

Regular Article - Theoretical Physics

Lund jet images from generative and cycle-consistent adversarial networks

Stefano Carrazza^{1,a}, Frédéric A. Dreyer²

- ¹ TIF Lab, Dipartimento di Fisica, Università degli Studi di Milano and INFN Milan, Via Celoria 16, 20133 Milan, Italy
- ² Rudolf Peierls Centre for Theoretical Physics, University of Oxford, Clarendon Laboratory, Parks Road, Oxford OX1 3PU, UK

Received: 8 September 2019 / Accepted: 19 November 2019 © The Author(s) 2019

Abstract We introduce a generative model to simulate radiation patterns within a jet using the Lund jet plane. We show that using an appropriate neural network architecture with a stochastic generation of images, it is possible to construct a generative model which retrieves the underlying two-dimensional distribution to within a few percent. We compare our model with several alternative state-of-the-art generative techniques. Finally, we show how a mapping can be created between different categories of jets, and use this method to retroactively change simulation settings or the underlying process on an existing sample. These results provide a framework for significantly reducing simulation times through fast inference of the neural network as well as for data augmentation of physical measurements.

1 Introduction

One of the most common objects emerging from hadron collisions at particle colliders such as the Large Hadron Collider (LHC) are jets. These are loosely interpreted as collimated bunches of energetic particles arising from the interactions of quarks and gluons, the fundamental constituents of the proton [1,2]. In practice, jets are usually defined through a sequential recombination algorithm mapping final-state particle momenta to jet momenta, with a free parameter *R* defining the radius up to which separate particles are clustered into a single jet [3–5].

Because of the high energies involved in the collisions at the LHC, heavy particles such as vector bosons or top quarks are frequently produced with very large transverse momenta. In this boosted regime, the decay products of these objects can become so collimated that they are reconstructed as a single jet. An active field of research is therefore dedicated to the theoretical understanding of radiation patterns within jets, notably to distinguish their physical origins and remove radiation unassociated with the hard process [6–26]. Furthermore, measurements of jet properties provide a unique opportunity for accurate comparisons between theoretical predictions and data, and can be used to tune simulation tools [27] or extract physical constants [28].

In recent years, there has also been considerable interest in applications of generative adversarial networks (GAN) [29] and variational autoencoders (VAE) [30] to particle physics, where such generative models can be used to significantly reduce the computing resources required to simulate realistic LHC data [31–40]. In this paper, we introduce a generative model to create new samples of the substructure of a jet from existing data. We use the Lund jet plane [22], shown in Fig. 1, as a visual representation of the clustering history of a jet. This provides an efficient encoding of a jets radiation patterns and can be directly measured experimentally [41]. The Lund jet image is used to train a Least Square GAN (LSGAN) [42] to reproduce simulated data within a few percent accuracy. We compare a range of alternative generative methods, and show good agreement between the original jets generated with Pythia v8.223 [43] using fast detector simulation with Delphes v3.4.1 particle flow [44] and samples provided by the different models [45]. Finally, we show how a cycleconsistent adversarial network (CycleGAN) [46] can be used to create mappings between different categories of jets. We apply this framework to retroactively change the parameters of the parton shower on an event, adding non-perturbative effects to an existing parton-level sample, and transforming quark and gluon jets to a boosted W sample.

These methods provide a systematic tool for data augmentation, as well as reductions of simulation time and storage space by several orders of magnitude, e.g. through a fast inference of the neural network with hardware architectures such as GPUs and field-programmable gate arrays (FPGA)

Published online: 27 November 2019



a e-mail: stefano.carrazza@cern.ch

979 Page 2 of 11 Eur. Phys. J. C (2019) 79:979

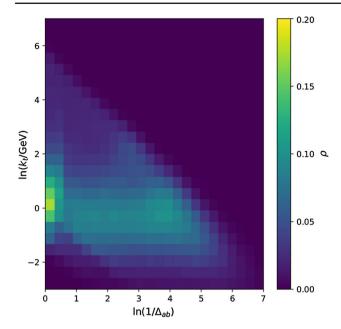


Fig. 1 Average Lund jet plane density for QCD jets simulated with Pythia v8.223 and Delphes v3.4.1

[47]. The code frameworks and data used in this work are available as open-source and published material in [48–50].

2 Generating jets

In this article we will construct a generative model, which we call gLund, to create new samples of radiation patterns of jets. We first introduce the basis used to describe a jet as an image, then construct a generative model which can be trained on these objects.

2.1 Encoding radiation patterns with Lund images

To describe the radiation patterns of a jet, we will use the primary Lund plane representation [22], which can be projected onto a two-dimensional image that serves as input to a neural network.

The Lund jet plane is constructed by reclustering a jet's constituents with the Cambridge-Aachen (C/A) algorithm [4, 51]. This algorithm sequentially recombines pairs of particles that have the minimal $\Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$ value, where y_i and ϕ_i are the rapidity and azimuth of particle i.

This clustering sequence can be used to construct an $n \times n$ pixel image describing the radiation patterns of the initial jet. We iterate in reverse through the clustering sequence, labelling the momenta of the two branches of a declustering as p_a and p_b , ordered in transverse momentum such that

¹ The codes are available at https://github.com/JetsGame/gLund and https://github.com/JetsGame/CycleJet.



 $p_{t,a} > p_{t,b}$. This procedure follows the harder branch a and at each step we activate the pixel on the image corresponding to the coordinates ($\ln \Delta_{ab}$, $\ln k_t$), where $k_t = p_{t,b} \Delta_{ab}$ is the transverse momentum of particle b relative to a.²

2.2 Input data

The data sample used in this article consists of 500k jets, generated using the dijet process in Pythia v8.223. Jets are clustered using the anti- k_t algorithm [5,52] with radius R=1.0, and are required to pass a selection cut, with transverse momentum $p_t>500$ GeV and rapidity |y|<2.5. Unless specified otherwise, results use the Delphes v3.4.1 fast detector simulation, with the delphes_card_CMS_NoFastJet.tcl card to simulate both detector effects and particle flow reconstruction.

