

LIOUVILLE-TYPE RESULTS IN EXTERIOR DOMAINS FOR RADIAL SOLUTIONS OF FULLY NONLINEAR EQUATIONS

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ABSTRACT. We give necessary and sufficient conditions for the existence of positive radial solutions for a class of fully nonlinear uniformly elliptic equations posed in the complement of a ball in \mathbb{R}^N , and equipped with homogeneous Dirichlet boundary conditions.

1. INTRODUCTION

Let B be any ball in \mathbb{R}^N and let $0 < \lambda \leq \Lambda$. The aim of this paper is to detect the optimal conditions on the exponent $p > 1$ for the existence of positive radial solutions of the fully nonlinear exterior Dirichlet problem

$$\begin{cases} -\mathcal{F}(D^2u) = u^p & \text{in } \mathbb{R}^N \setminus \bar{B}, \\ u = 0 & \text{on } \partial B, \end{cases} \quad (1.1)$$

where \mathcal{F} is either one of the Pucci's extremal operators $\mathcal{M}_{\lambda,\Lambda}^\pm$, defined respectively as

$$\begin{aligned} \mathcal{M}_{\lambda,\Lambda}^-(X) &:= \inf_{\lambda I \leq A \leq \Lambda I} \operatorname{tr}(AX) = \lambda \sum_{\mu_i \geq 0} \mu_i + \Lambda \sum_{\mu_i < 0} \mu_i \\ \mathcal{M}_{\lambda,\Lambda}^+(X) &:= \sup_{\lambda I \leq A \leq \Lambda I} \operatorname{tr}(AX) = \Lambda \sum_{\mu_i \geq 0} \mu_i + \lambda \sum_{\mu_i < 0} \mu_i \end{aligned}$$

μ_1, \dots, μ_N being the eigenvalues of any squared symmetric matrix X .

Pucci's extremal operators are the prototype of fully nonlinear uniformly elliptic operators. Acting as barriers in the whole class of operators with fixed ellipticity constants $\lambda \leq \Lambda$, they play a crucial role in the regularity theory for fully nonlinear elliptic equations, see [4]. Moreover, as sup/inf of linear operators, they frequently arise in the equations satisfied by the value function associated with stochastic optimal control problems, see e.g. [11, 18], with special application to mathematical finance problems.

We recall that if $\Lambda = \lambda$, both Pucci's operators reduce, up to a multiplicative factor, to the Laplace operator. Thus, they may be considered also as perturbations of the standard Laplacian, and the well-known Lane-Emden-Fowler equation $-\Delta u = u^p$ is included as a very special case of the problems we are considering.

In the semilinear case, the Lane-Emden-Fowler equation has been largely studied. The well known existence results, exhibiting the critical Sobolev exponent $p^* = \frac{N+2}{N-2}$ as threshold for the existence of entire solutions or solutions in bounded domains, are intimately related to the (lack of) compactness properties of the Sobolev embeddings. Moreover, the entire solutions existing in the critical case $p = p^*$ realize the best constant in the Sobolev inequality, see [19]. In other

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words, the critical nature of the exponent p^* may be largely interpreted in view of structural properties of the functional setting behind the equation.

As soon as $\Lambda > \lambda$, Pucci's operators loose the linear and variational structures. Nevertheless, Lane-Emden-Fowler type equations as (1.1) have been studied and, at least in the radial setting, some critical exponents p_{\pm}^* acting as thresholds for the existence of entire radial solutions or solutions in balls have been proved to exist, see [10]. In the fully nonlinear radial setting, the exponents p_{\pm}^* play the same role as the critical Sobolev exponent $p^* = \frac{N+2}{N-2}$ for the Laplacian. Though their appearance is motivated exclusively as threshold for the existence of entire radial solutions or solutions in balls, the recent results of [3], where some weighted energies associated with radial solutions of (1.1) are introduced and proved to be asymptotically preserved by almost critical solutions, suggest that the critical exponents reflect some intrinsic properties of the operators, maybe beyond the radial setting.

The Liouville-type results obtained in the present paper may be regarded as a further justification of the critical character of the exponents p_{\pm}^* , since they are proved to act as thresholds also for the existence of radial solutions to exterior Dirichlet problems. We observe that this result, while well expected for semilinear equations due to the duality between Dirichlet problems in balls and Dirichlet problems in exterior domains (via the Kelvin transform), in the fully non linear case it requires a direct proof. Moreover, the existence of positive solutions for exterior Dirichlet problems is also closely related to the existence or non-existence of sign-changing radial solutions in balls or in the whole space (see also Remark 5.2 in [14]). In particular, in the recent paper [13], the asymptotic behavior of sign changing solutions in balls is studied, and the obtained results strongly rely on the present theorems for exterior Dirichlet problems.

In order to describe the results, let us introduce the dimension like parameters

$$\begin{aligned}\tilde{N}_- &:= \frac{\Lambda}{\lambda}(N-1) + 1 && \text{for } \mathcal{M}_{\lambda,\Lambda}^-, \\ \tilde{N}_+ &:= \frac{\lambda}{\Lambda}(N-1) + 1 && \text{for } \mathcal{M}_{\lambda,\Lambda}^+, \end{aligned}$$

which have been proved in some previously studied cases to play a key role in existence results. Let us emphasize that one has always $\tilde{N}_+ \leq N \leq \tilde{N}_-$, and the equalities hold true if and only if $\Lambda = \lambda$.

The case of *entire supersolutions* has been considered in [5], where it has been proved that

$$\exists u > 0, -\mathcal{M}_{\lambda,\Lambda}^{\pm}(D^2u) \geq u^p \quad \text{in } \mathbb{R}^N \iff p > \frac{\tilde{N}_{\pm}}{\tilde{N}_{\pm} - 2},$$

meaning that, in particular, positive supersolutions never exist if $\tilde{N}_{\pm} \leq 2$. In the sequel, we will always assume that $\tilde{N}_{\pm} > 2$.

The same threshold has been proved in [2] to be optimal for the existence of *solutions in any exterior domain*, that is

$$\exists u > 0, -\mathcal{M}_{\lambda,\Lambda}^{\pm}(D^2u) = u^p \quad \text{in } \mathbb{R}^N \setminus K \iff p > \frac{\tilde{N}_{\pm}}{\tilde{N}_{\pm} - 2},$$

where $K \subset \mathbb{R}^N$ is any nonempty compact set. Let us also mention the results of [1], where more general nonlinearities $f(u)$ replacing u^p are considered. In the above results, supersolutions are meant in the viscosity sense and no symmetry property on u is required.

In the present paper we are concerned with *solutions* of the equation satisfying further *homogeneous Dirichlet boundary conditions*. By elliptic regularity theory, it is not restrictive to consider classical solutions of problem (1.1), that are C^2 functions satisfying pointwise the equation as well as the boundary condition. In its full generality, that is without assuming radial symmetry of the solutions, the problem is completely open, and also in the semilinear case (i.e.

when $\Lambda = \lambda$) few results are known for solutions of exterior Dirichlet problems, see e.g. [7] where solutions are constructed as perturbations of radial solutions. Our results, limited to radially symmetric solutions, may hopefully contribute to tackle the general fully nonlinear problem.

In the radial setting, the existence of *entire positive solutions* has been studied in [10], where it has been proved that there exist two critical exponents p_+^* and p_-^* associated with $\mathcal{M}_{\lambda,\Lambda}^+$ and $\mathcal{M}_{\lambda,\Lambda}^-$ respectively, such that

$$\exists u \text{ radial, } u > 0, \quad -\mathcal{M}_{\lambda,\Lambda}^\pm(D^2u) = u^p \quad \text{in } \mathbb{R}^N \iff p \geq p_\pm^*.$$

Unfortunately, the dependence of the radial critical exponents on the effective dimensions is not explicitly known. The radial critical exponents are proved in [10] to satisfy, when $\lambda < \Lambda$, the strict inequalities

$$\frac{\tilde{N}_- + 2}{\tilde{N}_- - 2} < p_-^* < \frac{N + 2}{N - 2} < p_+^* < \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}. \quad (1.2)$$

Note that inequalities in (1.2) become equalities when $\Lambda = \lambda$. Let us emphasize that inequalities (1.2) say that, with respect to the intrinsic dimensions, the exponent p_+^* is subcritical, whereas p_-^* is supercritical. Moreover, (1.2) show that the critical Sobolev exponent $p^* = \frac{N+2}{N-2}$ is not preserved even for small perturbations of the operator, since the critical exponents are instantaneously different as soon as $\Lambda > \lambda$. For an integral characterization of p_\pm^* , as well as for sharp estimates on entire critical solutions, we refer to [3].

Clearly, the analysis on the existence of entire positive radial solutions yields, as a by product, the dual result on the existence of positive solutions of Dirichlet problems in balls, namely

$$\exists u > 0, \quad -\mathcal{M}_{\lambda,\Lambda}^\pm(D^2u) = u^p \quad \text{in } B, \quad u = 0 \quad \text{on } \partial B \iff p < p_\pm^*.$$

Note that, in this case, the radial symmetry of the solutions is not a restriction, since, by [6], any positive solution in the ball is radial. The critical exponents p_\pm^* , therefore, give the optimal thresholds for the existence of positive solutions in balls and, as proved in [9], also in domains sufficiently close to balls.

On the other hand, as recently proved in [14], for annular domains (radial) solutions exist for any $p > 1$. More precisely, in [14] it has been proved that solutions of the initial value problems for the ODEs associated with the equations $-\mathcal{M}_{\lambda,\Lambda}^\pm(D^2u) = u^p$ give radial solutions in annular domains provided that they have sufficiently large initial slope. The results of the present paper, in a sense, complement the results of [14], since here we prove, in particular, that a sufficiently large initial velocity is needed for having radial solutions in annuli if and only if $p > p_\pm^*$.

The results of the subsequent sections are summarized in the following theorem.

Theorem 1.1. *There exist positive radial solutions of problem (1.1) if and only if $p > p_\pm^*$. Moreover, for any $p > p_\pm^*$, problem (1.1) has a unique positive radial solution u^* satisfying*

$$\lim_{r \rightarrow +\infty} r^{\tilde{N}_\pm - 2} u^*(r) = C > 0, \quad (1.3)$$

and infinitely many positive radial solutions u satisfying

$$\lim_{r \rightarrow +\infty} r^{\tilde{N}_\pm - 2} u(r) = +\infty. \quad (1.4)$$

Borrowing the terminology currently used, the solution u^* satisfying (1.3) will be referred to as the *fast decaying* solution. As far as solutions u satisfying (1.4) are concerned, they will be proved to satisfy either

$$\lim_{r \rightarrow +\infty} r^{\frac{2}{p-1}} u(r) = c > 0,$$

in which case they will be called *slow decaying* solutions, or

$$0 < \liminf_{r \rightarrow +\infty} r^{\frac{2}{p-1}} u(r) < \limsup_{r \rightarrow +\infty} r^{\frac{2}{p-1}} u(r) < +\infty,$$

in which case they will be named *pseudo-slow decaying* solutions (see [10]).

In the semilinear case, a proof of Theorem 1.1 can be found in [16]. For the fully nonlinear case, it will be a straightforward consequence of the results proved in the next sections, where we perform a careful analysis of the initial value problem for the ODE associated to radial solutions of problem (1.1). Recalling that the eigenvalues of the Hessian matrix D^2u of a smooth radial function $u = u(r)$ are nothing but $u''(r)$ and $\frac{u'(r)}{r}$ (with multiplicity at least $N - 1$), it is not difficult to write the ODE satisfied by a radial solution of problem (1.1). Nevertheless, since the coefficients of the operators $\mathcal{M}_{\lambda,\Lambda}^\pm$ depend on the sign of the eigenvalues of the Hessian matrix, we will obtain an ODE with discontinuous coefficients having jumps at the points where the solution u changes its monotonicity and/or concavity. This is a feature of the fully nonlinear problem which makes techniques previously developed for the semilinear case not directly applicable. In particular, the Kelvin transform which reduces a supercritical exterior Dirichlet problem to a subcritical Dirichlet problem in the punctured ball cannot be used.

We will essentially make use of the results of [10] for entire solutions, in particular of the fact that the critical exponents are the only exponents for which the entire solutions are fast decaying. Moreover, as in [10], we will take advantage of the Emden-Fowler transformed

$$x(t) = r^{\frac{2}{p-1}} u(r), \quad r = e^t,$$

which produces a new variable $x(t)$ satisfying an autonomous equation. Despite the fact that also the coefficients of the equation satisfied by x will have jumps at the points corresponding to the changes of monotonicity and concavity of u , the phase plane analysis of the trajectories associated to the solution x will be repeatedly used.

A particularly delicate step in the proof of Theorem 1.1 will be the proof of the non existence of solutions in the critical cases $p = p_\pm^*$, as well as of the uniqueness of the fast decaying solutions, which will be obtained by using different arguments for $\mathcal{M}_{\lambda,\Lambda}^+$ and $\mathcal{M}_{\lambda,\Lambda}^-$. For $\mathcal{M}_{\lambda,\Lambda}^-$, we heavily exploit the fact that $p_-^* > \frac{\tilde{N}_-+2}{\tilde{N}_--2}$ and the proof relies on some properties of the solutions of supercritical semilinear problems. For $\mathcal{M}_{\lambda,\Lambda}^+$, for which $p_+^* < \frac{\tilde{N}_++2}{\tilde{N}_+-2}$, a different proof will be obtained as an application of Gauss-Green Theorem in the phase plane.

