



Testbeam studies of pre-prototype silicon strip sensors for the LHCb UT upgrade project



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ARTICLE INFO

Article history:

Received 7 July 2015

Received in revised form

9 October 2015

Accepted 10 October 2015

Available online 17 October 2015

Keywords:

Silicon strip detector

Testbeam

Sensor irradiation

ABSTRACT

The LHCb experiment is preparing for a major upgrade in 2018–2019. One of the key components in the upgrade is a new silicon tracker situated upstream of the analysis magnet of the experiment. The Upstream Tracker (UT) will consist of four planes of silicon strip detectors, with each plane covering an area of about 2 m². An important consideration of these detectors is their performance after they have been exposed to a large radiation dose. In this paper we present test beam results of pre-prototype n-in-p and p-in-n sensors that have been irradiated with fluences up to 4.0×10^{14} n_{eq}/cm².

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1. Introduction

The Upstream Tracker (UT) detector (see Fig. 1) is a key part of the LHCb Upgrade, replacing the current TT tracking stations [1,2]. The UT improves over the current TT in that it (i) eliminates all gaps within the detector acceptance, (ii) has largely improved granularity to cope with the higher instantaneous luminosity expected in the LHCb upgrade, and (iii) improves the coverage close to the beam pipe by employing a circular cutout to match the beam pipe profile. The UT, like the TT, is situated just in front of LHCb's dipole analysis magnet. In this position, it provides a crucial link between segments reconstructed in the upgraded vertex detector [3] and the tracking chambers downstream [4] of the LHCb magnet. It provides a factor of 3 improvement in speed for the tracking in the fully-software-based trigger [5], reduces the rate of fake tracks being formed by a factor of 2–3, and improves the momentum resolution by about 25% relative to tracks not using UT hits. The factor of 3 in speed is enabled by enabling a very fast estimate of the momentum of charged particles, which can then be used to reduce the size of the hit search windows in the

downstream tracking stations. Due to the increased speed of the trigger and the higher purity of tracks considered, larger data sets with better signal-to-background can be acquired.

The UT detector consists of four silicon planes, each about 1.53 m in width and 1.34 m in height. Each plane is composed of 1.5 m long *staves* that are tiled with ~ 10 cm \times 10 cm silicon wafers. Consecutive wafers are mounted on opposite sides of the stave to ensure no gaps along the height, and adjacent staves are also overlapped to ensure no gaps in the horizontal direction. The majority of the detector area utilizes sensors with an approximate pitch of 190 μ m, however the inner region features sensors with half the pitch (95 μ m) to cope with higher occupancy. Both n-in-p and p-in-n type sensors are being considered for the outer region, but for the inner region, only n-in-p are being considered due to better radiation hardness. For the innermost region of the UT, the largest fluence expected, with a safety factor of about two, is about 5×10^{14} n_{eq}/cm². Outside this region, the fluence is not expected to exceed about 2×10^{13} n_{eq}/cm².

The primary goals of this test beam were to quantify the performance of several pre-prototype n-in-p mini-sensors from Micron Semiconductor, Ltd [6], after a high radiation fluence, and compare to corresponding results from similar unirradiated

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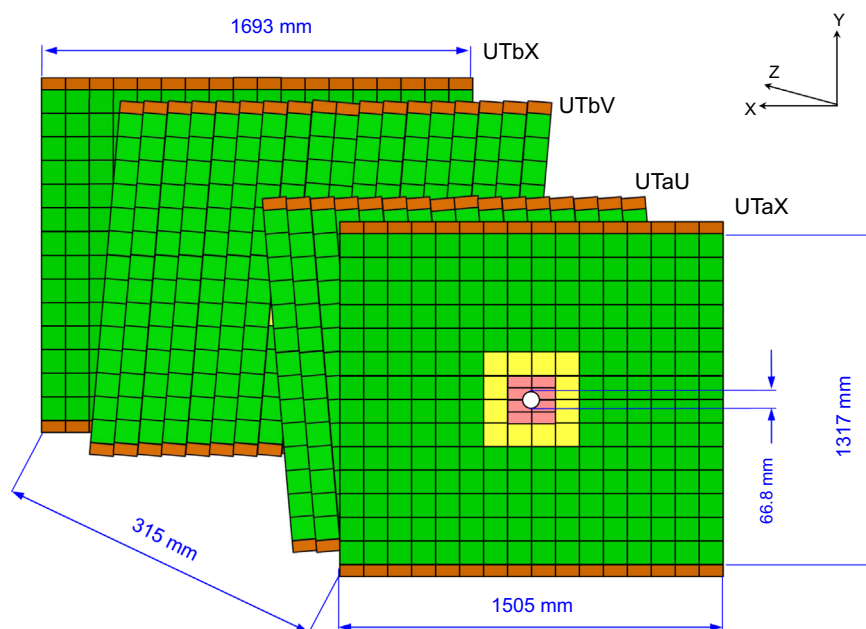


Fig. 1. Cartoon showing the layout and dimensions of the four UT planes.

detectors. The properties investigated include, but are not limited to:

- Landau distributions as a function of bias voltage;
- cluster size vs. bias voltage and angle of incidence;
- charge collected vs. interstrip position;
- resolution vs. angle; and
- characteristics of sensors near the quarter-circle region (these emulate the sensors surrounding the beam pipe).

One p-in-n mini-sensor irradiated to the maximum level expected in the outer region of the UT was also tested, but its study was not a primary focus of the test beam results presented here.

2. Experimental setup

The test beam discussed in this paper was conducted in October 2014 at the SPS at CERN. The beam consisted of positively charged hadrons with momentum of 180 GeV/c. The beam was delivered in spills at rate of about 4 spills/min, with each spill lasting about 4 s. For most of the data taking, the beam size was collimated down to about 0.5 cm in diameter and each spill provided a particle rate of order 1 MHz.

