

The zCOSMOS survey. The dependence of clustering on luminosity and stellar mass at $z = 0.2-1$ [★]

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ABSTRACT

Aims. We study the dependence of galaxy clustering on luminosity and stellar mass at redshifts $z \sim [0.2-1]$, using the first 10K redshifts from the zCOSMOS spectroscopic survey of the COSMOS field.

Methods. We measured the redshift-space correlation functions $\xi(r_p, \pi)$ and $\xi(s)$ and the projected function, $w_p(r_p)$ for subsamples covering different luminosity, mass, and redshift ranges. We explored and quantified in detail the observational selection biases from the flux-limited nature of the survey, using ensembles of realistic semi-analytic mock samples built from the Millennium simulation. We used the same mock data sets to carefully check our covariance and error estimate techniques, comparing the performances of methods based on the scatter in the mocks and on bootstrapping schemes. We finally compared our measurements to the cosmological model predictions from the mock surveys.

Results. At odds with other measurements at similar redshift and in the local Universe, we find a weak dependence of galaxy clustering on luminosity in all three redshift bins explored. A mild dependence on stellar mass is instead observed, in particular on small scales, which becomes particularly evident in the central redshift bin ($0.5 < z < 0.8$), where $w_p(r_p)$ shows strong excess power on scales $>1 h^{-1}$ Mpc. This is reflected in the shape of the full $\xi(r_p, \pi)$ that we interpret as produced by dominating structures almost perpendicular to the line of sight in the survey volume. Comparing to $z \sim 0$ measurements, we do not see any significant evolution with redshift of the amplitude of clustering for bright and/or massive galaxies.

Conclusions. This is consistent with previous results and the standard picture in which the bias evolves more rapidly for the most massive haloes, which in turn host the highest-stellar-mass galaxies. At the same time, however, the clustering measured in the zCOSMOS 10K data at $0.5 < z < 1$ for galaxies with $\log(M/M_\odot) \geq 10$ is only marginally consistent with the predictions from the mock surveys. On scales larger than $\sim 2 h^{-1}$ Mpc, the observed clustering amplitude is compatible only with $\sim 1\%$ of the mocks. Thus, if the power spectrum of matter is Λ CDM with standard normalisation and the bias has no “unnatural” scale-dependence, this result indicates that COSMOS has picked up a particularly rare, $\sim 2-3\sigma$ positive fluctuation in a volume of $\sim 10^6 h^{-1}$ Mpc³. These findings underline the need for larger surveys of the $z \sim 1$ Universe to appropriately characterise the level of structure at this epoch.

Key words. cosmology: observations – large-scale structure of Universe – surveys – Galaxy: evolution

1. Introduction

In the canonical scenario of galaxy formation, galaxies are thought to form through the cooling of baryonic gas within extended dark matter haloes (White & Rees 1978). The mass of the hosting halo is expected to play a significant role in the definition of the visible properties of the galaxy, as the total mass

in gas and stars, its luminosity, colour, star formation rate, and possibly, morphology.

Since it is the baryons that form the visible fabric of the Universe, a major challenge in testing the galaxy formation paradigm is to build clear connections between these observed properties and those of the hosting dark-matter haloes. This is a difficult task, as any direct connection initially existing between the dark-matter mass and the baryonic component cooling within the halo is modified by all subsequent dynamical processes affecting the halo-galaxy system, such as merging or dynamical friction. This is confirmed by simulations, which also show however that galaxy luminosity and stellar mass do in fact retain memory of the “original” (not actual) halo mass, i.e. before it experiences a major merger or is accreted by a larger halo (Conroy et al. 2006; Wang et al. 2006, 2007). This gives some hope that by measuring the dependence of the galaxy distribution on galaxy properties one is actually constraining the relationship between the dark and luminous components of galaxies.

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Measurements of first moments, such as the luminosity function (LF) or the stellar mass function, provide a way to understand how these are related to the total halo mass functions, which can be obtained from analytic predictions (e.g. Press & Schechter 1974) or n-body simulations (e.g. Warren et al. 2006). Similar investigations can be made on the second moment, i.e. the two-point correlation function (e.g. Springel et al. 2006). Studies of galaxy clustering in large local surveys have shown how clustering at $z \sim 0$ does depend significantly on several specific properties. These include luminosity (Hamilton 1988; Maurogordato & Lachieze-Rey 1991; Iovino et al. 1993; Benoist et al. 1996; Guzzo et al. 2000; Norberg et al. 2001; Norberg et al. 2002; Zehavi et al. 2005), colour or spectral type (Willmer et al. 1998; Norberg et al. 2002; Zehavi et al. 2002), morphology (Davis & Geller 1976; Giovanelli et al. 1986; Guzzo et al. 1997), stellar mass (Li et al. 2006), and environment (Abbas & Sheth 2006).

In recent years it has become possible to extend these investigations to high redshift, obtaining first indicative results on how these dependences evolve with time (Daddi et al. 2003; Coil et al. 2006; Phleps et al. 2006; Pollo et al. 2006; Meneux et al. 2006, 2008). The VIMOS-VLT Deep Survey (VVDS) (Pollo et al. 2006) and the DEEP2 survey (Coil et al. 2006) in particular, have provided new insights into the way galaxies of different luminosity cluster at $z \sim 1$. More specifically, Pollo et al. (2006) have shown that at these epochs galaxies already show a luminosity segregation, with more luminous galaxies being more clustered than faint objects. At the same time, however, a significant steepening with luminosity of the shape of their two-point correlation function for separations $< 1-2 h^{-1}$ Mpc, is observed. This behaviour is at variance with that at $z \sim 0$. A similar trend has been observed at the same redshift by the DEEP2 survey (Coil et al. 2006). In addition, Meneux et al. (2008) have shown a positive trend of clustering with stellar mass also at $z \sim 1$, with clear evidence of a stronger evolution of the bias factor for the most massive galaxies (see also Brown et al. 2008; Wake et al. 2008).

Interpreting the evolution in shape and amplitude of $w_p(r_p)$ with respect to luminosity and redshift is particularly interesting in the context of the halo model for galaxy formation. In this framework, the observed shape of $\xi(r)$ (or $w_p(r_p)$) is interpreted as being composed of the sum of two components: a) the 1-halo term, which dominates on small scales ($< 1-2 h^{-1}$ Mpc at the current epoch), where correlations are dominated by pairs of galaxies living within the same dark-matter halo (i.e. in a group or cluster); b) the 2-halo term on large scales, which is characterised by pairs of galaxies occupying different dark-matter haloes (see Cooray & Sheth (2002) for a review). Zheng et al. (2007) have modelled the luminosity-dependent $w_p(r_p)$ from both the DEEP2 (at $z \sim 1$) and SDSS (at $z \sim 0$) surveys, within such *Halo Occupation Distribution* (HOD) framework. In this way they establish evolutionary connections between galaxies and dark-matter haloes at these two epochs, providing a self-consistent scenario in which the growth of the stellar mass depends on the halo mass. Similar results have been obtained more recently in a combined analysis of the VVDS-Deep and SDSS data (Abbas et al. 2009).

In this paper we use the first 10000 redshifts from the zCOSMOS redshift survey (the “10K sample”) to further explore these high-redshift trends of clustering with luminosity and mass based on a new, independent sample. Although shallower than VVDS-Deep and DEEP2 ($I_{AB} < 22.5$ vs. 24 and 23.5, respectively), zCOSMOS covers a significantly larger area and samples a volume of $\sim 3 \times 10^6 h^{-1}$ Mpc to redshift $z = 1.2$. This should hopefully help reducing the effect of cosmic

variance (still strong for samples this size, Garilli et al. 2008; Stringer et al. 2009), while providing a better sampling of the high-end tail of the luminosity and mass functions. However, one main result of this analysis will be the explicit demonstration of the strength of the cosmic variance within volumes of this size. The clustering properties of the zCOSMOS sample in the volume contained within the redshift range 0.4–1 seem to lie at the extreme high end of the distribution of fluctuations on these scales, as already suggested by the angular clustering of the COSMOS data (McCracken et al. 2007). As we shall see, these results and those presented in the zCOSMOS series of clustering papers (de la Torre et al. 2009; Porciani et al., in prep.; Abbas et al., in prep.) indicate how cautious one should be in drawing far-reaching conclusions from the modelling of current clustering results from deep galaxy surveys.

A significant part of this paper is dedicated to discussing these cosmic-variance effects in detail, together with the impact of incompleteness on the derived results. This is particularly important when constructing mass-limited subsamples from a magnitude-limited survey, which introduces a mass incompleteness that depends on redshift and stellar mass. The intrinsic scatter in the galaxy mass-luminosity relation determines a progressive loss of faint galaxies with high mass-to-light ratio. We study the effect of this incompleteness on the measured clustering in detail using both the data themselves and mock samples built from the Millennium simulation. At the same time, we explore in quite some detail our ability to characterise measurement errors and the covariance matrix of our data, comparing estimates from the mock samples to those from bootstrap resamplings of the data.

The paper is organised as follows. In Sects. 2 and 3 we describe the zCOSMOS survey and the simulated mock samples used in the analysis, while in Sect. 4 we describe the selection of luminosity- and mass-limited subsamples, discussing extensively the incompleteness related to this operation. In Sect. 5 we describe our clustering estimators, while in Sect. 6 we discuss the observational biases and selection effects in detail, as well as how we account for them and what is their effect on the measured quantities; in Sect. 7 we explore the error budget and how to estimate the covariance properties of our measurements; in Sects. 8 and 9 we present our measurements of clustering as a function of luminosity and mass, respectively; in Sect. 10 we compare these results with those from other surveys and with simple model predictions; finally, in Sect. 11 we place these findings in a broader context and discuss future developments.

Throughout the paper we adopt a cosmology with $\Omega_m = 0.25$, $\Omega_\Lambda = 0.75$. When needed, we also adopt a value $\sigma_8 = 0.9$ for the normalisation of the matter power spectrum; this is chosen for consistency with the Millennium simulation, also used for comparison to model predictions. The Hubble constant is parameterised via $h = H_0/100$ to ease comparison with previous works. Stellar masses are quoted in unit of $h = 1$. All length values are quoted in co-moving coordinates.

2. The zCOSMOS survey data

The zCOSMOS survey (Lilly et al. 2007) is being performed with the VIMOS multi-object spectrograph at the ESO Very Large Telescope (Le Fèvre et al. 2003). Six hundred hours of observation have been allocated to this programme. These are being invested to measure spectra for galaxies in the COSMOS field (Scoville et al. 2007a), targeting: a) $\sim 20\,000$ galaxies brighter than $I \leq 22.5$ (*zCOSMOS Bright*); $\sim 10\,000$ sources at redshift $1.4 < z < 3.0$ pre-selected using colour-colour

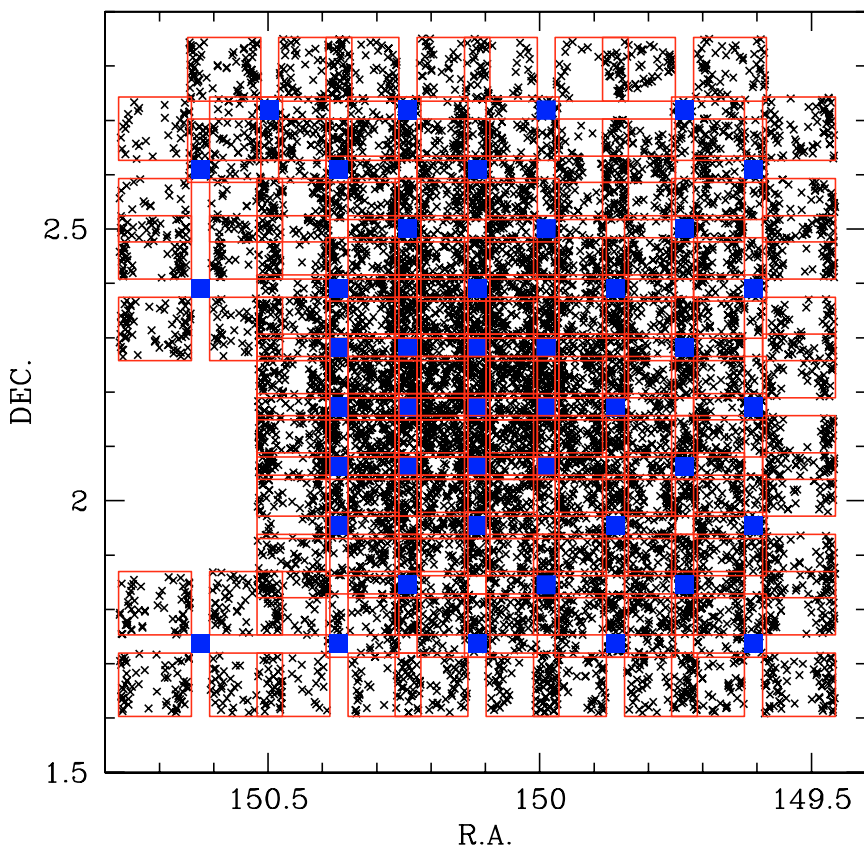


Fig. 1. Distribution on the sky of the $\sim 10\,000$ galaxies with measured redshift (crosses) forming the zCOSMOS “10K” sample. The large blue dots mark the centres of independent VIMOS pointings, each including four quadrants on the sky (as described by the red solid lines).

criteria (Lilly et al. 2007) (*zCOSMOS Faint*). So far, the survey has observed about half of the total “Bright” sample. This is the so-called “10K” sample used for the analysis presented in this paper, and is based on the observations of 83 VIMOS pointings over 44 distinct telescope positions on the sky (Lilly et al. 2009). These are shown in Fig. 1, where the footprint of VIMOS (4 quadrants of $\sim 7 \times 8$ arcmin² separated by a cross about 2 arcmin² wide) is evident. About every 3rd galaxy was observed in the field. The final “20K” zCOSMOS sample will be twice larger, reaching a sampling around 60–70%. The correction of the complex angular selection function is discussed later in the context of our galaxy clustering measurement.

Observations were performed using the medium-resolution RED grism, corresponding to $R \sim 600$ and covering the spectral range 5550–9650 Å. The average error on the redshift measurements was estimated from the repeated observations of 632 galaxies and found to be ~ 100 km s⁻¹ (Lilly et al. 2009). This corresponds roughly to a radial distance error of $1 h^{-1}$ Mpc. The reduction of the data to the redshift assignment was carried out independently at two institutes before a reconciliation process to solve discrepancies. The quality of each measured redshift was then quantified via a quality flag that provides us with a confidence level (see Lilly et al. 2007, 2009, for definition). For the present work, we only use redshifts with flags 1.5–4.5 and 9.3–9.5, corresponding to confidence levels greater than 98%.

