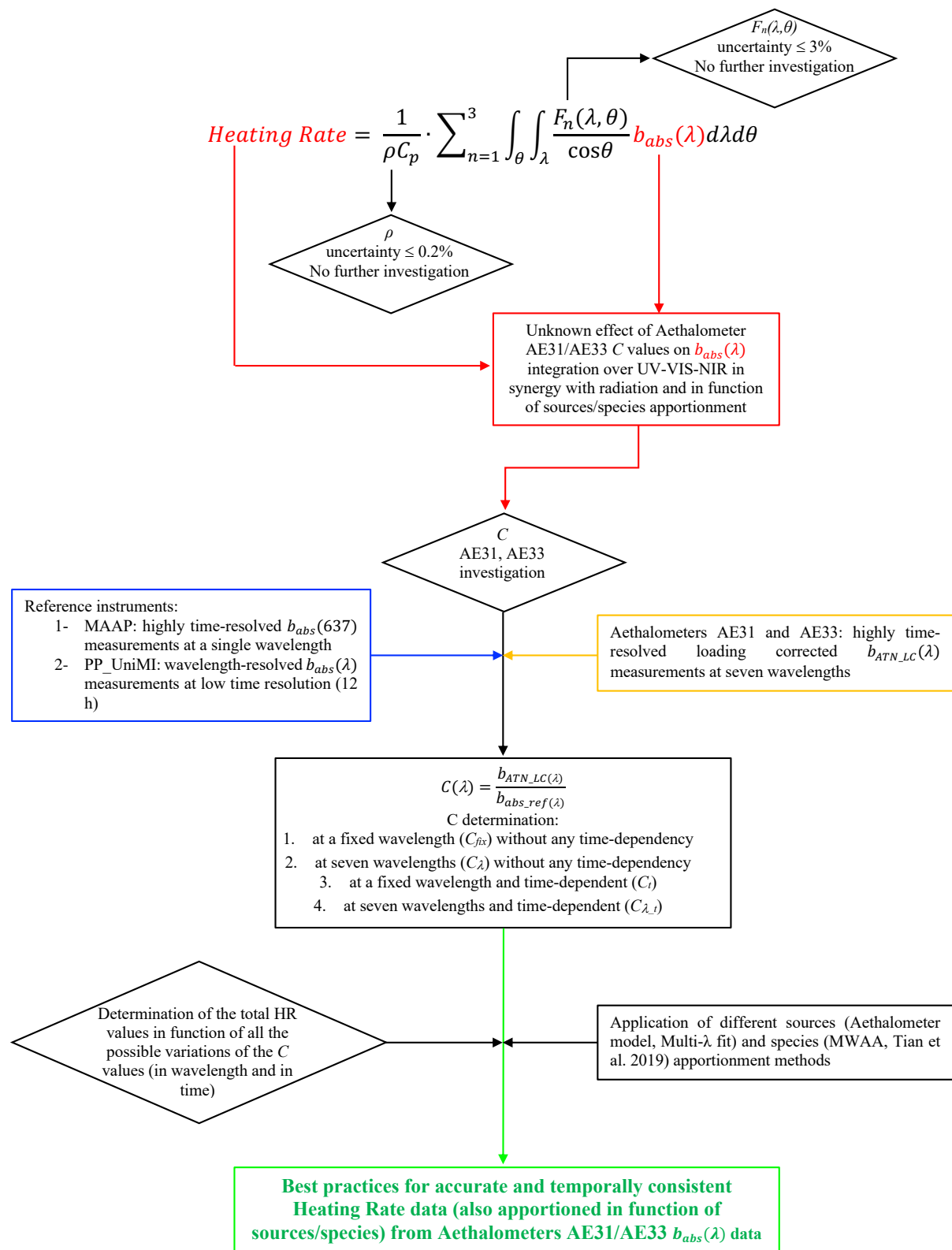


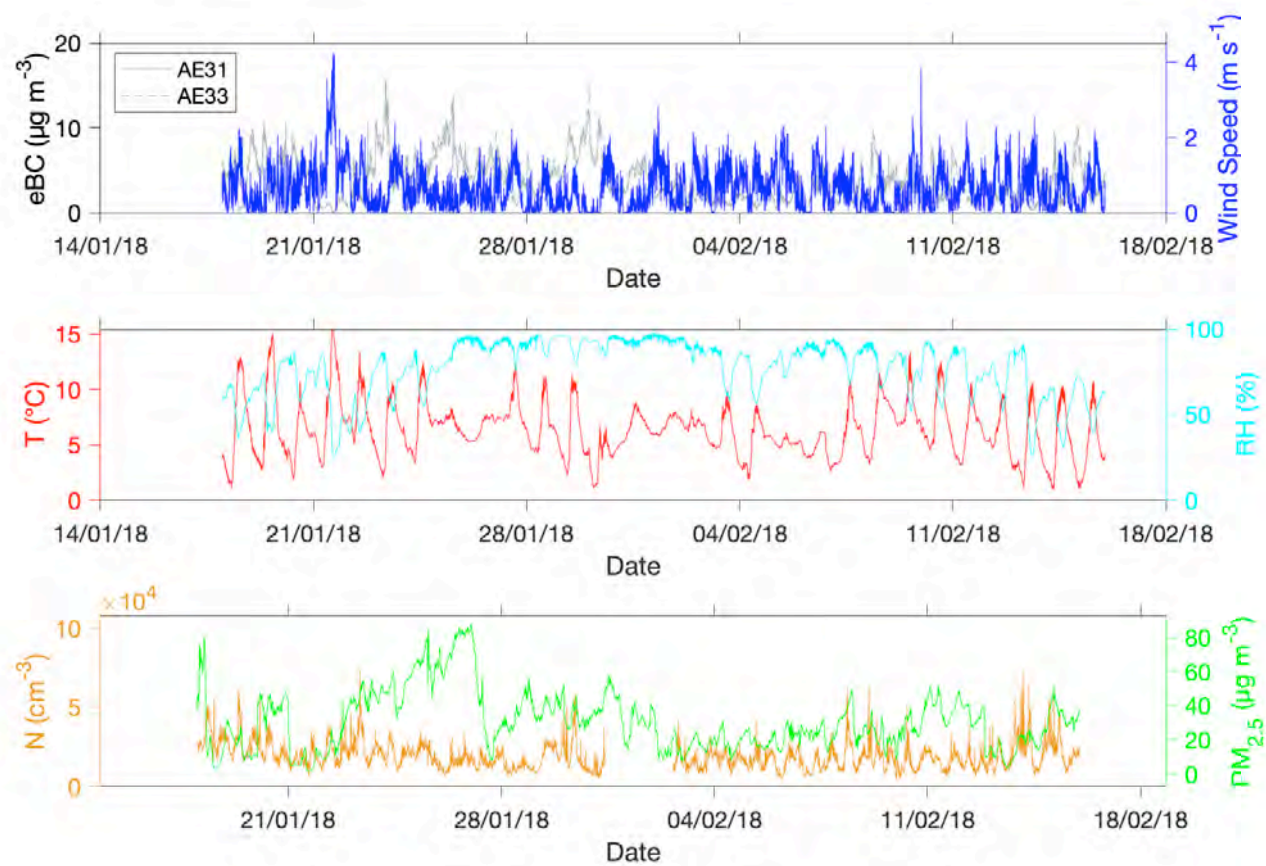
Referenc e	C determination	C values		Aethalometer Model		Multi-λ fit		Tian et al. (2019)		MWAA	
			HR	HR <sub>FF</sub>	HR <sub>BB</sub>	HR <sub>FF</sub>	HR <sub>BB</sub>	HR <sub>BC</sub>	HR <sub>BrC</sub>	HR <sub>BC</sub>	HR <sub>BrC</sub>
AE31	Weingartner et al. 2003	2.14	1.21±0.0 5	0.75±0.0 3	0.49±0.0 2	0.81±0.0 4	0.4±0.02	1.06±0.0 4	0.19±0.0 4	0.94±0.0 4	0.29±0.0 2
	MAAP high res (660 nm)	3.43±0.0 1	0.76±0.0 3	0.47±0.0 2	0.31±0.0 1	0.5±0.02	0.25±0.0 1	0.66±0.0 3	0.12±0.0 2	0.59±0.0 3	0.18±0.0 1