The simulated jets are then converted to Lund images with 24×24 pixels each using the procedure described in Sect. 2.1. A pixel is set to one if there is a corresponding $(\ln \Delta_{ab}, \ln k_t)$ primary declustering sequence, otherwise it is left at zero.

The full samples used in this article can be accessed online [50].

2.3 Probabilistic generation of jets

Generative adversarial networks [53] are one of the most successful unsupervised learning methods. They are constructed using both a generator G and discriminator D, which are competing against each other through a value function V(G, D).

In practice, we found improved performance when using a Least Square Generative Adversarial Network (LSGAN) [42], a specific class of GAN which uses a least squares loss function for the discriminator, and has objective functions defined as

$$\min_{D} V(D) = \frac{1}{2} \mathbb{E}_{x \sim p_{\text{data}}} [(D(x) - b)^{2}]
+ \frac{1}{2} \mathbb{E}_{z \sim p_{z}(z)} [(D(G(z)) - a)^{2}],$$
(1)

$$\min_{G} V(G) = \frac{1}{2} \mathbb{E}_{z \sim p_{z}(z)} [(D(G(z)) - c)^{2}], \tag{2}$$

where we defined $p_z(z)$ as a prior on input noise variables, and a, b and c are the labels for the fake, real and presumed fake data respectively. Thus D is trained in order to maximise the probability of correctly distinguishing the training examples and the samples from G, following Eq. (1), while

² For simplicity, we consider only whether a pixel is on or off, instead of counting the number of hits as in [22]. While these two definitions are equivalent only for large image resolutions, this limitation can easily be overcome e.g. by considering a separate channel for each activation level.

Eur. Phys. J. C (2019) 79:979 Page 3 of 11 979

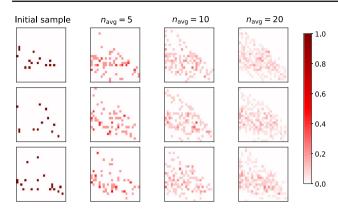


Fig. 2 Sample input images after averaging with $n_{\text{avg}} = 1, 5, 10$ and 20

the latter is trained to minimise Eq. (2). The generator's distribution p_g optimises Eq. (2) when $p_g = p_{\rm data}$, so that the generator learns how to generate new samples from z. The main advantage of the LSGAN over the original GAN framework is a more stable training process, due to an absence of vanishing gradients. In addition, we include a minibatch discrimination layer [54] to avoid collapse of the generator.

The LSGAN is trained on the full sample of QCD Lund jet images. In order to overcome the limitation of GANs due to the sparse and discrete nature of Lund images, we will use a probabilistic interpretation of the Lund images to train the model. To this end, we will first re-sample our initial data set into batches of n_{avg} and create a new set of 500k images, each consisting of the average of n_{avg} initial input images, as shown in Fig. 2. These images can be reinterpreted as physical events through a random sampling, where the pixel value is interpreted as the probability that the pixel is activated. The $n_{\rm avg}$ value is a parameter of the model, with a large value leading to increased variance in the generated images compared to the reference sample, while for too low values the model performs poorly due to the sparsity and discreteness of the data. A further data preprocessing step before training the LSGAN consists in rescaling the pixel intensities to be in the [-1, 1] range, and masking entries outside of the kinematic limit of the Lund plane. The images are then whitened using zero-phase components analysis (ZCA) whitening [55].

2.4 gLund model results

The optimal choice of hyperparameters, both for the LSGAN model architecture and for the image preprocessing, is determined using the distributed asynchronous hyperparameter optimisation library hyperopt [56].

The performance of each setup is evaluated by a loss function which compares the reference preprocessed Lund images to the artificial images generated by the LSGAN model. We define the loss function as

$$\mathcal{L}_h = I + 5 \cdot S \tag{3}$$

where I is the norm of the difference between the average of the images of the two samples and S is the absolute difference in structural similarity [57] values between 5000 random pairs of reference samples, and reference and generated samples.

We perform 1000 iterations and select the one for which the loss \mathcal{L}_h is minimal. In Fig. 3 we show some of the results obtained with the hyperopt library through the Treestructured Parzen Estimator (TPE) algorithm. The LSGAN is constructed from a generator and discriminator. The generator consists in three dense layers with 512, 1024 and 2048 units respectively using LeakyReLU [58] activation functions and batch normalisation layers, as well as a final layer matching the output dimension and using a hyperbolic tangent activation function. The discriminator is constructed from two dense layers with 768 and 384 units using a LeakyReLU activation function, followed by another 24dimensional dense layer connected to a minibatch discrimination layer, with a final fully connected layer with onedimensional output. The best parameters for this model are listed in Table 1. The loss of the generator and discriminator networks of the LSGAN is shown in Fig. 4 as a function of training epochs.

In Fig. 5, the first two images illustrate an example of input image before and after preprocessing while the last two images represent the raw output from the LSGAN model and the corresponding sampled Lund image.

A selection of preprocessed input images and images generated with the LSGAN model are shown in Fig. 6. The final averaged results for the Lund jet plane density are shown in Fig. 7 for the reference sample (left), the data set generated by the gLund model (centre) and the ratio between these two samples (right). We observe a good agreement between the reference and the artificial sample generated by the gLund model. The model is able to reproduce the underlying distribution to within a 3-5% accuracy in the bulk region of the Lund image. Larger discrepancies are visible at the boundaries of the Lund image and are due the vanishing pixel intensities. In practice this model provides a new approach to reduce Monte Carlo simulation time for jet substructure applications as well as a framework for data augmentation.