The paper is organized as follows: in Section 2 we recall some basic properties of radial solutions of problem (1.1) and their Emden-Fowler transform, whereas the existence of positive radial solutions for p supercritical is proved in Section 3. Section 4 will be then devoted to the proof of nonexistence of nontrivial solutions when p is strictly subcritical, while the critical cases $p = p_+^*$ and $p = p_-^*$ are addressed respectively in Section 5 and Section 6. Finally, in Section 7, we complete the proof of Theorem 1.1 by showing that problem (1.1) has a unique fast decaying solution and infinitely many slow or pseudo-slow decaying solutions, according to the initial slope.

2. RADIAL SOLUTIONS OF PROBLEM (1.1) AND THE EMDEN-FOWLER TRANSFORM

In order to study the existence of solutions for problem (1.1), we can assume, without loss of generality, that B is the unit ball of \mathbb{R}^N centered at the origin, by the invariance of the equation with respect to translations and to the scaling $\tilde{u}(x) := \gamma^{\frac{2}{p-1}} u(\gamma x)$ for any $\gamma > 0$.

For any $\alpha > 0$, let us introduce the initial value problem

$$\begin{cases} u''(r) = M^+ \left(-\frac{\lambda(N-1)}{r} K^+(u') - u^p \right) & \text{for } r > 1, \\ u(r) > 0 & \text{for } r > 1, \\ u(1) = 0, \quad u'(1) = \alpha, \end{cases} \quad (2.1)$$

where

$$M^+(s) := \begin{cases} s/\Lambda & \text{if } s \geq 0, \\ s/\lambda & \text{if } s < 0, \end{cases} \quad K^+(s) := \begin{cases} \frac{\Lambda}{\lambda} s & \text{if } s \geq 0, \\ s & \text{if } s < 0. \end{cases}$$

A direct computation shows that u is a radial solution of problem (1.1) with $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$ if and only if u satisfies (2.1) in $(1, +\infty)$ for some $\alpha > 0$.

Analogously, a radial solution of (1.1) with $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ is nothing but a global solution of

$$\begin{cases} u''(r) = M^- \left(-\frac{\Lambda(N-1)}{r} K^-(u') - u^p \right) & \text{for } r > 1, \\ u(r) > 0 & \text{for } r > 1, \\ u(1) = 0, \quad u'(1) = \alpha, \end{cases} \quad (2.2)$$

with

$$M^-(s) := \begin{cases} s/\lambda & \text{if } s \geq 0, \\ s/\Lambda & \text{if } s < 0, \end{cases} \quad K^-(s) := \begin{cases} \frac{\lambda}{\Lambda} s & \text{if } s \geq 0, \\ s & \text{if } s < 0. \end{cases}$$

For the rest of the section, let us focus on the case $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$. The same observations can be made for $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$, by just exchanging Λ with λ .

Let us denote by u_α the unique positive maximal solution of (2.1), defined on a maximal interval $[1, \rho_\alpha)$, with $\rho_\alpha \leq +\infty$. Some general properties of the function u_α have been established in [14]. In particular (see [14, Lemma 2.1]), it is known that, for any $\alpha > 0$, there exists a unique $\tau_\alpha \in (1, \rho_\alpha)$ such that $u'_\alpha(r) > 0$ for $r \in [1, \tau_\alpha)$ and $u'_\alpha(r) < 0$ for $r \in (\tau_\alpha, \rho_\alpha)$. Moreover, if $\rho_\alpha < +\infty$ one has $u_\alpha(\rho_\alpha) = 0$, whereas if $\rho_\alpha = +\infty$, then $\lim_{r \rightarrow +\infty} u_\alpha(r) = 0$ (see [14, Corollary 2.3]). We remark that if $\rho_\alpha < +\infty$, then u_α is a radial solution of the Dirichlet problem

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 u) = u^p & \text{in } A_{1, \rho_\alpha}, \\ u = 0 & \text{on } \partial A_{1, \rho_\alpha}, \end{cases}$$

where A_{1, ρ_α} is the annular domain with radii $1 < \rho_\alpha$, and $u'_\alpha(\rho_\alpha) < 0$ by the Hopf boundary maximum principle. On the contrary, if $\rho_\alpha = +\infty$, then u_α is a positive solution in the exterior domain $\mathbb{R}^N \setminus \overline{B_1}$ and, by the results of [2], this necessarily implies that $p > \frac{\tilde{N}_+}{\tilde{N}_+ - 2}$.

We further notice that the function u_α , which satisfies $u''_\alpha(1) < 0$, cannot be globally (or even definitely) concave in $(1, \rho_\alpha)$, since otherwise $\rho_\alpha < +\infty$ and

$$u_\alpha^p(\rho_\alpha) = -\lambda(u''_\alpha(\rho_\alpha) + (N-1)u'_\alpha(\rho_\alpha)/\rho_\alpha) > 0.$$

We will denote by $\sigma_\alpha \in (\tau_\alpha, \rho_\alpha)$ the first zero of u''_α . We will show below that σ_α actually is the first point where the function u_α changes its concavity.

A useful tool for studying the problem is the so called Emden-Fowler transform (see [12, 10]), which reduces the initial problem (2.1) to an autonomous equation. Setting

$$x(t) = r^{\frac{2}{p-1}} u(r), \quad r = e^t, \quad (2.3)$$

a direct computation shows that

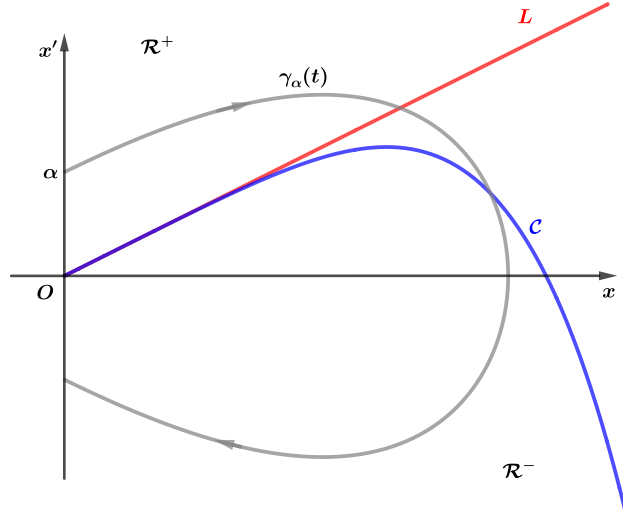
$$\begin{aligned} u' \geq 0 &\iff x' \geq \frac{2}{p-1}x, \\ u'' \leq 0 &\iff x' \geq \frac{2}{p-1}x - \frac{x^p}{\lambda(N-1)}. \end{aligned}$$

Hence, the new unknown x satisfies

$$\begin{cases} x'' = \tilde{a}_-x' + \tilde{b}_-x - \frac{x^p}{\lambda} & \text{if } x' \geq \frac{2}{p-1}x, \\ x'' = ax' + bx - \frac{x^p}{\lambda} & \text{if } \frac{2}{p-1}x \geq x' \geq \frac{2}{p-1}x - \frac{x^p}{\lambda(N-1)}, \\ x'' = \tilde{a}_+x' + \tilde{b}_+x - \frac{x^p}{\Lambda} & \text{if } x' \leq \frac{2}{p-1}x - \frac{x^p}{\lambda(N-1)}, \end{cases} \quad (2.4)$$

with coefficients defined respectively as

$$\begin{aligned} a &:= \frac{N+2-(N-2)p}{p-1}, & \tilde{a}_\pm &:= \frac{\tilde{N}_\pm + 2 - (\tilde{N}_\pm - 2)p}{p-1}, \\ b &:= \frac{2((N-2)p - N)}{(p-1)^2}, & \tilde{b}_\pm &:= \frac{2((\tilde{N}_\pm - 2)p - \tilde{N}_\pm)}{(p-1)^2}. \end{aligned} \quad (2.5)$$



Associated with a solution x of the above problem, one can consider the trajectory $\gamma(t) = (x(t), x'(t))$ in the phase plane. Thus, the coefficients of the autonomous equation satisfied by x are piecewise constant and have jumps whenever the trajectory γ crosses either the half-line $L = \left\{x' = \frac{2}{p-1}x\right\}$ or the curve $C = \left\{x' = \frac{2}{p-1}x - \frac{x^p}{\lambda(N-1)}\right\}$. We will denote by \mathcal{R}^+ and \mathcal{R}^- the open regions of the right half-plane lying respectively above and below the curve C .

In particular, the trajectory $\gamma_\alpha(t) = (x_\alpha(t), x'_\alpha(t))$ associated with the transform x_α of the solution u_α is defined on the interval $[0, \log \rho_\alpha)$, it lays in the right half plane for $t \in (0, \log \rho_\alpha)$ and it satisfies the initial condition $\gamma_\alpha(0) = (0, \alpha)$. We remind the reader that the function x_α

is increasing as long as the trajectory γ_α lays in the first quadrant, as well as x_α decreases when γ_α is in the fourth quadrant. If $\rho_\alpha < +\infty$, then the trajectory γ_α reaches the x' -axes at the finite time $t = \log \rho_\alpha$, with $\gamma_\alpha(\log \rho_\alpha) = \left(0, \rho_\alpha^{\frac{p+1}{p-1}} u'_\alpha(\rho_\alpha)\right)$ and $u'_\alpha(\rho_\alpha) < 0$. The typical behavior of a trajectory γ_α having $\rho_\alpha < +\infty$ is depicted in the figure above.

In the following result we collect some properties of the trajectory γ_α , partially observed in previous contributions (we refer in particular to [10, 3, 14]).

Lemma 2.1. *For each $\alpha > 0$, there exist unique $0 < t_\alpha < s_\alpha < \log \rho_\alpha$ such that*

- (i) $x'_\alpha(t) > \frac{2}{p-1}x_\alpha(t)$ for all $t \in [0, t_\alpha)$ and $x'_\alpha(t) < \frac{2}{p-1}x_\alpha(t)$ for all $t \in (t_\alpha, \log \rho_\alpha)$;
- (ii) $\frac{2}{p-1}x_\alpha(t) - \frac{x_\alpha^p(t)}{\lambda(N-1)} < x'_\alpha(t) < \frac{2}{p-1}x_\alpha(t)$ for all $t \in (t_\alpha, s_\alpha)$ and $\gamma_\alpha(s_\alpha) \in \mathcal{C}$;
- (iii) the trajectory γ_α , as any other trajectory γ , can intersect the curve \mathcal{C} from above only transversally and only at points satisfying $x^{p-1} > \frac{\lambda(N-1)}{p}$. In particular, $x_\alpha^{p-1}(s_\alpha) > \frac{\lambda(N-1)}{p}$ and $\gamma_\alpha(t) \in \mathcal{R}^-$ for t in a right neighborhood of s_α ;
- (iv) the trajectory γ_α , as any other trajectory γ , can intersect the curve \mathcal{C} from below only at points satisfying $x^{p-1} \leq \frac{\lambda(N-1)}{p}$ and intersections from below are always transversal if they occur at points satisfying $x^{p-1} < \frac{\lambda(N-1)}{p}$;
- (v) if either $\rho_\alpha < +\infty$ or $\rho_\alpha = +\infty$ and $\gamma_\alpha(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$, then the trajectory γ_α intersects (and crosses) the x -axis exactly once and $\gamma_\alpha(t) \in \mathcal{R}^-$ for all $t \in (s_\alpha, \log \rho_\alpha)$;
- (vi) the map $\gamma_\alpha(t)$ is bounded for $t \in [0, \log \rho_\alpha)$;
- (vii) if $\rho_\alpha = +\infty$, then, as $t \rightarrow +\infty$, only three possible cases can occur: either $\gamma_\alpha(t) \rightarrow (0, 0)$, or $\gamma_\alpha(t) \rightarrow \left((\tilde{b}_+ \Lambda)^{1/(p-1)}, 0\right)$, or $\gamma_\alpha(t)$ approaches the trajectory of a periodic orbit lying in the right half-plane and enclosing the equilibrium point $\left((\tilde{b}_+ \Lambda)^{1/(p-1)}, 0\right)$;
- (viii) if $\rho_\alpha = +\infty$ and $\gamma_\alpha(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$, then there exist non zero constants $c_{1,\alpha}$ and $c_{2,\alpha}$ such that

$$e^{-\lambda_1 t} \gamma_\alpha(t) \rightarrow (c_{1,\alpha}, c_{2,\alpha}) \quad \text{as } t \rightarrow +\infty$$

$$\text{where } \lambda_1 = -\frac{(\tilde{N}_+ - 2)p - \tilde{N}_+}{p-1};$$

- (ix) if $\rho_\alpha = +\infty$ and $\gamma_\alpha(t)$ approaches the trajectory of a periodic orbit as $t \rightarrow +\infty$, then $\gamma_\alpha(t)$ intersects the curve \mathcal{C} infinitely many times for $t \in [0, +\infty)$.