	PP all (660 nm)	4.30±0.0 7	0.60±0.0 2	0.37±0.0 2	0.24±0.0 1	0.4±0.02	0.2±0.01	0.53±0.0 2	0.09±0.0 2	0.47±0.0 2	0.14±0.0 1
	PP day (660 nm)	4.44±0.1 1	0.58±0.0 2	0.36±0.0 2	0.24±0.0 1	<b>0.39±0.0 2</b>	<b>0.19±0.0 1</b>	0.51±0.0 2	0.09±0.0 2	0.46±0.0 2	0.14±0.0 1
	PaM all (660 nm)	3.56±0.0 4	0.73±0.0 3	0.45±0.0 2	0.30±0.0 1	0.49±0.0 2	0.24±0.0 1	0.64±0.0 3	0.11±0.0 2	0.57±0.0 2	0.17±0.0 1
	PaM day (660 nm)	3.65±0.0 7	0.71±0.0 3	0.44±0.0 2	0.29±0.0 1	0.47±0.0 2	0.23±0.0 1	0.62±0.0 3	0.11±0.0 2	0.55±0.0 2	0.17±0.0 1
	PP day (370- 950 nm; Table 2)	min- max: 4.27- 4.63	<b>0.59±0.0 2</b>	<b>0.40±0.0 2</b>	<b>0.19±0.0 1</b>	//	//	<b>0.53±0.0 2</b>	<b>0.07±0</b>	//	//
AE33	Drinovec et al. 2015	1.57	1.34±0.0 5	0.98±0.0 4	0.37±0.0 2	0.85±0.0 4	0.46±0.0 2	1.21±0.0 5	0.12±0.0 1	1.15±0.0 5	0.19±0.0 1
	MAAP high res (660 nm)	2.64±0.0 1	0.80±0.0 3	0.58±0.0 2	0.22±0.0 1	0.51±0.0 2	0.28±0.0 1	0.72±0.0 3	0.07±0	0.68±0.0 3	0.12±0
	PP all (660 nm)	3.37±0.0 5	0.62±0.0 2	0.46±0.0 2	0.17±0.0 1	0.4±0.02	0.22±0.0 1	0.56±0.0 2	0.06±0	0.54±0.0 2	0.09±0