2.5 Comparisons with alternative methods

Let us now quantify the quality of the model described in Sect. 2.3 more concretely. As alternatives, we consider a variational autoencoder (VAE) [30,59,60] and a Wasserstein GAN [45,61].

A VAE is a latent variable model, with a probabilistic encoder $q_{\phi}(z|x)$, and a probabilistic decoder $p_{\theta}(x|z)$ to map a



979 Page 4 of 11 Eur. Phys. J. C (2019) 79:979

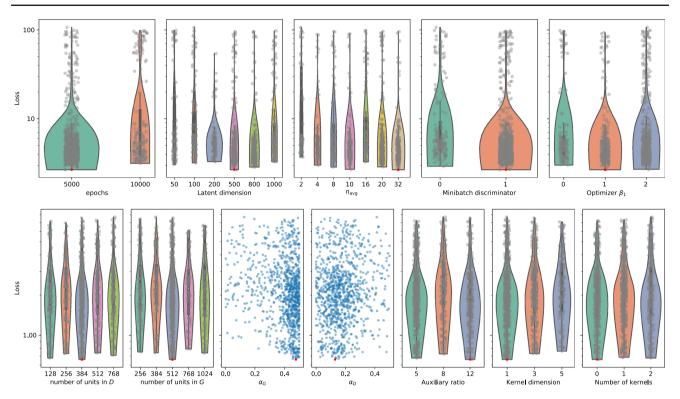


Fig. 3 Hyperparameter scan results obtained with the hyperopt library. The first row shows the scan over image and optimiser related parameters while the second row plots correspond to the final architecture scan

 $\textbf{Table 1} \ \ \text{Final parameters for the gLund model}$

Parameters	Value
Architecture	LSGAN
D units	384
G units	512
α_D	0.129
$lpha_G$	0.477
Aux ratio	12
Kernel dimension	1
Number of kernels	2
Minibatch discriminator	Yes
Epochs	5000
Batch size	32
Latent dimension	500
ZCA	Yes
n_{avg}	32
Learning rate	$6.5 \cdot 10^{-5}$
Decay β_1	$8 \cdot 10^{-9}$
Optimiser	Adagrad

representation from a prior distribution $p_{\theta}(z)$. The algorithm learns the marginal likelihood of the data in this generative process, which corresponds to maximising

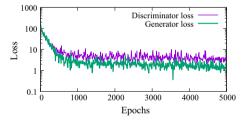


Fig. 4 Loss of the LSGAN discriminator and generator throughout the training stage

$$\mathcal{L}(\theta, \phi) = \mathbb{E}_{q_{\phi}(z|x)}[\log p_{\theta}(x|z)] - \beta D_{\text{KL}}(q_{\phi}(z|x)||p(z)), \tag{4}$$

where β is an adjustable hyperparameter controlling the disentanglement of the latent representation z. In our implementation, we will set $\beta=1$, which corresponds to the original VAE framework.

During the training of the VAE, we use KL cost annealing [62] to avoid a collapse of the VAE output to the prior distribution. This is a problem caused by the large value of the KL divergence term in the early stages of training, which is mitigated by adding a variable weight $w_{\rm KL}$ to the KL term in the cost function, expressed as

$$w_{\text{KL}}(n_{\text{step}}) = \min(1, 0.25 \cdot 1.05^{n_{\text{step}}}).$$
 (5)



Eur. Phys. J. C (2019) 79:979 Page 5 of 11 979

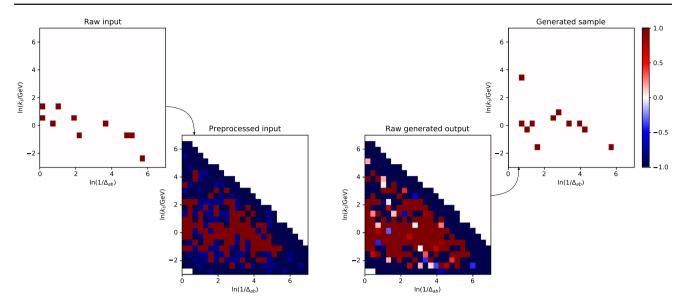


Fig. 5 Left two figures: sample input images before and after preprocessing. Right two: sample generated by the LSGAN and the corresponding Lund image

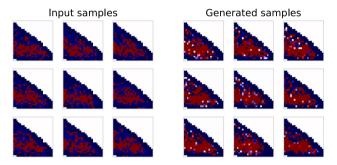


Fig. 6 A random selection of preprocessed input images (left), and of images generated with the LSGAN model (right). Axes and colour schemes are identical to Fig. 5

Finally, we will also consider a Wasserstein GAN with gradient penalty (WGAN-GP). WGANs [45] use the Wasserstein distance to construct the value function, but can suffer from undesirable behaviour due to the critic weight clipping. This can be mitigated through gradient penalty, where the norm of the gradient of the critic is penalised with respect to its input [61].

We determine the best hyperparameters for both of these models through a hyperopt parameter sweep, which is summarised in Appendix A. To train these models using Lund images, we then use the same preprocessing steps described in Sect. 2.3.

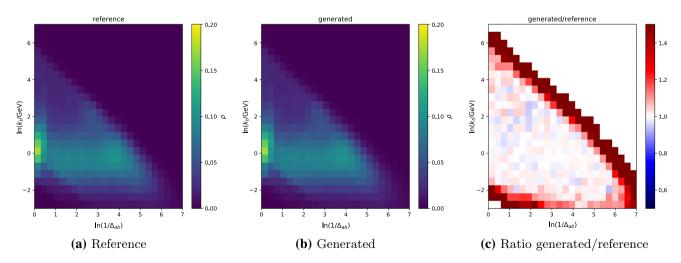


Fig. 7 Average Lund jet plane density for a the reference sample and b a data set generated by the gLund model. c The ratio between these two densities



979 Page 6 of 11 Eur. Phys. J. C (2019) 79:979

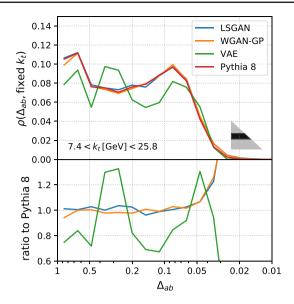


Fig. 8 Slice of the Lund plane along Δ_{ab} with 7.4 GeV $< k_t < 25.8 \; {\rm GeV}$

To compare our three models, we consider two slices of fixed k_t or Δ_{ab} size, cutting along the Lund jet plane horizontally or vertically respectively.

In Fig. 8, we show the k_t slice, with the reference sample in red. The lower panel gives the ratio of the different models to the reference Pythia 8 curve, showing very good performance for the LSGAN and WGAN-GP models, which are able to reproduce the data within a few percent. The VAE model also qualitatively reproduces the main features of the underlying distribution, however we were unable to improve the accuracy of the generated sample to more than 20% without avoiding the issue of posterior collapse. The same observations can be made in Fig. 9, which shows the Lund plane density as a function of k_t , for a fixed slice in Δ_{ab} .

In Fig. 10a we show the distribution of the number of activated pixels per image for the reference sample generated with Pythia 8 and the artificial images produced by the LSGAN, WGAN-GP and VAE models. All models except the VAE model provide a good description of the reference distribution.

We also use the Lund image to reconstruct the soft-drop multiplicity [63]. To this end, for a simpler correspondence between this observable and the Lund image, we retrained the generative models using $\ln(z\Delta)$ as y-axis. The soft-drop multiplicity can then be extracted from the final image, and is shown in Fig. 10b for each model using $z_{\rm cut}=0.007$ and $\beta=-1$. The dashed lines indicate the true reference distribution, as evaluated directly on the declustering sequence, and which differs slightly from the reconstructed curve due to the finite pixel and image size.

Finally, in Fig. 10c, we show the reconstructed mass of the groomed jet using the modified Mass Drop Tag-

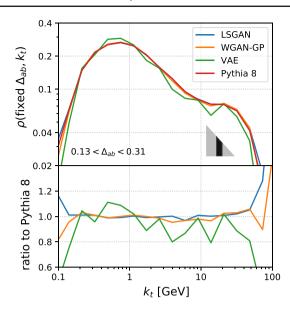


Fig. 9 Slice of the Lund plane along k_t with $0.13 < \Delta_{ab} < 0.31$

ger [17] with $z_{\text{cut}} = 0.1$, where we approximate the mass as

$$\rho = \frac{m^2}{R^2 p_t^2} \simeq \max_i \left[z^{(i)} \left(\Delta^{(i)} \right)^2 \right], \tag{6}$$

The dotted line shows the true mass distribution, evaluated with the left-hand side of Eq. (6) on the groomed jet. As in previous comparisons, we observe a very good agreement of the LSGAN and WGAN-GP models with the reference sample.

We note that while the WGAN-GP model is able to accurately reproduce the distribution of the training data, as discussed in Appendix A, the individual images themselves can differ quite notably from their real counterpart. For this reason, our preferred model in this paper is the LSGAN-based one.

3 Reinterpreting events using domain mappings

In this section, we will introduce a novel application of domain mappings to reinterpret existing event samples. To this end, we implement a cycle-consistent adversarial network (CycleGAN) [46], which is an unsupervised learning approach to create translations between images from a source domain to a target domain.

Using as input Lund images generated through different processes or generator settings, one can use this technique to create mappings between different types of jet. As examples, we will consider a mapping from parton-level to detectorlevel images, and a mapping from QCD images generated



Eur. Phys. J. C (2019) 79:979 Page 7 of 11 979

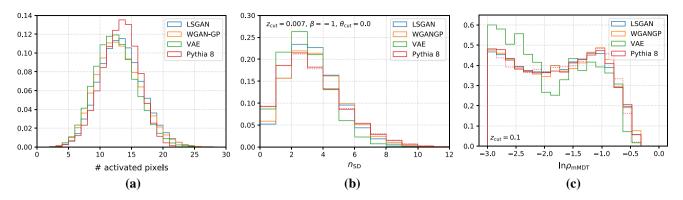


Fig. 10 Distribution of **a** the number of activated pixels per image, **b** the reconstructed soft-drop multiplicity for $z_{\text{cut}} = 0.007$, $\beta = -1$ and $\theta_{\text{cut}} = 0$, and **c** the jet mass after applying the modified Mass Drop Tagger with $z_{\text{cut}} = 0.1$

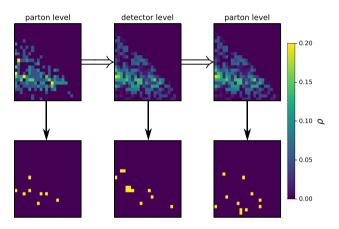


Fig. 11 Top: transition from parton-level to delphes-level and back using CycleJet. Bottom: corresponding sampled event

through Pythia 8's dijet process, to hadronically decaying W jets obtained from WW scattering.

The cycle obtained for a CycleGAN trained on parton and detector-level images is shown in Fig. 11, where an initial parton-level Lund image is transformed to a detector-level one, before being reverted again. The sampled image is shown in the bottom row.

3.1 CycleGANs and domain mappings

A CycleGAN learns mapping functions between two domains X and Y, using as input training samples from both domains. It creates an unpaired image-to-image translation by learning both a mapping $G: X \to Y$ and an inverse mapping $F: Y \to X$ which observes a forward cycle consistency $x \in X \to G(x) \to F(G(x)) \approx x$ as well as a backward cycle consistency $y \in Y \to F(y) \to G(F(y)) \approx y$. This behaviour is achieved through the implementation of a cycle consistency loss

$$\mathcal{L}_{\text{cyc}}(G, F) = \mathbb{E}_{x \sim p_{\text{data}}(x)}[\|F(G(x)) - x\|_{1}] + \mathbb{E}_{y \sim p_{\text{data}}(y)}[\|G(F(y)) - y\|_{1}], \tag{7}$$

Additionally, the full objective includes also adversarial losses to both mapping functions. For the mapping function $G: X \to Y$ and its corresponding discriminator D_Y , the objective is expressed as

$$\mathcal{L}_{GAN}(G, D_Y, X, Y) = \mathbb{E}_{y \sim p_{\text{data}}(y)}[\log D_Y(y)] + \mathbb{E}_{x \sim p_{\text{data}}(x)}[\log(1 - D_Y(G(x)))],$$
(8)

such that G is incentivized to generate images G(x) that resemble images from Y, while the discriminator D_Y attempts to distinguish between translated and original samples.

Thus, CycleGAN aims to find arguments solving

$$G^*, F^* = \arg\min_{G, F} \max_{D_X, D_Y} \mathcal{L}(G, F, D_X, D_Y),$$
 (9)

where \mathcal{L} is the full objective, given by

$$\mathcal{L}(G, F, D_X, D_Y) = \mathcal{L}_{GAN}(G, D_Y, X, Y) + \mathcal{L}_{GAN}(F, D_X, Y, X) + \lambda \mathcal{L}_{cvc}(G, F).$$
(10)

Here λ is parameter controlling the importance of the cycle consistency loss. We implemented a CycleGAN framework, labelled CycleJet, that can be used to create mappings between two domains of Lund images. By training a network on parton and detector-level images, this method can thus be used to retroactively add non-perturbative and detector effects to existing parton-level samples. Similarly, one can train a model using images generated through two different underlying processes, allowing for a mapping e.g. from QCD jets to W or top initiated jets.

³ CycleJet can also be used for similar practical purposes as DCTR [64], albeit it is of course limited to the Lund image representation.



979 Page 8 of 11 Eur. Phys. J. C (2019) 79:979

3.2 CycleJet model results

Following the pipeline presented in Sect. 2.4 we perform 1000 iterations of the hyperparameter scan using the hyperopt library and the loss function

$$\mathcal{L}_h = ||R_A - P_{B \to A}|| + ||R_B - P_{A \to B}|| \tag{11}$$

where A and B indexes refer to the desired input and output samples respectively so R_A and R_B are the average reference images before the CycleGAN transformation while $P_{B\rightarrow A}$ and $P_{A\rightarrow B}$ correspond to the average image after the transformation. Furthermore, for this model we noticed better results when preprocessing the pixel intensities with the standardisation procedure of removing the mean and scaling to unit variance, instead of a simpler rescaling in the [-1,1] range as done in Sect. 2.

The CycleJet model consists in two generators and two discriminators. The generators consist in a down-sampling module with three two-dimensional convolutional layers with 32, 64 and 128 filters respectively, and LeakyReLU activation function and instance normalisation [65], followed by an up-sampling with two two-dimensional convolutional layers with 64 and 32 filter. The last layer is a two-dimensional convolution with one filter and hyperbolic tangent activation function. The discriminators take three two-dimensional convolutional layers with 32, 64 and 128 filters and LeakyReLU activation. The first convolutional layer has additionally an instance normalisation layer and the final layer is a two-dimensional convolutional layer with one filter. The best parameters for the CycleJet model are shown in Table 2.

In the first row of Fig. 12 we show results for an initial average parton-level sample before (left) and after (right) applying the parton-to-detector mapping encoded by the CycleJet model, while in the second row of the same figure we perform the inverse operation by taking as input the average of

Table 2 Final parameters for the CycleJet model

Parameters	Value
D filters	32
G filters	32
λ cycle	10
λ identity factor	0.2
Epochs	3
Batch size	128
ZCA	Yes
n_{avg}	20
Learning rate	$6.7 \cdot 10^{-3}$
Decay β_1	0.7
Optimiser	Adam

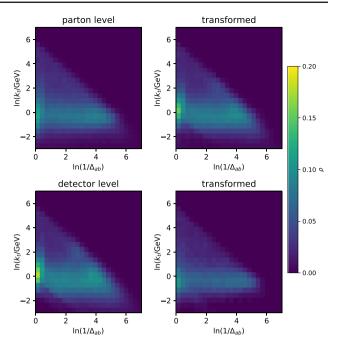


Fig. 12 Top: average of the parton-level sample before (left) and after (right) applying the parton-to-detector mapping. Bottom: average of the delphes-level sample before (left) and after (right) applying the detector-to-parton mapping

the dephes-level sample before (left) and after (right) applying the CycleJet detector-to-parton mapping. This example shows clearly the possibility to add non perturbative and detector effects to a parton level simulation within good accuracy. Similarly to the previous example, in Fig. 13 we present the mapping between QCD-to-W jets and vice-versa. Also in this case, the overall quality of the mapping is reasonable and provides and interesting successful test case for process remapping.

For both examples we observe a good level agreement for the respective mappings, highlighting the possibility to use such an approach to save CPU time for applying full detector simulations and non perturbative effects to parton level events. It is also possible to train the CycleJet model on Monte Carlo data and apply the corresponding mapping to real data.

4 Conclusions

We have conducted a careful study of generative models applied to jet substructure.

First, we trained a LSGAN model to generate new artificial samples of detector level Lund jet images. With this, we observed agreement to within a few percent accuracy in the bulk of the phase space with respect to the reference data. This new approach provides an efficient method for fast simulation of jet radiation patterns without requiring the long



Eur. Phys. J. C (2019) 79:979 Page 9 of 11 979

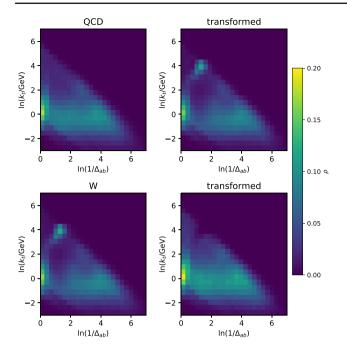


Fig. 13 Top: average of the QCD sample before (left) and after (right) applying the QCD-to-*W* mapping. Bottom: average of the *W* sample before (left) and after (right) applying the *W*-to-QCD mapping

runtime of full Monte Carlo event generators. Another advantage consists in the possibility of this method to be applied to real collider data to generate accurate physical samples, as well as making it possible to avoid the necessity for large storage space by generating realistic samples on-the-fly.

Secondly, a CycleGAN model was constructed to map different jet configurations, allowing for the conversion of existing events. This procedure can be used to change Monte Carlo parameters such as the underlying process or the shower parameters. As examples we show how to convert an existing sample of QCD jets into W jets and vice-versa, or how to add non perturbative and detector effects to a parton level simulation. As for the LSGAN, this method can be used to save CPU time by including full detector simulations and non perturbative effects to parton level events. Additionally, one could use CycleJet to transform real data using mappings trained on Monte Carlo samples or apply them to samples generated through gLund.

To achieve the results presented in this paper we have implemented a rather convolved preprocessing step which notably involved combining and resampling multiple images. This procedure was necessary to achieve accurate distributions but comes with the drawback of loosing information on correlations between emissions at wide angular and transverse momentum separation. Therefore, it is difficult to evaluate or improve the formal logarithmic accuracy of the generated samples. This limitation could be circumvented with an end-to-end GAN architecture more suited to sparse images.

We leave a more detailed study of this for future work. The full code and the pretrained models presented in this paper are available in [48,49].

Acknowledgements We thank Sydney Otten for discussions on β-VAEs. We also acknowledge the NVIDIA Corporation for the donation of a Titan Xp GPU used for this research. F.D. is supported by the Science and Technology Facilities Council (STFC) under Grant ST/P000770/1. S.C. is supported by the European Research Council under the European Union's Horizon 2020 research and innovation Programme (Grant agreement number 740006).

Data Availability Statement This manuscript has associated data in a data repository. [Authors' comment: The samples used for the training of models and figures in this article can be accessed online [50].]

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. Funded by SCOAP³.

Appendix A: VAE and WGAN-GP models

In this appendix we present the final parameters as well as generated event samples for the VAE and WGAN-GP models used in Sect. 2.5. These models are obtained after applying the hyperopt procedure described in Sect. 2.4.

The VAE encoder consists of a dense layer with 384 units with ReLU activation function connected to a latent space with 1000 dimensions. The decoder consists of a dense layer with 384 units with ReLU activation followed by an output layer which matches the shape of the images and has a hyperbolic tangent activation function. The reconstruction loss function used during training is taken to be the mean squared error. The best parameters for the VAE model obtained after

Table 3 Final parameters for the VAE model

Parameters	Value
Intermediate dimension	384
KL annealing rate	0.25
KL annealing factor	1.05
Minibatch discriminator	No
Epochs	50
Batch size	32
Latent dimension	1000
ZCA	Yes
n_{avg}	32
Learning rate	$4.2 \cdot 10^{-4}$
Decay β_1	0.9
Optimiser	Adam



979 Page 10 of 11 Eur. Phys. J. C (2019) 79:979

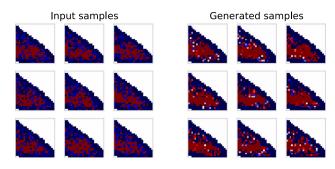


Fig. 14 A random selection of preprocessed input images (left), and of images generated with the VAE model (right). Axis and colour schemes are the same of Fig. 5

Table 4 Final parameters for the WGAN-GP model

Parameters	Value
D units	16
G units	4
α	0.3
Dropout	0.15
D momentum	0.7
G momentum	0.7
Minibatch discriminator	No
Epochs	300
Batch size	32
Latent dimension	800
ZCA	Yes
n_{avg}	32
Learning rate	$9.6 \cdot 10^{-5}$
Decay β_1	$2 \cdot 10^{-8}$
ρ	0.9
Optimiser	RMSprop

the hyperopt procedure are shown in Table 3. In Fig. 14 we show a random selection of preprocessed images generated through the VAE. From a qualitative point of view the images appear realistic on an event-by-event comparison however as highlighted in Sect. 2.5, the VAE model does not reproduce the underlying distribution accurately.

Finally, the WGAN-GP consists in a generator and discriminator. The generator architecture contains a dense layer with 1152 units with ReLU activation function followed by three sequential two-dimensional convolutional layers with a kernel size of 4 and respectively 32, 16 and 1 filters. Between these layers we apply batch normalisation and ReLU activation function while the final layer has a hyperbolic tangent activation function. On the other hand, the discriminator is composed by 4 two-dimensional convolutional layers with a kernel size of 3 and respectively 16, 32, 64, 128 and 128 filters. We apply batch normalisation for the last three layers and all of them LeakyReLU activation function with a

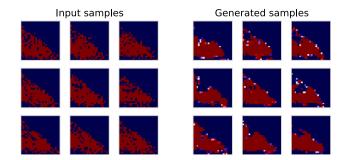


Fig. 15 A random selection of preprocessed input images (left), and of images generated with the WGAN-GP model (right). Axis and colour schemes are the same of Fig. 5

dropout layer. In Table 4 we provide the best parameters of the WGAN-GP model, always obtained through the hyperopt scan procedure. In Fig. 15 we show a random selection of preprocessed images generated through the WGAN-GP. Due to the convolutional filters of this model the preprocessing differs slightly from the description in Sect. 2.3 as we do not remove pixels outside the kinematic range resulting in images with non zero background pixels. While distributions presented in Sect. 2.5 are in good agreement with data, it is clear that for this WGAN-GP model the individual images look different from the input data.

References

- 1. G.F. Sterman, S. Weinberg, Phys. Rev. Lett. 39, 1436 (1977)
- 2. G.P. Salam, Eur. Phys. J. C 67, 637 (2010). arXiv:0906.1833
- S.D. Ellis, D.E. Soper, Phys. Rev. D 48, 3160 (1993). arXiv:hep-ph/9305266
- Y.L. Dokshitzer, G.D. Leder, S. Moretti, B.R. Webber, JHEP 08, 001 (1997). arXiv:hep-ph/9707323
- M. Cacciari, G.P. Salam, G. Soyez, JHEP 04, 063 (2008). arXiv:0802.1189
- 6. J. Thaler, L.T. Wang, JHEP 07, 092 (2008). arXiv:0806.0023
- D.E. Kaplan, K. Rehermann, M.D. Schwartz, B. Tweedie, Phys. Rev. Lett. 101, 142001 (2008). arXiv:0806.0848
- S.D. Ellis, C.K. Vermilion, J.R. Walsh, Phys. Rev. D 80, 051501 (2009). arXiv:0903.5081
- S.D. Ellis, C.K. Vermilion, J.R. Walsh, Phys. Rev. D 81, 094023 (2010). arXiv:0912.0033
- T. Plehn, G.P. Salam, M. Spannowsky, Phys. Rev. Lett. 104, 111801 (2010). arXiv:0910.5472
- 11. J. Thaler, K. Van Tilburg, JHEP 03, 015 (2011). arXiv:1011.2268
- A.J. Larkoski, G.P. Salam, J. Thaler, JHEP 06, 108 (2013). arXiv:1305.0007
- 13. Y.T. Chien, Phys. Rev. D 90, 054008 (2014). arXiv:1304.5240
- M. Cacciari, G.P. Salam, G. Soyez, Eur. Phys. J. C 75, 59 (2015). arXiv:1407.0408
- A.J. Larkoski, I. Moult, D. Neill, JHEP 12, 009 (2014). arXiv:1409.6298
- I. Moult, L. Necib, J. Thaler, JHEP 12, 153 (2016). arXiv:1609.07483
- M. Dasgupta, A. Fregoso, S. Marzani, G.P. Salam, JHEP 09, 029 (2013). arXiv:1307.0007



Eur. Phys. J. C (2019) 79:979 Page 11 of 11 979

 A.J. Larkoski, S. Marzani, G. Soyez, J. Thaler, JHEP 05, 146 (2014), arXiv:1402.2657

- P.T. Komiske, E.M. Metodiev, B. Nachman, M.D. Schwartz, JHEP 12, 051 (2017). arXiv:1707.08600
- P.T. Komiske, E.M. Metodiev, J. Thaler, JHEP 04, 013 (2018). arXiv:1712.07124
- F.A. Dreyer, L. Necib, G. Soyez, J. Thaler, JHEP 06, 093 (2018). arXiv:1804.03657
- F.A. Dreyer, G.P. Salam, G. Soyez, JHEP 12, 064 (2018). arXiv:1807.04758
- 23. A. Butter et al., SciPost Phys. 7, 014 (2019). arXiv:1902.09914
- 24. S. Carrazza, F.A. Dreyer (2019), arXiv:1903.09644
- P. Berta, L. Masetti, D.W. Miller, M. Spousta (2019), arXiv:1905.03470
- E.A. Moreno, O. Cerri, J.M. Duarte, H.B. Newman, T.Q. Nguyen, A. Periwal, M. Pierini, A. Serikova, M. Spiropulu, J.R. Vlimant (2019), arXiv:1908.05318
- N. Fischer, S. Gieseke, S. Plätzer, P. Skands, Eur. Phys. J. C 74, 2831 (2014). arXiv:1402.3186
- 28. Les Houches 2017: Physics at TeV Colliders Standard Model Working Group Report (2018), arXiv:1803.07977
- I. Goodfellow, J. Pouget-Abadie, M. Mirza, B. Xu, D. Warde-Farley, S. Ozair, A. Courville, Y. Bengio, Generative adversarial nets. In: Advances in neural information processing systems (2014), pp. 2672–2680
- D.P. Kingma, M. Welling, Auto-Encoding Variational Bayes, in 2nd International Conference on Learning Representations, ICLR 2014, Banff, AB, Canada, April 14–16, 2014, Conference Track Proceedings (2014). arXiv:1312.6114
- L. de Oliveira, M. Paganini, B. Nachman, Comput. Softw. Big Sci. 1, 4 (2017). arXiv:1701.05927
- M. Paganini, L. de Oliveira, B. Nachman, Phys. Rev. Lett. 120, 042003 (2018). arXiv:1705.02355
- M. Paganini, L. de Oliveira, B. Nachman, Phys. Rev. D 97, 014021 (2018). arXiv:1712.10321
- S. Otten, S. Caron, W. de Swart, M. van Beekveld, L. Hendriks,
 C. van Leeuwen, D. Podareanu, R. Ruiz de Austri, R. Verheyen
 (2019). arXiv:1901.00875
- P. Musella, F. Pandolfi, Comput. Softw. Big Sci. 2, 8 (2018). arXiv:1805.00850
- 36. K. Datta, D. Kar, D. Roy (2018). arXiv:1806.00433
- O. Cerri, T.Q. Nguyen, M. Pierini, M. Spiropulu, J.R. Vlimant, JHEP 05, 036 (2019). arXiv:1811.10276
- R. Di Sipio, M. Faucci Giannelli, S. Ketabchi Haghighat, S. Palazzo (2019). arXiv:1903.02433
- 39. A. Butter, T. Plehn, R. Winterhalder (2019). arXiv:1907.03764
- 40. (2019), arXiv:1909.04451
- 41. T.A. collaboration (ATLAS) (2019)
- 42. X. Mao, Q. Li, H. Xie, R.Y.K. Lau, Z. Wang, CoRR arXiv:1611.04076 (2016)
- T. Sjöstrand, S. Ask, J.R. Christiansen, R. Corke, N. Desai, P. Ilten, S. Mrenna, S. Prestel, C.O. Rasmussen, P.Z. Skands, Comput. Phys. Commun. 191, 159 (2015). arXiv:1410.3012
- J. de Favereau, C. Delaere, P. Demin, A. Giammanco, V. Lemaître,
 A. Mertens, M. Selvaggi, (DELPHES 3), JHEP 02, 057 (2014).
 arXiv:1307.6346

- M. Arjovsky, S. Chintala, L. Bottou, Wasserstein Generative Adversarial Networks, in Proceedings of the 34th International Conference on Machine Learning, ICML 2017, Sydney, NSW, Australia, 6-11 August 2017 (2017), pp. 214–223, http://proceedings.mlr.press/y70/arjovsky17a.html
- J.Y. Zhu, T. Park, P. Isola, A.A. Efros, Unpaired Image-to-Image Translation using Cycle-Consistent Adversarial Networkss, in Computer Vision (ICCV), 2017 IEEE International Conference on (2017)
- 47. J. Duarte et al., JINST 13, P07027 (2018). arXiv:1804.06913
- 48. F. Dreyer, S. Carrazza, *Jetsgame/glund v1.0.0* (2019), https://doi.org/10.5281/zenodo.3384920
- F. Dreyer, S. Carrazza, Jetsgame/cyclejet v1.0.0 (2019). https://doi. org/10.5281/zenodo.3384918
- S. Carrazza, F.A. Dreyer, *JetsGame/data v1.0.0* (2019), this repository is git-lfs. https://doi.org/10.5281/zenodo.2602514
- M. Wobisch, T. Wengler, Hadronization corrections to jet crosssections in deep inelastic scattering, in Monte Carlo generators for HERA physics. Proceedings, Workshop, Hamburg, Germany, 1998–1999, pp. 270–279 (1998). arXiv:hep-ph/9907280
- M. Cacciari, G.P. Salam, G. Soyez, Eur. Phys. J. C 72, 1896 (2012). arXiv:1111.6097
- I.J. Goodfellow, J. Pouget-Abadie, M. Mirza, B. Xu, D. Warde-Farley, S. Ozair, A. Courville, Y. Bengio, in *Proceedings of the 27th International Conference on Neural Information Processing Systems Volume 2* (MIT Press, Cambridge, MA, USA, 2014), NIPS'14, pp. 2672–2680. http://dl.acm.org/citation.cfm?id=2969033.2969125
- T. Salimans, I.J. Goodfellow, W. Zaremba, V. Cheung, A. Radford, X. Chen, CoRR. arXiv:1606.03498 (2016)
- 55. A.J. Bell, T.J. Sejnowski, Vision Research 37, 3327 (1997)
- J. Bergstra, D. Yamins, D.D. Cox, Making a Science of Model Search: Hyperparameter Optimization in Hundreds of Dimensions for Vision Architectures, in Proceedings of the 30th International Conference on International Conference on Machine Learning -Volume 28 (JMLR.org, 2013), ICML'13, pp. I–115–I–123
- 57. Zhou Wang, A.C. Bovik, H.R. Sheikh, E.P. Simoncelli, IEEE Transactions on Image Processing 13, 600 (2004)
- A.L. Maas, A.Y. Hannun, A.Y. Ng, Rectifier nonlinearities improve neural network acoustic models, in in ICML Workshop on Deep Learning for Audio, Speech and Language Processing (2013)
- I. Higgins, L. Matthey, X. Glorot, A. Pal, B. Uria, C. Blundell, S. Mohamed, A. Lerchner, CoRR arXiv:1606.05579 (2016)
- C.P. Burgess, I. Higgins, A. Pal, L. Matthey, N. Watters, G. Desjardins, A. Lerchner, CoRR arXiv:1804.03599 (2018)
- I. Gulrajani, F. Ahmed, M. Arjovsky, V. Dumoulin, A.C. Courville, CoRR arXiv:1704.00028 (2017)
- S.R. Bowman, L. Vilnis, O. Vinyals, A.M. Dai, R. Józefowicz, S. Bengio, CoRR arXiv:1511.06349 (2015)
- C. Frye, A.J. Larkoski, J. Thaler, K. Zhou, JHEP 09, 083 (2017).
 arXiv:1704.06266
- 64. A. Andreassen, B. Nachman (2019), arXiv:1907.08209
- D. Ulyanov, A. Vedaldi, V.S. Lempitsky, CoRR arXiv:1607.08022 (2016)