Proof. By setting $t_\alpha = \log \tau_\alpha$ and $s_\alpha = \log \sigma_\alpha$, properties (i) e (ii) immediately follow by the definition of τ_α and σ_α .

In order to prove (iii), let us assume that, for a time $s \in (0, \log \rho_\alpha)$, one has $\gamma_\alpha(s) \in \mathcal{C}$ and $\gamma_\alpha(t) \in \mathcal{R}^+$ for t in a left neighborhood of s . This means that

$$x'_\alpha(s) = \frac{2}{p-1}x_\alpha(s) - \frac{x_\alpha^p(s)}{\lambda(N-1)}$$

and

$$x'_\alpha(t) > \frac{2}{p-1}x_\alpha(t) - \frac{x_\alpha^p(t)}{\lambda(N-1)} \quad \text{for } t \in (s - \delta, s),$$

for some $\delta > 0$. It then follows that

$$x''_\alpha(s) \leq \left(\frac{2}{p-1} - \frac{p x_\alpha^{p-1}(s)}{\lambda(N-1)}\right) x'_\alpha(s) = \left(\frac{2}{p-1} - \frac{p x_\alpha^{p-1}(s)}{\lambda(N-1)}\right) \left(\frac{2}{p-1} - \frac{x_\alpha^{p-1}(s)}{\lambda(N-1)}\right) x_\alpha(s).$$

On the other hand, by (2.4) one has

$$x''_\alpha(s) = a x'_\alpha(s) + b x_\alpha(s) - \frac{x_\alpha^p(s)}{\lambda} = \left(\frac{4}{(p-1)^2} - \frac{p+3}{\lambda(N-1)(p-1)} x_\alpha^{p-1}(s) \right) x_\alpha(s).$$

Combining the above relationships, since $x_\alpha(s) > 0$, we then obtain

$$x_\alpha^{p-1}(s) \geq \frac{\lambda(N-1)}{p},$$

and the intersection is tangential if and only if $x_\alpha^{p-1}(s) = \frac{\lambda(N-1)}{p}$. But a second order analysis at the intersection point, as in the proof of Lemma 3.1 in [10], shows that intersections from above never occur at the point $\left(\left(\frac{\lambda(N-1)}{p} \right)^{\frac{1}{p-1}}, \frac{p+1}{p(p-1)} \left(\frac{\lambda(N-1)}{p} \right)^{\frac{1}{p-1}} \right)$. This proves statement (iii). The same argument, with reversed inequalities, proves claim (iv).

In order to prove (v), let γ_α be such that either $\rho_\alpha < +\infty$ or $\gamma_\alpha(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$. Since $\gamma_\alpha(0) = (0, \alpha)$, then the trajectory crosses at least once the x -axis. Let us call $t_0 \in (0, \log \rho_\alpha)$ the time of the first intersection between γ_α and the x -axis, that is $x'_\alpha(t_0) = 0$ and $x'_\alpha(t) > 0$ for all $t \in [0, t_0)$. Then, $x''_\alpha(t_0) \leq 0$. On the other hand, by (2.4), one has

$$x''_\alpha(t_0) = \begin{cases} x_\alpha(t_0) \left(b - \frac{x_\alpha(t_0)^{p-1}}{\lambda} \right) & \text{if } x_\alpha(t_0)^{p-1} \geq \frac{2\lambda(N-1)}{p-1}, \\ x_\alpha(t_0) \left(\tilde{b}_+ - \frac{x_\alpha(t_0)^{p-1}}{\Lambda} \right) & \text{if } x_\alpha(t_0)^{p-1} < \frac{2\lambda(N-1)}{p-1}. \end{cases}$$

This immediately implies that $x''_\alpha(t_0) < 0$ if $x_\alpha(t_0)^{p-1} \geq \frac{2\lambda(N-1)}{p-1}$, since $\frac{2\lambda(N-1)}{p-1} > \lambda b$, whereas it yields $x_\alpha(t_0)^{p-1} \geq \Lambda \tilde{b}_+$ in the case $x_\alpha(t_0)^{p-1} < \frac{2\lambda(N-1)}{p-1}$. But if $x_\alpha(t_0)^{p-1} = \Lambda \tilde{b}_+$, then $\tilde{b}_+ > 0$ and $\gamma_\alpha(t_0) = \left((\Lambda \tilde{b}_+)^{\frac{1}{p-1}}, 0 \right)$, which is impossible since $\left((\Lambda \tilde{b}_+)^{\frac{1}{p-1}}, 0 \right)$ is an equilibrium point. Hence, in both cases we obtain $x''_\alpha(t_0) < 0$ and, for $t > t_0$, the trajectory enters the fourth quadrant and $x_\alpha(t)$ starts to decrease. We then observe that for $t > t_0$ the trajectory never intersects the x -axis, since otherwise, arguing as above, it would cross it, entering again into the first quadrant, and, in order to avoid self-intersection, it could not reach any point of the x' -axes for subsequent times. This proves that γ_α intersects exactly once the x -axis. We notice further that, if $s_\alpha < t_0$, that is if the first intersection point between γ_α and \mathcal{C} lies in the first quadrant, then, for $t \in (s_\alpha, t_0]$ the trajectory lies in the first quadrant and $x_\alpha(t) > x_\alpha(s_\alpha)$, so that $\gamma_\alpha(t)$ cannot intersect \mathcal{C} by property (iv). Therefore, $\gamma_\alpha(t)$ stays in \mathcal{R}^- until it reaches the x -axis at the time t_0 , then it enters the fourth quadrant, $x_\alpha(t)$ becomes decreasing, and $\gamma_\alpha(t) \in \mathcal{R}^-$ also for $t \geq t_0$. The same argument can be applied in the case $s_\alpha \geq t_0$, since then $x'_\alpha(t) < 0$ for any $t \in (s_\alpha, \log \rho_\alpha)$ and, by property (iv), γ_α cannot intersect \mathcal{C} from below. This completely proves the claim.

Let us now prove (vi). The statement is trivial if $\rho_\alpha < +\infty$, so we assume $\rho_\alpha = +\infty$ and, thus, $p > \frac{\tilde{N}_+}{\tilde{N}_+ - 2}$.

From (iv) it follows that, if $\gamma_\alpha(t)$ intersects again the curve \mathcal{C} for $t > s_\alpha$, and if we call $\hat{s} > s_\alpha$ the first time such that $\gamma_\alpha(\hat{s}) \in \mathcal{C}$ and $\gamma_\alpha(t) \in \mathcal{R}^-$ for $t \in (s_\alpha, \hat{s})$, then necessarily $x'_\alpha(\hat{s}) > 0$. Hence, $\gamma_\alpha(t)$ intersects twice the x -axis for $t \in [0, \hat{s}]$ and therefore, for all $t \in [0, +\infty)$, one has

$$|\gamma_\alpha(t)| \leq \max_{s \in [0, \hat{s}]} |\gamma_\alpha(s)|.$$

It remains to prove the boundedness of the trajectories $\gamma_\alpha(t)$ satisfying $\gamma_\alpha(t) \in \mathcal{R}^-$ for all $t > s_\alpha$. In this case it is convenient to look back at the function u_α , which is convex and decreasing in

$(\sigma_\alpha, +\infty)$. By (2.1), u_α satisfies

$$u_\alpha'' + \frac{\tilde{N}_+ - 1}{r} u_\alpha' = -\frac{u_\alpha^p}{\Lambda} \quad \text{for } r \in [\sigma_\alpha, +\infty),$$

which can be written as

$$\left(r^{\tilde{N}_+ - 1} u_\alpha' \right)' = -r^{\tilde{N}_+ - 1} \frac{u_\alpha^p}{\Lambda} \quad \text{in } [\sigma_\alpha, +\infty). \quad (2.6)$$

Integrating between σ_α and any $r > \sigma_\alpha$, and using the decreasing monotonicity of u_α , we get

$$r^{\tilde{N}_+ - 1} u_\alpha'(r) \leq \int_{\sigma_\alpha}^r \left(s^{\tilde{N}_+ - 1} u_\alpha'(s) \right)' ds \leq -\frac{u_\alpha^p(r)}{\Lambda \tilde{N}_+} \left(r^{\tilde{N}_+} - \sigma_\alpha^{\tilde{N}_+} \right).$$

By integrating once again, we deduce

$$-\frac{u_\alpha^{1-p}(r)}{p-1} \leq \int_{\sigma_\alpha}^r \frac{u_\alpha'(s)}{u_\alpha^p(s)} ds \leq -\frac{1}{\Lambda \tilde{N}_+} \int_{\sigma_\alpha}^r \frac{s^{\tilde{N}_+} - \sigma_\alpha^{\tilde{N}_+}}{s^{\tilde{N}_+ - 1}} ds \leq -\frac{1}{2\Lambda \tilde{N}_+} \left(r^2 - \frac{\tilde{N}_+}{\tilde{N}_+ - 2} \sigma_\alpha^2 \right).$$

Hence, for $r \geq \hat{\sigma}_\alpha := \sqrt{\frac{2\tilde{N}_+}{\tilde{N}_+ - 2}} \sigma_\alpha$, we have

$$u_\alpha^{p-1}(r) \leq \frac{4\Lambda \tilde{N}_+}{(p-1)r^2} =: \frac{c_1}{r^2}. \quad (2.7)$$

Using (2.7) into (2.6), we also have

$$\left(r^{\tilde{N}_+ - 1} u_\alpha' \right)' \geq -\frac{c_1^{\frac{p}{p-1}}}{\Lambda} r^{\tilde{N}_+ - \frac{2p}{p-1} - 1} \quad \text{in } [\hat{\sigma}_\alpha, +\infty).$$

By integrating in $[\hat{\sigma}_\alpha, r]$ and recalling that $p > \frac{\tilde{N}_+}{\tilde{N}_+ - 2}$, we get

$$r^{\tilde{N}_+ - 1} u_\alpha'(r) \geq \hat{\sigma}_\alpha^{\tilde{N}_+ - 1} u_\alpha'(\hat{\sigma}_\alpha) - \frac{(p-1)c_1^{\frac{p}{p-1}}}{\Lambda \left(p(\tilde{N}_+ - 2) - \tilde{N}_+ \right)} r^{\frac{p(\tilde{N}_+ - 2) - \tilde{N}_+}{p-1}}.$$

Therefore, for $r \geq \hat{\sigma}_\alpha$, using again that $p > \frac{\tilde{N}_+}{\tilde{N}_+ - 2}$, we obtain the estimate

$$0 \leq -u_\alpha'(r) \leq \left(-\hat{\sigma}_\alpha^{\tilde{N}_+ - 1} u_\alpha'(\hat{\sigma}_\alpha) + \frac{(p-1)c_1^{\frac{p}{p-1}}}{\Lambda \left(p(\tilde{N}_+ - 2) - \tilde{N}_+ \right)} \right) \frac{1}{r^{\frac{p+1}{p-1}}} =: \frac{c_2}{r^{\frac{p+1}{p-1}}}. \quad (2.8)$$

Recalling the definition (2.3) of x_α , estimates (2.7) and (2.8) then yield

$$|\gamma_\alpha(t)| \leq c_3 \quad \text{for } t \geq \log \hat{\sigma}_\alpha,$$

for a positive constant c_3 depending on p, Λ, λ, N and α . This proves (vi).

Claim (vii) then follows by (vi) and the classical Poincaré-Bendixson Theorem (see e.g. [15]), after observing that the only equilibrium points of system (2.4) reachable by trajectories lying in the right half-plane are $(0, 0)$ and $\left((\tilde{b}_+ \Lambda)^{1/(p-1)}, 0 \right)$ if $p > \frac{\tilde{N}_+}{\tilde{N}_+ - 2}$ (i.e. $\tilde{b}_+ > 0$), which is always the case if $\rho_\alpha = +\infty$.

Also statement (viii) classically follows from the perturbed linear systems theory, λ_1 being the negative eigenvalue of the linear system obtained by linearizing around zero system (2.4). For a detailed proof, we refer e.g. to Lemma 3.3 in [17].

Let us finally prove (ix). We first claim that if a periodic orbit γ of system (2.4) is the ω -limit set of a trajectory γ_α , then γ intersects the curve \mathcal{C} . Indeed, the trajectory of such a periodic orbit, if any, is a closed curve γ contained in the right half-plane and winding around the equilibrium point $\left((\tilde{b}_+ \Lambda)^{1/(p-1)}, 0 \right)$. If, by contradiction, γ does not intersect \mathcal{C} , then $\gamma(t)$

lies entirely in the region \mathcal{R}^- . This implies that γ is the trajectory associated with a positive periodic solution $x(t)$ of the single equation

$$x'' = \tilde{a}_+ x' + \tilde{b}_+ x - \frac{x^p}{\Lambda}.$$

Hence, the energy function

$$E(x, y) = \frac{y^2}{2} - \tilde{b}_+ \frac{x^2}{2} + \frac{x^{p+1}}{\Lambda(p+1)}, \quad (2.9)$$

when evaluated along $\gamma(t)$, satisfies

$$\frac{d}{dt}(E(\gamma(t))) = \tilde{a}_+(x')^2.$$

In particular, if $\tilde{a}_+ \neq 0$, i.e. if $p \neq \frac{\tilde{N}_++2}{\tilde{N}_+-2}$, then the function $E(\gamma(t))$ is strictly monotone and periodic: a contradiction. On the other hand, if $\tilde{a}_+ = 0$, i.e. if $p = \frac{\tilde{N}_++2}{\tilde{N}_+-2}$, then $E(\gamma(t))$ is constant, that is γ is a level line of the function E contained in \mathcal{R}^- . Every level line of the function E contained in \mathcal{R}^- actually is a periodic orbit of system (2.4), and, since γ is contained in \mathcal{R}^- , by continuity there exists a larger level line of the function E still contained in \mathcal{R}^- and enclosing γ . This means that γ cannot be approached by any of the trajectories γ_α , again a contradiction to the assumptions.