	PP day (660 nm)	3.43±0.0 8	0.61±0.0 2	0.45±0.0 2	0.17±0.0 1	<b>0.39±0.0 2</b>	<b>0.21±0.0 1</b>	0.55±0.0 2	0.06±0	<b>0.53±0.0 2</b>	<b>0.09±0</b>
	PaM all (660 nm)	2.79±0.0 3	0.76±0.0 3	0.55±0.0 2	0.21±0.0 1	0.48±0.0 2	0.26±0.0 1	0.68±0.0 3	0.07±0	0.65±0.0 3	0.11±0
	PaM day (660 nm)	2.82±0.0 5	0.75±0.0 3	0.55±0.0 2	0.21±0.0 1	0.48±0.0 2	0.26±0.0 1	0.68±0.0 3	0.07±0	0.64±0.0 3	0.11±0

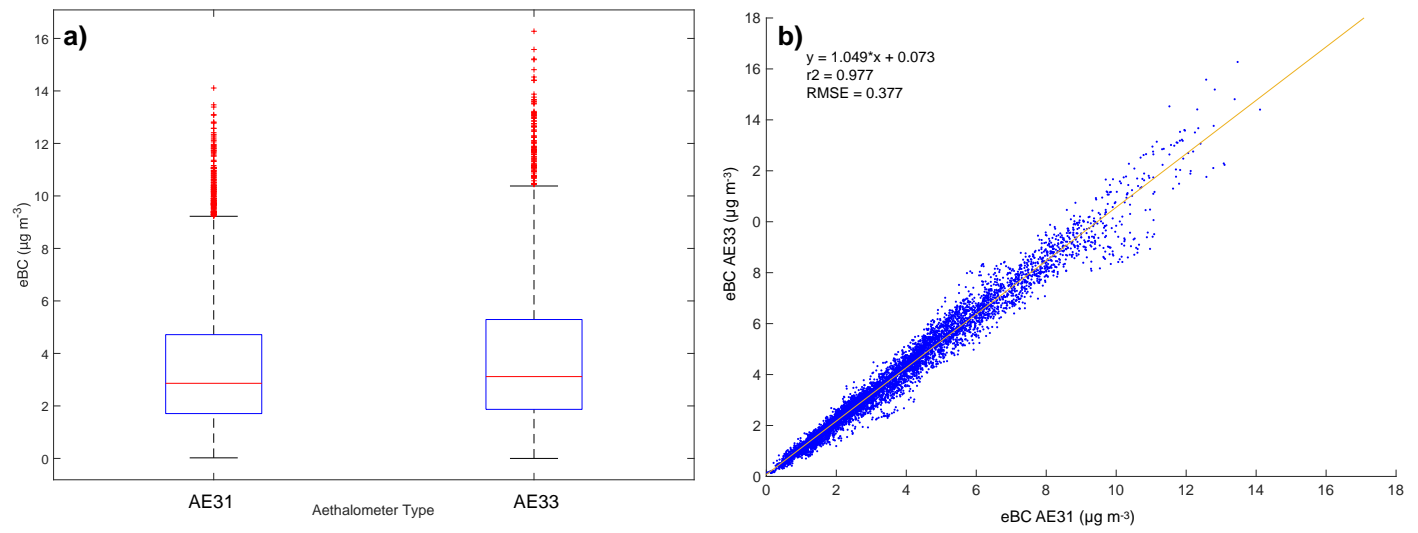
$\zeta$	PP day (370-950 nm; Table 2)	min-max: 3.41-3.78	<b>0.59±0.0</b> <b>2</b>	<b>0.42±0.0</b> <b>2</b>	<b>0.17±0.0</b> <b>1</b>	//	//	<b>0.54±0.0</b> <b>2</b>	<b>0.06±0</b>	//	//