2.1. Telescope description

The pre-prototype UT sensors, or detectors under test (DUT), were studied using the TimePix3 (TP3)-based telescope [7], composed of 8 pixel planes. Each pixel plane is about $1.4 \text{ cm} \times 1.4 \text{ cm}$ and has a pixel size of $55 \mu\text{m} \times 55 \mu\text{m}$. The planes are tilted in order to provide more charge sharing, and thus better position resolution. A cartoon of the telescope layout is shown in Fig. 2. With the high momentum beam of 180 GeV/c, the reconstructed tracks provide excellent pointing resolution of about $2 \mu\text{m}$ at the DUT. The telescope readout does not require an external trigger: hits are recorded continuously once a run is started. For each pixel hit, both position and a time-stamp with 1.56 ns precision is recorded. Tracks are then formed (offline) by combining hits that have compatible time values.

2.2. Detectors under test

The mini-sensor pre-prototypes tested in the October 2014 test beam were obtained from Micron Semiconductor, Ltd [6]. One of the sensors tested was a p-in-n, and six were n-in-p. The resistivities of the sensors, as determined from capacitance vs. voltage measurements, were about $0.90 \text{ k}\Omega \text{ cm}$ for the p-in-n sensor and about $2.8 \text{ k}\Omega \text{ cm}$ for the n-in-p. Each sensor was $1.115 \text{ cm} \times 1.125 \text{ cm}$ in size with a nominal thickness of $250 \mu\text{m}$, and had 128 strips with a strip pitch of $80 \mu\text{m}$ and a strip (implant) width of $30 \mu\text{m}$. Prior to irradiation, all of the sensors had a depletion voltage of about 180 V. A schematic of the mini-sensors is shown in Fig. 3. In this schematic, the strips run horizontally.

Of the six n-in-p sensors tested, three had the strips terminated such that they form a quarter-circle inactive region of the sensor. Two different guard-ring structures were implemented, as shown in Fig. 3. More details of the sensors under test are shown in Table 1. The MBP1 and MBP2 sensors are differentiated by the guard ring structures, and are shown on the left and right side of Fig. 3, respectively. The MBP1 (leftmost sensor in figure) differs from the MBP2 sensor (rightmost) in that it implements a stepped structure along the innermost guard ring to maintain an equal separation between the edge of the strip and the innermost guard ring, more like a conventional rectangle-shaped detector.

Six of the seven sensors were irradiated at the Massachusetts General Hospital (MGH) proton irradiation facility [8] in June 2014, using protons of kinetic energy equal to 226 meV and fluences ranging from $0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ to $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$. Between the time of the irradiation in June 2014, and the test beam in October 2014, the sensors were kept in a freezer, at a temperature below -10°C . The sensors were warmed up to room temperature for no more than 7 days to transport the sensors and for wirebonding.

The DUT readout for this testbeam was based on the Alibava DAQ system [9–12], which uses Beetle chips [13,14] as the front end ASIC. The main components of the Alibava system are a detector board to which sensors are mounted, a daughterboard that includes two Beetle chips (128 channels each), and a motherboard that manages the data flow to/from the Beetle chips and to/from the data acquisition PC.

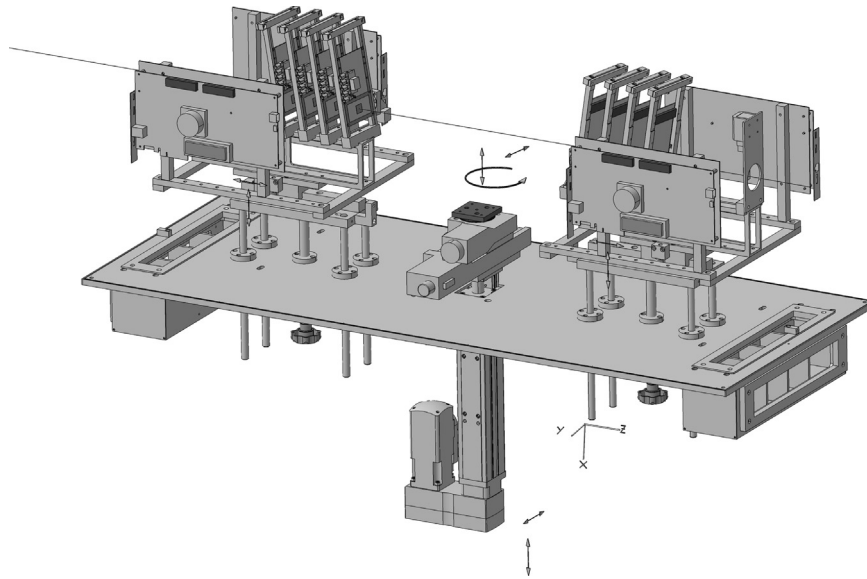


Fig. 2. Cartoon showing the TimePix3 telescope [7] used to study the UT sensors.

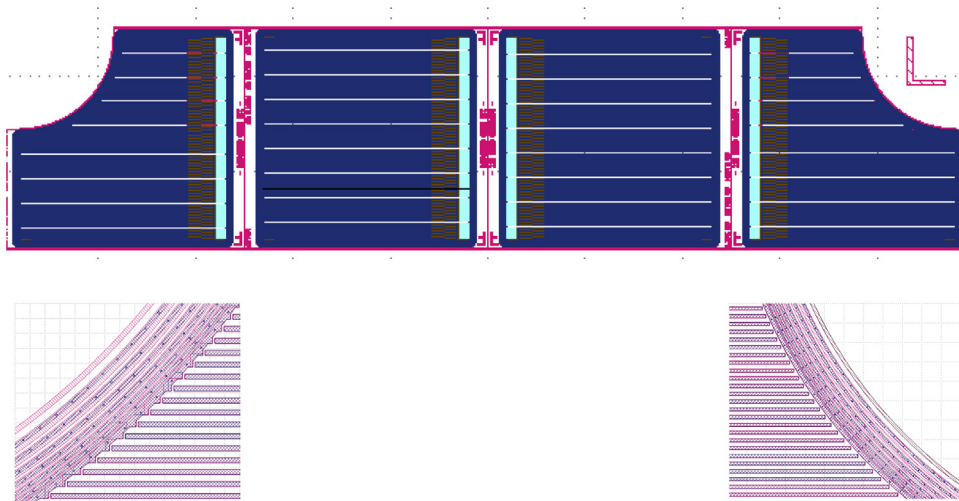


Fig. 3. Schematic of the mini-sensors tested during the October 2014 test beam. The two sensors in the middle are full mini-sensors, while the ones on the left and right have a quarter-circle cut out. The difference between the left-most (MBP1) and right-most sensors (MBP2) is in the guard-ring structure, as shown just below them.