Therefore, for any value of $p > \frac{\tilde{N}_++2}{\tilde{N}_+-2}$, the ω -limit γ intersects \mathcal{C} at least once. Hence, if γ_α is a trajectory winding around γ as $t \rightarrow +\infty$, then $\gamma_\alpha(t)$ intersects \mathcal{C} infinitely many times as t ranges in $[0, +\infty)$ (see also the proof of Proposition 3.3 in [10]). \square

By means of Lemma 2.1 (vii), the maximal solutions u_α of problem (2.1) for which $\rho_\alpha = +\infty$ can be classified in terms of the asymptotic behavior of the associated trajectory γ_α . Precisely, a solution u_α is called fast, slow or pseudo-slow decaying as $r \rightarrow +\infty$ provided that $\gamma_\alpha(t)$ converges as $t \rightarrow +\infty$ respectively to $(0, 0)$, $\left((\tilde{b}_+\Lambda)^{1/(p-1)}, 0\right)$ or to a periodic orbit. Furthermore, thanks to Lemma 2.1 (viii), a fast decaying solution may be equivalently defined as a solution $u_\alpha : [1, +\infty) \rightarrow \mathbb{R}$ satisfying the asymptotic condition (1.3).

Remark 2.2. Positive entire radial solutions of equations

$$-\mathcal{M}_{\lambda, \Lambda}^\pm(D^2u) = u^p \quad \text{in } \mathbb{R}^N$$

have been studied in detail in [10]. In this case the initial value problems for the ODE to be considered are

$$\begin{cases} v''(r) = M^\pm \left(-\frac{\lambda(N-1)}{r} K^\pm(v') - v^p \right) & \text{for } r > 0, \\ v(r) > 0 & \text{for } r > 0, \\ v(0) = \alpha, \quad v'(0) = 0, \end{cases}$$

and for the maximal solutions v_α , defined on $[0, r_\alpha)$ for some $0 < r_\alpha \leq +\infty$, one has $r_\alpha < +\infty$ if and only if $p < p_\pm^*$. We recall that, by the homogeneity property of the equation, for any $\alpha, \alpha' > 0$ one has

$$v_{\alpha'}(r) = \frac{\alpha'}{\alpha} v_\alpha \left(\left(\frac{\alpha'}{\alpha} \right)^{\frac{p-1}{2}} r \right), \quad (2.10)$$

that is changing the initial condition just reflects in scaling the solution. Therefore, the property $r_\alpha < +\infty$ or $r_\alpha = +\infty$ only depends on p .

The trajectory $\Gamma_\alpha(t) = (\xi_\alpha(t), \xi'_\alpha(t))$ associated with the Emden-Fowler transform $\xi_\alpha(t)$ of a solution $v_\alpha(r)$, defined through (2.3), is defined for $t \in (-\infty, \log r_\alpha)$ and satisfies

$$\lim_{t \rightarrow -\infty} \Gamma_\alpha(t) = (0, 0).$$

Relationship (2.10) implies that

$$\Gamma_{\alpha'}(t) = \Gamma_\alpha \left(t + \frac{p-1}{2} \log \frac{\alpha'}{\alpha} \right),$$

and thus the initial value α for the solution $v_\alpha(r)$ does not affect the support of the trajectory $\Gamma_\alpha(t)$, which is then denoted simply by $\Gamma(t)$. In the phase plane, the trajectory $\Gamma(t)$ is a curve exiting from the origin, lying always below the line L , and staying above the curve \mathcal{C} for $t \in (-\infty, S_\alpha)$, for some $S_\alpha < \log r_\alpha$. We recall from [10] that $\Gamma(t)$ satisfies properties (iii)-(ix) of Lemma 2.1, and fast, slow or pseudo-slow decaying solutions v_α are accordingly defined. A crucial result proved in [10] is that, independently of α , v_α is a fast decaying solution if and only if $p = p_\pm^*$.

3. EXISTENCE IN THE SUPERCRITICAL CASE $p > p_\pm^*$ FOR $\mathcal{M}_{\lambda, \Lambda}^\pm$

This section is devoted to the proof of the existence statement in Theorem 1.1.

Theorem 3.1. *If $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$ and $p > p_+^*$ then (1.1) has a nontrivial radial solution. If $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ and $p > p_-^*$ then (1.1) has a nontrivial radial solution.*

Proof. We write the proof for $\mathcal{M}_{\lambda, \Lambda}^+$. Unless otherwise said, the proof for $\mathcal{M}_{\lambda, \Lambda}^-$ is obtained just by exchanging λ with Λ .

Assume by contradiction that the thesis is false. Then, for any $\alpha > 0$, the unique maximal solution $u_\alpha = u_\alpha(r)$ of the initial value problem (2.1) is defined on the finite interval $[1, \rho_\alpha]$. We recall that there exists a unique $\tau_\alpha \in (1, \rho_\alpha)$ such that $u'_\alpha(\tau_\alpha) = 0$, with $u'_\alpha(r) > 0$ for $r \in [1, \tau_\alpha)$ and $u'_\alpha(r) < 0$ for $r \in (\tau_\alpha, \rho_\alpha]$. Moreover, by Lemma 2.1 (ii)-(iii)-(v), there exists a unique $\sigma_\alpha \in (\tau_\alpha, \rho_\alpha)$ such that $u''_\alpha(\sigma_\alpha) = 0$, with $u''_\alpha(r) < 0$ for $r \in [1, \sigma_\alpha)$ and $u''_\alpha > 0$ for $r \in (\sigma_\alpha, \rho_\alpha]$.

Next, we perform the asymptotic analysis as $\alpha \rightarrow 0^+$. From [14, Lemma 3.1] we know that $\tau_\alpha \rightarrow +\infty$, $u_\alpha(\tau_\alpha) \rightarrow 0$. Hence, we have also $\sigma_\alpha \rightarrow +\infty$, $\rho_\alpha \rightarrow +\infty$ as $\alpha \rightarrow 0^+$. Setting $m_\alpha := u_\alpha(\tau_\alpha)$, we consider the rescaled function

$$\tilde{u}_\alpha(r) := \frac{1}{m_\alpha} u_\alpha \left(\frac{r}{m_\alpha^{\frac{p-1}{2}}} \right), \quad r \in \left[m_\alpha^{\frac{p-1}{2}}, m_\alpha^{\frac{p-1}{2}} \rho_\alpha \right].$$

For notational convenience, let us set $\tilde{\tau}_\alpha := m_\alpha^{\frac{p-1}{2}} \tau_\alpha$, $\tilde{\sigma}_\alpha := m_\alpha^{\frac{p-1}{2}} \sigma_\alpha$, $\tilde{\rho}_\alpha := m_\alpha^{\frac{p-1}{2}} \rho_\alpha$. Note that, by construction, \tilde{u}_α is a radial solution to

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 u) = u^p & \text{in } \tilde{A}_\alpha, \\ u > 0 & \text{in } \tilde{A}_\alpha, \\ u = 0 & \text{on } \partial \tilde{A}_\alpha, \end{cases} \quad (3.1)$$

where $\tilde{A}_\alpha = \left\{ x \in \mathbb{R}^N : m_\alpha^{\frac{p-1}{2}} < |x| < \tilde{\rho}_\alpha \right\}$, and one has $1 = \tilde{u}_\alpha(\tilde{\tau}_\alpha) = \max_{\tilde{A}_\alpha} \tilde{u}_\alpha$.

Step 1: As $\alpha \rightarrow 0^+$, we have $\tilde{\tau}_\alpha \rightarrow 0^+$.

Let us consider the functional $E_1 : [1, \tau_\alpha] \rightarrow \mathbb{R}$ defined by

$$E_1(r) := \frac{(u'_\alpha(r))^2}{2} + \frac{u_\alpha^{p+1}(r)}{\lambda(p+1)}.$$

A direct computation, as in the proof of [14, Lemma 3.1], shows that $E_1'(r) \leq 0$ for $r \in [1, \tau_\alpha]$. Thus, from $E_1(r) \geq E_1(\tau_\alpha)$ and $u_\alpha'(r) \geq 0$ for $r \in [1, \tau_\alpha]$, we infer that

$$u_\alpha'(r) \geq \sqrt{\frac{2}{\lambda(p+1)}} \left(m_\alpha^{p+1} - u_\alpha^{p+1}(r) \right) \quad \forall r \in [1, \tau_\alpha].$$

Integrating between 1 and τ_α we get that

$$\int_1^{\tau_\alpha} \frac{u_\alpha'(r)}{\sqrt{m_\alpha^{p+1} - u_\alpha^{p+1}(r)}} dr \geq \sqrt{\frac{2}{\lambda(p+1)}} (\tau_\alpha - 1). \quad (3.2)$$

Performing the change of variable $t = \frac{u_\alpha(r)}{m_\alpha}$, we rewrite (3.2) as

$$\frac{1}{m_\alpha^{\frac{p-1}{2}}} \int_0^1 \frac{1}{\sqrt{1-t^{p+1}}} dt \geq \sqrt{\frac{2}{\lambda(p+1)}} (\tau_\alpha - 1). \quad (3.3)$$

Since $\int_0^1 \frac{1}{\sqrt{1-t^{p+1}}} dt \leq \int_0^1 \frac{1}{\sqrt{1-t^2}} dt = \frac{\pi}{2}$ and $m_\alpha \rightarrow 0$ as $\alpha \rightarrow 0^+$, from (3.3) we deduce that

$$\limsup_{\alpha \rightarrow 0^+} \tilde{\tau}_\alpha \leq \frac{\pi}{2} \sqrt{\frac{\lambda(p+1)}{2}}.$$

Assume now, by contradiction, that $l := \limsup_{\alpha \rightarrow 0^+} \tilde{\tau}_\alpha > 0$ and let us select a sequence, still denoted by α , such that $\tilde{\tau}_\alpha \rightarrow l$ as $\alpha \rightarrow 0^+$.

By (2.1) and the concavity/monotonicity properties of \tilde{u}_α , it follows that \tilde{u}_α in particular satisfies

$$\tilde{u}_\alpha'' + \frac{\tilde{N}_- - 1}{r} \tilde{u}_\alpha' = -\frac{\tilde{u}_\alpha^p}{\lambda} \quad \text{in } [m_\alpha^{\frac{p-1}{2}}, \tilde{\tau}_\alpha].$$

Since $0 \leq \tilde{u}_\alpha \leq 1$, it follows that, up to a further subsequence if necessary, $\tilde{u}_\alpha(r) \rightarrow \tilde{u}(r)$ in $C_{loc}^2((0, l))$, for some \tilde{u} satisfying $0 \leq \tilde{u} \leq 1$, $\tilde{u}' \geq 0$ in $(0, l]$, $\tilde{u}(l) = 1$ and

$$\tilde{u}'' + \frac{\tilde{N}_- - 1}{r} \tilde{u}' = -\frac{\tilde{u}^p}{\lambda} \quad \text{in } (0, l].$$

We observe that \tilde{u} cannot be identically equal to 1, since the constant function 1 does not satisfy the above equation. Hence, there exists $r_0 \in (0, l)$ such that $\tilde{u}'(r_0) > 0$. Observing further that

$$(\tilde{u}' r^{\tilde{N}_- - 1})' = -\frac{1}{\lambda} \tilde{u}^p r^{\tilde{N}_- - 1} \leq 0,$$

we infer that, for all $r \in (0, r_0]$,

$$\tilde{u}'(r) \geq \tilde{u}'(r_0) \frac{r_0^{\tilde{N}_- - 1}}{r^{\tilde{N}_- - 1}}.$$

For any $\varepsilon \in (0, r_0)$, let us integrate the above inequality from ε to $r \in [\varepsilon, r_0]$. We obtain

$$1 \geq \tilde{u}(r) - \tilde{u}(\varepsilon) \geq \frac{\tilde{u}'(r_0) r_0^{\tilde{N}_- - 1}}{\tilde{N}_- - 2} \left(\frac{1}{\varepsilon^{\tilde{N}_- - 2}} - \frac{1}{r^{\tilde{N}_- - 2}} \right),$$

which yields a contradiction as $\varepsilon \rightarrow 0^+$. The proof of Step 1 is complete.

Step 2: Up to a sequence $\alpha \rightarrow 0^+$, we have $\tilde{\sigma}_\alpha \rightarrow l$ for some $l \in (0, +\infty)$.

We first observe that, by the definition of σ_α and of x_α given by (2.3), from Lemma 2.1-(iii) it follows that

$$\sigma_\alpha^2 u_\alpha^{p-1}(\sigma_\alpha) > \frac{\lambda(N-1)}{p}.$$

Recalling the definition of $\tilde{\sigma}_\alpha$ and observing that $u_\alpha(\sigma_\alpha) \leq m_\alpha$, this readily implies that, for every $\alpha > 0$,

$$\tilde{\sigma}_\alpha > \sqrt{\frac{\lambda(N-1)}{p}}. \quad (3.4)$$

Therefore, in order to prove Step 2, it is enough to show that $\limsup_{\alpha \rightarrow 0^+} \tilde{\sigma}_\alpha < +\infty$.

Assume by contradiction that for some sequence $\alpha \rightarrow 0^+$ one has $\tilde{\sigma}_\alpha \rightarrow +\infty$. Since \tilde{u}_α is positive and concave in $[\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha]$, from $\tilde{u}_\alpha(\tilde{\tau}_\alpha) = 1$ we infer that

$$\tilde{u}_\alpha(r) \geq \frac{\tilde{\sigma}_\alpha - r}{\tilde{\sigma}_\alpha - \tilde{\tau}_\alpha} \quad \text{for any } r \in [\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha]. \quad (3.5)$$

On the other hand, in the present assumption the limit domain of \tilde{u}_α is $(0, +\infty)$, and, by standard elliptic estimates, up to a further subsequence, $\tilde{u}_\alpha \rightarrow \tilde{u}$ in $C_{loc}^2((0, +\infty))$, where \tilde{u} is a radial solution of

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 w) = w^p & \text{in } \mathbb{R}^N \setminus \{0\}, \\ w \geq 0 & \text{in } \mathbb{R}^N \setminus \{0\}. \end{cases} \quad (3.6)$$

Passing to the limit as $\alpha \rightarrow 0^+$ in (3.5) and using Step 1, we deduce that

$$\tilde{u}(r) \geq 1 \quad \text{for any } r \in (0, +\infty).$$

But $\tilde{u} \leq 1$ by construction, and, therefore, $\tilde{u} \equiv 1$, which is absurd because \tilde{u} solves (3.6). The proof of Step 2 is complete.

Step 3: As $\alpha \rightarrow 0^+$, we have $\tilde{\rho}_\alpha \rightarrow +\infty$.

Let us assume, by contradiction, that $k := \liminf_{\alpha \rightarrow 0^+} \tilde{\rho}_\alpha < +\infty$. Then, there exists a sequence still denoted by $\alpha \rightarrow 0^+$ such that $\tilde{\rho}_\alpha \rightarrow k$. By (3.4), we have that $k > 0$. Recall that, by Step 1, $\tilde{\tau}_\alpha \rightarrow 0$ and, possibly considering a further subsequence, $\tilde{\sigma}_\alpha \rightarrow l_1$ for some $0 < l_1 \leq k$. By construction and standard elliptic estimates, it follows that, possibly up to a further subsequence, $\tilde{u}_\alpha \rightarrow \tilde{u}$ in $C_{loc}^2((0, k))$, where \tilde{u} is in this case a radial solution of

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 w) = w^p & \text{in } B_k \setminus \{0\}, \\ w \geq 0 & \text{in } B_k \setminus \{0\}, \\ w = 0 & \text{on } \partial B_k. \end{cases} \quad (3.7)$$

We claim that \tilde{u} can be extended to a smooth non-trivial radial solution of (3.7) in the whole ball B_k . Indeed, by (2.1), in $[\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha]$ the function \tilde{u}_α satisfies

$$\tilde{u}_\alpha''(r) + \frac{N-1}{r} \tilde{u}_\alpha'(r) = -\frac{\tilde{u}_\alpha^p(r)}{\lambda},$$

and, therefore,

$$(r^{N-1} \tilde{u}_\alpha'(r))' = -\frac{1}{\lambda} r^{N-1} \tilde{u}_\alpha^p(r) \geq -\frac{1}{\lambda} r^{N-1}. \quad (3.8)$$

Integrating between $\tilde{\tau}_\alpha$ and $r \in (\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha)$, we get

$$r^{N-1} \tilde{u}_\alpha'(r) \geq -\frac{1}{\lambda N} (r^N - \tilde{\tau}_\alpha^N),$$

which yields

$$\tilde{u}_\alpha'(r) \geq -\frac{1}{\lambda N} r \quad \text{for } r \in (\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha). \quad (3.9)$$

Integrating again between $\tilde{\tau}_\alpha$ and $r \in (\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha)$ and taking into account that $\tilde{u}_\alpha(\tilde{\tau}_\alpha) = 1$, we have

$$\tilde{u}_\alpha(r) \geq 1 - \frac{1}{2\lambda N} r^2 \quad \text{for } r \in (\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha). \quad (3.10)$$

Therefore, for any fixed $r \in (0, l_1)$, letting $\alpha \rightarrow 0^+$ we infer that

$$1 \geq \tilde{u}(r) \geq 1 - \frac{1}{2\lambda N} r^2,$$

which implies that

$$\lim_{r \rightarrow 0} \tilde{u}(r) = 1. \quad (3.11)$$

Hence \tilde{u} can be extended by continuity in 0 by the value $\tilde{u}(0) = 1$. Moreover, since $\tilde{u}'_\alpha(r) \leq 0$ in $(\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha)$, from (3.9) we get that

$$|\tilde{u}'_\alpha(r)| \leq \frac{r}{\lambda N} \quad \text{for any } r \in (\tilde{\tau}_\alpha, \tilde{\sigma}_\alpha).$$

Therefore, fixing $r \in (0, l_1)$ and passing to the limit as $\alpha \rightarrow 0^+$, we get

$$|\tilde{u}'(r)| \leq \frac{r}{\lambda N} \quad (3.12)$$

which gives $\lim_{r \rightarrow 0} \tilde{u}'(r) = 0$. Thus, \tilde{u} in particular satisfies

$$\begin{cases} \tilde{u}'' + \frac{N-1}{r} \tilde{u}' = -\frac{\tilde{u}^p}{\lambda} & \text{in } (0, l_1), \\ \tilde{u}(0) = 1, \quad \tilde{u}'(0) = 0, \end{cases}$$

and this implies that it is of class C^2 up to 0.

Summing up, \tilde{u} is a smooth nontrivial radial solution to

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 w) = w^p & \text{in } B_k, \\ w > 0 & \text{in } B_k, \\ w = 0 & \text{on } \partial B_k. \end{cases}$$

Since $p > p_+^*$, this is a contradiction to the results of [6, 10]. The proof of Step 3 is complete.

Conclusion: By the previous steps, we deduce that, up to a sequence $\alpha \rightarrow 0^+$, $\tilde{u}_\alpha \rightarrow \tilde{u}$ in $C_{loc}^2((0, +\infty))$, where \tilde{u} can be extended at the origin as a smooth positive radial solution of

$$-\mathcal{M}_{\lambda, \Lambda}^+(D^2 u) = u^p \quad \text{in } \mathbb{R}^N. \quad (3.13)$$

Moreover, by construction, \tilde{u} satisfies $\tilde{u}(0) = 1$, it is radially decreasing in $[0, +\infty)$ and it is convex in $[l, +\infty)$.

If $p \in]p_+^*, \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}]$, then it is known (see [10]) that the only positive radial solutions of (3.13) are pseudo-slow decaying solutions, which are functions changing concavity infinite times (see the proof of Lemma 2.1 (ix), or the proof of Proposition 3.3 in [10], and Remark 2.2). This clearly gives a contradiction because \tilde{u} changes concavity exactly once. We point out that this case is considered only for $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$, since for $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ we have $p_-^* > \frac{\tilde{N}_- + 2}{\tilde{N}_- - 2}$.

On the other hand, if $p > \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$, again by the results of [10] the only positive radial solutions of (3.13) are slow decaying functions, i.e. solutions for which the trajectory $\tilde{\gamma}(t) = (\tilde{x}(t), \tilde{x}'(t))$, associated with the Emden-Fowler transform $\tilde{x}(t)$ defined through (2.3), satisfies

$$\lim_{t \rightarrow +\infty} \tilde{\gamma}(t) = (c^*, 0), \quad (3.14)$$

where $c^* = (\Lambda \tilde{b}_+)^{\frac{1}{p-1}}$, with \tilde{b}_+ given by (2.5).

Let us denote by $\tilde{x}_\alpha(t)$ the Emden-Fowler transform (2.3) of $\tilde{u}_\alpha(r)$, and let $\tilde{\gamma}_\alpha(t) = (\tilde{x}_\alpha(t), \tilde{x}'_\alpha(t))$ be the associated trajectory, both defined for $t \in [\frac{p-1}{2} \log m_\alpha, \log \tilde{\rho}_\alpha]$. Then, $\tilde{\gamma}_\alpha \rightarrow \tilde{\gamma}$ in $C_{loc}^1(\mathbb{R})$.

We observe that for $t \in [\log \tilde{\sigma}_\alpha, \log \tilde{\rho}_\alpha]$, the trajectory $\tilde{\gamma}_\alpha(t)$ belongs to \mathcal{R}^- by Lemma 2.1 (v), and, therefore, the energy function E defined in (2.9) satisfies, by (2.4),

$$\frac{d}{dt}(E(\tilde{\gamma}_\alpha(t))) = \tilde{a}_+ \tilde{x}_\alpha'^2.$$

Now, noticing that the condition $p > \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$ is equivalent to $\tilde{a}_+ < 0$, we infer that $E(\tilde{\gamma}_\alpha(t))$ is monotone decreasing for $t \in [\log \tilde{\sigma}_\alpha, \log \tilde{\rho}_\alpha]$. Hence,

$$E(\tilde{\gamma}_\alpha(t)) \geq E(\tilde{\gamma}_\alpha(\log \tilde{\rho}_\alpha)) = \frac{[\tilde{x}'_\alpha(\log \tilde{\rho}_\alpha)]^2}{2} > 0.$$

Letting $\alpha \rightarrow 0^+$, we deduce

$$E(\tilde{\gamma}(t)) \geq 0 \quad \text{for all } t \geq \log l.$$

On the other hand, by (3.14), we also have

$$\lim_{t \rightarrow +\infty} E(\tilde{\gamma}(t)) = E(c^*, 0) = -\frac{(\Lambda \tilde{b}_+)^{\frac{2}{p-1}} \tilde{b}_+(p-1)}{2(p+1)} < 0,$$

and we reach a contradiction in this case as well. \square

4. NON-EXISTENCE IN THE SUBCRITICAL CASE $p < p_\pm^+$ FOR $\mathcal{M}_{\lambda, \Lambda}^\pm$

In this and the next two sections we are concerned with the non-existence part of Theorem 1.1.

Theorem 4.1. *Problem (1.1) does not have any positive radial solution when $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$ and $1 < p < p_+^*$, or when $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ and $1 < p < p_-^*$.*

Proof. In view of [2], it is sufficient to prove the result for $p > \frac{N_\pm}{N_\pm - 2}$. Let us consider $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$, the proof for $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ being completely analogous.

We will show that for any $\alpha > 0$ the unique maximal solution $u_\alpha = u_\alpha(r)$ of (2.1) vanishes at some $\rho_\alpha \in (1, +\infty)$. Consider the unique positive radial solution $u = u(r)$ of

$$\begin{cases} -\mathcal{M}_{\lambda, \Lambda}^+(D^2 u) = u^p & \text{in } B, \\ u = 0 & \text{on } \partial B, \end{cases} \quad (4.1)$$

whose existence is guaranteed by the assumption $p < p_+^*$ (see [10, Theorem 5.1]). Performing the Emden-Fowler change of variable (2.3) for u , we obtain in the phase plane a trajectory $\gamma(t) = (x(t), x'(t))$ defined for all $t \leq 0$ and satisfying

$$\lim_{t \rightarrow -\infty} \gamma(t) = (0, 0), \quad \gamma(0) = (0, u'(1)).$$

By Lemma 2.1 (v), there exists a unique time $t_0 < 0$ such that $\gamma(t_0) = (x(t_0), 0)$ and $x(t_0) > c^* = (\Lambda \tilde{b}_+)^{\frac{1}{p-1}}$, $(c^*, 0)$ being the unique equilibrium point.

We then deduce that the trajectory $\gamma(t)$ and the x' -axis bound a closed region containing the equilibrium points $(0, 0)$ and $(c^*, 0)$. Consider now the trajectory $\gamma_\alpha(t) = (x_\alpha(t), x'_\alpha(t))$ relative to u_α . Clearly, γ_α cannot intersect γ . Moreover it cannot approach neither the equilibrium points nor a periodic orbit, since otherwise it should cross γ . By Lemma 2.1 (vii), it follows that $\gamma_\alpha(t)$ must leave the fourth quadrant in finite time, i.e. there exists $\rho_\alpha \in (1, +\infty)$ such that $u_\alpha(\rho_\alpha) = 0$ as desired. \square

5. NON-EXISTENCE AT THE CRITICAL LEVEL $p = p_+^*$ FOR $\mathcal{M}_{\lambda,\Lambda}^+$

Theorem 5.1. *If $\mathcal{F} = \mathcal{M}_{\lambda,\Lambda}^+$ and $p = p_+^*$ then (1.1) does not have any positive radial solution.*

Proof. Let us assume by contradiction that (1.1) admits a positive radial solution. Then there exists $u_\alpha = u_\alpha(r)$ solution of (2.1) defined in $[1, +\infty)$, and the trajectory $\gamma_\alpha(t) := (x_\alpha(t), x'_\alpha(t))$, associated with the Emden-Fowler transform (2.3) of u_α , is defined for $t \in [0, +\infty)$.