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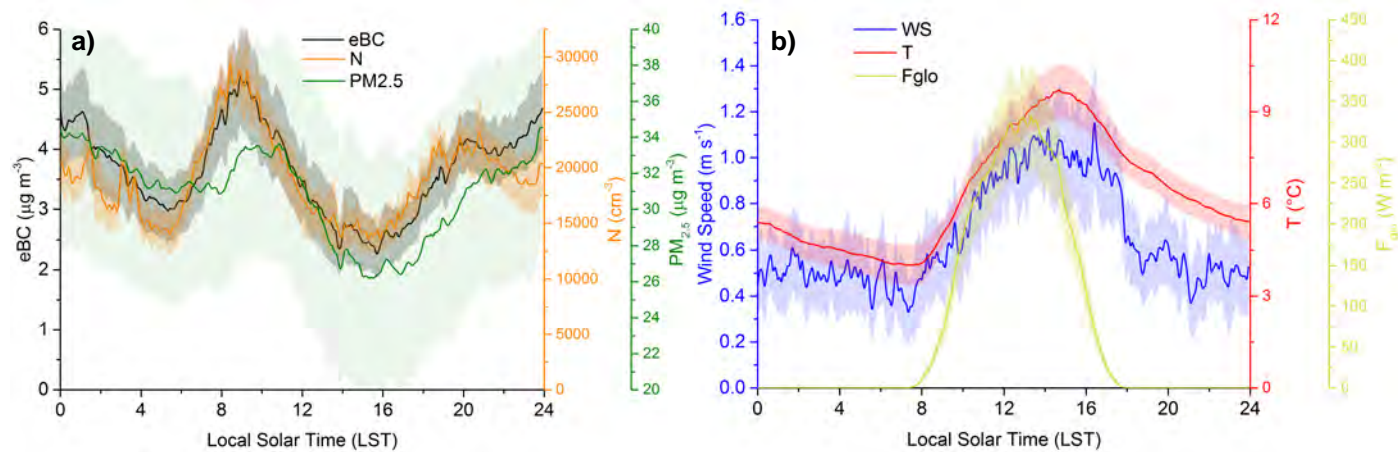
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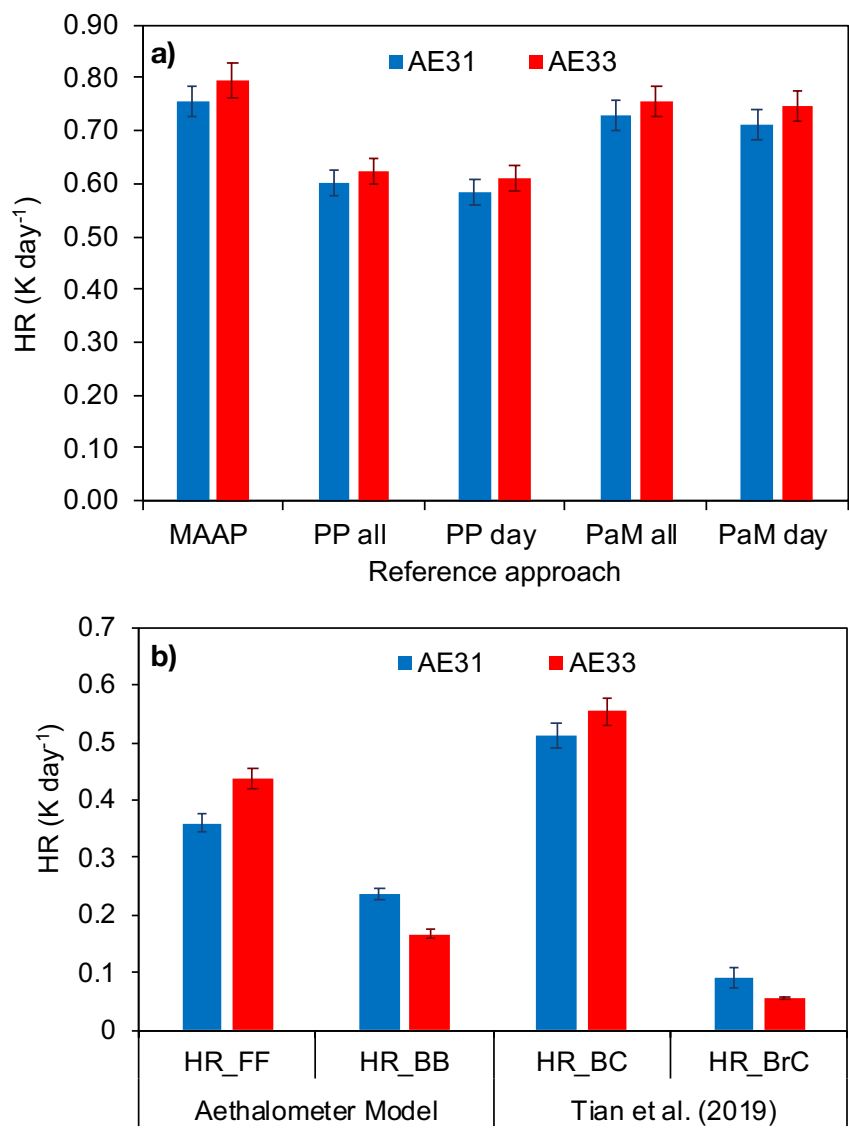
$\lambda$ (nm)	C daytime values	
	AE31	AE33
370	4.63	3.78
470	4.42	3.56
520	4.38	3.57
590	4.38	3.55
660	4.44	3.43
880	4.34	3.41
950	4.27	3.57



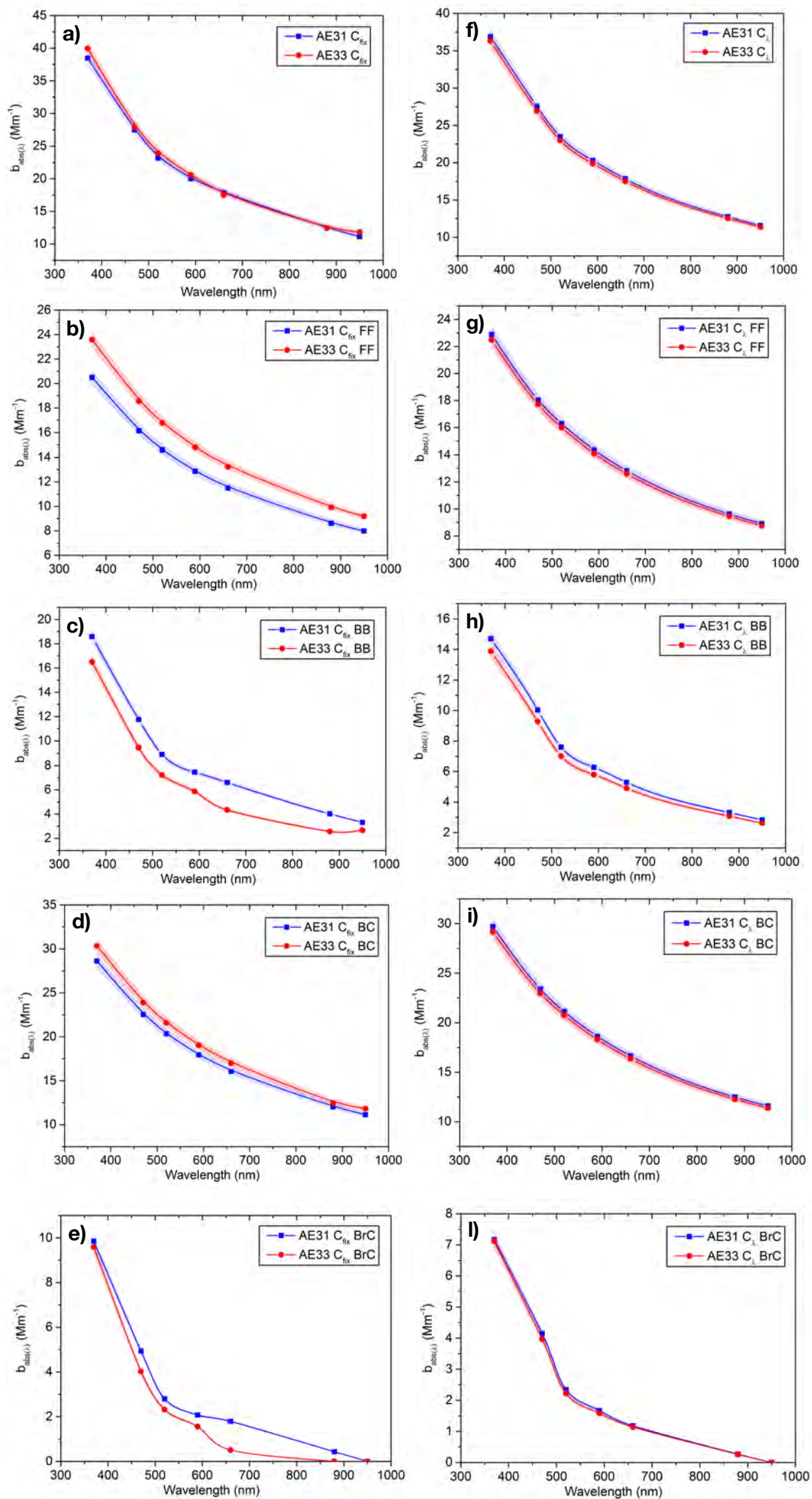


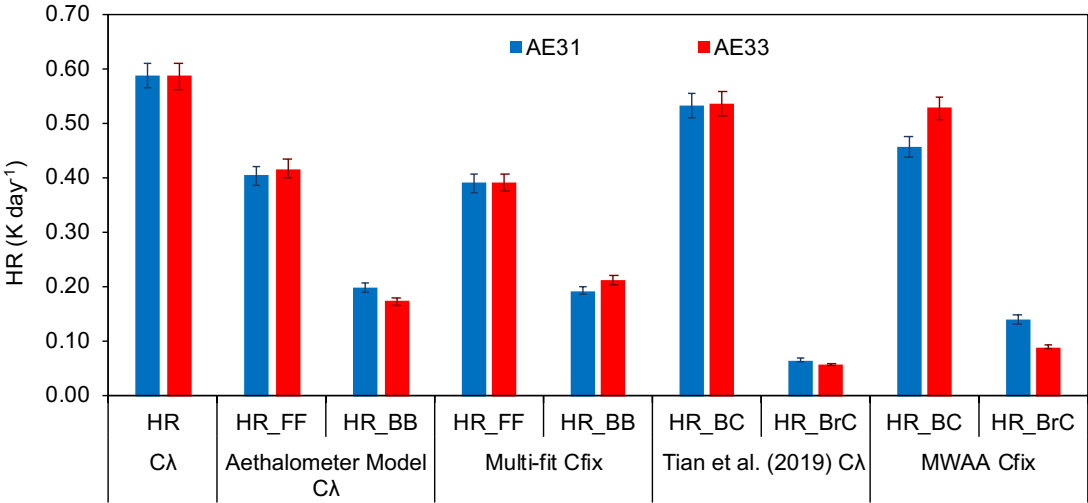


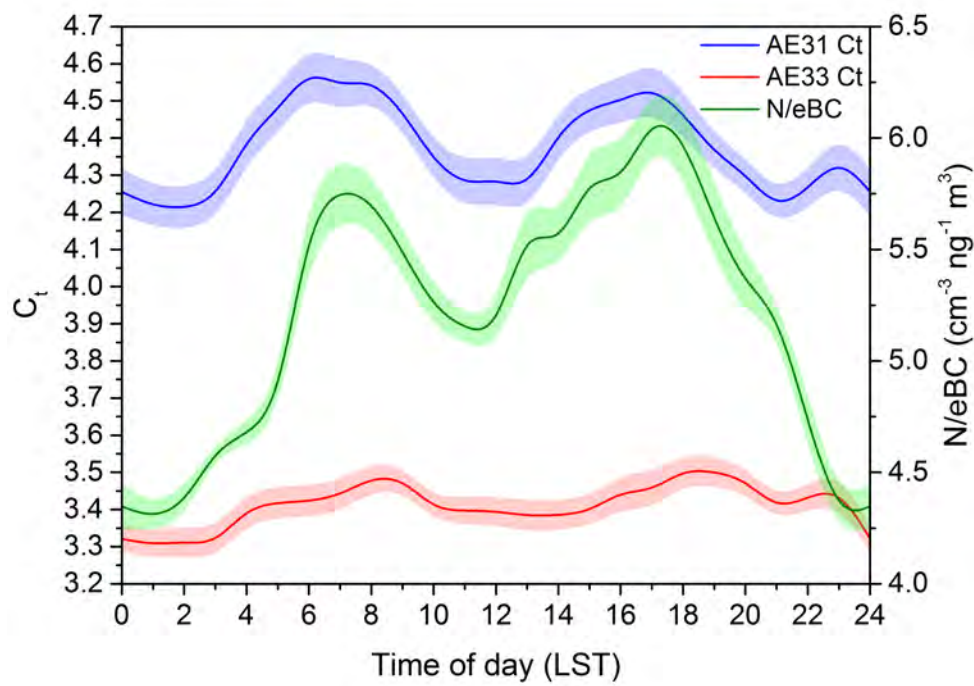