The DUT was housed in an aluminum box to provide both shielding and a light-tight environment. A Peltier device was used to cool the sensor to about $-13\text{ }^{\circ}\text{C}$, and a continuous nitrogen flow maintained the relative humidity close to zero. The entire box was mounted on a stage that allowed for horizontal and vertical translations, as well as rotations about the vertical axis to allow for studies of the detector performance vs. particle incident angle. The system was equipped with both temperature and relative humidity monitoring. A photograph of the telescope with the UT DUT being installed is shown in Fig. 4. Various components are indicated.

2.3. Trigger and synchronization

The trigger was formed from the coincidence of two scintillators, one just upstream and a second just downstream of the TimePix3 telescope (see Fig. 2) and a *not-busy* signal from the Alibava motherboard. This trigger signal was fanned out to both the TimePix3 telescope and the Alibava motherboard. Because the data acquisition systems of the TimePix3 telescope and the DUT were independent of one another, the common trigger signal was used to synchronize the two systems. The signal sent to the

Alibava system initiated a readout of the DUT, while the signal sent to the TimePix3 system produced a timestamp in the TimePix3 event record. This timestamp used the same clock used to time stamp the pixel hits, and in this way, pixel hits could be associated with a specific trigger. Since the number of time stamps in the TimePix3 event record was identical to the number of triggers sent to the Alibava system, the synchronization only required matching the first event in the Alibava system with the pixel hits/tracks associated to the first trigger time stamp, and then the second event matched to the second trigger time stamp, and so on. In this way, the beam tracks reconstructed in the TimePix3 system were properly matched to the corresponding hits produced in the DUT.

3. Corrections and calibrations

Before being able to quantify the performance of the sensors, various calibrations and corrections were applied. These corrections included pedestal and common-mode noise suppression, a cross-talk correction, an out-of-time correction, and are detailed below.

Table 1

Specific details of the sensors tested in the October 2014 test beam. Here, ρ refers to the resistivity of the silicon and $N_{e,h}$ indicates either the electron or hole concentration. For sensor 3091-10-MBP2, CV measurements were not available.

Sensor ID	Type	Fluence (10^{14} n_{eq}/cm^2)	$N_{e,h}$ (10^{12} cm^{-3})	ρ ($k\Omega$ cm)	Other spec.
3092-1-MS2	p-in-n	0.27	5.1	0.87	250 μ m
3091-8-MS1	n-in-p	1.0	5.0	2.79	242 μ m, p-spray
3091-10-MBP2	n-in-p	0.0	n/a	n/a	238 μ m, p-spray
3091-7-MS2	n-in-p	4.0	5.0	2.76	254 μ m, p-spray
3091-7-MS1	n-in-p	1.8	4.9	2.85	254 μ m, p-spray
3091-8-MBP2	n-in-p	4.0	5.2	2.69	242 μ m, p-spray
3091-8-MBP1	n-in-p	4.0	5.0	2.79	242 μ m, p-spray

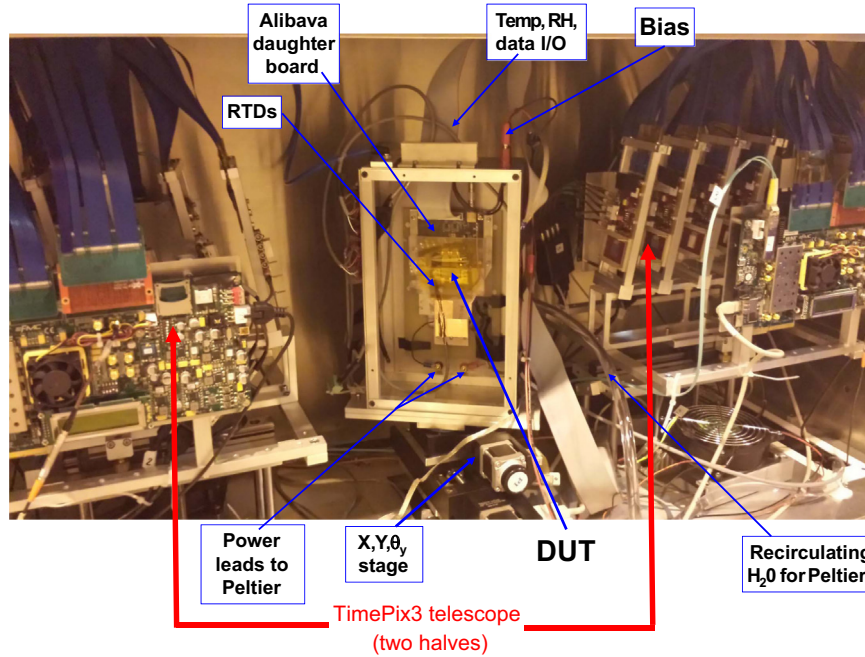


Fig. 4. Photograph of the TimePix3 telescope and the DUT during installation at the SPS, together with a description of the different components.

3.1. Calibration scans

Calibration scans, which were taken periodically during the test beam, showed that 1 ADC was equivalent to about $275e^-$. The exact value varied by a few percent, depending on the sensor/readout board. Since the calibration circuit has its own inherent parasitic capacitances, we estimate that there is an inherent uncertainty in the overall calibration not larger than about 5%. Charges presented throughout this paper are given in terms of ADC counts, and must be multiplied by 275 to obtain the corresponding charge in number of electrons.