The zCOSMOS survey benefits the large multi-wavelength coverage of the COSMOS field (Capak et al. 2007), which with the latest additions now comprises 30 photometric bands (Ilbert et al. 2009) extending well into the infrared. These include, in particular, accurate *K*-band and Spitzer-IRAC photometry over the whole area, which have allowed us to derive relevant physical

properties as rest-frame luminosity and stellar mass with unprecedented accuracy (Bolzonella et al. 2009; Pozzetti et al. 2009; Zucca et al. 2009).

3. Mock survey catalogues

In this paper we make intense use of mock surveys constructed from the Millennium simulation (Springel et al. 2005). This was done a) to understand the effect of our selection criteria on the measured quantities (Sect. 6.3) and b) to estimate the measurement errors and covariance of the data (Sect. 7).

We used two sets of light cones, constructed as explained in Kitzbichler & White (2007) and Blaizot et al. (2005) by combining dark-matter halo trees from the Millennium run to the Munich semi-analytic model of galaxy formation (De Lucia & Blaizot 2007). The two sets contain 24 1.4×1.4 deg² mocks built by Kitzbichler & White (2007) and 40 1×1 deg² mocks built by De Lucia & Blaizot (2007), which we name KW24 and DLB40, respectively. The main difference between the two sets, in addition to the different survey area, is that the DLB40 set contains all galaxies irrespective of any criteria down to the simulation limit that corresponds roughly to $10^8 M_{\odot}$, up to redshift $z = 1.7$, whereas the KW24 set only contains galaxies brighter than $I \leq 22.5$. This implies that the DLB40 set allowed us to select in stellar mass down to very low masses and to test selection effects. The observing strategy of the zCOSMOS 10K sample was only applied to the KW24 set, allowing us to do careful error analysis of our measurements.

The Millennium run contains $N = 2160^3$ particles of mass $8.6 \times 10^8 h^{-1} M_{\odot}$ in a cubic box of size $500 h^{-1}$ Mpc. The simulation was built with a Λ CDM cosmological model with $\Omega_m = 0.25$, $\Omega_{\Lambda} = 0.75$, $\sigma_8 = 0.9$ and $H_0 = 73$ km s⁻¹ Mpc⁻¹.

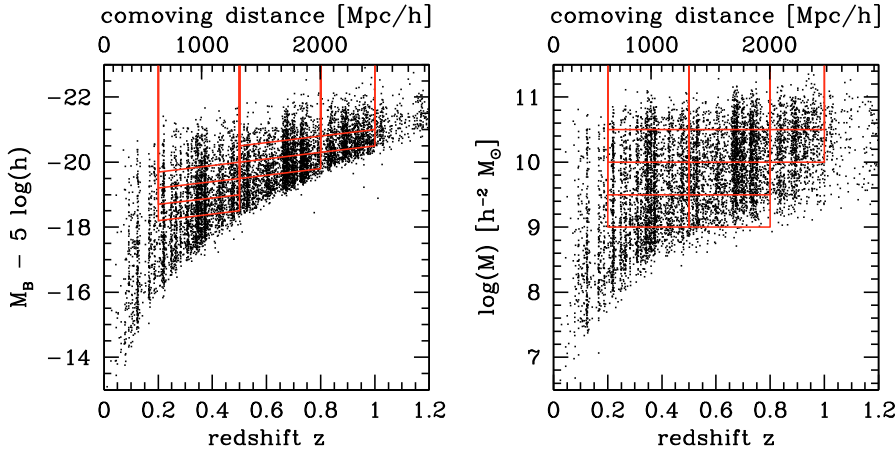


Fig. 2. Selection boundaries of the different subsamples of the zCOSMOS 10K survey used in this paper. *Left:* luminosity-redshift selection, which accounts for the average luminosity evolution of galaxies; *Right:* mass-redshift selection.

4. Luminosity- and mass-selected subsamples

4.1. Luminosity selection

Absolute magnitudes were derived for the 10K galaxies using the code ALF (Ilbert et al. 2005; Zucca et al. 2009), which is based on fitting a spectral energy distribution (SED) to the observed multi-band photometry. There are various sources of uncertainties to take into account (errors on apparent magnitudes, number of available photometric bands, method used, etc.). A direct comparison with absolute magnitudes derived with the independent code ZEBRA (Feldmann et al. 2006) shows consistent estimates with a small dispersion of $\sigma \sim 0.05$ magnitudes, in particular in the B -band. This can reasonably be considered as the typical error on our absolute magnitudes.

For our analysis, the goal is to define luminosity-limited samples that are as close as possible to truly volume-limited samples, i.e. with a constant number density. This might be done within a few independent redshift ranges. The size of the redshift slices in which to split the sample has to be chosen as a compromise between two aspects: a) large enough to have sufficient statistics and provide a good measurement of clustering and b) not too large to avoid significant evolution within each redshift bin.

We know, however, that the luminosity of galaxies evolves through the redshift range covered by the zCOSMOS survey ($0.2 < z < 1.1$), with a clear change in the characteristic parameters of the LF (Ilbert et al. 2005). This evolution does depend on the morphological/spectral type of the galaxy considered. To be able to select a nearly volume-limited sample within a given redshift interval, we need to take the corresponding evolution into account. This can only be done realistically in a statistical way by looking at the population-averaged evolution of the global LF.

We therefore considered the observed LF measured from the same data (Zucca et al. 2009) and modelled its change with redshift as a pure luminosity evolution (i.e. keeping a constant slope α and normalisation factor Φ^*), which is a fair description of the observed behaviour. We find that the characteristic absolute magnitude $M^*(z)$ evolves with redshift as

$$M^*(z) = M_0^* + A z, \quad (1)$$

where $A \sim -1$. In the companion paper, de la Torre et al. (2009) split the zCOSMOS galaxy samples into 3 morphological classes. They observed different luminosity evolutions for elliptical, spiral and irregular galaxies, with A varying from ~ -0.7 for to ~ -1.2 but with large uncertainties making $A = -1$ compatible for all classes. Porciani et al. (in prep.) reach similar

Table 1. Properties of the luminosity-selected samples.

Sample	Redshift range	Mean redshift	$M_{B,\text{cut}} (z=0)$	Number of galaxies
L1.1	0.2–0.5	0.37	−18.00	1892
L1.2	0.2–0.5	0.37	−18.50	1311
L1.3	0.2–0.5	0.37	−19.00	811
L1.4	0.2–0.5	0.37	−19.50	469
L2.1	0.5–0.8	0.67	−19.00	1848
L2.2	0.5–0.8	0.67	−19.50	1025
L2.3	0.5–0.8	0.67	−20.00	441
L3.1	0.8–1.0	0.91	−19.50	971
L3.2	0.8–1.0	0.91	−20.00	447

Table 2. Properties of the mass-selected samples.

Sample	Redshift range	Mean redshift	$\log(M/M_\odot)$		Number of galaxies
			Range	Median	
M1.1	0.2–0.5	0.36	≥ 9.0	9.80	2159
M1.2	0.2–0.5	0.37	≥ 9.5	10.09	1445
M1.3	0.2–0.5	0.36	≥ 10.0	10.36	827
M1.4	0.2–0.5	0.37	≥ 10.5	10.66	275
M2.1	0.5–0.8	0.66	≥ 9.0	9.97	2831
M2.2	0.5–0.8	0.66	≥ 9.5	10.12	2276
M2.3	0.5–0.8	0.67	≥ 10.0	10.38	1366
M2.4	0.5–0.8	0.67	≥ 10.5	10.68	477
M3.1	0.8–1.0	0.90	≥ 10.0	10.46	755
M3.2	0.8–1.0	0.90	≥ 10.5	10.73	344

conclusions when dividing the zCOSMOS 10K sample into three colour classes.

We therefore defined our luminosity-limited samples by an effective absolute magnitude cut at $z = 0$, $M_{B,\text{cut}}$, including all galaxies with $M_B(z) - 5 \log(h) \leq M_{B,\text{cut}} - z$. The resulting selection loci for different values of $M_{B,\text{cut}}$ are plotted over the data in the luminosity-redshift plane in the left panel of Fig. 2. As is evident in the figure, the faintest allowed threshold $M_{B,\text{cut}}$ depends on the redshift range considered, i.e. $z = [0.2-0.5]$, $z = [0.5-0.8]$, and $z = [0.8-1.0]$. The details of the resulting samples are described in Table 1.

4.2. Mass selection

Stellar mass has become a quantity routinely measured in recent years, thanks to surveys with multi-wavelength photometry, extending to the near-infrared (e.g. Rettura et al. 2006), although some uncertainties related to the detailed modelling of stellar

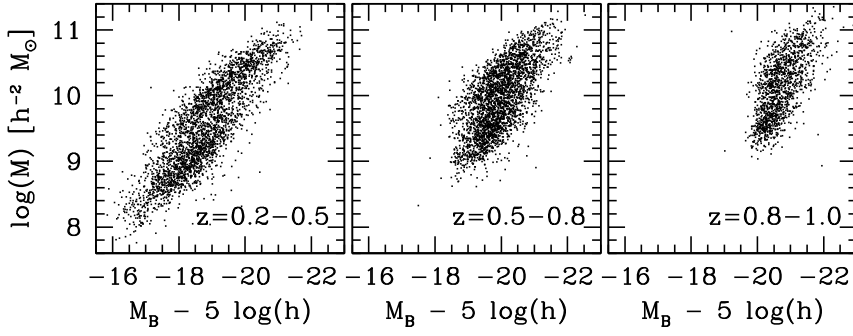


Fig. 3. The observed relationship between stellar mass and luminosity for galaxies in the 10K sample, within the three redshift ranges studied in this paper. The left panel shows an aspect of the galaxy bimodality, with red galaxies more massive and brighter than blue ones.

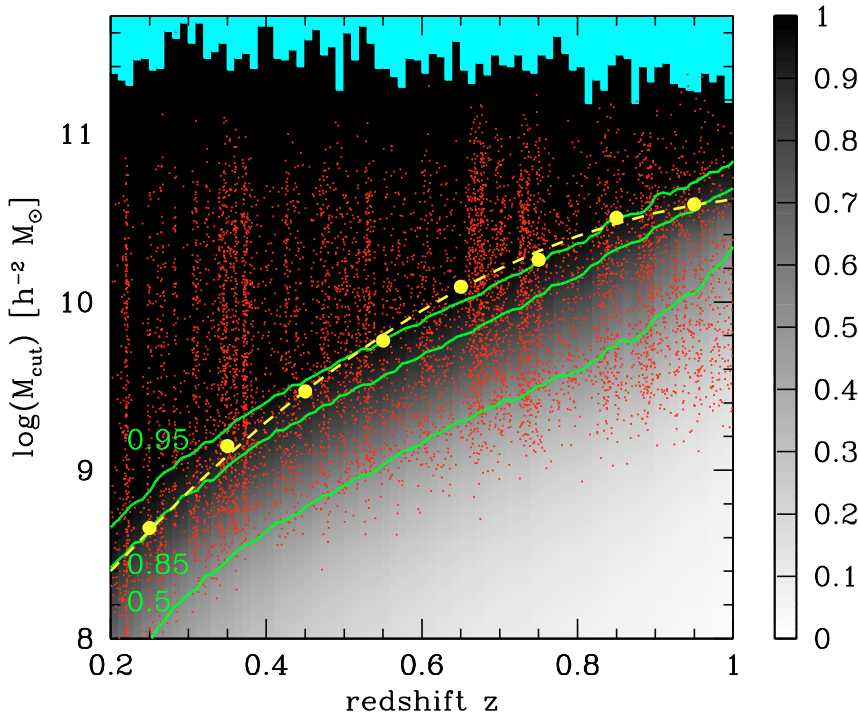


Fig. 4. Estimate of how the completeness in stellar mass changes as a function of redshift, due to the survey flux limit ($I_{AB} < 22.5$). The shaded grey area and green contours describe the loci of constant completeness. They are derived from the DLB40 mock samples of $1 \times 1 \text{ deg}^2$ and defined as the fraction of observed ($I_{AB} \leq 22.5$) galaxies over the total number in a given cell with size $\Delta z = 0.01$ and $\log(M) \geq \log(M_{\text{cut}})$. The superimposed red points correspond to the actual data of the 10K sample. The yellow points and dashed line show the 95% M/L ratio completeness level derived independently by Pozzetti et al. (2009) (see text), directly from the observed data.

evolution remain (Pozzetti et al. 2007). This has made studies of clustering as a function of stellar mass possible for large statistical samples. We used stellar masses estimated by fitting the SED, as sampled by the large multi-band photometry, with a library of stellar population models based on Bruzual & Charlot (2003). We used the code *Hyperzmass*, a modified version of the photometric redshift code *Hyperz* (Bolzonella et al. 2000). The typical error on stellar masses is ~ 0.2 dex. The method and accuracy of these measurements are fully described in Bolzonella et al. (2009) and Pozzetti et al. (2009).

We have thus constructed a set of mass-selected samples, containing galaxies more massive than a given threshold. We chose the same redshift ranges as used for the luminosity-selected samples. The properties of the selected subsamples are summarised in Table 2 and represented in Fig. 2.

4.3. Mass completeness

The flux-limited nature of surveys like zCOSMOS ($I_{AB} < 22.5$) mean that the lowest-mass samples are affected to varying degrees by incompleteness related to the scatter in the mass-luminosity relation (Fig. 3). This introduces a bias against

objects that would be massive enough to enter the mass-selected samples, but too faint to fulfil the apparent-magnitude limit of the survey. These missed high mass-to-light ratio galaxies will be those dominated by low-luminosity stars, i.e. the red and faint objects. Clearly, if this is not accounted for in some way, it would inevitably affect the estimated clustering properties, with respect to a truly complete, mass-selected sample (Meneux et al. 2008). It is therefore necessary to understand the effective completeness level in detail in the stellar mass of the samples that we defined for our analysis.

Meneux et al. (2008) have used two different methods to explore and quantify the completeness limit in stellar mass as a function of redshift. The first is based on the observed scatter in the mass-luminosity relation, obtained from the data themselves and extrapolated to fainter fluxes. The second instead makes use of mock survey samples, under the hypothesis that they provide a realistic description of the mass-luminosity relation and its scatter: the DLB40 set of mock survey catalogues that are complete in stellar mass are “observed” under the same conditions as the real data, i.e. selected at $I \leq 22.5$. The completeness is then simply defined, for a given redshift range and mass threshold, as the ratio of the number of galaxies brighter than the zCOSMOS flux

Table 3. The completeness in stellar mass of mass-selected mock subsamples reproducing the properties and selection criteria of our 10K data samples.

Sample	Redshift range	Stellar mass ($\log(M/M_\odot)$) range	Completeness
M1.1	0.2–0.5	≥ 9.0	0.783
M1.2	0.2–0.5	≥ 9.5	0.972
M1.3	0.2–0.5	≥ 10.0	1.000
M1.4	0.2–0.5	≥ 10.5	1.000
M2.1	0.5–0.8	≥ 9.0	0.349
M2.2	0.5–0.8	≥ 9.5	0.652
M2.3	0.5–0.8	≥ 10.0	0.919
M2.4	0.5–0.8	≥ 10.5	0.996
M3.1	0.8–1.0	≥ 10.0	0.571
M3.2	0.8–1.0	≥ 10.5	0.882

limit over those at any flux. Interestingly, even if this method is model-dependent (in particular, on the prescription of galaxy formation used in the semi-analytic models), this approach leads to similar completeness limits to the first one. The results of this second exercise are shown as a function of redshift and mass threshold and for a flux limit $I \leq 22.5$, in Fig. 4. Completeness is estimated in narrow redshift ranges ($\Delta z = 0.01$) for different mass thresholds M_{cut} increasing from 10^8 to $10^{11.7} h^{-2} M_\odot$ with a step of $10^{0.01} h^{-2} M_\odot$. A large fraction of low-mass objects are clearly missed at high redshift.

We also add in Fig. 4 the completeness limit estimated from the observed scatter in the M/L relation of the data, and defined at each redshift as the lower boundary, $M_{\text{min}}(z)$, including above it 95% of the mass distribution (Pozzetti et al. 2009). It is very encouraging to notice the very good agreement between this independent estimation from the data and that based on the DLB40 set of mock catalogues. This adds confidence in the use of the simulated samples. Table 3 summarises the completeness estimates derived from these mock catalogues for each of the 10 zCOSMOS galaxy samples defined in Table 2. The sample M2.1 shows the strongest incompleteness: 65.1% of the galaxies more massive than $10^9 h^{-2} M_\odot$ are fainter than $I = 22.5$ at $z = [0.5–0.8]$ and then, not included in our sample. In Sect. 6.3 we discuss the effects of this incompleteness on the galaxy clustering measurement.

5. Estimating the two-point correlation function

The two-point correlation function is the simplest estimator for quantifying galaxy clustering, because it is related to the second moment of the galaxy distribution, i.e. its variance. In practice, it describes the excess probability $\xi(r)$ of observing a pair of galaxies at a given separation r , with respect to that of a random distribution (Peebles 1980). Here we estimate the redshift-space correlation function $\xi(r_p, \pi)$, which allows one to incorporate the effect of peculiar motions on the pure Hubble recession velocity. In this case, galaxy separations are split into the tangential and radial components, r_p and π (Davis & Peebles 1983; Fisher et al. 1994).

The real-space correlation function $\xi_R(r)$ can be recovered by projecting $\xi(r_p, \pi)$ along the line of sight, as

$$w_p(r_p) \equiv 2 \int_0^\infty \xi(r_p, \pi) d\pi = 2 \int_0^\infty \xi_R \left[(r_p^2 + y^2)^{1/2} \right] dy. \quad (2)$$

For a power-law correlation function, $\xi_R(r) = (r/r_0)^{-\gamma}$, this integral can be solved analytically and fitted to the observed

$w_p(r_p)$ to find the best-fitting values of the correlation length r_0 and slope γ (e.g. Davis & Peebles 1983). In computing $w_p(r_p)$, a finite upper integration limit has to be chosen in practice. Its value has to be high enough as to include most of the clustering signal dispersed along the line of sight by peculiar motion. However, it must not be too high to avoid adding only noise, which is dominant above a certain π . Previous works (Pollo et al. 2005) have shown that, for similar data, the best results are obtained with an integration limit π_{max} between 20 and 40 h^{-1} Mpc. Our tests show that the scatter in the recovered $w_p(r_p)$ is obtained using the lowest value in this range. This can introduce a 5–10% underestimate in the recovered large-scale amplitude, which can be accounted for when fitting a model to $w_p(r_p)$. In the following, we in general use $\pi_{\text{max}} = 20 h^{-1}$ Mpc and show examples of how the amplitude is biased by this choice for the real data.

To estimate $\xi(r_p, \pi)$ from each galaxy sample, we used the standard estimator of Landy & Szalay (1993):

$$\xi(r_p, \pi) = \frac{N_R(N_R - 1) GG(r_p, \pi)}{N_G(N_G - 1) RR(r_p, \pi)} - \frac{N_R - 1}{N_G} \frac{GR(r_p, \pi)}{RR(r_p, \pi)} + 1, \quad (3)$$

where N_G is the mean galaxy density (or, equivalently, the total number of objects) in the sample, N_R the mean density of a catalogue of random points distributed within the same survey volume and with the same selection function as the data, $GG(r_p, \pi)$ the number of independent galaxy-galaxy pairs with separation between r_p and $r_p + dr_p$ and between π and $\pi + d\pi$, $RR(r_p, \pi)$ the number of independent random-random pairs within the same interval of separations, and $GR(r_p, \pi)$ represents the number of galaxy-random cross pairs.

6. Observational biases and selection effects

6.1. Correction of VIMOS angular footprint and varying sampling

To properly estimate the correlation function from the 10K zCOSMOS data, we need to correct for its spatial sampling rate, which is on average $\sim 30\%$, but which varies with the position on the sky due to the VIMOS footprint and the superposition of multiple passes (see Fig. 1). The correction scheme used here is an evolution of the one discussed in Pollo et al. (2005), but with a simplified weighting scheme. The main differences between this sample and the VVDS-Deep data used by Pollo et al. (2005) are that: a) this sample is 1.5 magnitudes brighter; b) the spectra are taken with higher resolution, which produces longer spectra and thus fewer objects along the dispersion direction; and c) there are as many as eight repeated observations (“passes”) covering each point on the sky in the central area of the COSMOS field. The net result of these differences can be appreciated in Fig. 1, where the sample is characterised by a well-sampled central region, but also by rather sparsely sampled VIMOS pointings in the outskirts of the field. In particular, these external pointings clearly show target galaxies concentrated along rows. This effect is produced by the significant length of the spectra on the CCD in medium-resolution mode: not more than two spectra can be aligned on top of each other on the detector in each quadrant, which results in the observed two “stripes”. This is significantly different from what happens in the low-resolution observing mode (such as in VVDS-Deep, Le Fèvre et al. 2005a), where spectra are shorter, and up to four of them can be packed along the same column on the CCD.

We tested three different algorithms to correct for the angular selection function of the survey and obtained comparable results. Other weighting schemes use in particular the angular correlation

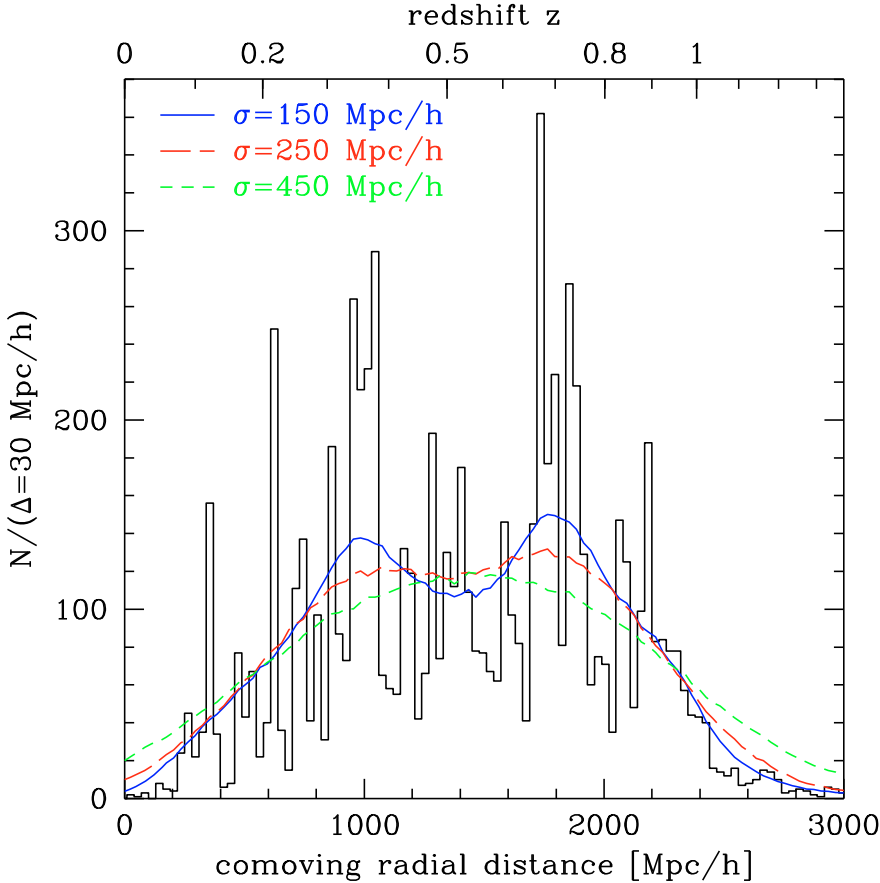


Fig. 5. Overall radial distribution of the zCOSMOS 10K sample, compared to three different smoothed distributions. These are obtained by filtering the observed data with a Gaussian kernel of increasing $\sigma = 150, 250,$ and $450 h^{-1}$ Mpc. The first two smoothed curves retain information on the two large structures located at ~ 1000 and $\sim 1800 h^{-1}$ Mpc along the line of sight, while the third one overestimates the number density of galaxies at low and high redshifts.

functions of the 10K sample and the photometric catalogue to correct for the nonuniform spatial sampling rate. These methods are discussed in the parallel clustering analyses by de la Torre et al. (2009) and Porciani et al. (in preparation). In the latter paper in particular, comparative tests of the three algorithms are presented.

Since the subsamples analysed in this work are essentially volume-limited (above the luminosity/mass completeness limits), we did not need to apply any further minimum-variance weighting scheme (as e.g. the J_3 weighting, Fisher et al. 1994). This is normally necessary for purely flux-limited surveys in which the selection function varies significantly as a function of redshift, such that different parts of the volume are sampled by galaxies with different luminosities and number densities (e.g. Li et al. 2006).

6.2. Construction of reference random samples

A significant source of uncertainty that we encountered in estimating two-point functions from our 10K subsamples is related to the construction of the random sample and in particular to its redshift distribution. We soon realised that the strongly clustered nature of the COSMOS field along the line of sight, with several dominating structures at different redshifts, required some particular care so as not to generate systematic biases in the random sample. These superclusters are already evident as vertical stripes in Fig. 2 and even more clearly so in the redshift histogram of Fig. 5. We point out the big “walls” at $z = 0.35, 0.75,$ and 0.9 , which are also clearly identified by the density field reconstruction of Kovač et al. (2009).

A standard way to generate a random redshift coordinate accounting for the radial selection function of the data uses a

Gaussian-filtered version of the data themselves. This is normally obtained using smoothing kernels with a dispersion σ (in co-moving coordinates) in the range $150\text{--}250 h^{-1}$ Mpc. The results of applying this technique to the current 10K data are shown in Fig. 5. One notes how for smoothing scales of 150 and $250 h^{-1}$ Mpc, the curves still retain memory of the two largest galaxy fluctuations. These are only erased when a very strong smoothing filter ($450 h^{-1}$ Mpc) is adopted. However, in this case the smoothed curve is unable to correctly follow the global shape of the distribution, overestimating the number density in the lowest and highest redshift ranges. The situation for our specific analysis, however, is somewhat simpler than this general case. Our luminosity-limited or mass-limited samples are in principle “volume-limited”, i.e. samples that – if properly selected – should have a constant density within the specific redshift bin. One such case is shown in the zoom of Fig. 6, where the redshift distribution in the range $z = [0.5, 0.8]$ is plotted.

An alternative way to generate the radial distribution of the random sample is to integrate the galaxy LF in steps along the redshift direction, computing at each step a value for the density of galaxies expected at that redshifts. Ideally, the LF can be measured from the sample itself and would include any detected evolution of its parameters. This is what we did here, using the evolving LF parameters presented in the companion dedicated paper (Zucca et al. 2009). The dashed red line in Fig. 6 shows the result one obtains if smoothing with a kernel of $450 h^{-1}$ Mpc, compared to the one obtained from the integration of the LF. The latter is fully consistent with what is expected from a truly volume-limited sample with the given selection criteria, with the number of objects increasing as the square of the radial co-moving distance.

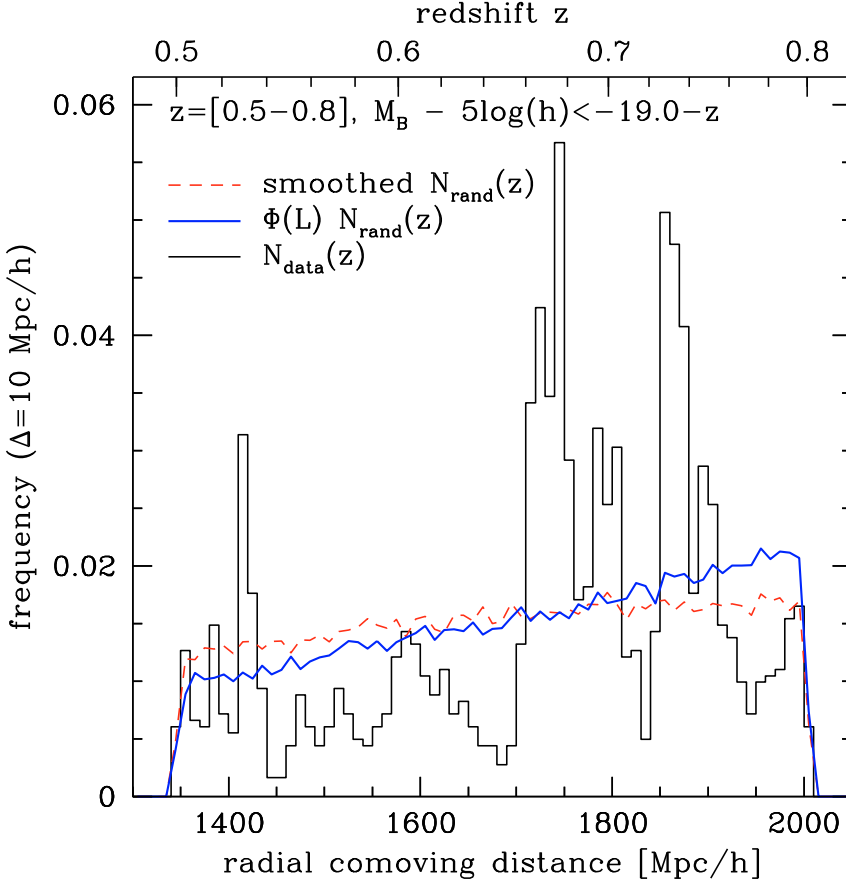


Fig. 6. The radial distribution of the luminosity-selected sample L2.1 compared to a smoothed curved (with a kernel of $\sigma = 450 h^{-1} \text{ Mpc}$ – dashed curve) and a radial distribution generated by integrating the LF (solid curve). The latter is consistent with what is expected for such galaxy sample.

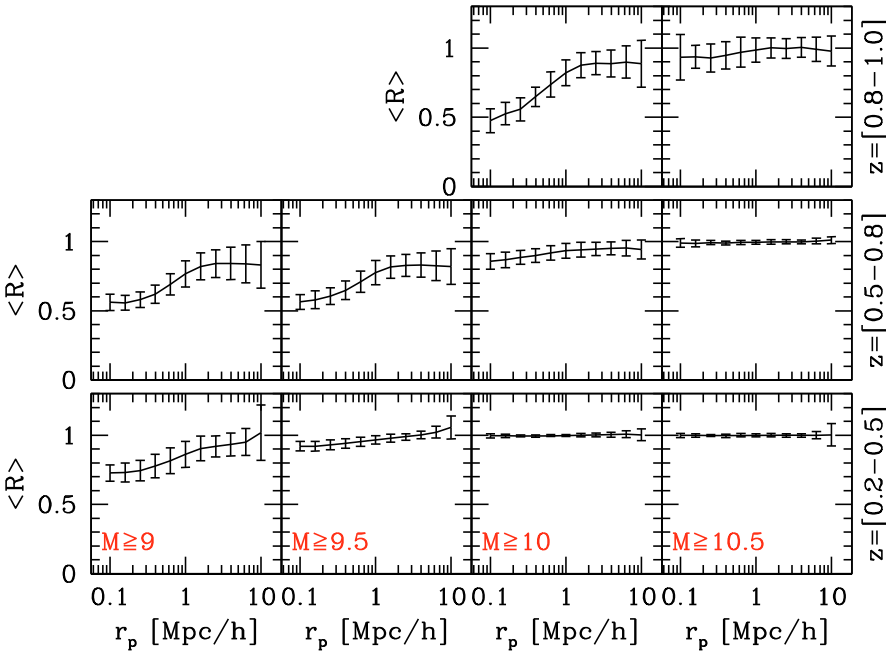


Fig. 7. The effect of stellar mass incompleteness on the measured $w_p(r_p)$, estimated from the $1 \times 1 \text{ deg}^2$ Millennium mock samples. The figure shows the average of the quantity R over 40 mock samples as a function of r_p . R is defined as the ratio of the estimates of $w_p(r_p)$ with and without the flux cut at $I_{AB} = 22.5$, i.e. for a sample mimicking the 10K selection and a sample 100% complete in stellar mass.

6.3. Effect of mass incompleteness on $w_p(r_p)$

As discussed in Sect. 4.3 when we constructed our mass-limited samples, a fraction of the galaxies more massive than the formal mass threshold are in fact lost because of the limiting $I_{AB} < 22.5$ flux cut of the survey. This becomes more and more important with increasing redshift. As said, this population of missing

galaxies is inevitably dominated by red objects with high mass-to-light ratio (Meneux et al. 2008), which are known to cluster more strongly than the average population (Meneux et al. 2006; Coil et al. 2008; McCracken et al. 2008). Having defined our clustering tools, we can now further extend the analysis of Sect. 4.3 and use the DLB40 mock samples to directly quantify the effect this has on the measured $w_p(r_p)$. We thus computed

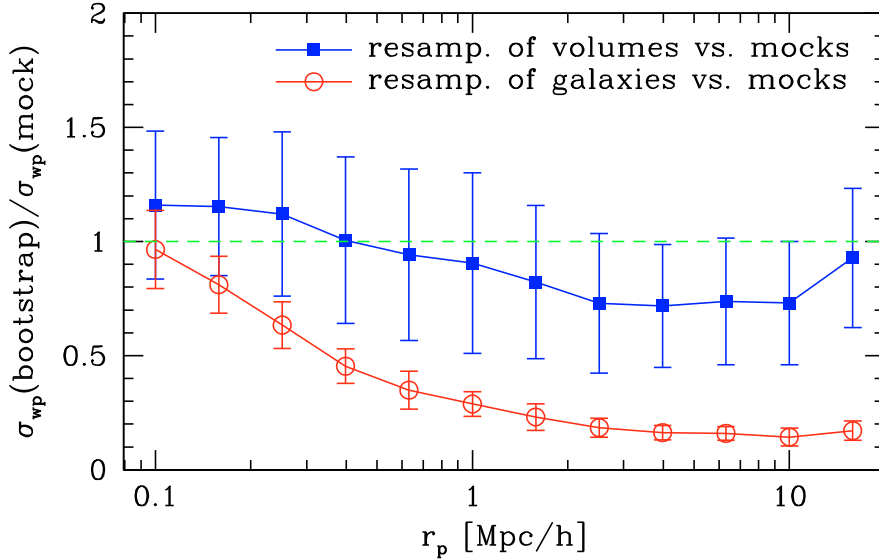


Fig. 8. Ratio of the diagonal errors on $w_p(r_p)$, obtained through the *bootstrap* resampling method to the “true” ones obtained from the variance of 24 mock catalogues. Filled and open symbols correspond to two different bootstrapping techniques, resampling, respectively, subvolumes of the survey or single galaxies. The former technique clearly provides a standard deviation which is closer to the “true” one obtained from repeated measurements.

the statistics for each of the mocks, which are complete down to very small masses ($\sim 10^8 M_\odot$), with and without applying the apparent-magnitude cut. Clearly, we are making a very strong hypothesis here, i.e. that the simulated samples have intrinsic clustering properties (and their relation to the galaxy’s M/L ratio), which are similar to those of real data.

The ratio of these two estimates (“true” over “observed”) averaged over the 40 mock catalogues is shown in Fig. 7. For a mass selection that is 100% complete within the given redshift bin, we would measure $\langle R \rangle = 1$ at all separations. We can see that the only mass range for which this is strictly happening at any redshift is the one with $\log(M/M_\odot) \geq 10.5$. For lower mass samples, we see a clear reduction of the clustering amplitude. However, we can also see that, for most samples, the shape of $w_p(r_p)$ is distorted mainly only below $< 1 h^{-1}$ Mpc. Above this scale, the mass incompleteness introduces an amplitude reduction up to $\sim 20\%$ in the worst cases. This will have to be considered when comparing our measurements with models (although keeping in mind that these estimates come from simulated data, not from real observations). For general comparisons, however, the amount of amplitude reduction of $w_p(r_p)$ is typically negligible on scales larger than $\sim 1 h^{-1}$ Mpc, given the statistical errors of the data measurements.

7. Systematic and statistical errors on correlation estimates

The derivation of realistic errors on the galaxy correlation function has been the subject of debate since its early measurements (see e.g. Bernstein 1994). In particular, it is well known that the measured values of the two-point correlation function are not independent on different scales. This means that, the bins of $w_p(r_p)$ have a degree of correlation among them, which needs to be taken into account when fitting a model to the observed values. This can be done if we are able to reconstruct the $N \times N$ covariance (or correlation) matrix of the N bins (Fisher et al. 1994).

In a recent paper, Norberg et al. (2009) compare in detail three different methods for estimating the covariance matrix of a given set of measurements. These use a) the ensemble variance from a set of mock catalogues, reproducing as accurately as possible the clustering properties and selection

function of the real data; b) a set of *bootstrap* resamplings of the volume containing the data; and c) a so-called *jack-knife* subset of volumes of the survey. In this latest case, the survey volume is divided into N_V subvolumes and the statistics under study recomputed each time excluding one of the subparts. In the “block-wise” incarnation of the bootstrap technique (Porciani & Giavalisco 2002, method “b”), instead, N subvolumes are selected each time *with repetition*, i.e. excluding some of them, but counting two or more times some others as to always get a global sample with the same total volume. We note, however, that there are historically two possible ways of resampling internally the data set. The classical “old” bootstrap (Ling et al. 1986) entailed bootstrapping the sample “galaxy-by-galaxy”. This means each time randomly picking a sample of N_G galaxies among our data set of N_G galaxies, allowing repetitions. In this way, within one bootstrap realisation a galaxy can be selected more than once, while some others are never selected. This technique has been shown to generally lead to some underestimation of the diagonal errors (Fisher et al. 1994). Here we also directly test this aspect.

The advantage of using mock samples is that, under the assumption that these are a realistic realisation of the real data, they allow us to obtain a true ensemble average and standard deviation from samples with the same size as the data sample, including both Poissonian noise and cosmic variance. Unfortunately, the covariance properties derived from mock samples are not necessarily a good description of those of the real data, thus making the use of the derived covariance matrix (e.g. in model fitting) doubtful. Conversely, depending on the sample size, *jack-knife* or volume-bootstrap covariance matrices can exacerbate the peculiarities of some subregions, again not adequately representing the true covariance properties of the data.

For the present investigation, we put considerable effort in understanding how to best estimate a sensible covariance matrix for our $w_p(r_p)$ measurements. The available mock samples were crucial for allowing us to perform direct comparisons of the performances of the different techniques. After some initial attempts, we excluded the *jack-knife* method because of the limited size of the survey volume. We then directly compared the covariance matrices derived through the bootstrap technique and from the KW24 mock catalogues. For the bootstrap method, we decided

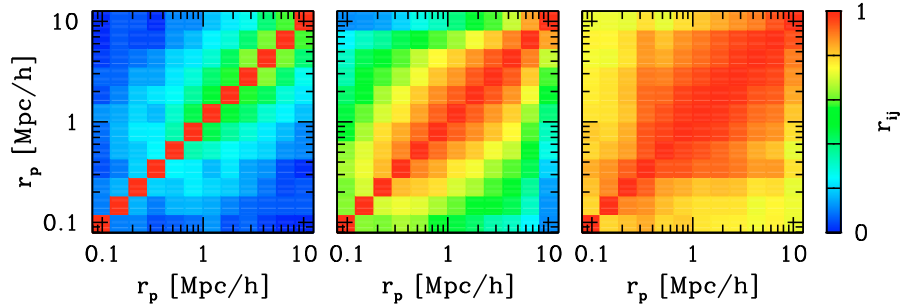


Fig. 9. Mean of the 24 correlation matrix derived resampling the galaxies of each KW24 mocks (*left*), or resampling 8 equal subvolumes (*center*). These are compared to the correlation matrix derived directly from the KW24 mocks (*right*). The redshift range considered here is $z = [0.5-0.8]$. The averaging over the 24 realisations of the 2 left matrices suppress the negative off-diagonal terms, which are sometimes present for a given mock catalogues. Correlation coefficients are then colour-coded from 0 to 1.

to directly test how galaxy- and volume-bootstrap were performing. We concentrated on the redshift range $z = [0.5-0.8]$ by selecting simulated galaxies brighter than $M_B - 5 \log(h) \leq -19.5 - z$. After computing the correlation function $w_p(r_p)$ for all 24 mock samples, we constructed for each of them a) 100 galaxy-galaxy-bootstrap samples and b) 100 volume-volume-bootstrap samples. In the latter case, we considered 8 equal subvolumes, defined as redshift slices within the redshift range considered. The number of subvolumes was chosen as the best compromise between having enough of them and not having volumes that were too small. With this choice, their volume is $\sim 1.4 \times 10^5 h^{-3} \text{ Mpc}^3$ for the samples with $z = [0.5-0.8]$ and $z = [0.8-1.0]$ and $\sim 0.6 \times 10^5 h^{-3} \text{ Mpc}^3$ for $z = [0.2-0.5]$. The two bootstrap techniques led to a total of 4800 samples and corresponding estimates of $w_p(r_p)$. We then calculated the covariance (and correlation) matrices for each of these two cases, along with the one derived from the correlation function of the 24 mocks themselves.

In Fig. 8 we compare the standard deviations derived from the two bootstrap techniques, to those derived from the 24 mocks. In each case, these values correspond by definition to the square root of the diagonal elements of the covariance matrix. In the plot we show the mean (over the 24 mocks) of the ratio of σ_{w_p} from the bootstrap to the “true” one from the ensemble of mock surveys. This shows clearly how the rms values obtained with the single-galaxy-bootstrap grossly underestimate the true variance, up to one order of magnitude on large scales. Bootstrapping by volumes produces a better result, providing a realistic estimate of σ_{w_p} between 0.1 and $1 h^{-1} \text{ Mpc}$, and a 20–25% underestimate on larger scales.

Each element of the the correlation matrix r_{ij} is obtained from the corresponding element of the covariance matrix σ_{ij} as $r_{ij} = \sigma_{ij} / \sqrt{\sigma_{ii}\sigma_{jj}}$. By definition, the off-diagonal terms of the correlation matrix will then range between -1 and 1 , indicating the degree of correlation between different scales of the function $w_p(r_p)$. Considering the redshift range $z = [0.5-0.8]$, we show in Fig. 9 the mean of the 24 correlation matrices derived by resampling the galaxies (left panel) or by resampling 8 equal subvolumes (centre), 100 times each. These are compared to the correlation matrix directly derived from the 24 mocks (right panel). The first case shows a mainly diagonal correlation matrix where off-diagonal terms are mostly noise. In the second case they instead decrease smoothly from 1 to 0 as a function of bin separation. The matrix derived from the 24 mocks shows high correlation on all scales.

In order to directly compare the properties of the correlation matrices derived with the 3 methods, we compute the principal

Table 4. The five main eigenvalues of the correlation matrix derived with the bootstrap resampling of galaxies (first column) and subvolumes (Col. 2), and from the ensemble variance of the 24 mocks (Col. 3).

Eigenvalue	Bootstrap		Mocks catalogues
	Galaxies	Volumes	
λ_1	3.87784	8.28193	11.82062
λ_2	1.95834	2.17866	0.15708
λ_3	1.36841	0.71635	0.01794
λ_4	1.11316	0.38252	0.00436
λ_5	0.91683	0.20340	0.00000

For the two latest cases, each mock sample is used in turn as “data” and the reported eigenvalues are the obtained as the average over the 24 mocks.

components and the amplitudes of the corresponding eigenvalues λ_i ($i = 1-12$) for each of the $24 + 24 + 1$ matrices. The sum of the eigenvalues of a correlation matrix is always equal to its dimension, i.e. 12 in our case. We report in Table 4 the values of the five main eigenvalues obtained with the 24 mocks (first column) compared to the averages over the 24 mocks of those obtained with the two resampling methods. The numbers show that the correlation matrix derived from the 24 mocks essentially contains four principal components and is mostly dominated by one of them. This indicates a strong correlation in the data. The bootstrap matrices, instead, show more than five non-negligible components, with the fifth one the same order of magnitude as the second in the mock matrix. This implies a lower correlation. We note, however, that volume resampling tends to produce a matrix whose structure is closer to that of the mocks, with 1–2 dominant components. This is another indication of how volume-bootstrapting, although not perfectly reproducing the intrinsic covariance properties of the sample, better estimates the variance and correlation in the data than does a galaxy-galaxy-bootstrap.

These experiments are extended and further discussed in our parallel accompanying papers, in particular Porciani et al. (in preparation). The bottom-line result of our extensive investigations is that a volume-bootstrap provides a good enough reconstruction of the intrinsic covariance matrix of the data set, if enough resamplings are used. This is obtained at the expense of a slightly less accurate account of cosmic variance on large scale than what can be obtained from the scatter among mock samples, where wavelengths longer than the survey size can be sampled. However, we have shown (Fig. 8) that this effect on scales $\sim 10 h^{-1} \text{ Mpc}$ is limited to $\sim 20\%$.

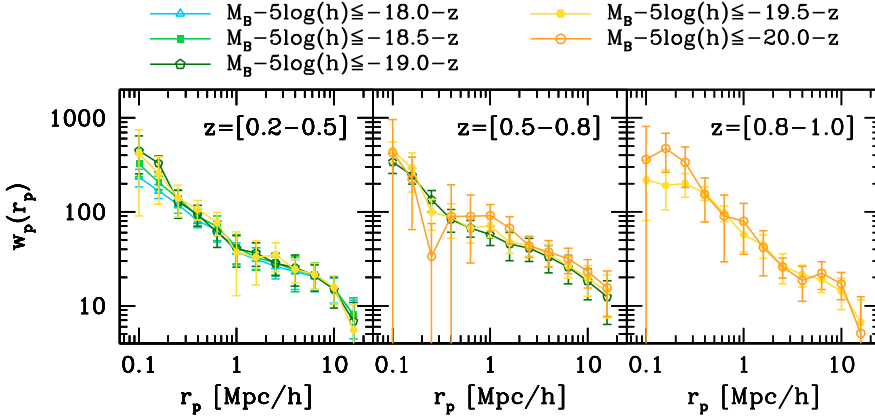


Fig. 10. Projected correlation function $w_p(r_p)$ measured as a function of galaxy luminosity within three redshift ranges. No significant dependence on luminosity is observed within the explored ranges.

8. Dependence of galaxy clustering on luminosity

8.1. Luminosity dependence at fixed redshift

Figure 10 shows the projected correlation function $w_p(r_p)$ estimated for our nine luminosity-selected subsamples at different redshifts. Error bars correspond to the 1σ dispersion provided by 200 volume-bootstrap resamplings, as extensively discussed in Sect. 7.

No clear dependence on luminosity is observed within any of the three redshift ranges. Also, in the shape of $w_p(r_p)$, there is some hint of the usual “shoulder”, i.e. a change in slope around $1 h^{-1}$ Mpc, but no clear separation between the expected 1-halo term on small scales and the 2-halo component above this scale (see the Introduction for definitions). In particular, in the intermediate-redshift bin, all subsamples show a rather flat large-scale slope, with no evidence of the usual breakdown above $\sim 2 h^{-1}$ Mpc.

To try to understand the origin of the observed flat shape, it is interesting to look directly at the contour plots of the bidimensional redshift-space correlation function $\xi(r_p, \pi)$. These are shown in Fig. 11 for the three luminosity-selected subsamples L1.4, L2.2 and L3.1 (see Table 1 for definitions), which include galaxies brighter than $M_B - 5 \log(h) \leq -19.5 - z$. The three contour plots show some interesting features. First, one clearly notices the much stronger distortion along the line of sight π , in the central panel. At the same time, a much more extended signal is also observed along the perpendicular direction r_p in the same redshift range. It is tempting to interpret both these effects as produced in some way by the two dominating structures that we showed in Fig. 6. The excess signal along the line of sight is very plausibly due to the distortions by “Fingers of God”, due to an anomalous number of virialised systems (groups and clusters) within these structures. At the same time, the extension along r_p indicates that there is also an excess of pairs perpendicular to the line of sight, with respect to an isotropic distribution. In fact, we know (Guzzo et al. 2007; Scoville et al. 2007b) that the large-scale structure at $z \simeq 0.73$ extends over a large part of the COSMOS area. This evidently biases the observed number of pairs along r_p , for simple geometrical reasons. We cannot exclude that part of the large-scale compression observed in $\xi(r_p, \pi)$ is also generated by an excess of galaxy infall onto this structure, thus producing what is known as the Kaiser effect (Kaiser 1987). This effect is proportional to the growth of structure (see e.g. Guzzo et al. 2008, for a recent direct estimate at similar redshift) and can be extracted when the underlying clustering can be assumed to be isotropic. In this case it is in practice impossible to disentangle this dynamical distortion from

the geometrical anisotropy generated by having one dominating structure elongated perpendicular to the line of sight.

The flatter shape in $w_p(r_p)$ in Fig. 10 is also consistent with the overdense samples of Abbas & Sheth (2007), who notice not only a higher amplitude for the most overdense (10% and 30%) samples of mock and SDSS galaxies, but also a flattening in the correlation function compared to the full sample. This is another line of evidence favouring the hypothesis that the zCOSMOS field is centred on an overdensity.

The plots of Fig. 11 also explicitly show the reasons for our choice of $\pi_{\max} = 20 h^{-1}$ Mpc as the upper integration limit in computing $w_p(r_p)$, a value that provides a reasonable compromise between including most of the signal and excluding the noisiest regions in the upper part the diagrams. In the central redshift bin, however, some real clustering power may still be present above this scale, for small r_p 's. In Fig. 12 we show directly how $w_p(r_p)$ changes, when π_{\max} is extended from 20 to $30 h^{-1}$ Mpc. We see that, somewhat counter intuitively, below $1 h^{-1}$ Mpc, no extra amplitude is gained, while – as indicated by the mock experiments (see Sect. 5) – the scatter is increased. Conversely, one can see the slight scale-dependent bias on the amplitude at larger separations, which gets up to $\sim 10\%$ at $15 h^{-1}$ Mpc when increasing π_{\max} .

8.2. Redshift evolution at fixed (evolving) luminosity

In Sect. 4.1 we discussed how our luminosity selection was devised such as to account for the average evolution in the luminosity of galaxies, assuming this to be the dominant effect in modifying the mean density of objects at any given luminosity. Under this assumption, it is then interesting to test how $w_p(r_p)$ changes with redshift for galaxies within the same (de-evolved) luminosity interval. This is shown in Fig. 13 for $M_B - 5 \log(h) \leq -19.5 - z$, i.e. for the same three samples for which $\xi(r_p, \pi)$ is plotted in Fig. 11. No coherent evolution of the amplitude and shape of the projected correlation function is observed among the three samples, characterised by mean redshifts 0.37, 0.61, and 0.91. The three curves are consistent within the error bars, with – as expected – the intermediate-redshift bin (triangles) showing a systematically higher amplitude than the other two. Again, this is easily interpreted as a local effect, resulting from the extreme large-scale clustering observed in this redshift range. On the other hand, the overall lack of apparent amplitude evolution of luminous galaxies with redshift is consistent with previous results from the VVDS-Deep (Pollo et al. 2006), DEEP2 (Coil et al. 2006), and SDSS (Zehavi et al. 2005) surveys for galaxies brighter than $\sim M^*$. The only evolutionary effect shown in

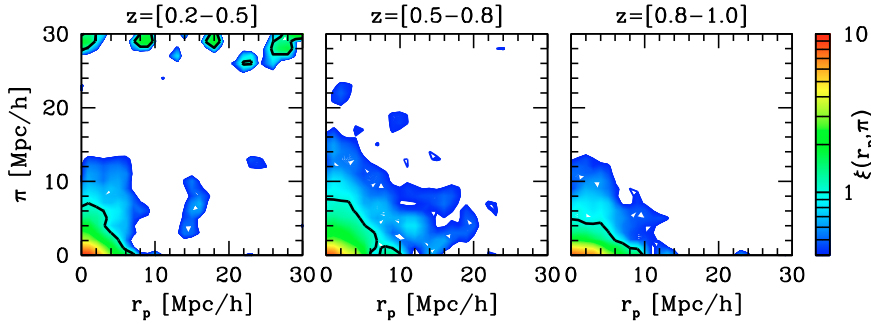


Fig. 11. Iso-correlation contours of $\xi(r_p, \pi)$ (here smoothed with a Gaussian kernel) for galaxies brighter than $M_B - 5 \log(h) \leq -19.5 - z$, computed in 3 redshift ranges. The amplitude is colour-coded according to the side bar, while the black contour corresponds to $\xi(r_p, \pi) = 1$. White values correspond to $\xi(r_p, \pi) < 0.4$.

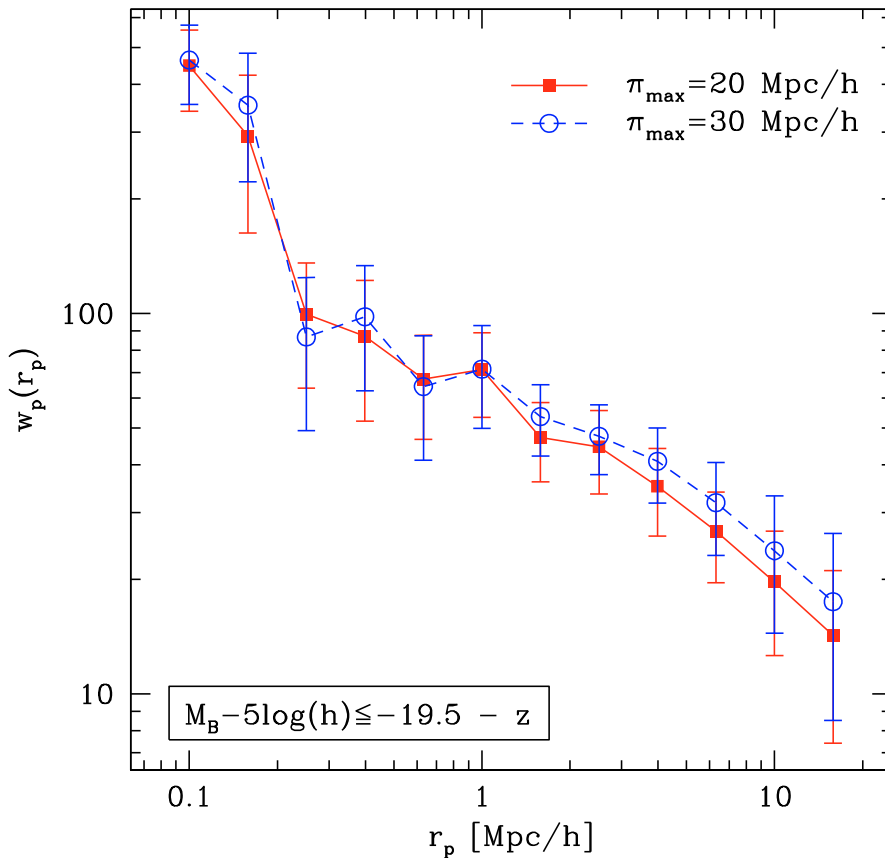


Fig. 12. Sensitivity of the projected function $w_p(r_p)$ to the upper integration limit π_{\max} , for one luminosity-selected sample in the central redshift bin.

particular by the VVDS data is a steepening of the small-scale slope of $w_p(r_p)$ (i.e. the 1-halo term) for high-redshift luminous galaxies (Pollo et al. 2006). McCracken et al. (2008) also show a lack of clustering amplitude evolution for a large sample of luminous ($-22 \leq M_B - 5 \log(h) \leq -19$) galaxies based on accurate photometric redshifts in the CFHTLS Deep fields; interestingly, they show that such invariance is also maintained when splitting the sample into early- and late-type galaxies.

9. Dependence of galaxy clustering on stellar mass

The relation of clustering properties to galaxy stellar masses is in principle more informative and straightforward for interpreting because stellar mass is a more fundamental physical parameter than luminosity.

9.1. Mass dependence at fixed redshift

Also in the case of stellar mass dependence, it is interesting to look at the shape of the full correlation function $\xi(r_p, \pi)$ in the three redshift ranges. We show in Fig. 14 the result obtained for the 3 samples M1.3, M2.3, and M3.2 (see Table 2), which include galaxies more massive than $10^{10} h^{-2} M_\odot$. Also in this case, the central panel is significantly different from the other two, with $\xi(r_p, \pi)$ remaining positive out to much larger scales in both r_p and π directions. We note that the small-scale stretching along π seems to be less extended than the one obtained for the corresponding luminosity-selected sample, although it is hard to say whether this difference is significant.

Figure 15 shows the projected correlation function $w_p(r_p)$ of the 10 mass-selected samples. The plotted points are not corrected for the residual stellar mass incompleteness (see Sect. 6.3). Errors are estimated as in the luminosity case using 200 bootstrap resamplings of 8 equal subvolumes of each data

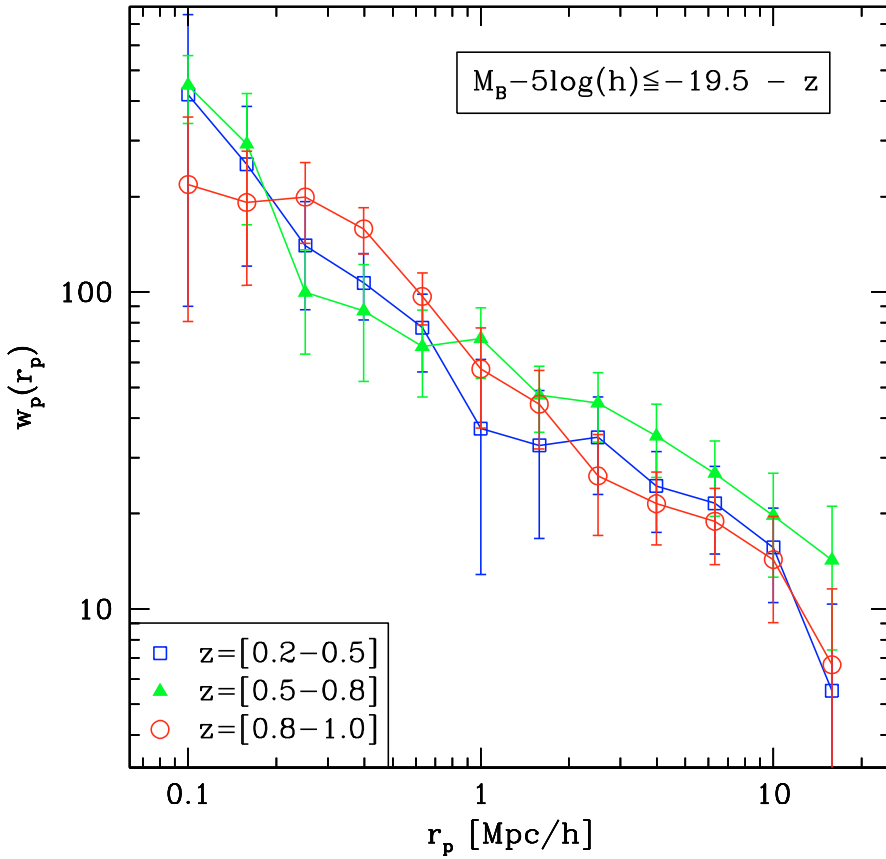


Fig. 13. Evolution of the projected function $w_p(r_p)$ of galaxies with $M_B - 5 \log(h) \leq -19.5 - z$ between redshift $z = 0.2$ and $z = 1$.

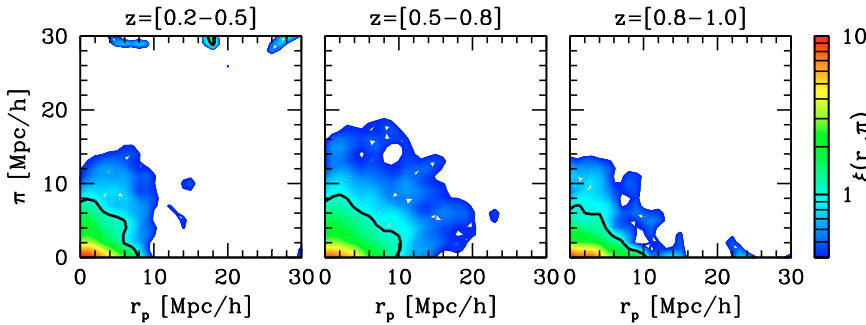


Fig. 14. Example of full redshift-space correlation function $\xi(r_p, \pi)$ (here smoothed with a Gaussian kernel) for the subsample containing galaxies more massive than $10^{10} h^{-2} M_\odot$, computed in the same 3 redshift ranges as in Fig. 11. $\xi(r_p, \pi)$ is computed in cells of $1 h^{-1}$ Mpc side in both r_p and π and the iso-correlation levels are colour-coded according to the side bar. The thick black contour corresponds to $\xi(r_p, \pi) = 1$ and the white values to $\xi(r_p, \pi) < 0.4$.

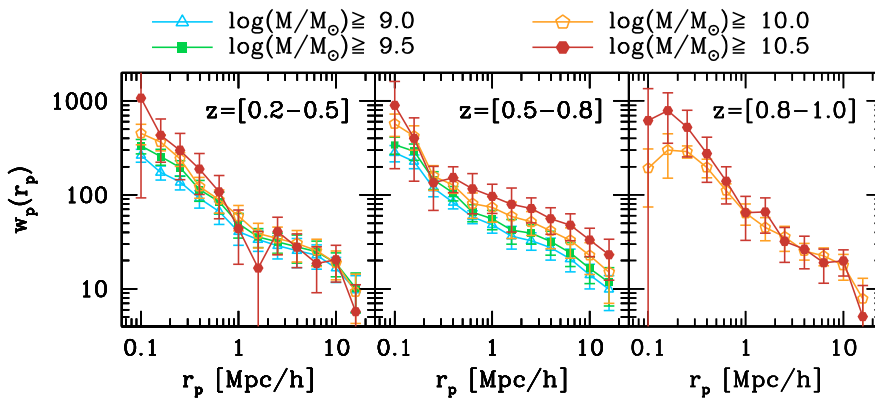


Fig. 15. The projected correlation function $w_p(r_p)$ as a function of galaxy stellar masses in the zCOSMOS 10K sample, in three redshift ranges.

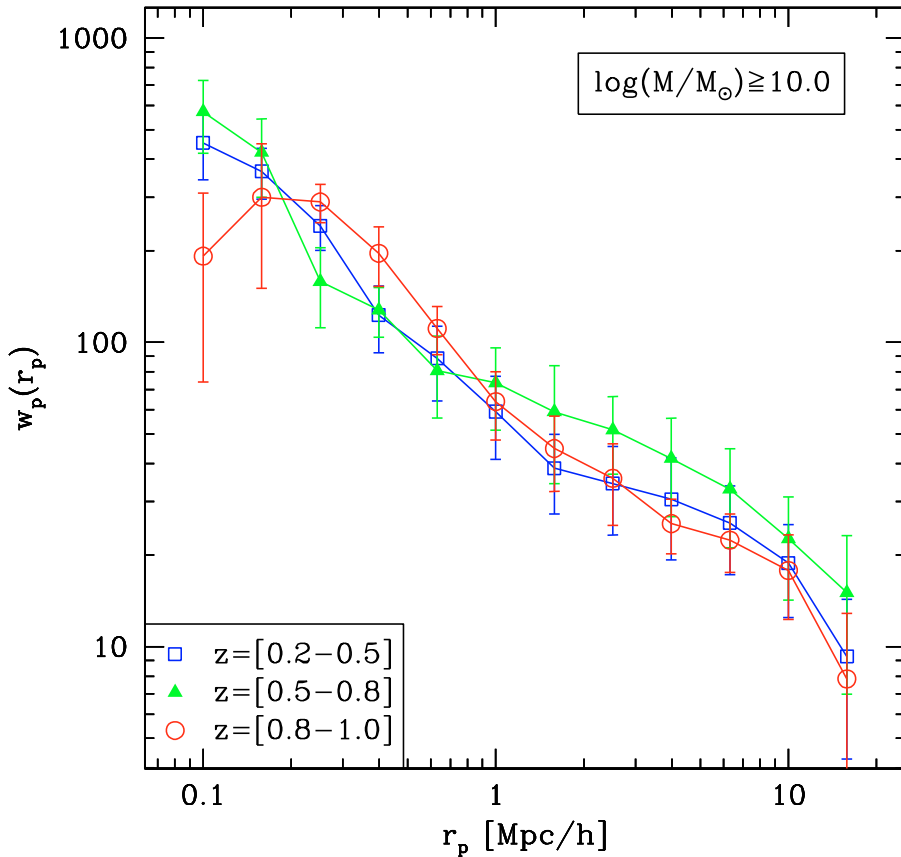


Fig. 16. Evolution of the projected function $w_p(r_p)$ of galaxies with $\log(M/M_\odot) \geq 10$ between redshift $z = 0.2$ and $z = 1$.

set. The figure shows a weak mass dependence of clustering in the low- and high-redshift bins, in particular at small separations. At the same time, a strong dependence at all separations is evident in the intermediate-redshift slice. There, the slope of $w_p(r_p)$ remains extremely flat out to the largest explored scales, even more strongly than in the luminosity-selected cases. Finally, in the low- and high-redshift bins there is evidence of a steeper “1-halo term” contribution at $r_p < 1 h^{-1}$ Mpc (with no clear indication of an evolution in redshift of the transition scale to the 2-halo term). Conversely, the central redshift range seems to be characterised by the same, flat power-law shape down to $0.2 h^{-1}$ Mpc, where a sudden steepening is then observed. The slope below $0.2 h^{-1}$ Mpc seems to depend directly on the limiting mass, with more massive galaxies showing a steeper correlation function. In summary, no clear overall trend can be shown among the three redshift ranges, with the central volume again displaying peculiar clustering properties that apparently dominate any possible cosmological effect.

9.2. Clustering evolution at fixed stellar mass

It is also interesting to directly compare the evolution of $w_p(r_p)$ with redshift, when a specific class of stellar mass is selected. As mentioned earlier in this section, this is particularly interesting because it should not require accounting for strong galaxy evolutionary trends as in the case of luminosity. We are assuming here that stellar mass does not grow significantly between $z \sim 1$ and $z \sim 0.2$, which we know is only partially true. A factor of ~ 2 growth in stellar mass is in fact expected on average between $z = 1$ and $z = 0$, which would have little effect on the estimated correlation function, however, if taken into account.

In Fig. 16 we show $w_p(r_p)$ computed for the same three samples with $\log(M/M_\odot) \geq 10$ of Fig. 14. Similarly to the luminosity samples, we do not see any clear evolution with redshift of the amplitude and shape of the projected correlation function. The three curves are consistent with each other within 1σ .

10. Comparison with independent measurements and models

10.1. Redshift evolution of $w_p(r_p)$

An accurate $z \sim 0$ reference measurement of $w_p(r_p)$ as a function of stellar mass has been obtained by the SDSS (Li et al. 2006). Meneux et al. (2008) do find evidence for evolution of the amplitude of $w_p(r_p)$ for galaxies less massive than $10^{10.5} h^{-2} M_\odot$, when comparing this to the measurements from VVDS-Deep at $z \sim 0.9$. The SDSS and zCOSMOS stellar masses were derived with the same initial mass function (Chabrier 2003) and normalised to $h = 1$. They are directly comparable. The SDSS clustering measurements were obtained within differential stellar mass ranges (Li et al. 2006), while ours correspond to galaxies that are more massive than a given threshold. However, from Fig. 2 we see that the zCOSMOS sample includes a very small number of galaxies that are more massive than $10^{11} h^{-2} M_\odot$, due to the much smaller volume than for the SDSS. Any of our mass-selected samples has therefore, in practice, an upper bound at this value of mass. This implies that we can coherently compare two of the SDSS measurements of $w_p(r_p)$ (for their galaxy samples with stellar masses in the ranges $[10.0-10.5]$ and $[10.5-11.0]$) to those from our samples M3.1 and M3.2, which include galaxies more massive than 10^{10} and $10^{10.5} h^{-2} M_\odot$,

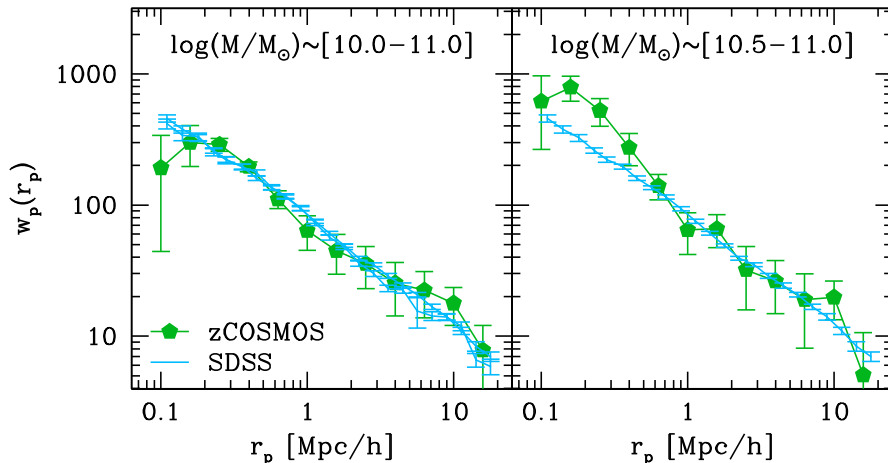


Fig. 17. The measured projected function $w_p(r_p)$ for galaxies at $z = [0.8-1.0]$ with $\log(M/M_\odot) \sim [10-11]$ (left) and $\log(M/M_\odot) \sim [10.5-11]$ (right) from the zCOSMOS survey (filled diamonds), compared to the $z \sim 0.1$ estimate by the SDSS (Li et al. 2006) (blue curve).

respectively, within the redshift range $z = [0.8-1.0]$. This comparison (Fig. 17) does not show a clear evolution with redshift. For both samples, the large-scale amplitude of $w_p(r_p)$ is virtually the same as in the local SDSS samples. Considering a simple evolution of structures, this implies that the bias for galaxies more massive than $10^{10} h^{-2} M_\odot$ has evolved significantly between $z \sim 1$ and today, as to keep their apparent clustering amplitude substantially unchanged. This implies in practice that the bias $b(z)$ must evolve in such a way that $b(0)D(0) \simeq b(z)D(z)$, where $D(z)$ is the linear growth rate of density fluctuations. In the standard model, this implies that, at the approximate mean redshifts of our redshift bins, $z = 0.35, 0.75, 0.9$, the bias of massive galaxies must have been, respectively, 1.2, 1.44, and 1.53 times its value at the current epoch. Meneux et al. (2008) observed the same effect at $\langle z \rangle \sim 0.8$ in the VVDS data but only for galaxies more massive than $\sim 10^{10.5} h^{-2} M_\odot$, with lower-mass objects showing a weaker bias evolution. A lack of any evolution of the clustering of the most massive galaxies was also noticed in the NDWFS (Brown et al. 2008) and 2SLAQ surveys (Wake et al. 2008).

10.2. Observed and predicted shape of $w_p(r_p)$ at $0.5 < z < 1$

The only available measurement of clustering as a function of stellar mass at redshifts comparable to those explored by our sample comes from the VVDS-Deep survey (Meneux et al. 2008) at $0.5 < z < 1.2$. VVDS-Deep goes 1.5 magnitude deeper (although over a smaller area of $\sim 0.5 \text{ deg}^2$), which allows the analysis to be extended beyond $z = 1$. To provide a qualitative, yet meaningful comparison of these two data sets, we can recompute the correlation function for the 10K data within the widest usable redshift range overlapping with the VVDS interval, i.e. $[0.5-1.0]$. We applied the same stellar mass selection limits, keeping in mind the residual incompleteness that will affect the highest redshift part of the sample. The result is shown in Fig. 18, where the VVDS and zCOSMOS mass-selected samples are directly compared. The difference in shape and amplitude in the $w_p(r_p)$ derived from the two data sets is rather striking. The zCOSMOS points show in general a much flatter relation than those from the VVDS. The amplitudes for a given mass selection also seem to be incompatible at several standard deviations between the two samples, especially above $1 h^{-1} \text{ Mpc}$.

10.3. Comparison to analytic and semi-analytic models

At this point, it is relevant to compare the available observations from zCOSMOS and VVDS with model predictions in a standard Λ CDM scenario. We can do this in two ways. We first used the *HALOFIT* public code (Smith et al. 2003), which uses the halo model to compute the expected nonlinearly evolved power spectrum at $z = 0.8$, which we take as a reasonable mean redshift for the two samples. Our conclusions would not differ at all if predictions for $z = 0.7$ or 0.9 were used. The corresponding projected function $w_p(r_p)$ is then computed by Fourier-transforming the power spectrum and projecting the resulting real-space correlation function. The result gives the expected $w_p(r_p)$ of the mass density field at $z = 0.8$. Second, we can compute the expectation value and the scatter expected in the same redshift range ($0.5 < z < 1$) for $w_p(r_p)$ using the available semi-analytic mock surveys built from the Millennium run. To this end, we used the DLB40 mocks for which we have full control over stellar masses, selecting simulated galaxies with $I_{AB} \leq 22.5$ and $\log(M/M_\odot) \geq 10$ and reproducing the sampling rate of the 10K data. In Fig. 19 we plot the *HALOFIT* prediction both for the dark matter and for an arbitrarily biased population of haloes with $b = \sqrt{2.6} \sim 1.61$, together with the DLB40 mean $w_p(r_p)$ and the corresponding 1σ and 3σ scatter corridor from the 40 mock surveys. As a consistency check, we note the rather good agreement between the analytic *HALOFIT* result and the expected value from the full N-body plus semi-analytic simulation. On these model predictions, we overplot the corresponding zCOSMOS and VVDS estimates. The zCOSMOS points agree with the models at better than 68% confidence in both shape and amplitude on scales smaller than $1 h^{-1} \text{ Mpc}$. On larger scales, however, the observed $w_p(r_p)$ would require a strongly scale-dependent bias to be compatible with the model predictions. This scale-dependence would also behave contrary to what models and very general considerations suggest, implying a bias that grows with scale, rather than declines. From the plot we see in fact that the 10K data are compatible with $b \simeq 1.6$ on small scales, but would require $b \simeq 2.45$ on $10 h^{-1} \text{ Mpc}$ scales. The shaded area shows that this large-scale excess is marginally compatible with the model predictions, representing a very strong positive fluctuation. A few percent of volumes this size would show this high clustering amplitude (on these scales and for this kind of galaxies), in a Λ CDM Universe.

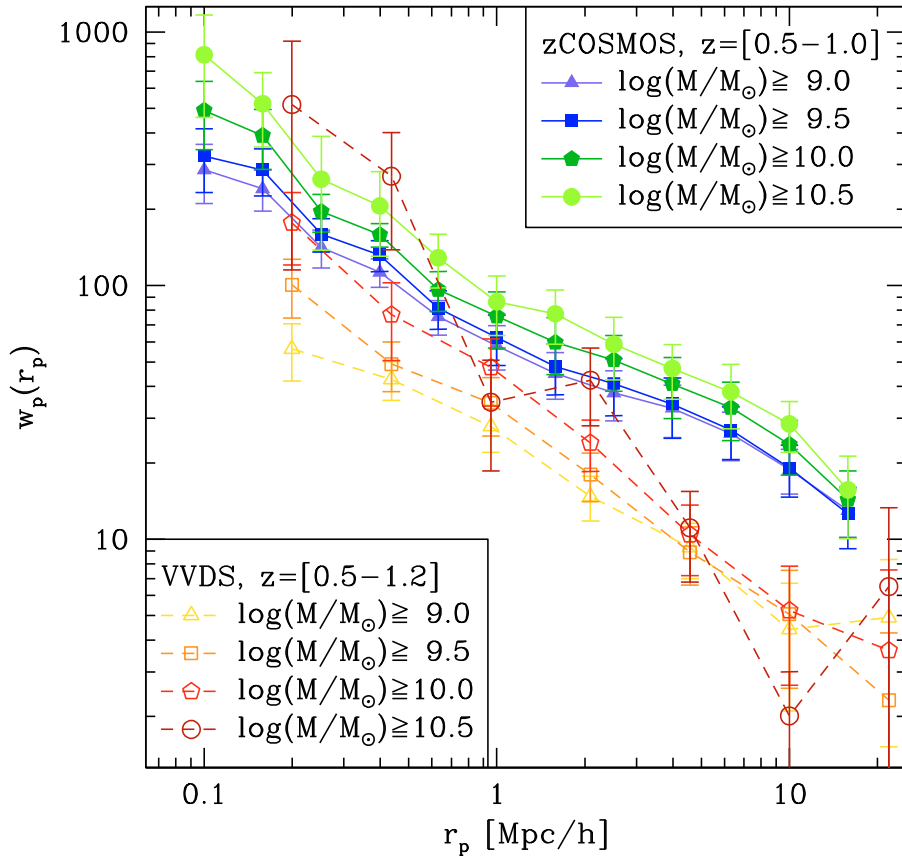


Fig. 18. Direct comparison of the dependence of clustering on stellar mass in the VVDS-Deep and zCOSMOS samples, over a similar redshift range.

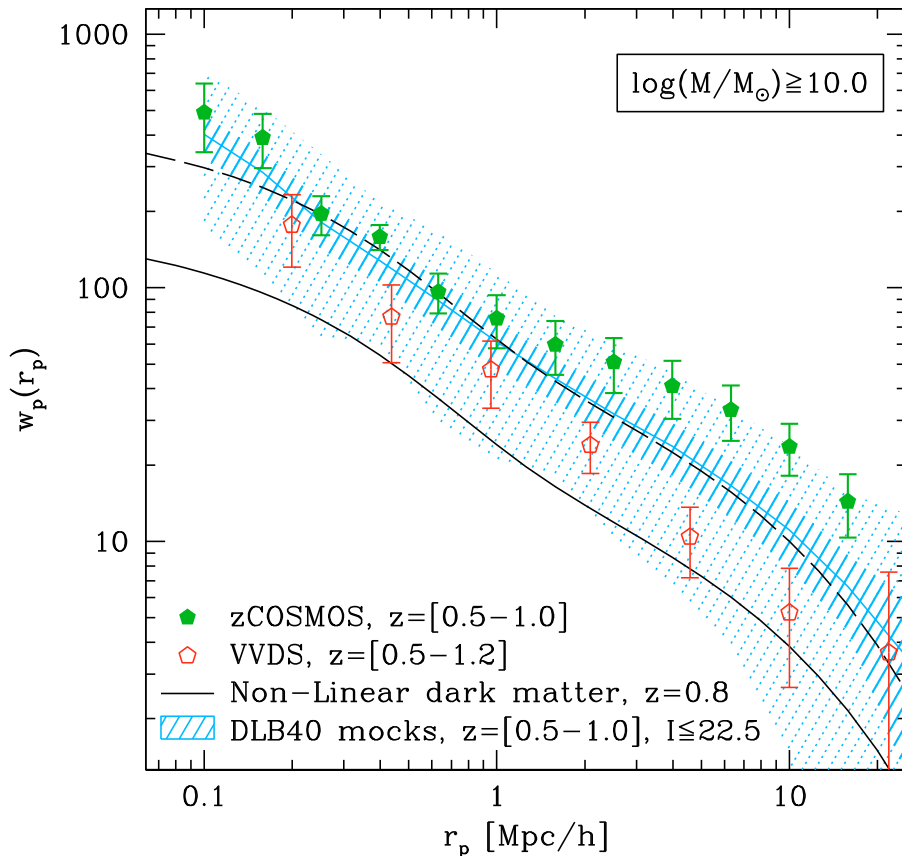


Fig. 19. The zCOSMOS (solid circles) and VVDS-Deep (open circles) $w_p(r_p)$ of galaxies with mass larger than $10^{10} h^{-2} M_\odot$ and $\langle z \rangle \sim 0.8$, compared to model predictions. These include the nonlinear mass projected correlation function computed using *HALOFIT* (lower solid line, Smith et al. 2003), and, for reference, the corresponding $w_p(r_p)$ for a population of haloes with bias $b^2 = 2.6$ (dashed curve). The latter curve is a very good description of the full nonlinear $w_p(r_p)$ (light solid blue line), obtained averaging the 40 DLB40 mocks from the Millennium run after applying the same sampling, magnitude and mass selections of the 10K sample. The shaded areas (thick and thin shaded) give, respectively, the 1σ and 3σ confidence corridors around the mean.

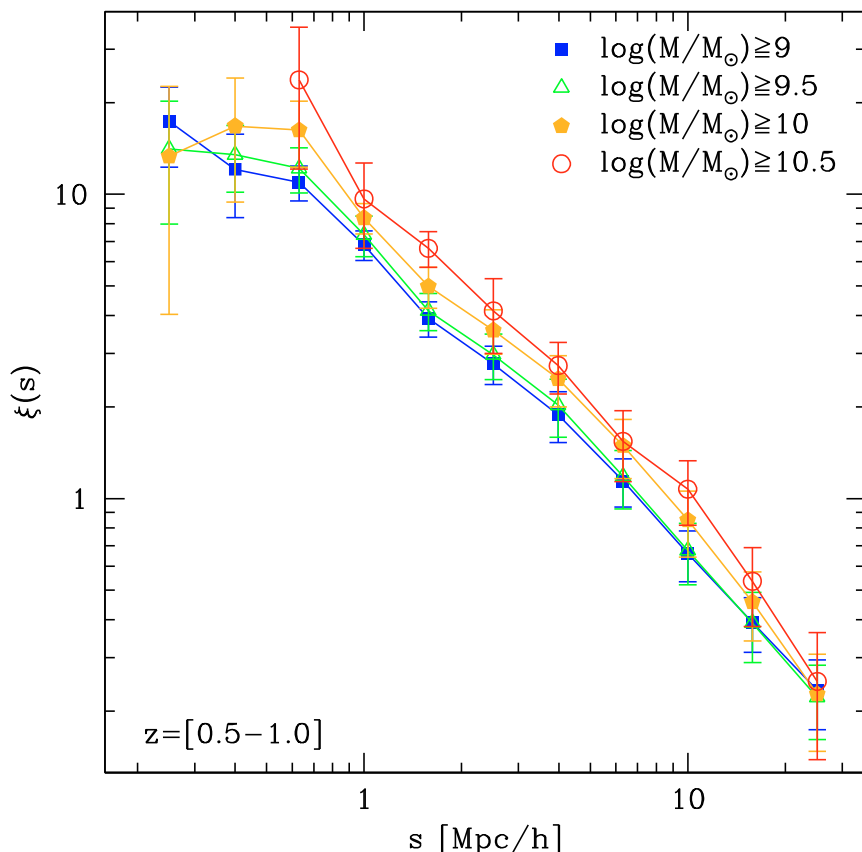


Fig. 20. The redshift-space, angle-averaged correlation function $\xi(s)$ of the mass-selected samples in the redshift bin $[0.5, 1]$. A mild systematic mass dependence is visible.

It is interesting to note, at the same time, how the VVDS measurements lie on the opposite side of the distribution, at about $1.5\text{--}2\sigma$ from the mean, but with a shape that is compatible with the model prediction over the whole range (corresponding to a linear bias $b \sim 1.2$). Based on the results of Abbas & Sheth (2007), one would also conclude that the central volume of the COSMOS survey is dominated by overdense regions, while these should be slightly under-represented in the corresponding volume of the VVDS survey.

These results show how a full HOD model fitting to the $w_p(r_p)$ measured from the 10K data – originally planned for this paper – would add no meaningful information to the current analysis. Our first experiments with HOD models based on the universal halo mass function indicate that rather unrealistic sets of parameters are required to reproduce the observed function. An interesting possibility would be to use in such modelling a halo mass function that depends on local environment (e.g. Abbas & Sheth 2005, 2006), to consider the evidence that a large part of this sample is dominated by an overdensity. We plan to explore this possibility using the larger 20K zCOSMOS sample that is now nearly complete.

11. Discussion

Together with previous analyses (McCracken et al. 2007; Kovač et al. 2009), these results suggest that a significant fraction of the volume of Universe bounded by the COSMOS field is indeed characterised by particularly extreme density fluctuations. We have seen how, in statistical terms, these seem to lie at the 3σ limit of the distribution of amplitudes expected in volumes of a few $10^6 h^{-3} \text{Mpc}^3$. We should consider, however, that these conclusions are drawn from measurements that are strongly

affected by the angular distribution of structure. The McCracken et al. (2007) result is based on the angular correlation function, while here we studied the projected function $w_p(r_p)$. Although making use of the redshift information, the latter is in practice a clustering measure dominated by galaxy pairs lying almost perpendicular to the line of sight. The underlying assumption when measuring $w_p(r_p)$ is that the geometrical distribution of structures within the sample being analysed is completely isotropic, in other words, that there are superclusters aligned along several directions, such that the only remaining radial signal is produced by galaxy-peculiar velocities. The very reason for using $w_p(r_p)$ is indeed to get rid of the distortions introduced in the shape of $\xi(s)$ (the redshift-space, angle-averaged correlation function) by galaxy motions. If this is true, and only in this case, then $w_p(r_p)$ is fully equivalent to an integral over the real-space correlation function $\xi(r)$, so it carries the same cosmological information. However, if, as in the case we have encountered here, there is one or more dominating structures extending preferentially along one direction, then the use of $w_p(r_p)$ to infer cosmological information is inappropriate.

One may thus wonder whether more robust cosmological information could instead be inferred by looking directly at the simplest, angle-averaged redshift-space correlation function $\xi(s)$. The expectation is that the average over all directions reduces the weight of the excess pair counts produced by just a few structures oriented along one preferred direction. In such case any analytic modelling (e.g. with HOD models) should also include an appropriate model for the linear and nonlinear redshift distortions (Scoccimarro 2004; Tinker et al. 2007). More simply, we can use the available mock samples in redshift-space to compute the nonlinear redshift-space $\xi(s)$ and its variance and compare it to the data, as we did for $w_p(r_p)$. In

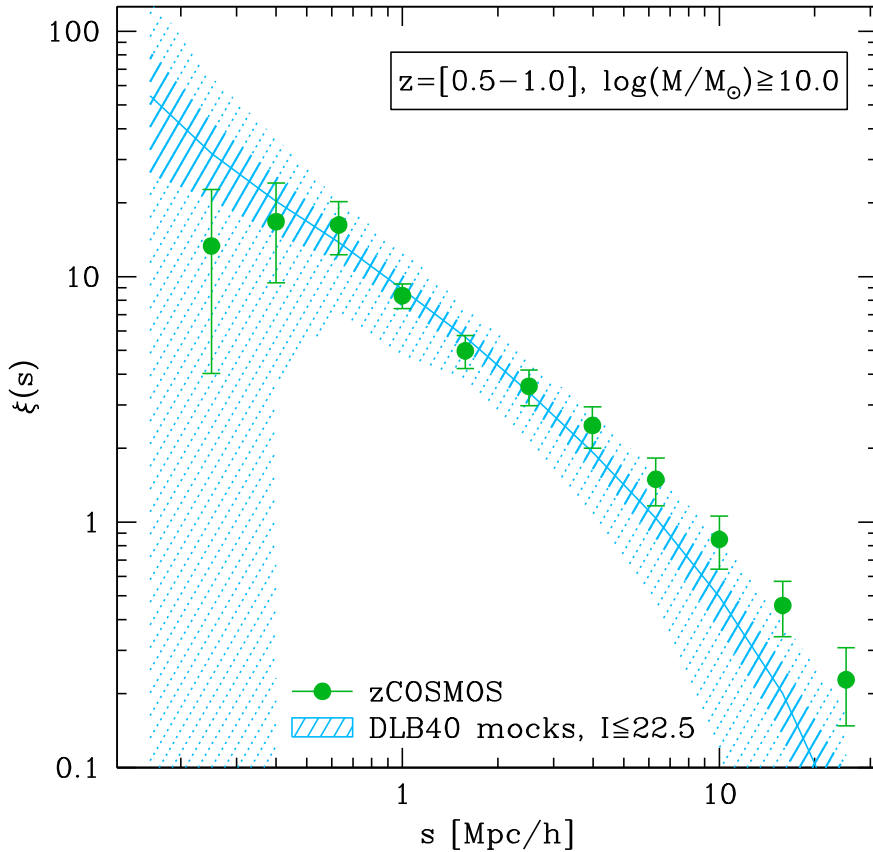


Fig. 21. Comparison of the redshift-space correlation function for galaxies with $\log(M/M_{\odot}) \geq 10$ at $0.5 < z < 1$ in the zCOSMOS sample (filled circles), with the model predictions from the DLB40 mock samples. The solid line gives the average of the 40 mocks, with the dashed areas corresponding to the 1σ and 3σ confidence error corridor. Similarly to what we found for $w_p(r_p)$, the agreement between the models and the zCOSMOS measurement at $1 h^{-1}$ Mpc is remarkable. Nevertheless, the large-scale shape of the zCOSMOS $\xi(s)$ is at $\sim 3\sigma$ the mean amplitude of the mock catalogues.

Fig. 20 we first plot $\xi(s)$ for the 10K sample, computed for the usual four mass ranges in the broad redshift range $0.5 < z < 1$. The four data sets show a smooth power-law behaviour, with some evidence of a mass dependence of the clustering amplitude, in particular at the upper mass limit. The overall shape is well described by a rather flat power-law $\xi(s) \sim (s_0/s)^\gamma$, with slope $\gamma \sim 1$ and a correlation length s_0 between 6 and $\sim 10 h^{-1}$ Mpc. These values for the shape and amplitude of $\xi(s)$ are similar to those measured for luminous red galaxies at $z = 0.55$ in the 2SLAQ survey (see Fig. 7 in Ross et al. 2007) and for luminous early-type galaxies in the 2dFGRS (Norberg et al. 2002). This is consistent with the most massive objects in the 10K sample being predominantly red, early-type galaxies that show moderate or no evolution in the overall clustering amplitude with redshift.

In Fig. 21, instead, we compare $\xi(s)$ of our “reference” sample with $\log(M/M_{\odot}) \geq 10$, with the mean and scatter (at 1σ and 3σ confidence, respectively) of the similarly-selected set of DLB40 mocks. Despite the angular average, we note a behaviour that is similar to what is observed in $w_p(r_p)$, although now the agreement extends to slightly larger scales. The observed clustering is compatible with the predictions of the standard model (to better than the 68% level) on scales smaller than $\sim 2 h^{-1}$ Mpc. On larger scales, $\xi(s)$ also shows excess power with respect to the models, which places the zCOSMOS volume at the upper 3σ limit of the statistical distribution obtained from the mocks. This exercise shows that, even after angle-averaging our clustering estimator, the amount of structure present in this specific volume of the Universe remains outstanding in comparison to the model expectations. The conclusion can only be that we have either been very unlucky in the selection of the COSMOS field and

picked up a fluctuation with a probability of $\sim 1\%$ to be found in such a volume or fluctuations with this amplitude are in reality more common than what the standard cosmology predicts.

12. Summary

We used the 10K zCOSMOS spectroscopic sample to study galaxy clustering as a function of galaxy luminosity and stellar mass, in the range of redshift $[0.2, 1]$. To this end, we built luminosity and mass-selected samples from the 10K catalogue sampling three separate redshift ranges. We used mock catalogues to quantify the effect of stellar mass incompleteness on the measured clustering, as a function of redshift. We carefully checked our covariance and error estimate techniques, comparing the performances of methods based on the scatter in the mocks and on bootstrapping schemes. We adopted the latter, based on 200 resamplings of 8 subvolumes of the survey, as the most appropriate description of the covariance properties of the data.

By measuring the redshift-space correlation functions $\xi(s)$ and $\xi(r_p, \pi)$ and the projected function $w_p(r_p)$ for these subsamples, we found the following results.

- Surprisingly, we do not see any clear dependence on luminosity of the correlation function at all redshifts. This is at odds with results in the local Universe by the 2dFGRS and with measurements at similar redshift by the VVDS and DEEP2 surveys, which found a significant steepening of $w_p(r_p)$ with luminosity.
- We find a mildly more evident (although not striking) dependence of $w_p(r_p)$ on stellar mass, especially on small scales. The central redshift bin ($0.5 < z < 0.8$) displays in general a more evident effect, with a very flat shape of $w_p(r_p)$

on scales $r_p = [1-10] h^{-1}$ Mpc. The overall shape of the corresponding map of $\xi(r_p, \pi)$ shows strong distortions that we interpret as the effect of dominant structure extending preferentially perpendicular to the line of sight.

- From comparison to the SDSS measurements at $z \sim 0$, we do not see any significant evolution with redshift of the amplitude of clustering for bright and/or massive galaxies. Together with previous results from VVDS, this is consistent with a more rapid evolution of the linear bias for the most massive objects with respect to the general population. In the zCOSMOS sample, this invariance in the clustering amplitude between $z \sim 1$ and $z \sim 0$ seems to remain valid down to lower masses than in the VVDS, an effect easily explained by the overall larger clustering amplitude observed in general for $z > 0.5$ in this sample. This is shown by a much flatter shape (higher amplitude) of $w_p(r_p)$ of zCOSMOS galaxies with respect to VVDS galaxies, when selected with the same criteria.
- This particularly high level of structure is confirmed by comparing the measured $w_p(r_p)$ and $\xi(s)$ at $0.5 < z < 1$ with model predictions, concentrating on the sample with $\log(M/M_\odot) \geq 10$. On scales smaller than $\sim 1-2 h^{-1}$ Mpc, the observations agree very well with the model expectation values in the standard Λ CDM scenario for a linear bias $b \sim 1.6$. On these scales, the measured values are compatible to better than 68% with the DLB40 mocks. On larger scales, however, the observed clustering amplitude is reproduced in only a few percent of the mocks. In other words, if the shape of the power spectrum is that of Λ CDM and the bias has no “unnatural” scale-dependence, COSMOS has picked up a volume of the Universe that is rare, $2-3\sigma$ positive fluctuation. This conclusion is also corroborated by comparison with the VVDS measurements, which on the other hand lie on the lower side of the distribution, at about $1.5-2\sigma$.

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References

Abbas, U., & Sheth, R. K. 2005, MNRAS, 364, 1327
 Abbas, U., & Sheth, R. K. 2006, MNRAS, 372, 1749
 Abbas, U., & Sheth, R. K. 2007, MNRAS, 378, 641
 Abbas, U., de la Torre, S., Le Fèvre, O., et al. 2009, A&A, submitted
 Arnouts, S., Walcher, C. J., Le Fèvre, O., et al. 2007, A&A, 476, 137
 Bell, E. F., Naab, T., McIntosh, D. H., et al. 2006, ApJ, 640, 241
 Benoist, C., Maurogordato, S., da Costa, L. N., Cappi, A., & Schaeffer, R. 1996, ApJ, 472, 452
 Bernstein, G. M. 1994, ApJ, 424, 569
 Bolzonella, M., Miralles, J.-M., & Pelló, R. 2000, A&A, 363, 476
 Bolzonella, M., Kovač, K., Pozzetti, L., et al. 2009, A&A, submitted [arXiv:0907.0013]
 Blaizot, J., Wadadekar, Y., Guiderdoni, B. et al. 2005, MNRAS, 360, 159
 Brown, M. J. I., Zheng, Z., White, M., et al. 2008, ApJ, 682, 937
 Bruzual, G., & Charlot, S. 2003, MNRAS, 344, 1000
 Capak, P., Aussel, H., Ajiki, M., et al. 2007, ApJS, 172, 99
 Chabrier, G. 2003, PASP, 115, 763
 Coil, A. L., Newman, J. A., Cooper, M. C., et al. 2006, ApJ, 644, 671
 Coil, A. L., Newman, J. A., Croton, D., et al. 2008, ApJ, 672, 153
 Cooray, A., & Sheth, R. 2002, Phys. Rep., 372, 1
 Conroy, C., Wechsler, R. H., & Kravtsov, A. V. 2006, ApJ, 647, 201
 Daddi, E., Röttgering, H. J. A., Labbé, I., et al. 2003, ApJ, 588, 50
 Davis, M., & Geller, M. J. 1976, ApJ, 208, 13
 Davis, M., & Peebles, J. P. E. 1983, ApJ, 267, 465
 de la Torre, S., Le Fèvre, O., et al. 2009, A&A, submitted
 De Lucia, G., & Blaizot, J. 2007, MNRAS, 375, 2
 Feldmann, R., Carollo, C. M., Porciani, C., et al. 2006, MNRAS, 372, 565

Fisher, K. B., Davis, M., Strauss, M. A., Yahil, A., & Huchra, J. 1994, MNRAS, 266, 50
 Garilli, B., Le Fèvre, O., Guzzo, L., et al. 2008, A&A, 486, 683
 Giacomini, R., Zirm, A., Wang, J., et al. 2002, ApJS, 139, 369
 Giovanelli, R., Haynes, M. P., & Chincarini, G. L. 1986, ApJ, 300, 77
 Guzzo, L., Strauss, M. A., Fisher, K. B., Giovanelli, R., & Haynes, M. P. 1997, ApJ, 489, 37
 Guzzo, L., Bartlett, J. G., Cappi, A., et al. 2000, A&A, 355, 1
 Guzzo, L., Cassata, P., Finoguenov, A., et al. 2007, ApJS, 172, 254
 Guzzo, L., Pierleoni, M., Meneux, B., et al. 2008, Nature, 451, 541
 Hamilton, A. J. S. 1988, ApJ, 331, L59
 Iovino, A., Giovanelli, R., Haynes, M., Chincarini, G., & Guzzo, L. 1993, MNRAS, 265, 21
 Ilbert, O., Tresse, L., Zucca, E., et al. 2005, A&A, 439, 863
 Ilbert, O., Capak, P., Salvato, M., et al. 2009, ApJ, 690, 1236
 Kaiser, N. 1987, MNRAS, 227, 1
 Kitzbichler, M. G., & White, S. D. M. 2007, MNRAS, 376, 2
 Kovač, K., Lilly, S. J., Cucciati, O., et al. 2009, ApJ, submitted [arXiv:0903.3409]
 Landy, S. D., & Szalay, A. S. 1993, ApJ, 412, 64
 Le Fèvre, O., Saisse, M., Mancini, D., et al. 2003, Proc. SPIE, 4841, 1670
 Le Fèvre, O., Vettolani, G., Garilli, B., et al. 2005a, A&A, 439, 845
 Le Fèvre, O., Guzzo, L., Meneux, B., et al. 2005b, A&A, 439, 877
 Li, C., Kauffmann, G., Jing, Y. P., et al. 2006, MNRAS, 368, 21
 Lilly, S. J., Le Fèvre, O., Renzini, A., et al. 2007, ApJS, 172, 70
 Lilly, S. J., Le Brun, V., Maier, C., et al. 2009, ApJS, submitted
 Ling, E. N., Barrow, J. D., & Frenk, C. S. 1986, MNRAS, 223, 21
 Maurogordato, S., & Lachieze-Rey, M. 1991, ApJ, 369, 30
 McCracken, H. J., Peacock, J. A., Guzzo, L., et al. 2007, ApJS, 172, 314
 McCracken, H. J., Ilbert, O., Mellier, Y., et al. 2008, A&A, 479, 321
 Meneux, B., Le Fèvre, O., Guzzo, L., et al. 2006, A&A, 452, 387
 Meneux, B., Guzzo, L., Garilli, B., et al. 2008, A&A, 478, 299
 Norberg, P., Baugh, C. M., Hawkins, E., et al. 2001, MNRAS, 328, 64
 Norberg, P., Baugh, C. M., Hawkins, E., et al. 2002, MNRAS, 332, 827
 Norberg, P., Baugh, C. M., Gaztanaga, E., & Croton, D. J. 2009, MNRAS, 396, 19
 Peebles, P. J. E. 1980, The Large Scale Structure of the Universe (Princeton: Princeton University Press)
 Phleps, S., Peacock, J. A., Meisenheimer, K., & Wolf, C., 2006, A&A, 457, 145
 Pollo, A., Meneux, B., Guzzo, L., et al. 2005, A&A, 439, 887
 Pollo, A., Guzzo, L., Le Fèvre, O., et al. 2006, A&A, 451, 409
 Porciani, C., & Giavalisco, M. 2002, ApJ, 565, 24
 Pozzetti, L., Bolzonella, M., Lamareille, F., et al. 2007, A&A, 474, 443
 Pozzetti, L., Bolzonella, M., Zucca, E., et al. 2009, A&A, submitted
 Press, W. H., & Schechter, P. 1974, ApJ, 187, 425
 Rettura, A., Rosati, P., Strazzullo, V., et al. 2006, A&A, 458, 717
 Ross, N. P., da Ángela, J., Shanks, T., et al. 2007, MNRAS, 381, 573
 Scoccimarro, R. 2004, Phys. Rev. D, 70, 083007
 Scoville, N., Abraham, R. G., Aussel, H., et al. 2007a, ApJS, 172, 38
 Scoville, N., Aussel, H., Benson, A., et al. 2007b, ApJS, 172, 150
 Skibba, R. A., Sheth, R. K., & Martino, M. C. 2007, MNRAS, 382, 1940
 Smith, R. E., Peacock, J. A., Jenkins, A., et al. 2003, MNRAS, 341, 1311
 Springel, V., White, S. D. M., Jenkins, A., et al. 2005, Nature, 435, 629
 Springel, V., Frenk, C. S., & White, S. D. M. 2006, Nature, 440, 1137
 Stringer, M. J., Benson, A. J., Bundy, K., Ellis, R. S., & Quetin, E. L. 2009, MNRAS, 143
 Tinker, J. L., Norberg, P., Weinberg, D. H., & Warren, M. S. 2007, ApJ, 659, 877
 Wake, D. A., Sheth, R. K., Nichol, R. C., et al. 2008, MNRAS, 387, 1045
 Wang, L., Li, C., Kauffmann, G., & De Lucia, G. 2006, MNRAS, 371, 537
 Wang, L., Li, C., Kauffmann, G., & De Lucia, G. 2007, MNRAS, 377, 1419
 Warren, M. S., Abazajian, K., Holz, D. E., & Teodoro, L. 2006, ApJ, 646, 881
 Willmer, C. N. A., da Costa, L. N., & Pellegrini, P. S. 1998, AJ, 115, 869
 White, S. D. M., & Rees, M. J. 1978, MNRAS, 183, 341
 Zehavi, I., Blanton, M. R., Frieman, J. A., et al. 2002, ApJ, 571, 172
 Zehavi, I., Zheng, Z., Weinberg, D. H., et al. 2005, ApJ, 630, 1
 Zheng, Z., Coil, A., & Zehavi, I. 2007, ApJ, 667, 760
 Zucca, E., Ilbert, O., Bardelli, S., et al. 2006, A&A, 455, 879
 Zucca, E., Bardelli, S., Bolzonella, M., et al. 2009, A&A, submitted

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