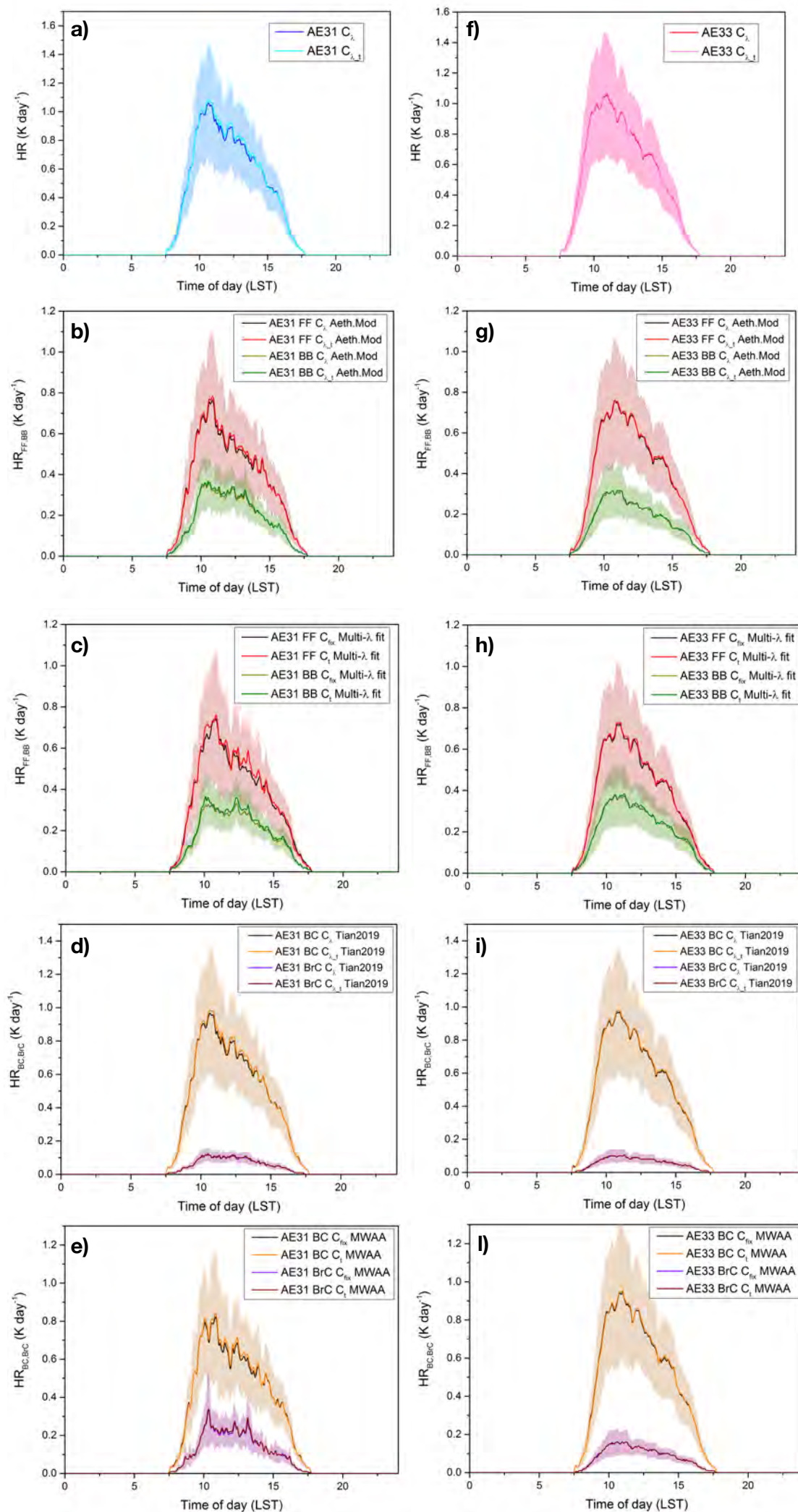




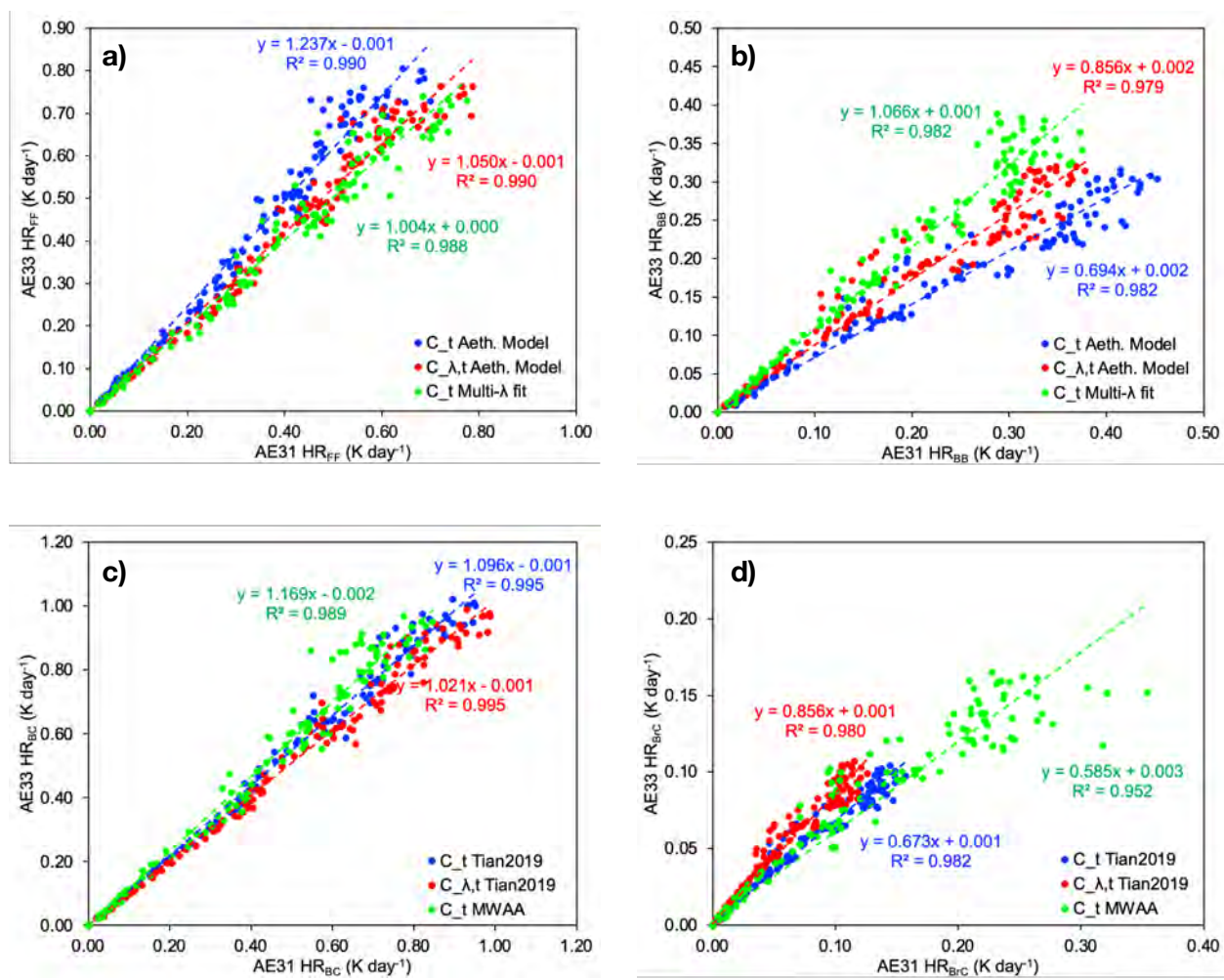


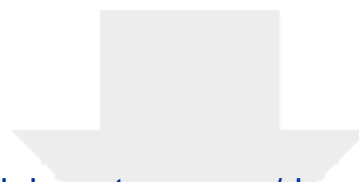






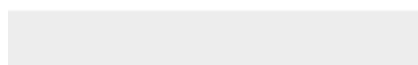
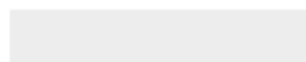






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## Author contribution

**Conceptualization:** L.F., V.B., E.B., G.V., and R.V. designed and organized the sampling campaign to meet the final goals. L.F. performed the data analysis for heating rate determination and its apportionment in function of the optimized multiple-scattering enhancement parameter and the multi- $\lambda$  fit approach applied the Aethalometer model. **Data curation:** L.F. and V.B. validated and assembled the final database. **Formal analysis:** L.F., V.B., L.S., S.C., G.M., A.B., N.L., S.V., G.V., D.M and P.P. collaborated to data analysis and reduction. **Methodology:** L.F., V.B., S.C., G.V., R.V. realized the sampling campaign. A.G., M.R., G.M. gave support for Aethalometers set-up. F.S., S.V. performed PP\_UniMI measurements. S.C. performed spectroradiometer measurements. **Software:** L.F. developed the software for the heating rate determination and apportionment in function of LAA sources and species. **Supervision:** E.B. and L.F. supervised all the scientific activity. **Writing – original draft:** L.F. and L.S. wrote the original draft. **Writing – review & editing:** all co-authors commented and contributed to the final version of the paper.

**Declaration of interests**

☒ The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

☐The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: