

# Effective Visible Light Exploitation by Copper Molybdo-tungstate Photoanodes

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## ABSTRACT

The need for stable oxide-based semiconductors with narrow band gap, able to maximize the exploitation of the visible light portion of the solar spectrum, is a challenging issue for photoelectrocatalytic (PEC) applications. In the present work  $\text{CuW}_{1-x}\text{Mo}_x\text{O}_4$  ( $E_g = 2.0$  eV for  $x = 0.5$ ), exhibiting a significantly reduced optical band gap  $E_g$  compared to isostructural  $\text{CuWO}_4$  ( $E_g = 2.3$  eV), was investigated as photoactive material for the preparation of photoanodes.  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes with different thickness (80 – 530 nm), prepared by a simple solution-based process in the form of multilayer films, effectively exhibit visible light photoactivity up to 650 nm, i.e., extended compared to  $\text{CuWO}_4$  photoanodes prepared by the same way. Furthermore, the systematic investigation on the effects on photoactivity of the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  layer thickness evidenced that long wavelength photons can better be exploited by thicker electrodes. PEC measurements in the presence of  $\text{NaNO}_2$ , acting as a suitable hole scavenger ensuring enhanced photocurrent generation compared to that of water oxidation while minimizing dark currents, allowed us to elucidate the role that molybdenum incorporation plays on the charge separation efficiency in the bulk and on the charge injection efficiency at the photoanode surface. The adopted Mo for W substitution increases the visible light photoactivity of copper tungstate towards improved exploitation and storage of visible light into chemical energy *via* photoelectrocatalysis.

**KEYWORDS:** *Mo modified  $\text{CuWO}_4$ , band gap reduction, extended visible light photoactivity, photoelectrocatalysis, photoanode, charge separation*

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## 1. Introduction

Our need of storable forms of energy is continuously growing,<sup>1,2</sup> as well as the attention towards renewable energy sources, in order to prevent crucial fossil fuel-related issues such as climate change and air pollution. The sun is the primary source of energy for our planet and the evolution towards a solar energy society, already envisaged long time ago, has now become an inescapable choice.<sup>3</sup> Mimicking the action of bacteria and plants which are able to convert solar energy into nutrients *via* photosynthesis,<sup>4</sup> effective strategies were foreseen to convert solar light into energy-containing chemicals.<sup>1,4-6</sup> Among them, the photoelectrocatalytic (PEC) conversion of solar light into highly energetic chemical fuels was proposed, in thermodynamically uphill reactions such as hydrogen production from water splitting,<sup>7,8</sup> and intense efforts were made in identifying photoactive materials for the fabrication of efficient photoelectrodes, in particular photoanodes for water oxidation, as this reaction represents the kinetic bottleneck of the overall PEC water splitting process.<sup>9</sup>

Ternary oxides, such as  $\text{BiVO}_4$ ,  $\text{MFe}_2\text{O}_4$  ( $\text{M} = \text{Cu}, \text{Mg}, \text{Zn}$ ),  $\text{InTaO}_4$  and  $\text{CuWO}_4$ , emerged as promising candidate materials for this application,<sup>10-17</sup> because of their high stability towards the harsh oxidative conditions required for oxygen evolution,<sup>18</sup> their valence band edge position being lower in energy with respect to the water oxidation potential of 1.23 V *vs.* RHE.<sup>19</sup> Commonly employed  $\text{BiVO}_4$  and  $\text{MFe}_2\text{O}_4$  oxides with a relatively narrow band gap may allow up to 8 and 10% harvesting of the solar radiation, respectively, which is still not suitable for industrial applications.<sup>13,19</sup> However, their efficient use as well as that of other ternary oxides such as  $\text{InTaO}_4$  and  $\text{CuWO}_4$  (having 2.6 eV<sup>10</sup> and 2.3 eV<sup>20</sup> band gap, respectively) requires their modification to narrow their band gap, for an effective solar energy conversion. Attempts of doping these materials with elements such as Ni,<sup>21</sup> N,<sup>22</sup> Cr,<sup>22-24</sup> and Mo<sup>25,26</sup> proved to increase their visible light absorption capability. In particular, partial substitution of  $\text{W}^{6+}$  with  $\text{Mo}^{6+}$  in the  $\text{CuWO}_4$  structure results in a band gap reduction of *ca.* 0.3 eV (from 2.3 to 2.0 eV)<sup>25</sup> corresponding to a significant red shift of the absorption onset of the ternary oxide material.<sup>25-27</sup>

Aiming at ascertaining the role that molybdenum for tungsten substitution has in increasing the PEC performance of intrinsically poorly performing  $\text{CuWO}_4$ , in relation to the transport properties of the charge carriers photogenerated in differently thick photoactive layers and to the charge transfer efficiency at the oxide-solution interface, in the present work we investigate the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  semiconductor oxide obtained by 50%  $\text{Mo}^{6+}$  for  $\text{W}^{6+}$  substitution, which is quite easy to achieve, the two ions having very similar radii.<sup>28</sup> This material was prepared by a facile solution-based synthetic route and is composed of a single phase. Its thorough PEC characterization allowed us to demonstrate that the attained reduction in band gap energy directly results in a more efficient exploitation of longer wavelength photons compared to pure  $\text{CuWO}_4$  photoanodes. The role that molybdenum incorporation has on the efficiencies of photoproduced charge separation in the bulk and charge injection at the electrode/electrolyte interface has been clarified by investigating how the thickness of the photoactive layer in photoanodes affects their wavelength-dependent PEC performance, also in the presence of a suitable hole scavenger.

## **2. Experimental**

### **2.1. Chemicals and materials**

The following chemicals, all purchased from Sigma Aldrich, were employed as supplied: copper(II) nitrate trihydrate (99%,  $\text{Cu}(\text{NO}_3)_2 \cdot 3\text{H}_2\text{O}$ ), ammonium metatungstate hydrate (99%,  $(\text{NH}_4)_6\text{H}_2\text{W}_{12}\text{O}_{40} \cdot x\text{H}_2\text{O}$ ), citric acid (99%), boric acid (99%) and ethanol (99%). Molybdenum(VI) oxide bis(2,4-pentanedionate) (99%,  $\text{C}_{10}\text{H}_{14}\text{MoO}_6$ ) was an Alfa Aesar product. Fluorine-doped 2 mm thick tin oxide (FTO) glass was purchased from Pilkington Glass (TEC-7).

### **2.2. Photoelectrodes preparation**

A 0.5 M solution of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  was prepared as follows. 0.270 g of citric acid, 0.122 g of copper nitrate, 0.062 g of ammonium metatungstate and 0.082 g of molybdenum oxide bis-pentanedionate were added to 1.0 mL of an ethanol-water 2:1 solution. Complete dissolution of the metal precursors,

corresponding to a 1:1 W:Mo molar ratio, was ensured in the chosen solvent composition and any phase segregation was excluded in the resulting film. The precursors were dissolved by keeping the solution under constant stirring for 45 min at 80 °C. The so obtained green paste is stable for several weeks. The photoelectrodes were prepared by spin coating it onto a  $2.5 \times 2.5 \text{ cm}^2$  FTO glass at 4000 rpm for 30 s. Prior to deposition, the FTO glass was cleaned by 30 min-long sonication in a soap solution, followed by careful washing, sonication in ethanol for 30 min and drying in air. The clean glass slices then underwent a 15 min-long UV-cleaner ozone treatment to remove any organic species deposited onto the FTO surface. They were finally soaked in isopropanol for a few seconds right before the spin coating deposition, to increase the FTO affinity for the metal oxide precursor solution and reduce the light scattering of the resulting films in the long wavelengths spectral region.

After deposition, the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  films were dried at 250 °C for 10 min, followed by annealing at 550 °C for 1 h. Pure  $\text{CuWO}_4$  films were prepared by a similar procedure without adding the molybdenum precursor to the initial solution. Multilayer films (up to 5 layers), labeled as nL ( $n = 1-5$ )  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$ , were prepared by coating the so deposited films with the same precursor solution, followed by annealing at 550°C after each deposition step.

### **2.3. Optical, structural, morphological and photoelectrochemical tests**

UV–visible absorption spectra were recorded in the transmission mode using a Jasco V-670 spectrophotometer. The crystalline phase of the materials was investigated through X-ray powder diffraction (XRPD) analysis using a Philips PW1820 diffractometer, equipped with a Cu sealed tube that provided  $\text{K}\alpha$  radiation at 40 mA and 40 kV. A model LEO 1430 scanning electron microscope operating at a 10 kV accelerating voltage and at 8 mm working distance was used to acquire the top view and cross-section images of the films, up to three deposited layers. A Dektak XT Bruker profilometer was employed to measure the thickness of the thicker 4L and 5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  films.

Linear Sweep Voltammetry (LSV) measurements were carried out using a three electrode cell equipped with two quartz windows. The FTO/CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> film was used as working electrode, an Ag/AgCl (3.0 M NaCl) as reference electrode and a platinum gauze as counter electrode. The electrical bias was swept at 10 mV s<sup>-1</sup> using an Autolab PGSTAT 12, controlled by the NOVA software. The light source was an Oriel, Model 81172 solar simulator equipped with an AM 1.5 G filter. The light intensity, measured by means of a Thorlabs PM200 power meter equipped with a S130VC power head with Si detector, was 100 mW cm<sup>-2</sup>. During a typical LSV test, 5 consecutive *J-V* scans were performed with each electrode, preceded by Fermi level equilibrium under irradiation up to open circuit potential (OCP). Identical OCP values were recorded prior to the beginning of each scan with each mono- or multi-layer electrode and all of them provided stable and reproducible photocurrent from the first to the last scan.

The investigated films were tested as photoanodes under both back (through the FTO side) and front (through the deposited film side) irradiation configuration, in contact with a 0.1 M K<sub>3</sub>BO<sub>3</sub> aqueous solution at pH 9.<sup>29</sup> The buffer borate solution was prepared by adding KOH to aqueous boric acid up to the desired pH. Furthermore, a series of different sacrificial agents, acting as electron donor species, was employed in LSV measurements. 0.1 M aqueous solutions of either H<sub>2</sub>O<sub>2</sub>, NH<sub>3</sub> or NaNO<sub>2</sub> were buffered at pH 9 in a K<sub>3</sub>BO<sub>3</sub> solution, while the 0.1 M Na<sub>2</sub>SO<sub>3</sub> aqueous solution did not contain K<sub>3</sub>BO<sub>3</sub> as it is naturally at pH 9. The potential values *vs.* Ag/AgCl were converted into the RHE scale using the following equation:  $E_{\text{RHE}} = E_{\text{AgCl}} + 0.059 \text{ pH} + E^{\circ}_{\text{AgCl}}$ , with  $E^{\circ}_{\text{AgCl}} (3.0 \text{ M NaCl}) = 0.210 \text{ V}$  at 25 °C.

Incident photon to current efficiency (IPCE) measurements were carried out at 1.23, 1.5 and 1.7 V *vs.* RHE under irradiation with a 300 W Lot-Oriel Xe lamp equipped with a Lot-Oriel Omni-λ 150 monochromator and a Thorlabs SC10 automatic shutter, in the above described single-compartment three-electrode cell containing the K<sub>3</sub>BO<sub>3</sub> buffered solution, to which the NaNO<sub>2</sub> hole scavenger was eventually added. The IPCE values were calculated using the following equation:

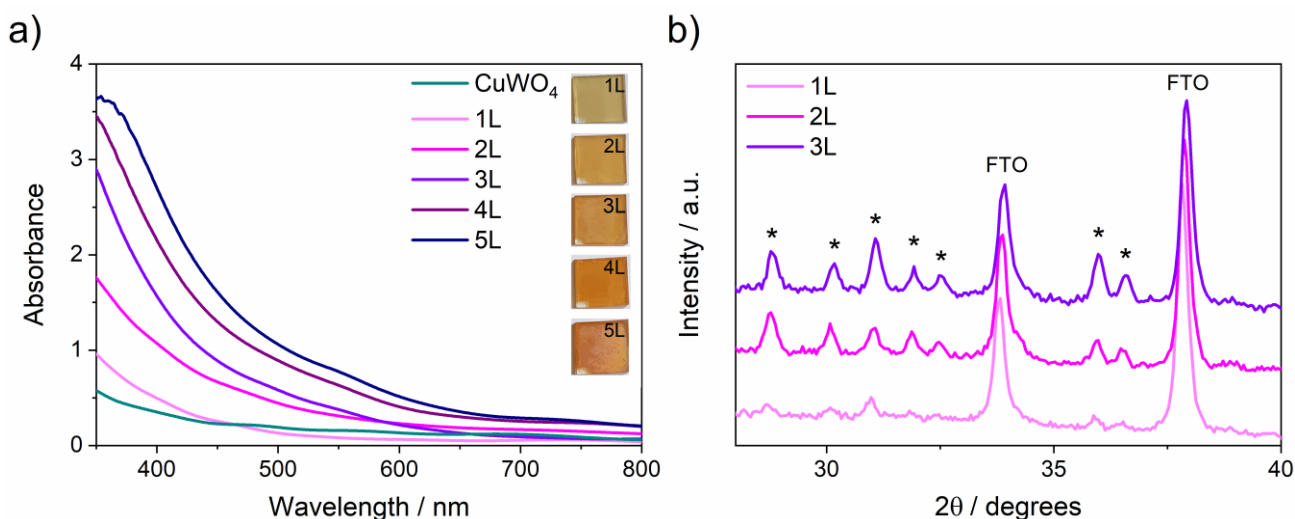
$$\text{IPCE} = \frac{1240 \times J}{P_{\lambda} \times \lambda} \times 100 \quad (1)$$

where  $J$  is the photocurrent density ( $\text{mA cm}^{-2}$ ) and  $P_{\lambda}$  ( $\text{mW cm}^{-2}$ ) is the power measured at each specific wavelength  $\lambda$  (nm).

Chopped chronoamperometric scans at different wavelengths were recorded in back side configuration at 1.23 V *vs.* RHE within the 300-650 nm wavelength range, with a 10 nm step. A 420 nm filter was employed at  $\lambda > 500$  nm, to avoid any contribution from the high-order harmonics originated from the monochromator.

### 3. Results and discussion

#### 3.1. Photoelectrodes characterization



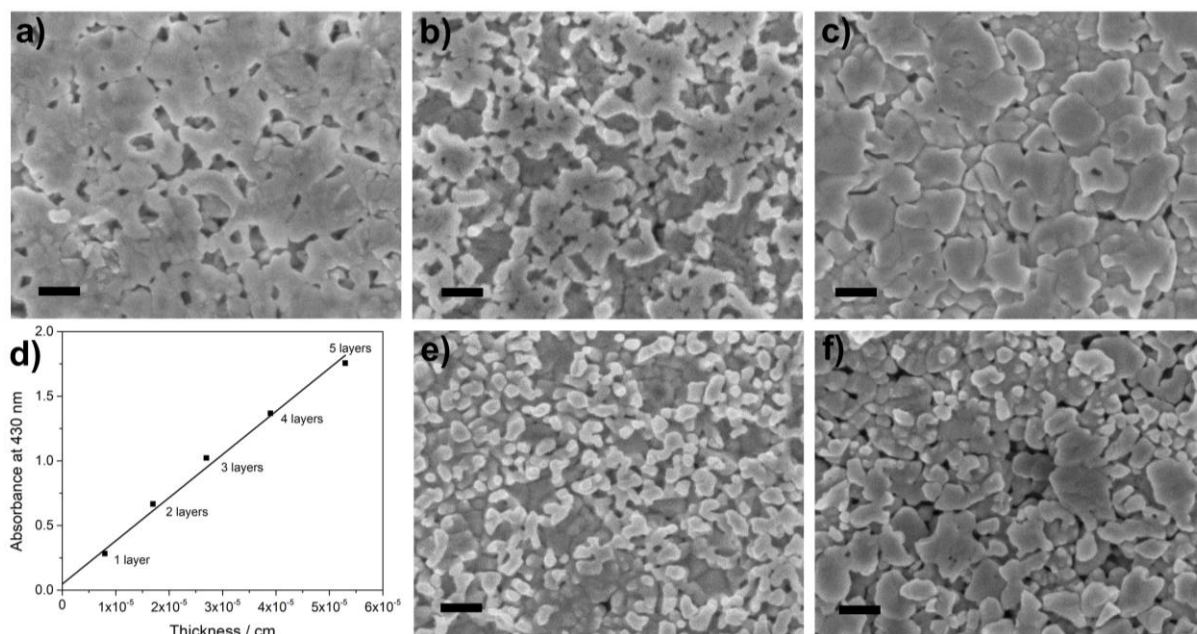
**Figure 1.** (a) Absorption spectra of multilayer (1L–5L)  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  photoelectrodes. Inset: picture of the photoelectrodes. (b) X-ray powder diffraction (XRPD) patterns of 1L, 2L and 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes, in the 28 – 40 degrees  $2\theta$  range. The asterisks mark the reflections typical of wolframite.

As shown in Figure 1a, the multilayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes exhibit an absorption increasing with increasing number of deposited layers, with an absorption tail due to light scattering observed with the thickest films. The gradual colour increase can be appreciated also from their pictures shown in the inset

of Figure 1a. The absorption onset of the material is at *ca.* 650 nm, corresponding to a band gap of *ca.* 1.9 eV, in agreement with the *ca.* 2.0 eV estimated band gap reported for a 50% degree of Mo for W substitution.<sup>25</sup> Thus, the absorption onset of pure copper tungstate, having a band gap energy of 2.3 eV,<sup>30</sup> is effectively extended towards the visible light region upon 50% Mo for W substitution.

The XRPD patterns relative to the 1L, 2L and 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  films, reported in Figure 1b, show the reflections characteristic of pure wolframite,<sup>25,26</sup> which become more intense with increasing the material thickness. Thus, 50% Mo for W substitution has no effects on the crystalline phase formation, within the detection limits of the XRPD technique. The presence of Mo is confirmed by the change of lattice constants and the consequent peak shifts observed in the XRPD patterns of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  in comparison with  $\text{CuWO}_4$  (see a detailed view of Figure 1b reported in Figure S1 of the Supporting Information), which are almost identical to those recently reported<sup>26</sup> for a  $\text{CuW}_{1-x}\text{Mo}_x\text{O}_4$  sample with the same nominal Mo:W molar ratio.

The top view scanning electron microscopy (SEM) images of 1L, 2L and 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes (Figure 2a–c) reveal a structure composed of densely aggregated crystallites, which may facilitate the electrolyte diffusion across the film by increasing the contact area between the electrolyte solution and the photoactive material. In the case of the thickest film a decrease in the agglomerates size can be observed in the top view SEM images acquired with the same electrode after PEC tests (see Figure 2e–f), resulting from an extended contact between the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  film and the electrolyte.



**Figure 2.** Top-view SEM images of (a) 1L, (b) 2L and (c) 3L CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrodes. (d) Absorption at 430 nm of multilayer photoanodes vs. their thickness. SEM images recorded after PEC tests of (e) 2L and (f) 3L CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrodes. The scalebar is 200 nm.

The thickness values of the multilayer CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrodes, evaluated either by means of SEM cross section images (Figure S2) or by profilometry, are reported in Table 1 and indicate that the average increase in film thickness is *ca.* 110 nm upon each deposited layer. From these values, together with the absorption spectrum of the films, the absorption coefficients of CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> at different wavelengths were calculated according to the Lambert Beer law, as the slopes of the lines obtained by plotting the absorbance of the film at a selected wavelength against the film thickness of multilayer CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrodes (see Figures 2d and S3). The best fitting was obtained for  $\lambda = 430$  nm, the corresponding absorption coefficient being  $\alpha_{430\text{nm}} = (3.41 \pm 0.10) \times 10^4 \text{ cm}^{-1}$ .



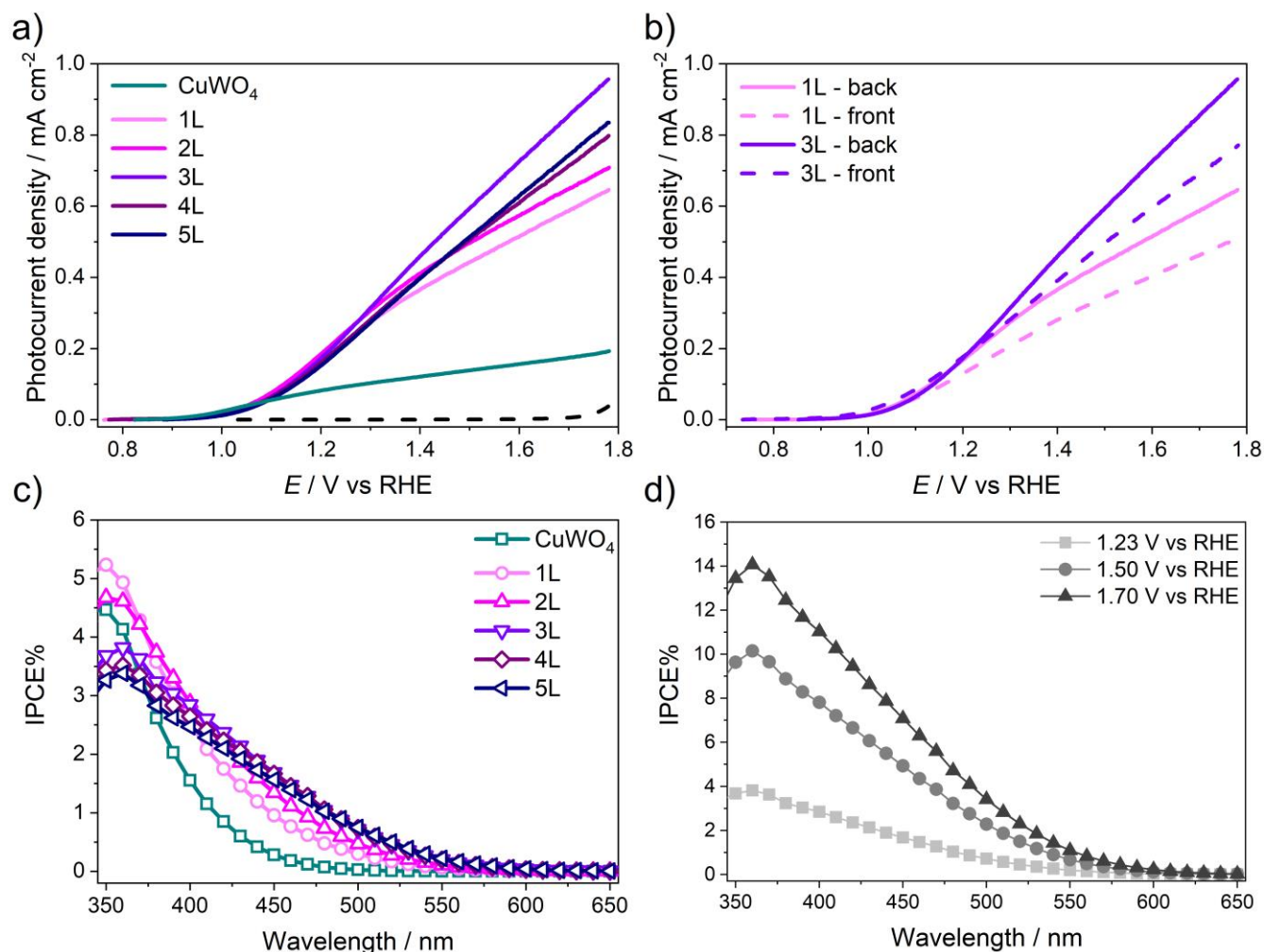
**Table 1.** Thickness of multilayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes estimated from SEM cross-section images and by profilometry.

Sample	Thickness / nm
1L	$80 \pm 20$
2L	$170 \pm 10$
3L	$270 \pm 10$
4L	$410 \pm 20$
5L	$530 \pm 30$

### 3.2. PEC performances

During a typical LSV test under simulated solar light irradiation, 5 consecutive photocurrent density vs. applied potential ( $J$ - $V$ ) scans were performed with each electrode. The  $J$ - $V$  curves reported in Fig. 3ab correspond to the fifth scan performed with each electrode. Each mono- or multi-layer electrode provided stable and reproducible photocurrent from the first to the last scan, with no difference in photocurrent, as shown in Figure S4. Thus high repeatability was obtained with all films, in line with the high photostability of  $\text{CuWO}_4$ -based materials.

Figure 3a shows the  $J$ - $V$  plots recorded with the multilayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes and a monolayer  $\text{CuWO}_4$  electrode under back side solar simulated light (*i.e.* the electrodes were irradiated through the FTO glass), in contact with a 0.1 M  $\text{K}_3\text{BO}_3$  solution at pH 9. First of all, the introduction of Mo into  $\text{CuWO}_4$  produces a more than 3-fold improvement in photocurrent density at 1.7 V vs. RHE. Moreover, no significant difference in the LSV curves of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  films could be noted up to an applied potential of *ca.* 1.3 V vs. RHE, while progressively higher photocurrent density values were recorded at higher applied potentials with increasing number of deposited layers up to the 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrode, which was found to be the most active one.



**Figure 3.** Linear Sweep Voltammetry (LSV) of (a) 1L–5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes and  $\text{CuWO}_4$  electrode in back configuration and (b) 1L and 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes in back and front configuration; AM 1.5 G solar simulated irradiation, scan rate  $10 \text{ mV s}^{-1}$ . The current in the absence of irradiation is also shown in panel (a) (black dashed line). Incident photon to current efficiency (IPCE) of (c)  $\text{CuWO}_4$  monolayer and 1L–5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes at 1.23 V vs. RHE and (d) the 3L electrode at different applied potentials.

Consequently, the highest photocurrent density of *ca.*  $1 \text{ mA cm}^{-2}$  was recorded at  $1.7 \text{ V vs. RHE}$  for the 3L film, while lower performances were obtained with thicker 4L and 5L electrodes, despite they practically absorb 100% of the incident light, also at long wavelengths. This provides a first indication of electron mobility issues occurring in films composed of more than three layers, *i.e.* thicker than *ca.*

270 nm. On the other hand, the slightly higher activity resulting for the 5L film compared to the 4L film could be related to the contribution to photocurrent of the longest wavelength photons, the exploitation of which increases with increasing film thickness, *i.e.* with increasing absorption of long wavelengths radiation (see Figure S5, showing a magnification of the IPCE profile of Figure 3c).

The  $J$ - $V$  curves of Figure 3a were compared with the results of LSV analyses performed under front-side illumination (Figure S6), *i.e.* by irradiating the material through the electrode/electrolyte interface. This allows to get more insight into the charge carriers mobility within the material and to assess if the internal charge transport is limited by either the minority or the majority charge carriers,<sup>31</sup> *i.e.*, by the holes or the electrons, respectively, in a n-type semiconductor material.<sup>32</sup> The LSV plots recorded with the 1L and 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes (*ca.* 80 and 270 nm thick, respectively) under back and front side irradiation are compared in Figure 3b. Since for both selected films the photocurrent recorded under front side irradiation is lower than that recorded under back side irradiation, poor electron transport appears to be the limiting factor for this material. The 3L electrode exhibits the best PEC performance among the tested films either under front or back irradiation (compare Figure 3b and Figure S6). Thus a *ca.* 270 nm film thickness seems to correspond to the best balance between photon absorption in the photoactive layer and electron transport efficiency within it.

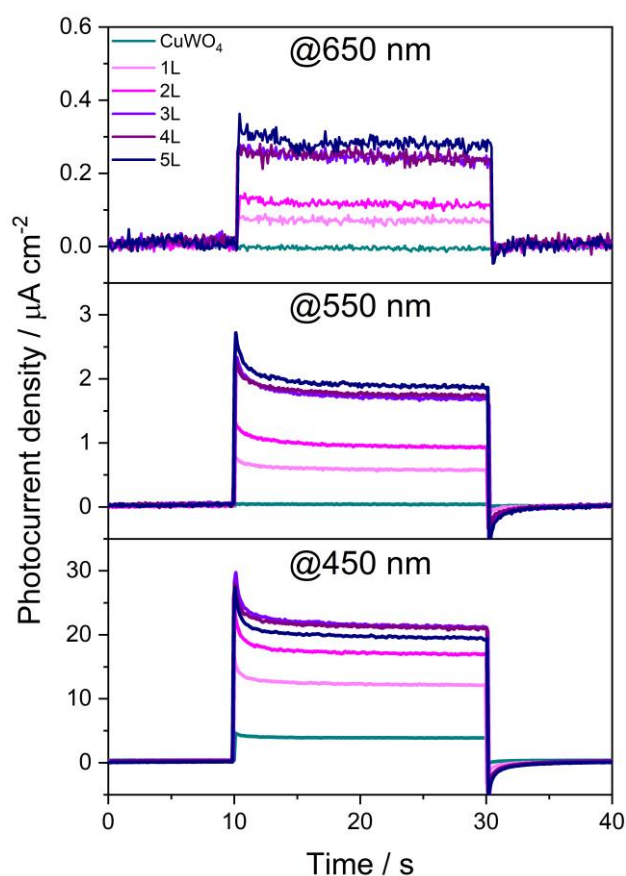
The IPCE plots recorded under back side irradiation at 1.23 V *vs.* RHE are reported in Figure 3c for the complete series of multilayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes and for a single layer  $\text{CuWO}_4$  photoanode. By comparing the curves obtained with the two monolayer  $\text{CuWO}_4$  and  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  photoanodes evidence is obtained of the higher efficiency of the composite molybdenum-containing material with respect to pure  $\text{CuWO}_4$ , over the entire investigated wavelengths range. This demonstrates that the increased visible light absorption capability (see Figure 1a) consequent to Mo for W substitution results in a higher number of photogenerated charge couples, which are effectively responsible for the higher visible light activity of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  with respect to  $\text{CuWO}_4$ .

Furthermore, by comparing the IPCE curves obtained with multilayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes (Figure 3c), the thinnest 1L and 2L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  films are most efficient in the UV region up to 390 nm, whereas thicker films exhibit the best photoactivity in the visible light region. The high absorption coefficient value of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  in the UV region (the extinction coefficient at 400 nm is three times greater than at 500 nm, see Figure S7) implies that UV radiation is absorbed almost quantitatively even by the thinner films, which results in high photocurrent values. On the other hand, in thicker electrodes (thicker than 2L) and under back side irradiation, most of the UV photogenerated holes need to travel across the whole film to reach their extraction sites at the film/electrolyte interface. The probability that they recombine with photopromoted electrons increases with increasing film thickness leading to lower IPCE values in the UV region with respect to thinner films. Concerning the relatively higher IPCE values of the thicker films in the visible region, owing to the lower absorption capability of the material at longer wavelengths (Figure S7), quantitative light harvesting can be attained only in thicker films, which are consequently able to better exploit visible light. The best compromise between efficient hole transport across the film and maximum visible light exploitation under back side irradiation is attained with the *ca.* 270 nm thick film (3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrode).

The recorded IPCE values are also in good agreement with the performances under full lamp irradiation recorded in LSV scans. In fact, as shown in Table S1, there is a good matching between the photocurrent density values at 1.23 V *vs.* RHE in *J-V* curves and the photocurrent density calculated by integrating the product between the IPCE curves and the standard AM 1.5 G solar spectrum<sup>33</sup> over the whole investigated wavelengths range (300-650 nm for the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  multilayer films).

Besides at 1.23 V *vs.* RHE, the IPCE curve for the best performing 3L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrode was recorded in back side configuration also at 1.50 and 1.70 V *vs.* RHE (Figure 3d), in order to evaluate the conversion efficiency of the material upon increasing charge separation. As shown in Figure 3d, the conversion efficiency gets higher with increasing applied potential over the whole investigated

wavelengths, with the IPCE value recorded at 400 nm showing a 2.5-fold and 3.5-fold enhancement as the potential increases from 1.23 to 1.5 and to 1.7 V *vs.* RHE, respectively. The observed behavior is in line with the photocurrent density enhancement recorded in LSV scans under simulated solar light irradiation. In particular, the IPCE almost linearly increases with increasing applied potential (see Figure 3d), because the holes generated by high energy photons, which are mainly confined in the proximity of the FTO back contact, benefit of the enhanced charge carrier separation due to the progressively higher external bias and have larger probability to reach the film/electrolyte interface where O<sub>2</sub> evolution occurs.



**Figure 4.** Chopped chronoamperometry at 1.23 V *vs.* RHE of the 1L–5L CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrodes and of the CuWO<sub>4</sub> electrode under irradiation at 650 nm (top), 550 nm (mid) and 450 nm (bottom panel).

Finally, the chopped photocurrent measurements under monochromatic irradiation at different wavelengths performed with our photoanodes at 1.23 V *vs.* RHE provide uncontroversial evidence that

all  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  – based photoanodes absorb light and exhibit photoactivity up to 650 nm, while  $\text{CuWO}_4$  is photoactive only below 550 nm. In fact, as shown in Figure 4, whereas all tested photoanodes, including the  $\text{CuWO}_4$  – based one, originate a photocurrent signal under irradiation at 450 nm (Figure 4, bottom panel), the photocurrent signal is zero for  $\text{CuWO}_4$  under irradiation at 550 nm (Figure 4, mid panel), while all 1L–5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  photoanodes remain photoactive at this wavelength and maintain their photoactivity up to 650 nm (Figure 4, top panel). This demonstrates that Mo for W substitution in  $\text{CuWO}_4$  largely extends the photoactivity of this material into the visible region, at least up to 650 nm.

### 3.3. Sacrificial agents

Two are the main limiting factors for PEC water oxidation performed by semiconductor-based photoanodes, *i.e.*, *i*) the transport of photopromoted electrons and photogenerated holes through the bulk material to the FTO back contact and to the film/electrolyte interface, respectively, and *ii*) the interfacial hole injection kinetics, determining the oxygen evolution efficiency. Decoupling these two contributions is therefore required in order to shed light on the intrinsic properties of the investigated photoactive material. In particular, the use of a suitable hole scavenger-containing electrolyte, acting as electron donor, allows one to study the intrinsic properties of the bulk material, under the assumption that hole injection at the material/electrolyte interface is not rate limiting under such conditions.<sup>34</sup>

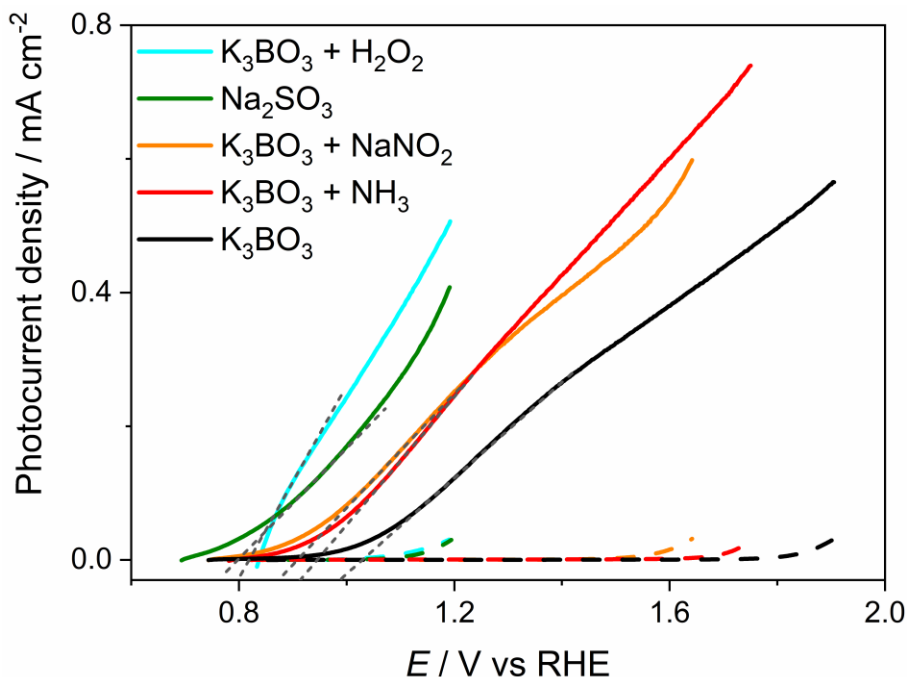
Previous studies pointed out that significant dark current is observed if  $\text{CuWO}_4$  electrodes are in contact with  $\text{Na}_2\text{SO}_3$ - and  $\text{H}_2\text{O}_2$ -containing solutions.<sup>35</sup> Therefore, the performance of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  photoanodes was investigated in the presence of various hole scavengers, *i.e.*  $\text{Na}_2\text{SO}_3$ ,  $\text{H}_2\text{O}_2$ ,  $\text{NaNO}_2$  and  $\text{NH}_3$ , to check the existence of dark currents and verify if and how the response of this material is affected by the redox potential of the donor species. The concentration of all sacrificial agents was kept constant at 0.1 M within 0.1 M  $\text{K}_3\text{BO}_3$  buffer solutions at pH 9, except for the 0.1 M  $\text{Na}_2\text{SO}_3$  solution, which was at pH 9 in the absence of buffer.

The dark current onsets in the presence of different sacrificial agents was measured first using both the 1L CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> film and a glassy carbon as working electrode, in order to attain an apparent redox scale for the employed species (see Figure S8). For the electrode immersed in 0.1 M K<sub>3</sub>BO<sub>3</sub> in the absence of any hole scavenger, the dark current potential onset is the highest, *i.e.*, around 1.8 V *vs.* RHE. On the other hand, the most easily oxidizable species is H<sub>2</sub>O<sub>2</sub>, immediately followed by Na<sub>2</sub>SO<sub>3</sub>, with a dark current onset located at *ca.* 1.0 V *vs.* RHE. More positive onset potentials, *i.e.*, *ca.* 1.4 and 1.6 V *vs.* RHE, were found for NaNO<sub>2</sub> and NH<sub>3</sub>, respectively. Thus, moving from the most stable to the most oxidizable species, the following scale results: H<sub>2</sub>O (*i.e.* K<sub>3</sub>BO<sub>3</sub> aqueous solution) > NH<sub>3</sub> > NaNO<sub>2</sub> >> Na<sub>2</sub>SO<sub>3</sub> > H<sub>2</sub>O<sub>2</sub>.

The LSV curves obtained with a monolayer CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> electrode in contact with the hole scavenger-containing solutions under back side simulated solar irradiation are shown in Figure 5. The LSV curve obtained with the 0.1 M K<sub>3</sub>BO<sub>3</sub> aqueous solution is also reported, as reference. In the absence of any electron donor species the photocurrent onset potential was *ca.* 1.0 V *vs.* RHE, *i.e.*, 800 mV lower than under dark conditions. The photocurrent onset potentials in the presence of the hole scavengers, extrapolated from the curves shown in Figure 5, are reported in Table S2. Compared to the pure K<sub>3</sub>BO<sub>3</sub> electrolyte solution, the use of Na<sub>2</sub>SO<sub>3</sub> and H<sub>2</sub>O<sub>2</sub> as hole scavengers led to the largest decrease of photocurrent onset potential, which was around 0.8 V *vs.* RHE in both cases, in good agreement with the conduction band edge value of *ca.* 0.7 V *vs.* RHE.<sup>25</sup> However, owing to the significant dark current generated in the presence of these sacrificial agents at a potential as low as 1.2 V *vs.* RHE, Na<sub>2</sub>SO<sub>3</sub> and H<sub>2</sub>O<sub>2</sub> were discarded as hole scavengers in our PEC investigation.

On the other hand, in the presence of NH<sub>3</sub> and NaNO<sub>2</sub> hole scavengers, a lower negative shift of the photocurrent onset potential (located at *ca.* 0.9 V *vs.* RHE) was observed and negligible dark current up to 1.75 and 1.6 V *vs.* RHE, respectively. This allows the exploitation of a wider applied potential window in PEC measurements. Because of the high volatility of NH<sub>3</sub>, which may result in uncontrollable

concentration variation of its aqueous solutions, we finally selected  $\text{NaNO}_2$  as the most suitable sacrificial agent for our PEC tests. In its presence the photocurrent density at 1.23 V vs. RHE is double compared to that measured in pure  $\text{K}_3\text{BO}_3$  (0.30 vs. 0.15  $\text{mA cm}^{-2}$ , see Figure 5), with negligible dark current at such applied bias.



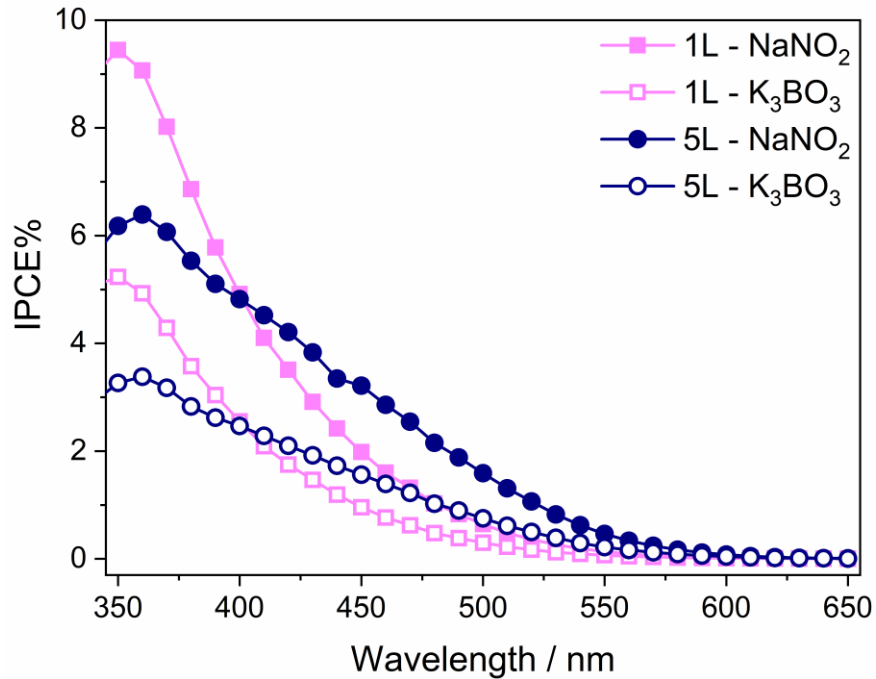
**Figure 5.** LSV curves recorded with a monolayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrode in contact with a 0.1 M  $\text{K}_3\text{BO}_3$  solution, 0.1 M  $\text{K}_3\text{BO}_3$  solutions containing different hole scavengers at a 0.1 M concentration or 0.1 M  $\text{Na}_2\text{SO}_3$ , under dark (dashed lines) or back side AM 1.5 G irradiation conditions (continuous lines). Scan rate 10  $\text{mV s}^{-1}$ . All solutions were at pH 9. The photocurrent onset potential was calculated by extrapolation of each LSV line (black dashed lines).

Anyway, regardless of the employed electron donor, the photocurrent just moderately increased with respect to water oxidation (in  $\text{K}_3\text{BO}_3$ ). This is in contrast with the considerable increase typically observed with semiconductors the PEC performance of which is limited by surface hole accumulation and slow water oxidation kinetics. For instance,  $\text{BiVO}_4$  photoanodes generate up to 100 times higher current in contact with a hole scavenger containing solution,<sup>36</sup> whereas  $\text{Fe}_2\text{O}_3$  electrodes in contact with



$\text{H}_2\text{O}_2$  allow the complete consumption of surface holes.<sup>37</sup> This suggests that the transfer of surface holes to water or electron donor species has minor limiting effects on PEC efficiency of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$ .

The IPCE curves obtained at 1.23 V vs. RHE with the 1L and 5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes in the presence of the  $\text{NaNO}_2$  hole scavenger (Figure 6) evidence the extended visible light photoactivity of the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  material, which is maximum for the thickest 5L photoanode.



**Figure 6.** Incident photon to current efficiency (IPCE) of 1L and 5L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrodes at 1.23 V vs. RHE in  $\text{K}_3\text{BO}_3$  both in the absence (void symbols) and in the presence (full symbols) of  $\text{NaNO}_2$ .

### 3.4. Charge separation and charge injection efficiencies

PEC measurements in the presence of  $\text{NaNO}_2$  as hole scavenger allowed us to evaluate the charge separation and charge injection efficiencies and thus assess if the performance of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  photoanodes is limited by either the material bulk properties or the surface hole injection kinetics.

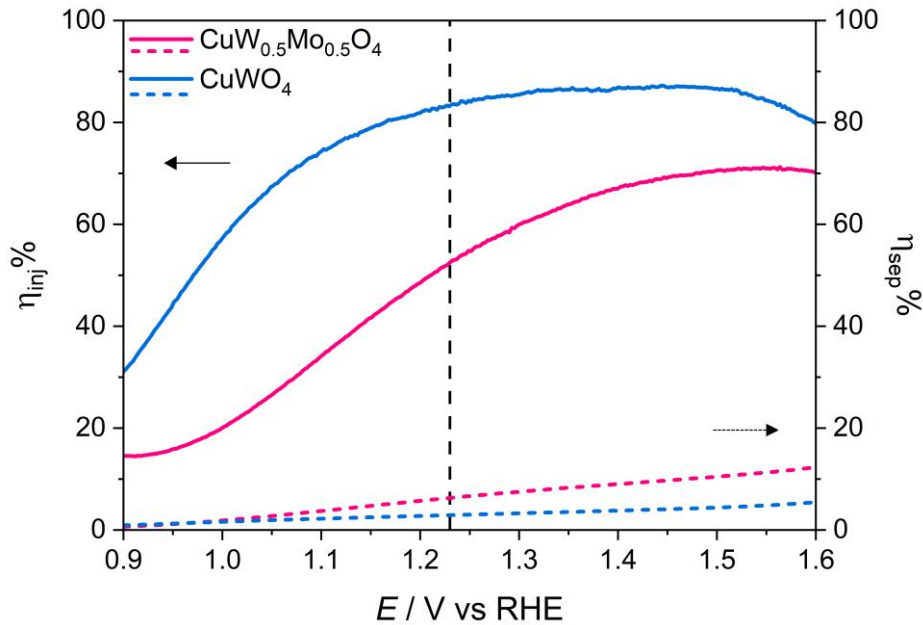
The photocurrent density  $J$  can be expressed as the combination of different contributions, according to equation (2):<sup>37,38</sup>

$$J = J_{\text{abs}} \eta_{\text{sep}} \eta_{\text{inj}} \quad (2)$$

where  $J_{\text{abs}}$  is the theoretical maximum photocurrent density obtained if all absorbed photons are converted into electricity, which was estimated by converting into current the integral of the product between the standard AM 1.5 G solar spectrum and the absorption spectrum of the photoelectrode over the 300–650 nm range ( $4.39 \text{ mA cm}^{-2}$  for our 1L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  material),<sup>15</sup>  $\eta_{\text{sep}}$  is the charge separation efficiency and  $\eta_{\text{inj}}$  is the interfacial charge injection efficiency. From this equation the charge separation and charge injection efficiencies can be easily calculated by taking into account the photocurrent densities recorded in the presence ( $J_{\text{NaNO}_2}$ ) and absence ( $J_{\text{K}_3\text{BO}_3}$ ) of the hole scavenger, as follows:

$$\eta_{\text{inj}} = J_{\text{K}_3\text{BO}_3} / J_{\text{NaNO}_2} \quad (3)$$

$$\eta_{\text{sep}} = J_{\text{NaNO}_2} / J_{\text{abs}} \quad (4)$$



**Figure 7.** Charge injection efficiency  $\eta_{\text{inj}}$  (continuous lines) and charge separation efficiency  $\eta_{\text{sep}}$  (dashed lines) calculated for the monolayer  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  (fuchsia) and  $\text{CuWO}_4$  (light blue) electrodes vs. the applied potential.

The calculated  $\eta_{\text{inj}}$  and  $\eta_{\text{sep}}$  values for the 1L  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  electrode and for the  $\text{CuWO}_4$  electrode are shown in Figure 7 as a function of the applied potential. At 1.23 V vs. RHE  $\eta_{\text{inj}}$  and  $\eta_{\text{sep}}$  are 52% and

6%, respectively, for 1L CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub>, which clearly indicate that charge separation in the bulk material is the main limiting factor for the PEC performance of our investigated material, in line with the limited photoactivity improvement attained in the presence of electron donor species.

For our CuWO<sub>4</sub> photoanode (with  $J_{\text{abs}}$  amounting to 3.60 mA cm<sup>-2</sup> in the 300–550 nm range)  $\eta_{\text{inj}} = 83\%$  and a  $\eta_{\text{sep}} = 3\%$  at 1.23 V vs. RHE. This implies that Mo incorporation into CuWO<sub>4</sub> induces a two-fold increase of the charge separation properties in the bulk, in line with recent reports,<sup>27</sup> while it negatively affects the hole injection efficiency at the CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> film/electrolyte interface, with a  $\eta_{\text{inj}}$  decrease.

The increased charge separation efficiency can result from the increased majority carrier (*i.e.*, electron) concentration in Mo-doped CuWO<sub>4</sub> with respect to CuWO<sub>4</sub> evidenced through Mott-Schottky plots in previous studies on Mo-doped CuWO<sub>4</sub>.<sup>25-27</sup> This would improve the material conductivity and thus reduce charge recombination in CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> films. On the other hand, the reduction in charge injection efficiency may be ascribed to a Mo-induced increase of the surface trap states of CuWO<sub>4</sub>,<sup>12,29,35,39,40</sup> which would decrease the efficiency of water oxidation especially at relatively low applied potentials, by inducing Fermi level pinning at the semiconductor–liquid junction. An indirect evidence of the increase of surface trap states in CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> films is the modest delay in photocurrent onset of the LSV curves recorded with CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub>-based photoanodes with respect to the CuWO<sub>4</sub> photoanode appearing in Figure 3a and better evidenced in the enlargement of this figure at low potentials shown in Figure S9. As recently outlined,<sup>27</sup> such delay is compatible with an increase in surface trap states upon Mo incorporation, which hampers the photocurrent onset at low overpotentials, *i.e.*, when the driving force to separate photoproduced electron-hole couples is small. Taken together, the two-fold improvement in charge separation and the decrease in charge injection efficiency result in the overall enhanced PEC activity of CuW<sub>0.5</sub>Mo<sub>0.5</sub>O<sub>4</sub> with respect to CuWO<sub>4</sub>.

#### **4. Conclusions**

Substitution of Mo for W into the  $\text{CuWO}_4$  structure is responsible for the conduction band edge energy lowering and consequent band gap narrowing. Thus, the increased PEC performance of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  with respect to  $\text{CuWO}_4$  mainly results from the extended visible light absorption ability and photoactivity of the  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  material up to 650 nm. This effect is maximum for the highest film thickness explored in this study ensuring the exploitation of long wavelength photons in the visible region, which are efficiently harvested for longer optical paths within the photoactive films. However, Mo for W substitution also increases the efficiency of photoproduced charge separation, due to an increased conductivity of the material, while the lower charge injection efficiency of  $\text{CuW}_{0.5}\text{Mo}_{0.5}\text{O}_4$  with respect to pure  $\text{CuWO}_4$  may result from an increased amount of surface trap states. With its 2.0 eV band gap, copper molybdo-tungstate shows promise as photoanode material for practical water splitting. Further improvements are needed and might be attained through nanostructuring, the development of efficient oxygen evolution co-catalysts for this class of materials and the fine modification of the electronic structure by doping with other elements.

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## **Supporting Information**

Details of XRPD spectra, cross-section SEM images, absorption coefficient and band gap (Tauc plot) determination, magnification of IPCE data, LSV under front side irradiation, dark current onset with different electrolytes, experimental and calculated (from IPCE integration) photocurrents.

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