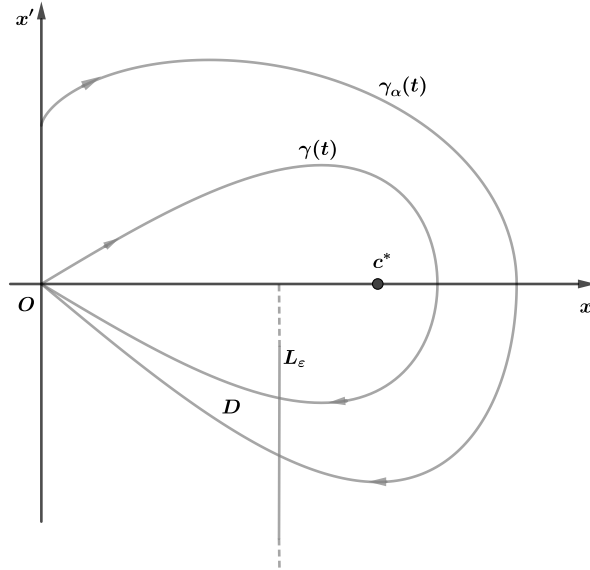
Let us also consider the trajectory $\gamma(t) = (x(t), x'(t))$, $t \in (-\infty, +\infty)$, associated with the unique (up to scaling) solution of

$$-\mathcal{M}_{\lambda,\Lambda}^+(D^2u) = u^{p_+^*} \quad \text{in } \mathbb{R}^N.$$

By [10] (see also Remark 2.2), we know that $\lim_{t \rightarrow \pm\infty} \gamma(t) = (0, 0)$ and the support of $\gamma(t)$ bounds a compact region including the equilibrium point $(c^*, 0)$ of the dynamical system (2.4) for $p = p_+^*$, with $c^* = (\Lambda \tilde{b}_+)^{\frac{1}{p_+^*-1}}$. Since γ_α and γ cannot intersect (apart from the origin), we infer that $\gamma_\alpha(t)$ cannot approach as $t \rightarrow +\infty$ neither a periodic orbit nor the equilibrium point $(c^*, 0)$. Thus, from Lemma 2.1 (vii), it follows that $\gamma_\alpha(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$ as well. Moreover, Lemma 2.1 (v), which also applies to γ , yields that both $\gamma_\alpha(t)$ and $\gamma(t)$ intersect (and cross) the x -axis exactly once. Hence, for $\varepsilon > 0$ small enough, the vertical line $L_\varepsilon = \{x = \varepsilon\}$ intersects exactly once γ_α and γ in the fourth quadrant. Therefore, there exist unique t_ε and $t_{\alpha,\varepsilon}$ such that

$$x(t_\varepsilon) = \varepsilon = x_\alpha(t_{\alpha,\varepsilon}), \quad x'_\alpha(t_{\alpha,\varepsilon}) < x'(t_\varepsilon) < 0.$$

Moreover, again by Lemma 2.1 (v), we can assume that $\varepsilon > 0$ is so small that $\gamma(t)$, $\gamma_\alpha(t) \in \mathcal{R}^-$ respectively for $t \geq t_\varepsilon$ and $t \geq t_{\alpha,\varepsilon}$, see the picture below.



Consider now the compact region D bounded by the closed piecewise smooth curve formed by the trajectories $\gamma(t)$ for $t \geq t_\varepsilon$, $\gamma_\alpha(t)$ for $t \geq t_{\alpha,\varepsilon}$ and the vertical segment joining $\gamma(t_\varepsilon)$ and $\gamma_\alpha(t_{\alpha,\varepsilon})$.

Denoting by $|D|$ the area of D and applying the Gauss-Green Theorem (with ∂D clock-wise oriented), we infer that

$$|D| = \int_{\partial D^-} y \, dx = \int_{t_{\alpha,\varepsilon}}^{+\infty} [x'_\alpha(t)]^2 \, dt + \int_{+\infty}^{t_\varepsilon} [x'(t)]^2 \, dt. \quad (5.1)$$

On the other hand, considering the energy function E defined in (2.9), we obtain that

$$[x'_\alpha(t)]^2 = \frac{1}{\tilde{a}_+} \frac{d}{dt} [E(\gamma_\alpha(t))], \quad [x'(t)]^2 = \frac{1}{\tilde{a}_+} \frac{d}{dt} [E(\gamma(t))].$$

Hence, using also that $x(t_\varepsilon) = x_\alpha(t_{\alpha,\varepsilon})$, we deduce

$$|D| = \frac{1}{\tilde{a}_+} [E(\gamma(t_\varepsilon)) - E(\gamma_\alpha(t_{\alpha,\varepsilon}))] = \frac{1}{2\tilde{a}_+} \left[[x'(t_\varepsilon)]^2 - [x'_\alpha(t_{\alpha,\varepsilon})]^2 \right]. \quad (5.2)$$

Now, since $p_+^* < \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$, it holds that $\tilde{a}_+ = \tilde{a}_+(p_+^*) > 0$, and since $x'_\alpha(t_{\alpha,\varepsilon}) < x'(t_\varepsilon) < 0$, from (5.2) it then follows that $|D| < 0$, which clearly is a contradiction. \square

Remark 5.2. The previous proof does not work for $\mathcal{F} = \mathcal{M}_{\lambda,\Lambda}^-$ and $p = p_-^*$ because in this case $p_-^* > \frac{\tilde{N}_- + 2}{\tilde{N}_- - 2}$ and thus $\tilde{a}_-(p_-^*) < 0$. In particular (5.2) does not give a contradiction.

6. NON-EXISTENCE AT THE CRITICAL LEVEL $p = p_-^*$ FOR $\mathcal{M}_{\lambda,\Lambda}^-$

Theorem 6.1. *If $\mathcal{F} = \mathcal{M}_{\lambda,\Lambda}^-$ and $p = p_-^*$ then (1.1) does not have any positive radial solution.*

The proof of Theorem 6.1 will be given at the end of this section after some preliminary results on solutions of semilinear ODEs.

Let $\nu \in (2, +\infty)$. For any $p > 1$ and $\alpha > 0$, let $u_\alpha = u_\alpha(r)$ be the positive maximal solution of the initial value problem

$$\begin{cases} u''(r) + (\nu - 1) \frac{u'(r)}{r} = -\frac{u^p(r)}{\lambda} & \text{for } r > 1 \\ u(1) = 0, \quad u'(1) = \alpha. \end{cases} \quad (6.1)$$

The function u_α is defined on a maximal interval $[1, \rho_\alpha)$ with $\rho_\alpha \leq +\infty$. If $\rho_\alpha < +\infty$ then $u_\alpha(\rho_\alpha) = 0$, otherwise $u_\alpha(r) > 0$ for any $r > 1$ and $u_\alpha(r) \rightarrow 0$ as $r \rightarrow +\infty$. Set

$$D = \{\alpha \in (0, +\infty) : \rho_\alpha < +\infty\}.$$

The set D is nonempty (see e.g. [14, Proposition 3.2]) and open by the continuous dependence on the initial data for (6.1). Let

$$\alpha^* = \alpha^*(p) := \inf D. \quad (6.2)$$

In the next proposition, the behavior as $r \rightarrow +\infty$ of the solutions u_α is analyzed. In particular, when $\nu = N$ is an integer, we recover the well known results (see e.g. [16]) about radial solutions of the semilinear exterior Dirichlet problem

$$\begin{cases} -\Delta u = u^p & \text{in } \mathbb{R}^N \setminus \bar{B}, \\ u = 0 & \text{on } \partial B, \end{cases}$$

in the case they exist, i.e. in the supercritical regime $p > \frac{N+2}{N-2}$. For the sake of completeness, we include a (different) proof. Let us preliminarily observe that Lemma 2.1 clearly applies to the trajectory $\gamma_\alpha(t) = (x_\alpha(t), x'_\alpha(t))$, where $x_\alpha(t)$ is the Emden-Fowler transform (2.3) of $u_\alpha(r)$, since problem (6.1) is a special case of (2.2) with $\Lambda = \lambda$ and $N = \nu$.

Proposition 6.2. *Let $p > \frac{\nu+2}{\nu-2}$. Then*

- (i) $\alpha^* > 0$;
- (ii) if $\alpha > \alpha^*$ then $\rho_\alpha < +\infty$;
- (iii) $\rho_{\alpha^*} = +\infty$ and $\lim_{r \rightarrow +\infty} r^{\nu-2} u_{\alpha^*}(r) = C$, for some positive constant C ;
- (iv) if $\alpha < \alpha^*$ then $\rho_\alpha = +\infty$ and $\lim_{r \rightarrow +\infty} r^{\frac{2}{p-1}} u_\alpha(r) = \left(\frac{2\lambda((\nu-2)p-\nu)}{(p-1)^2} \right)^{\frac{1}{p-1}}$.

Proof. (i). By contradiction let us assume $\alpha^* = 0$. Then there is a sequence $\alpha_n \in D$ such that $\alpha_n \rightarrow 0$ and $\rho_{\alpha_n} \rightarrow +\infty$ as $n \rightarrow +\infty$, in view of [14, Lemma 3.1].

Let us consider the rescaled function

$$\tilde{u}_{\alpha_n}(r) = \frac{1}{m_{\alpha_n}} u_{\alpha_n} \left(\frac{r}{m_{\alpha_n}^{\frac{p-1}{2}}} \right), \quad r > m_{\alpha_n}^{\frac{p-1}{2}},$$

with $m_{\alpha_n} = u_{\alpha_n}(\tau_{\alpha_n}) = \max_{[1, \rho_{\alpha_n}]} u_{\alpha_n}$. Following the arguments of Step 1 and Step 2 of Theorem 3.1 we infer that $\tilde{\tau}_{\alpha_n} = m_{\alpha_n}^{\frac{p-1}{2}} \tau_{\alpha_n} \rightarrow 0$ and $\tilde{\rho}_{\alpha_n} = m_{\alpha_n}^{\frac{p-1}{2}} \rho_{\alpha_n}$ is bounded from below by a positive constant. Let us show that $\tilde{\rho}_{\alpha_n} \rightarrow +\infty$. If not, up to a subsequence, $\tilde{\rho}_{\alpha_n} \rightarrow k \in (0, +\infty)$ and, again as in the proof of Theorem 3.1, $\tilde{u}_{\alpha_n} \rightarrow \tilde{u}$ in $C_{loc}^2(0, k)$ with

$$\begin{cases} \tilde{u}''(r) + (\nu-1) \frac{\tilde{u}'(r)}{r} = -\frac{\tilde{u}^p(r)}{\lambda} & \text{for } r > 0 \\ \tilde{u}(0) = 1, \quad \tilde{u}'(0) = 0, \quad \tilde{u}(k) = 0. \end{cases} \quad (6.3)$$

Multiplying the ODE in (6.3) by $\tilde{u}(r)r^{\nu-1}$ we have

$$(\tilde{u}'(r)r^{\nu-1})' \tilde{u}(r) + \frac{1}{\lambda} \tilde{u}^{p+1}(r)r^{\nu-1} = 0.$$

Then, integrating by parts from 0 to r , this yields

$$\int_0^r (\tilde{u}')^2 s^{\nu-1} ds = \frac{1}{\lambda} \int_0^r \tilde{u}^{p+1} s^{\nu-1} ds + \tilde{u}(r) \tilde{u}'(r) r^{\nu-1}. \quad (6.4)$$

Similarly, multiplying (6.3) by $\tilde{u}'(r)r^\nu$ and integrating, we obtain

$$\left(\frac{\nu-2}{2} \right) \int_0^r (\tilde{u}')^2 s^{\nu-1} ds = -\frac{(\tilde{u}'(r))^2}{2} r^\nu - \frac{\tilde{u}^{p+1}(r)}{\lambda(p+1)} r^\nu + \frac{\nu}{\lambda(p+1)} \int_0^r \tilde{u}^{p+1} s^{\nu-1} ds. \quad (6.5)$$

From (6.4) and (6.5) we deduce the following Pohozaev-like identity

$$r^\nu \frac{(\tilde{u}'(r))^2}{2} + r^\nu \frac{\tilde{u}^{p+1}(r)}{\lambda(p+1)} + \frac{\nu-2}{2} r^{\nu-1} \tilde{u}(r) \tilde{u}'(r) = \frac{1}{\lambda} \left(\frac{\nu}{p+1} - \frac{\nu-2}{2} \right) \int_0^r s^{\nu-1} \tilde{u}^{p+1} ds,$$

which leads to a contradiction for $r = k$, since $\frac{\nu}{p+1} - \frac{\nu-2}{2} < 0$ in view of the assumption $p > \frac{\nu+2}{\nu-2}$. Hence, $\tilde{\rho}_{\alpha_n} \rightarrow +\infty$ and \tilde{u}_{α_n} converges in $C_{loc}^2((0, +\infty))$ to \tilde{u} , with

$$\begin{cases} \tilde{u}''(r) + (\nu-1) \frac{\tilde{u}'(r)}{r} = -\frac{\tilde{u}^p(r)}{\lambda} & \text{for } r > 0 \\ \tilde{u}(0) = 1, \quad \tilde{u}'(0) = 0 \end{cases} \quad (6.6)$$

Denoting, respectively, by \tilde{x}_{α_n} , \tilde{x} the Emden-Fowler transform (2.3) of \tilde{u}_{α_n} , \tilde{u} , and by $\tilde{\gamma}_{\alpha_n} = (\tilde{x}_{\alpha_n}, \tilde{x}'_{\alpha_n})$, $\tilde{\gamma} = (\tilde{x}, \tilde{x}')$ the associated trajectories, in the phase plane $\tilde{\gamma}_{\alpha_n}(t) \rightarrow \tilde{\gamma}(t)$ as $n \rightarrow +\infty$,

the convergence being $C_{loc}^1(\mathbb{R})$. We observe that in this case \tilde{x}_{α_n} and \tilde{x} are solution, respectively for $t \in (\frac{p-1}{2} \log m_{\alpha_n}, \log \tilde{\rho}_{\alpha_n})$ and $t \in \mathbb{R}$, of

$$x''(t) = ax'(t) + bx(t) - \frac{x^p(t)}{\lambda}, \quad (6.7)$$

where now $a = \frac{\nu+2-(\nu-2)p}{p-1}$ and $b = \frac{2((\nu-2)p-\nu)}{(p-1)^2}$. Since $p > \frac{\nu+2}{\nu-2}$, then $a < 0$ and, by a direct computation, the energy function

$$E(x, y) = \frac{y^2}{2} - b \frac{x^2}{2} + \frac{x^{p+1}}{\lambda(p+1)} \quad (6.8)$$

decreases along $\tilde{\gamma}_{\alpha_n}$ and $\tilde{\gamma}$. Hence

$$E(\tilde{\gamma}_{\alpha_n}(t)) > E(\tilde{\gamma}_{\alpha_n}(\log \tilde{\rho}_{\alpha_n})) = \frac{(\tilde{x}'_{\alpha_n})^2(\log \tilde{\rho}_{\alpha_n})}{2} > 0 \quad \forall t \in \left(\frac{p-1}{2} \log m_{\alpha_n}, \log \tilde{\rho}_{\alpha_n} \right)$$

and

$$E(\tilde{\gamma}(t)) < \lim_{t \rightarrow -\infty} E(\tilde{\gamma}(t)) = E(0, 0) = 0 \quad \forall t \in (-\infty, +\infty).$$

Since $E(\tilde{\gamma}_{\alpha_n}(t)) \rightarrow E(\tilde{\gamma}(t))$ for any fixed $t \in \mathbb{R}$, this leads to a contradiction.

(ii). Let $\alpha > \alpha^*$. Take $\bar{\alpha} \in D$ such that $\bar{\alpha} \in (\alpha^*, \alpha)$ and consider in the phase plane the trajectory $\gamma_{\bar{\alpha}}(t) = (x_{\bar{\alpha}}(t), x'_{\bar{\alpha}}(t))$, where $x_{\bar{\alpha}}$ is the Emden-Fowler transform (2.3) associated with $u_{\bar{\alpha}}$. By construction, $\gamma_{\bar{\alpha}}$ is defined for $t \in [0, \log \rho_{\bar{\alpha}}]$ and satisfies $\gamma_{\bar{\alpha}}(0) = (0, \bar{\alpha})$, $\gamma_{\bar{\alpha}}(\log \rho_{\bar{\alpha}}) = (0, x'_{\bar{\alpha}}(\log \rho_{\bar{\alpha}}))$, with $x'_{\bar{\alpha}}(\log \rho_{\bar{\alpha}}) < 0$, and $x_{\bar{\alpha}}(t) > 0$ for $t \in (0, \log \rho_{\bar{\alpha}})$. By Lemma 2.1 (v), there exists a unique $T \in (0, \log \rho_{\bar{\alpha}})$ such that $x'_{\bar{\alpha}}(T) = 0$ and, moreover, $x_{\bar{\alpha}}(T) > c^*$, where $c^* = (\lambda b)^{\frac{1}{p-1}}$ and $(c^*, 0)$ is the only equilibrium point in the right half plane for equation (6.7).

Let us consider now the trajectory $\gamma_{\alpha}(t) = (x_{\alpha}(t), x'_{\alpha}(t))$. If, by contradiction, $\rho_{\alpha} = +\infty$, then γ_{α} satisfies Lemma 2.1 (vii). On the other hand, in this case equation (6.7) does not admit periodic solutions, due to the monotonicity along trajectories of the energy function (6.8). Hence, $\gamma_{\alpha}(t)$ must approach as $t \rightarrow \infty$ either $(0, 0)$ or $(c^*, 0)$. In both cases, it intersects the trajectory $\gamma_{\bar{\alpha}}$, which is impossible since $\alpha > \bar{\alpha}$. Thus, the only possibility is that γ_{α} leaves the right half plane in finite time, i.e. $\rho_{\alpha} < +\infty$ as desired.

(iii). Since $\alpha^* > 0$ by (i) and the set D is open, then necessarily $\alpha^* \notin D$, namely $\rho_{\alpha^*} = +\infty$.

In view of (ii), for any $\alpha > \alpha^*$ the trajectory $\gamma_{\alpha}(t) = (x_{\alpha}(t), x'_{\alpha}(t))$ crosses the x' -axis at $t = \log \rho_{\alpha}$, with $\log \rho_{\alpha} \rightarrow +\infty$ as $\alpha \rightarrow \alpha^*$ due to the locally uniform convergence of γ_{α} to γ_{α^*} , γ_{α^*} being the trajectory related to u_{α^*} . Moreover, by the monotonicity of the energy function (6.8) along trajectories, one has $E(\gamma_{\alpha}(t)) > 0$ for any $t \in [0, \log \rho_{\alpha}]$, so that $E(\gamma_{\alpha^*}(t)) \geq 0$ for any $t \geq 0$.

Arguing as in the proof of (ii), we have that, as $t \rightarrow +\infty$, the trajectory $\gamma_{\alpha^*}(t)$ converges either to $(0, 0)$ or to $(c^*, 0)$, with $c^* = (\lambda b)^{\frac{1}{p-1}}$. Since $E(c^*, 0) = -\frac{(c^*)^2 b(p-1)}{2(p+1)} < 0$, the latter case is excluded and, by Lemma 2.1 (viii), the conclusion follows.

(iv). Let $\alpha < \alpha^*$. By (6.2), we have $\rho_{\alpha} = +\infty$. Then, arguing as above, either $\gamma_{\alpha}(t) \rightarrow (0, 0)$ or $\gamma_{\alpha}(t) \rightarrow (c^*, 0)$ as $t \rightarrow +\infty$. By contradiction, let us suppose that $\gamma_{\alpha}(t) \rightarrow (0, 0)$. By Lemma 2.1 (viii), this means that the maximal solution u_{α} is of (6.1) satisfies $\lim_{r \rightarrow +\infty} r^{\nu-2} u_{\alpha}(r) = C$ and $\lim_{r \rightarrow \infty} r^{\nu-1} u'_{\alpha}(r) = -(\nu-2)C$, for a positive constant C .

Let us now consider the Kelvin transform of u_{α} , defined by

$$\bar{u}_{\alpha}(r) := \begin{cases} r^{2-\nu} u_{\alpha}\left(\frac{1}{r}\right) & \text{if } r \in (0, 1] \\ C & \text{if } r = 0. \end{cases} \quad (6.9)$$

By a straightforward computation, one checks that \bar{u}_α is a bounded solution of

$$\begin{cases} u''(r) + (\nu - 1)\frac{u'(r)}{r} = -\frac{1}{\lambda}r^{p(\nu-2)-\nu-2}u^p(r) & \text{for } r \in (0, 1), \\ u(1) = 0. \end{cases} \quad (6.10)$$

Moreover, one has $\frac{\bar{u}'_\alpha(r)}{r} \rightarrow 0$ as $r \rightarrow 0^+$, so that \bar{u}_α can be extended to a solution of (6.10) for $r \in [0, 1)$. Indeed, using the identity

$$(r^{\nu-1}\bar{u}'_\alpha)' = -\frac{1}{\lambda}r^{p(\nu-2)-3}\bar{u}_\alpha^p, \quad r \in (0, 1), \quad (6.11)$$

we first infer that $r^{\nu-1}\bar{u}'_\alpha$ is monotone decreasing. Furthermore, by definition (6.9) and by the decaying property as $r \rightarrow +\infty$ of $u_\alpha(r)$, $u'_\alpha(r)$, we have

$$\lim_{r \rightarrow 0^+} r^{\nu-1}\bar{u}'_\alpha(r) = -\lim_{r \rightarrow +\infty} (u'_\alpha(r)r + (\nu - 2)u_\alpha(r)) = 0.$$

Integrating (6.11), then we get

$$r^{\nu-1}\bar{u}'_\alpha(r) = -\frac{1}{\lambda} \int_0^r s^{p(\nu-2)-3}\bar{u}_\alpha^p(s) ds$$

and

$$\left| \frac{\bar{u}'_\alpha(r)}{r} \right| = \frac{1}{\lambda r^\nu} \int_0^r s^{p(\nu-2)-3}\bar{u}_\alpha^p(s) ds \leq \frac{\|\bar{u}_\alpha^p\|_\infty}{\lambda(p(\nu-2)-2)} r^{p(\nu-2)-\nu-2} \rightarrow 0 \quad \text{as } r \rightarrow 0^+.$$

Then \bar{u}_α is solution of (6.10) also for $r \rightarrow 0^+$.

By exactly the same argument, the Kelvin transform \bar{u}_{α^*} of u_{α^*} also is a solution of (6.10) for $r \in [0, 1)$. But the proof of [8, Proposition 5.2] shows that (6.10) has a unique positive solution in $[0, 1)$. Indeed, the proof of [8] is purely based on the analysis of the ODE problem (6.10), and the used arguments do not depend on the first order coefficient to be natural or a real number larger than 1. Hence, by uniqueness, we have $u_\alpha(r) = u_{\alpha^*}(r)$ for all $r \geq 1$ and this is a contradiction since $\alpha < \alpha^*$. \square

Proof of Theorem 6.1. Assume by contradiction that problem (1.1), with $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$ and $p = p_-^*$, admits a radial solution $u_\alpha = u_\alpha(r)$ with $u'_\alpha(1) = \alpha > 0$. As in the proof of Theorem 5.1, let us consider the unique (up to scaling) entire radial solution u_* of $-\mathcal{M}_{\lambda, \Lambda}^-(D^2u) = u^{p_-^*}$ in \mathbb{R}^N (see [10]). After the Emden-Fowler transform (2.3), in the phase plane u_α and u_* correspond to the trajectories $\gamma_\alpha(t) = (x_\alpha(t), x'_\alpha(t))$ and $\gamma_*(t) = (x_*(t), x'_*(t))$, defined respectively for $t \in [0, +\infty)$ and $t \in \mathbb{R}$. Note that $\gamma_\alpha(t)$ and $\gamma_*(t)$ cannot intersect and, as in the proof of Theorem 5.1, we obtain that both $\gamma_\alpha(t)$ and $\gamma_*(t)$ converge to $(0, 0)$ as $t \rightarrow +\infty$.

By Lemma 2.1 (iii) and (v), we have that there exist a unique $s_\alpha > 0$ and a unique $s_* \in \mathbb{R}$ such that $\gamma_\alpha(s_\alpha), \gamma_*(s_*) \in \mathcal{C}$. Moreover, $\gamma_\alpha(t)$ and $\gamma_*(t)$ cross transversally from above the curve \mathcal{C} for $t = s_\alpha$ and $t = s_*$ respectively, with $x_{\alpha^-}^{p_-^*-1}(s_\alpha), x_{*^-}^{p_-^*-1}(s_*) > \frac{\Lambda(N-1)}{\lambda p_-^*}$, and, by (2.4) with $p = p_-^*$ (and λ and Λ interchanged), $x_\alpha(t)$ and $x_*(t)$ satisfy, respectively for $t \geq s_\alpha$ and $t \geq s_*$,

$$x''(t) = \tilde{a}_-x'(t) + \tilde{b}_-x(t) - \frac{x^{p_-^*}(t)}{\lambda}, \quad (6.12)$$

where \tilde{a}_-, \tilde{b}_- are defined in (2.5). We recall that along any trajectory $(x(t), x'(t))$ associated with a solution x of (6.12), the energy function E defined by

$$E(x, y) = \frac{y^2}{2} - \tilde{b}_- \frac{x^2}{2} + \frac{x^{p_-^*+1}}{\lambda(p_-^*+1)} \quad (6.13)$$

is monotone decreasing, since $\tilde{a}_-(p_-^*) < 0$. Thus, equation (6.12) does not admit periodic orbits and its only equilibrium points in the phase plane are $(0, 0)$ and $(c^*, 0)$, with $c^* = (\tilde{b}_-\lambda)^{\frac{1}{p_-^*-1}}$.

Let us now consider the maximal positive solution $\bar{x}_\alpha = \bar{x}_\alpha(t)$ of the semilinear initial value problem

$$\begin{cases} x''(t) = \tilde{a}_-x'(t) + \tilde{b}_-x(t) - \frac{x^{p_-^*}(t)}{\lambda}, \\ x(s_\alpha) = x_\alpha(s_\alpha), \quad x'(s_\alpha) = x'_\alpha(s_\alpha). \end{cases}$$

Similarly, let $\bar{x}_* = \bar{x}_*(t)$ be the maximal positive solution of

$$\begin{cases} x''(t) = \tilde{a}_-x'(t) + \tilde{b}_-x(t) - \frac{x^{p_-^*}(t)}{\lambda}, \\ x(s_*) = x_*(s_*), \quad x'(s_*) = x'_*(s_*). \end{cases}$$

In other words, we look at the problems with initial conditions given by $\gamma_\alpha(s_\alpha)$ and $\gamma_*(s_*)$, and with an equation having fixed coefficients given by \tilde{a}_- , \tilde{b}_- and $1/\lambda$.

Clearly, $\bar{\gamma}_\alpha(t) := (\bar{x}_\alpha(t), \bar{x}'_\alpha(t)) = \gamma_\alpha(t)$ for $t \geq s_\alpha$ and $\bar{\gamma}_*(t) := (\bar{x}_*(t), \bar{x}'_*(t)) = \gamma_*(t)$ for $t \geq s_*$. Let us analyze the behavior of $\bar{\gamma}_\alpha(t)$ and $\bar{\gamma}_*(t)$ for $t < s_\alpha$ and $t < s_*$ respectively. We will show that there exist finite times $\bar{s}_\alpha < s_\alpha$ and $\bar{s}_* < s_*$ such that $\bar{\gamma}_\alpha(t)$ and $\bar{\gamma}_*(t)$ reach the x' -axis for $t = \bar{s}_\alpha$ and $t = \bar{s}_*$ respectively.

Let us consider only the trajectory $\bar{\gamma}_*(t)$, the arguments for $\bar{\gamma}_\alpha$ being exactly the same.

We observe that, by the transversality of the intersection with \mathcal{C} at the point $\bar{\gamma}_*(s_*)$, the trajectory $\bar{\gamma}_*(t)$ belongs to the region \mathcal{R}^+ for t in a left neighborhood of s_* . Let us first show that $\bar{\gamma}_*(t) \in \mathcal{R}^+$ for $t < s_*$. Indeed, assume, by contradiction, that there exists $t_* < s_*$ such that $\bar{\gamma}_*(t_*) \in \mathcal{C}$ and $\bar{\gamma}_*(t) \in \mathcal{R}^+$ for $t \in (t_*, s_*)$. By Lemma 2.1 (iii) and (iv) applied to $\bar{\gamma}_*$, we obtain that $\bar{x}_*^{p_-^*-1}(t_*) \leq \frac{\lambda(N-1)}{p_-^*}$ and that either $\bar{\gamma}_*(t) \in \mathcal{R}^-$ for all $t < t_*$ or $\bar{\gamma}_*$ intersects again \mathcal{C} at some point on the right of $\frac{\lambda(N-1)}{p_-^*}$. In any case, $\bar{\gamma}_*(t)$ is bounded as $t \rightarrow -\infty$, and the Poincaré-Bendixson Theorem implies that either $\bar{\gamma}_*(t) \rightarrow (0, 0)$ or $\bar{\gamma}_*(t) \rightarrow (c^*, 0)$ as $t \rightarrow -\infty$. In both cases we reach a contradiction, since at the equilibrium points the energy (6.13) is non positive, whereas along $\bar{\gamma}_*(t)$ it is decreasing and it converges to 0 as $t \rightarrow +\infty$. This proves the claim.

In order to show that $\bar{\gamma}_*(t)$ leaves the right half-plane in a finite time, let us use relation (2.3) to switch from $\bar{x}_*(t)$ to $\bar{u}_*(r)$, which is the maximal positive solution of the initial value problem

$$\begin{cases} u'' + (\tilde{N}_- - 1)\frac{u'}{r} = -\frac{u^{p_-^*}}{\lambda} \\ u(\sigma_*) = u_*(\sigma_*), \quad u'(\sigma_*) = u'_*(\sigma_*) \end{cases}$$

where $\sigma_* = e^{s_*}$ is the only zero of $u''(r)$. The solution \bar{u}_* is defined on a maximal interval $(\bar{r}_*, +\infty)$, with $0 \leq \bar{r}_* < \sigma_*$.

The fact that $\bar{\gamma}_*(t) \in \mathcal{R}^+$ for all $t < s_*$ means that \bar{u}_* is a concave function in $(\bar{r}_*, \sigma_*]$. Then, \bar{u}_* is bounded. If $(\bar{u}_*)'(r) < 0$ for all $r \in (\bar{r}_*, \sigma_*]$, then $(\bar{u}_*)'$ also is bounded, since it is decreasing, and, therefore, $\bar{\gamma}_*(t)$ is bounded for $t < \bar{s}_*$. Moreover, for $t < \bar{s}_*$, $\bar{\gamma}_*(t)$ lays in the portion of \mathcal{R}^+ below the line L and, again by Poincaré-Bendixson Theorem, it follows that $\bar{\gamma}_*(t) \rightarrow (0, 0)$ as $t \rightarrow -\infty$, which is impossible. Therefore, there exists $\bar{r}_* \in (\bar{r}_*, \sigma_*)$ such that $(\bar{u}_*)'(\bar{r}_*) = 0$ and, since $(\bar{u}_*)''(\bar{r}_*) = -\bar{u}_*^{p_-^*}(\bar{r}_*)/\lambda < 0$, it follows that $(\bar{u}_*)'(r) > 0$ for $r \in (\bar{r}_*, \bar{r}_*)$.

By writing, as usual, the equation satisfied by \bar{u}_* in the form

$$\left(r^{\tilde{N}_- - 1} \bar{u}'_* \right)'(r) = -\frac{r^{\tilde{N}_- - 1} \bar{u}_*^{p_-^*}(r)}{\lambda} < 0, \quad (6.14)$$

we further notice that if, $\bar{r}_* = 0$, then, for any fixed $r_0 \in (0, \bar{r}_*)$, one has

$$(\bar{u}_*)'(r) \geq \frac{r_0^{\tilde{N}_- - 1} (\bar{u}_*)'(r_0)}{r^{\tilde{N}_- - 1}} \quad \text{for all } r \in (0, r_0],$$

and this contradicts the boundedness of \bar{u}_* . Thus $\bar{r}_* > 0$, and from (6.14), it necessarily follows that

$$\bar{u}_*(\bar{r}_*) = 0, \quad (\bar{u}_*)'(\bar{r}_*) =: \bar{\alpha}_* \in (0, +\infty).$$

The same arguments show that $\bar{x}_\alpha(t)$ also is the Emden-Fowler transform of a function $\bar{u}_\alpha(r)$ which is the maximal positive solution of

$$\begin{cases} u'' + (\tilde{N}_- - 1) \frac{u'}{r} = -\frac{u^{p^*}}{\lambda} & \text{for } r > \bar{r} \\ u(\bar{r}) = 0, \quad u'(\bar{r}) = \bar{\alpha} \end{cases}$$

for some $\bar{r}, \bar{\alpha} > 0$. Then, the two functions \bar{u}_* and \bar{u}_α , suitably rescaled, are two distinct fast decaying solutions of problem (6.1) with $\nu = \tilde{N}_-$ and $p = p^* > \frac{\tilde{N}_- + 2}{\tilde{N}_- - 2}$, in contradiction with Proposition 6.2. \square

7. EXISTENCE AND UNIQUENESS OF FAST DECAYING SOLUTIONS

In this section we complete the proof of Theorem 1.1, by showing in particular the existence and the uniqueness of fast decaying solutions of problems (1.1) for $p > p_\pm^*$.

For any $\alpha > 0$, let u_α be the unique maximal solution of either (2.1) or (2.2), defined on the maximal interval $[1, \rho_\alpha)$ with $\rho_\alpha \leq +\infty$. As for the proof of Proposition 6.2, let us set

$$D = \{\alpha \in (0, +\infty) : \rho_\alpha < +\infty\},$$

and

$$\alpha^* = \alpha^*(p) = \inf D.$$

Arguing as in the proof of Proposition 6.2, one has $D = (\alpha^*, +\infty)$. Moreover, by Theorems 3.1, 4.1, 5.1 and 6.1, we have that $\alpha^*(p) > 0$ if and only if $p > p_\pm^*$.

Thus, for $p > p_\pm^*$ and $0 < \alpha \leq \alpha^*$, one has $\rho_\alpha = +\infty$. Let $\gamma_\alpha(t) = (x_\alpha(t), x'_\alpha(t))$ be the trajectory associated with the Emden-Fowler transform $x_\alpha(t)$ of u_α , given by (2.3).

By continuous dependence on the initial data, $\gamma_\alpha(t) \rightarrow \gamma_{\alpha^*}(t)$ in $C_{loc}^1([0, +\infty))$ as $\alpha \rightarrow \alpha^*$, and, since $\gamma_\alpha(t)$ reaches the x' -axis in a finite time for $\alpha > \alpha^*$, it follows necessarily that $\gamma_{\alpha^*}(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$. Indeed, by Lemma 2.1 and the locally uniform convergence, it follows that $s_\alpha \rightarrow s_{\alpha^*}$ and $\gamma_{\alpha^*}(t) \in \mathcal{R}^-$ for all $t > s_{\alpha^*}$. This excludes the convergence of $\gamma_{\alpha^*}(t)$ to a periodic orbit as $t \rightarrow +\infty$, and the convergence to the equilibrium point $(c^*, 0)$ (here $c^* = (\tilde{b}_+ \Lambda)^{1/(p-1)}$ if $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^+$ and $c^* = (\tilde{b}_- \lambda)^{1/(p-1)}$ if $\mathcal{F} = \mathcal{M}_{\lambda, \Lambda}^-$) is excluded as well, by arguing as in the conclusion of the proof of Theorem 3.1 and taking into account also that γ_{α^*} cannot intersect the trajectory γ associated with entire solutions. Therefore, $u_{\alpha^*}(r)$ is a fast decaying solution.

Next, in order to prove that u_{α^*} is the only fast decaying solution, let us consider separately the cases of $\mathcal{M}_{\lambda, \Lambda}^+$ and $\mathcal{M}_{\lambda, \Lambda}^-$.

Assume first that u_α solves (2.1) and, by contradiction, suppose that $\gamma_\alpha(t) \rightarrow (0, 0)$ as $t \rightarrow +\infty$ for some $\alpha < \alpha^*$. If $p_+^* < p < \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$, then we can apply the argument of the proof of Theorem 5.1 to the two trajectories $\gamma_{\alpha^*}(t)$ and $\gamma_\alpha(t)$, reaching a contradiction. On the other hand, if $p = \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$, then the trajectory associated with any fast decaying solution lies for sufficiently large t on the zero level set of the energy function E given by (2.9). Hence, $\gamma_{\alpha^*}(t)$ and $\gamma_\alpha(t)$ definitively coincide, again a contradiction. Finally, if $p > \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$, then we can argue as in the

proof of Theorem 6.1 and, thanks to Proposition 6.2 for $\nu = \tilde{N}_+$, we obtain a contradiction as well.

For the operator $\mathcal{M}_{\lambda,\Lambda}^-$, the uniqueness of the fast decaying solution follows by the same proof of Theorem 6.1.

Therefore, in both cases, we obtain that, for $\alpha < \alpha^*$, the trajectories $\gamma_\alpha(t)$ cannot approach the origin as $t \rightarrow +\infty$. Furthermore, since, for any α , $\gamma_\alpha(t)$ cannot intersect the trajectory $\gamma(t)$ associated with any entire solution u of $-\mathcal{M}_{\lambda,\Lambda}^\pm(D^2u) = u^p$, and since, for $p > p_\pm^*$ and $\alpha < \alpha^*$, neither u nor u_α is a fast decaying solution, it follows that $\gamma(t)$ and $\gamma_\alpha(t)$ must have the same behavior as $t \rightarrow +\infty$.

According to Theorem 1.1 of [10], we then obtain the following result.

Theorem 7.1. *Let $p > p_+^*$ and let u_α denote the maximal solution of (2.1). Then:*

- (i) u_{α^*} is a fast decaying solution;
- (ii) for any $\alpha < \alpha^*$, u_α is a pseudo-slow decaying solution if $p_+^* < p \leq \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$;
- (iii) for any $\alpha < \alpha^*$, u_α is a slow decaying solution if $p > \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2}$.

Analogously, by applying Theorem 1.2 of [10], we deduce the following theorem.

Theorem 7.2. *Let $p > p_-^*$ and let u_α denote the maximal solution of (2.2). Then:*

- (i) u_{α^*} is a fast decaying solution;
- (ii) for any $\alpha < \alpha^*$, u_α is either a pseudo-slow or a slow decaying solution if $p_-^* < p \leq \frac{N+2}{N-2}$;
- (iii) for any $\alpha < \alpha^*$, u_α is a slow decaying solution if $p > \frac{N+2}{N-2}$.

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