3.2. Pedestal and common mode noise suppression

Pedestals are calculated from datasets of 10^4 – 10^5 events recorded during beam stops and are then subtracted from the beam-on runs recorded under the same conditions. (The events are self-triggered through the Alibava motherboard.) Dead or noisy strips are excluded from all parts of the analysis. The pedestal of a given Beetle channel is evaluated as the mean of the raw ADC counts of that channel over the whole pedestal run, after excluding anomalously large or small ADC values. After subtracting the pedestals, the common mode noise is computed (event-by-event) as the average ADC value in a given chip, and is then subtracted from each channel. After the pedestal and common mode

noise suppression, the noise level is about 3.5 ADC counts, or about $1000e^-$.

3.3. Cross-talk correction

From analysis of the data, it was found that there was cross-talk between the N th and the $(N+1)$ th channel. The effect was easily seen by studying the charge asymmetry between the $(N-1)$ th and $(N+1)$ th channel about the peak strip (N) in a cluster with tracks at normal incidence. Due to capacitive coupling some charge sharing is expected, but it should not be asymmetric. The asymmetry is studied for even and odd-numbered peak strips separately. The resulting asymmetries are shown in Fig. 5 (top) before the correction and (bottom) after the linear correction for one of the DUTs. The even channels show a larger cross-talk than the odd channels. An empirical correction based on the linear fit shown is applied, and the asymmetry after the correction is shown. This is believed to be due to a sub-optimal timing of the ADC sampling phase, which was a fixed setting in the firmware of the motherboard. The Beetle chip itself is known to have a small amount of cross-talk [14], but the level is below 2%; thus most of the effect is attributable to the readout. All DUTs have similar odd–even cross-talk corrections. The even/odd cross-talk correction was similar for all of the DUTs, but not identical. Therefore, each sensor had a unique cross-talk correction to remove this bias.

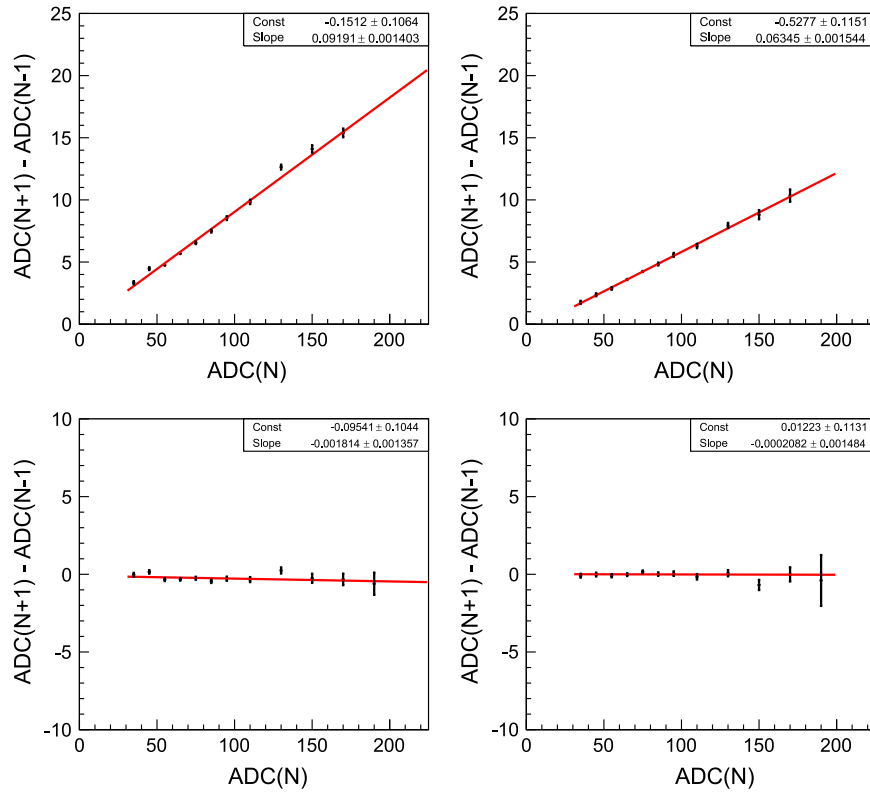


Fig. 5. Mean of the difference in charge detected, $\langle \text{ADC}(N+1) - \text{ADC}(N-1) \rangle$ between the $(N+1)$ th and $(N-1)$ th strips, as a function of the charge seen in the N th strip, for sensor 3092-1-MS2. The top pair of plots show the values for even and odd channels before the correction, the bottom pair after the correction. Plots with an even peak strip numbers are shown on the left, while plots with an odd peak strip number are shown on the right. The red lines are linear fits to the data. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

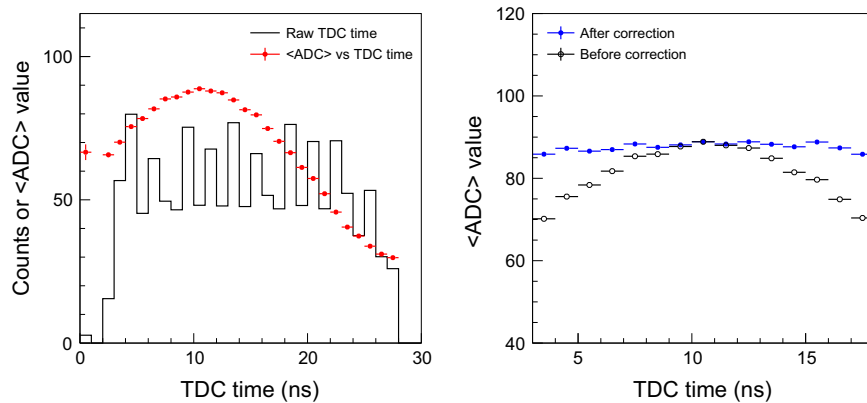


Fig. 6. (Left) Raw TDC time spectrum, overlaid with the average ADC vs. TDC time for DUT signal clusters. (Right) Average ADC vs. TDC time for DUT signal clusters before (black points) and after (blue points) the time correction. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

3.4. Out-of-time correction

The Alibava system provides a 40 MHz clock to sample the signals from DUT in the front end of the Beetle chip. However, beam particles arrive asynchronously with respect to this clock. The Alibava system stores the TDC time between the signals from the scintillators and the edge of the 40 MHz clock. In the offline analysis, the sampled pulse heights are sorted with the TDC values. When timing in the system, we adjusted the latency such that the signal from a beam particle arrives at about 10 ns. That is, the maximum charge is collected from the Beetle when the TDC time is 10 ns. To make use of the data not precisely at the peak, we include and correct the measured charge for signals within \pm

3.5 ns of the peak. The correction is obtained by fitting the average ADC value for signal clusters as a function of the TDC time to a Gaussian function. In this limited range, the correction is no more than about 5–6%. Fig. 6 (left) shows the raw time spectrum of all triggered events, overlaid with the average ADC for signal clusters as a function of the TDC time. The right panel shows the average ADC vs. time before and after the correction, where only the region used in the fit is shown.

3.5. Alignment

The coordinate system is defined such that X is horizontal, Y is vertical, and Z is parallel to the beam axis. Sensors are aligned with

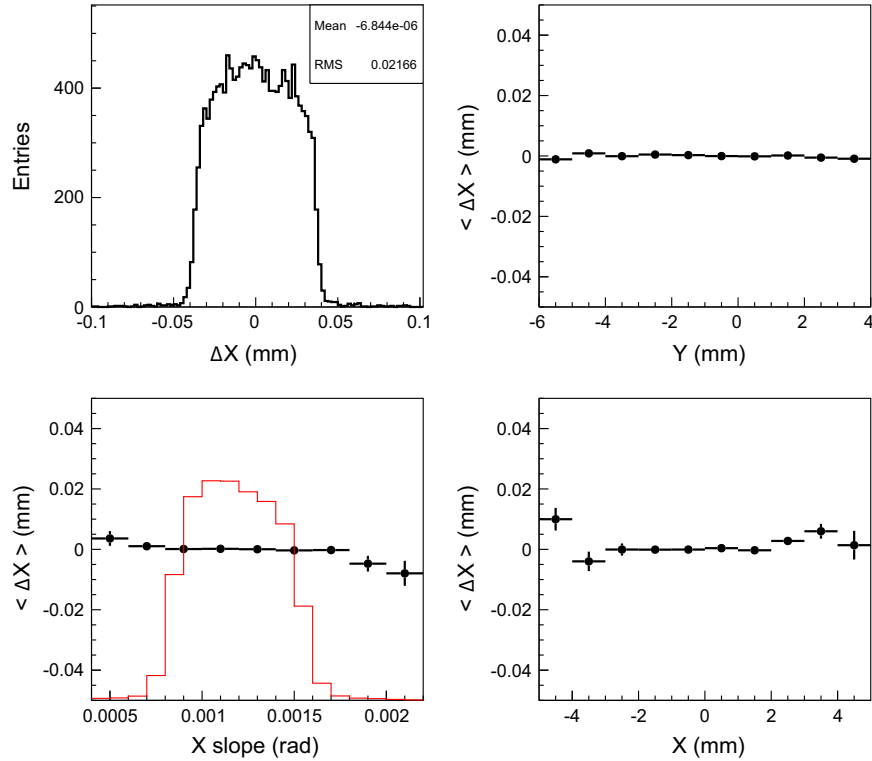


Fig. 7. Alignment plots for sensor 3092-1-MS2, after alignment, showing (top left) X residual, (top right) average X residual vs. Y position, (bottom left) average X residual vs. X slope, and (bottom right) average X residual vs. X position, after the alignment is performed. The X and Y positions and X slope are obtained from the matching TimePix3 track. Superimposed on the bottom left plot (red line) is the slope distribution of the tracks. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

respect to tracks using the *residual* distribution, which is defined as the difference between the reconstructed X position of the hit and the X position of the track extrapolated to the position of the sensor. The degrees of freedom considered are: (i) offsets in the X direction, (ii) offsets along the Z-axis, (iii) rotations around the Z-axis, and (iv) rotations around the Y-axis. A typical set of alignment plots are shown in Fig. 7, after the alignment is done. The profile plots are close to flat at zero. Small deviations are seen in the edges of the bottom pair of distributions, but is likely an artifact of limited numbers of tracks and a possible bias due to the steep droppoff of the beam profile near the edges.

4. Results

The results of the test beam are presented below, starting with the results obtained from particles striking the detector at normal incidence, and continuing with the dependence on the incidence angle.

4.1. Cluster finding in the DUT

Clusters in the DUT are built up by searching for a *seed strip* that has a collected charge more than 20 ADC counts. Moving away from the seed strip, we add adjacent *side strips* having at least 11 ADC counts (about $3 \times \sigma_{\text{noise}}$). The cluster is terminated when a side strip charge is below 11 ADC threshold. Thus, by definition, a cluster has between 1 and 5 strips included; each strip has charge of at least 11 ADC counts, and there is at least one strip with at least 20 ADC counts. Therefore the minimum charge of any cluster is 20 ADC counts. When there are multiple strips in the cluster, the position is computed using linear charge weighting, namely $\bar{x} = \sum x_i q_i / \sum q_i$, where x_i and q_i are the positions and charges of

the strips in the cluster. Better resolution can be achieved with a non-linear correction, but optimizing the resolution on these pre-prototype sensors is not a central goal of this test beam study.

4.2. Tracking information

Tracks were reconstructed using the TimePix3 telescope. About 95% of tracks had 8-pixel hits; the remainder was 7-hit tracks. Only tracks with good fit quality were used, by requiring that the $\chi^2/\text{ndf} < 4$, where ndf is the number of degrees of freedom. Fig. 8 shows the beam profile (Y vs. X) for tracks that have a DUT hit within 100 μm from the track. The quarter-circle regions in board 2 chip 1, and board 4 chips 0 and 1 are evident. The sharp vertical edges are due to the collimators in the beamline, and the boundaries along Y are fiducial cuts used in the analysis. The beam intensity was tuned so that most of the triggered events had a single beam track, although about 10% of events had more than one track. Multi-track events are excluded from the analysis, to ensure that the reconstructed track is the one which produced the trigger.

4.3. Single event displays

A few typical events, after all corrections are applied, are shown in Fig. 9. Here, no requirement is made on the number of tracks. Strips with large ADC counts are indicative of the passage of a beam particle through the detector. The other channels show roughly Gaussian fluctuations about zero, typical of incoherent detector noise. Signals generally stand out significantly above the noise.

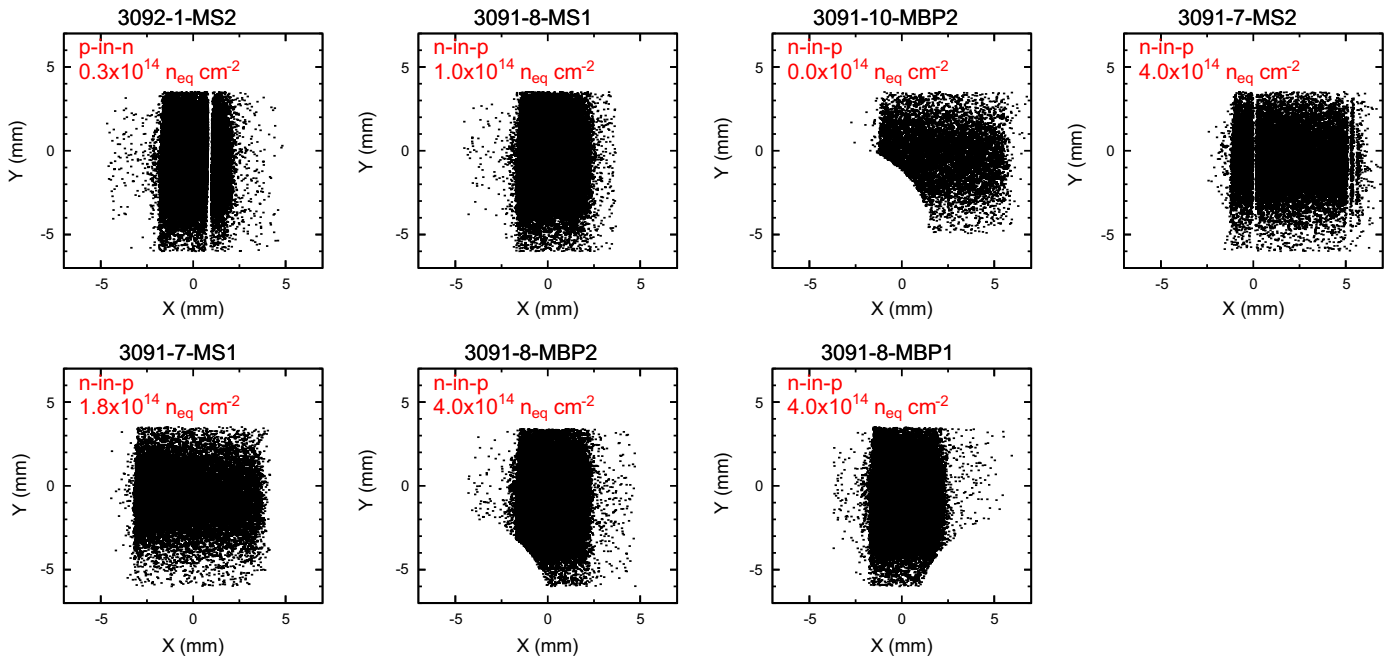


Fig. 8. Profile of the beam (Y vs. X) at the DUT for each of the 7 sensors for a typical set of runs. The 3 DUTs with the quarter-circle edge are evident in these distributions.

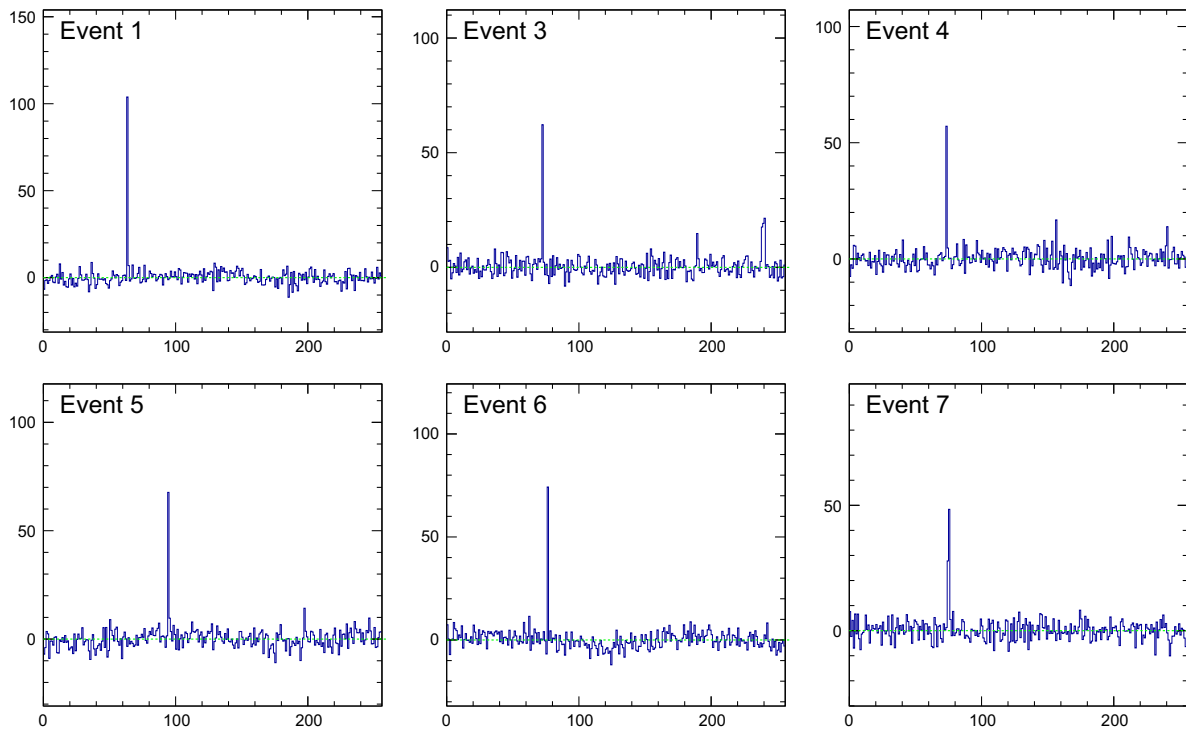


Fig. 9. A few examples of single events for sensor 3092-1-MS2 (p-in-n), with a bias voltage of 250 V.

4.4. Noise

To investigate the noise more quantitatively, the distribution of ADC counts on all strips for a typical run of each sensor is shown in Fig. 10. The runs are chosen to be in a range where each detector is fully depleted. While these are beam-on data, the signal contribution to the shown distributions is negligible, since clusters typically have charge ~ 70 – 80 ADC counts on average, and are therefore off the plot. In addition, most events have only one cluster among the 255 channels. For each distribution, the core of

the distribution is fit to a Gaussian function and show the width, which is found to be in the range of about 3.3–3.5 ADC counts. Since the charge calibration is about $275e^-/\text{ADC}$, this corresponds to a noise level of about $1000e^-$. Pedestal data with the beam off give compatible results.

4.5. Landau distributions vs. bias voltage

The most important aspect of these test beam studies is to validate that the sensors still have a high signal-to-noise ratio

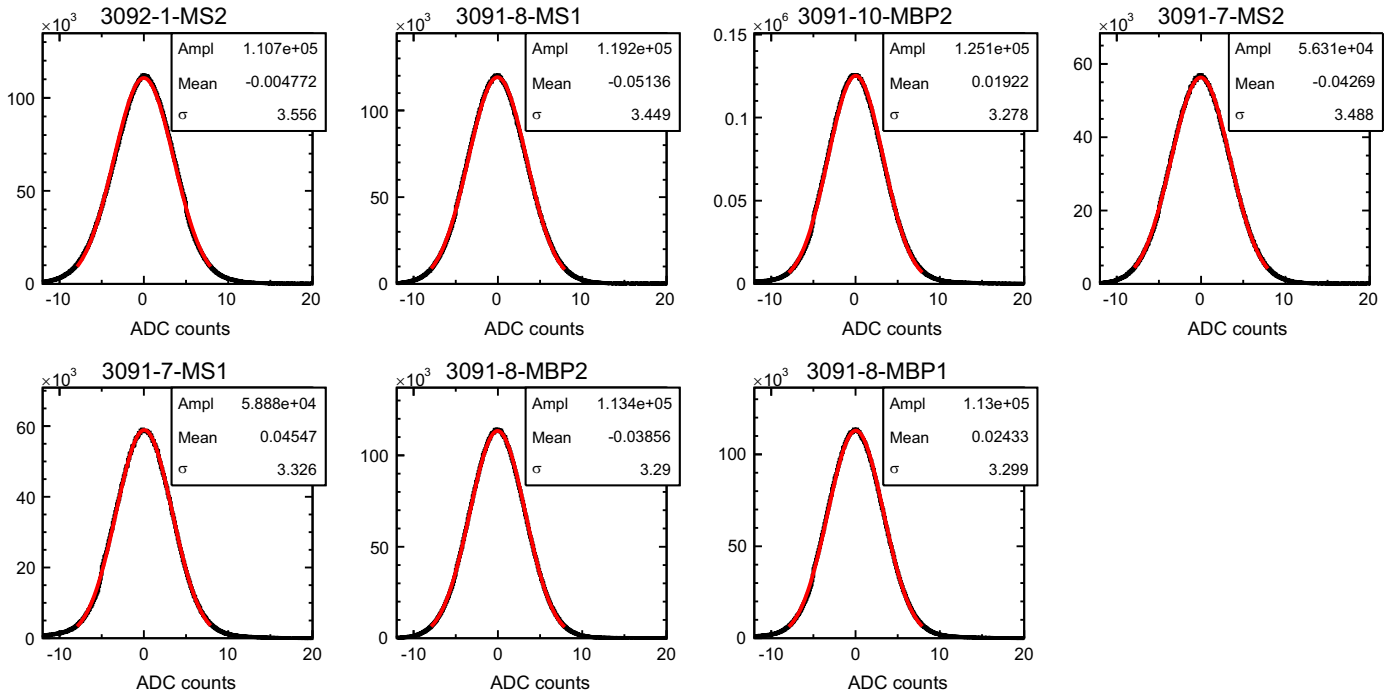


Fig. 10. Noise distributions (in ADC counts) for the DUTs after all corrections during beam-on runs. The red curves show a Gaussian fit to the core of the distribution. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

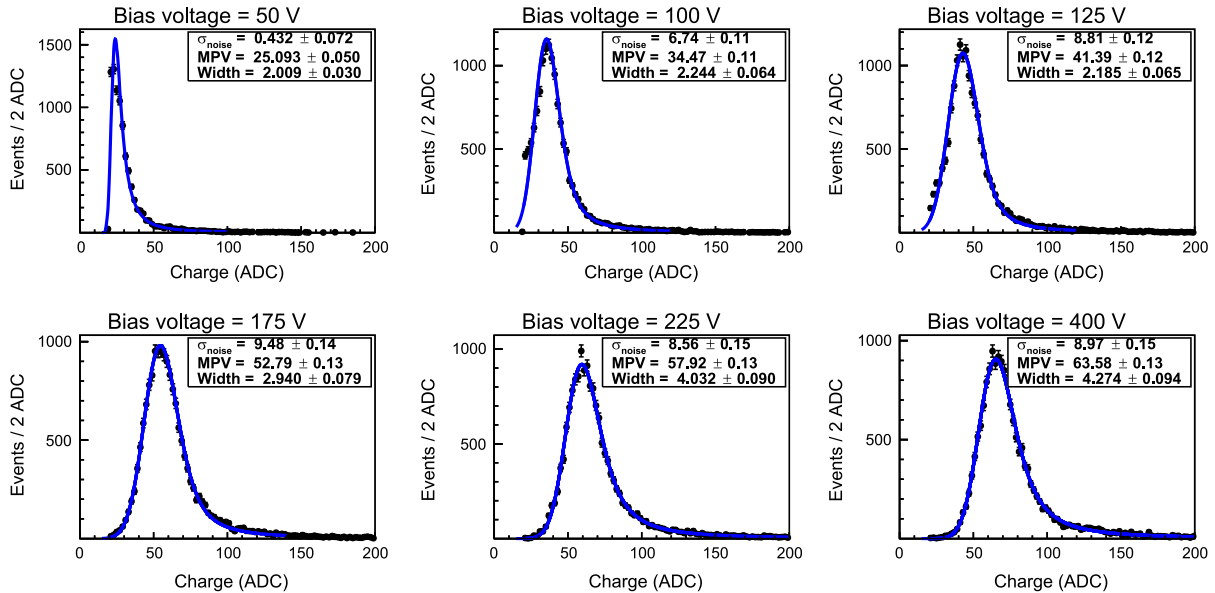


Fig. 11. Signal distributions (in ADC counts) for sensor 3091-8-MBP2, which is an n-in-p type sensor, irradiated to $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$, for several bias voltages. Each distribution is fitted to a Landau convoluted with a Gaussian resolution function and is overlaid.

(S/N) at the highest expected radiation fluence. For each sensor, data was recorded at bias voltages ranging from 50 V (well below full depletion) to 500 V (well above it). A few of the resulting distributions of cluster charge for sensor 3091-8-MBP2 (n-in-p, $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$) are shown in Fig. 11. The cluster charge distributions are each fit to a Landau convoluted with a Gaussian function, and the fit parameters are indicated. For bias voltages smaller than 75 V, the fits are biased because of the 20 ADC threshold requirement on the seed strip in the cluster. As the bias voltage increases, so does the collected charge.

The results for all DUTs tested are summarized in Fig. 12, which shows the most probable value (MPV) of the Landau fit as a

function of the bias voltage, for each of the sensors. The p-in-n sensor is shown with a dashed curve and the n-in-p sensors are represented by solid curves. All of the curves exhibit a plateau in the 300–400 V range, although there is a clear trend that the sensors with the higher accumulated fluence exhibit a lower MPV in the plateau region. The loss is about 10% at a fluence of $4 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$, compared to no irradiation. Similar loss of collected charge in highly irradiated detectors has been seen previously (see for example Refs. [15–19]). The three highest fluence ($4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$) sensors show a similar behavior. A small difference in the MPV in the plateau is seen for one of the sensors, and this is most likely indicative of small differences in its charge

calibration relative to the other two. Since the noise is ~ 3.5 ADC per channel, the S/N is still quite high, about 18 for $V > 400$ V, even in the most irradiated detectors. The S/N depends on the detector capacitance, and hence we would expect smaller S/N in full size sensors.

4.6. Cluster size

Fig. 13 shows the distribution of cluster sizes for the DUTs after all corrections are applied, all at normal incidence, and at bias voltages in the plateau region of Fig. 12. In the p-in-n sensor, about 85% of the clusters are single hit clusters. On the zero fluence n-in-

p sensor, the single hit cluster fraction is a bit lower, about 82%. The irradiated n-in-p sensors show a substantially larger fraction of 2-strip clusters, of about 40%. While there is a loss of total charge collected in the highly irradiated detectors, a larger amount of charge sharing is observed.

4.7. Cluster charge vs. cluster size and interstrip position

In this section, we investigate the dependence of the collected charge on the cluster size. With zero thresholds, and no radiation effects, one would expect that the collected charge to be independent of the cluster size. However, radiation effects and non-zero thresholds can lead to a difference. Fig. 14 shows the Landau distributions in the plateau region of the sensors, and overlaid are the separate contributions from 1 and 2-strip clusters. The 1-strip and 2-strip clusters have similar MPVs, although not identical. Due to the threshold bias mentioned above, the $0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ p-in-n and the zero fluence n-in-p sensors have larger cluster charge for 2-strip clusters compared to the 1-strip clusters. However, for the irradiated sensors, the opposite is seen, that is, 2-strip clusters have a lower charge collected than 1-strip clusters. From these observations, we conclude that the n-in-p sensors that received the radiation fluence (i) have reduced total charge, (ii) have a larger fraction of two-strip clusters, and (iii) have less charge collected in two-strip clusters compared to 1-strip clusters.

Additional information can be gained by investigating the charge collection as a function of the relative position between two strips. For each track, the interstrip position is defined to be the position at the DUT mapped onto the range from [0, 1], where 0 is the center of the N th strip and 1 is the center of the $(N+1)$ th strip. The distributions of the average cluster charge as a function of the interstrip position are shown in Fig. 15 for each of the DUTs. It is seen that the heavily irradiated n-in-p sensors exhibit a steady decrease in charge collection as the track approaches the middle region between two strips. At the center, a loss of approximately 15% in collected charge is observed. The p-in-n and unirradiated n-in-p show a small loss in the middle between two strips, but the effect is not more than a few percent. Similar loss of collected

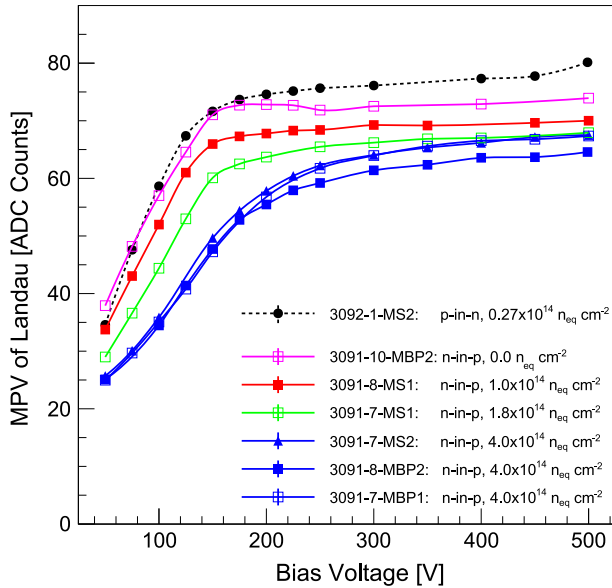


Fig. 12. The most probable value of the collected charge for all cluster sizes from the Landau fit as a function of bias voltage for each of the DUT. Tracks are at normal incidence to the sensor.

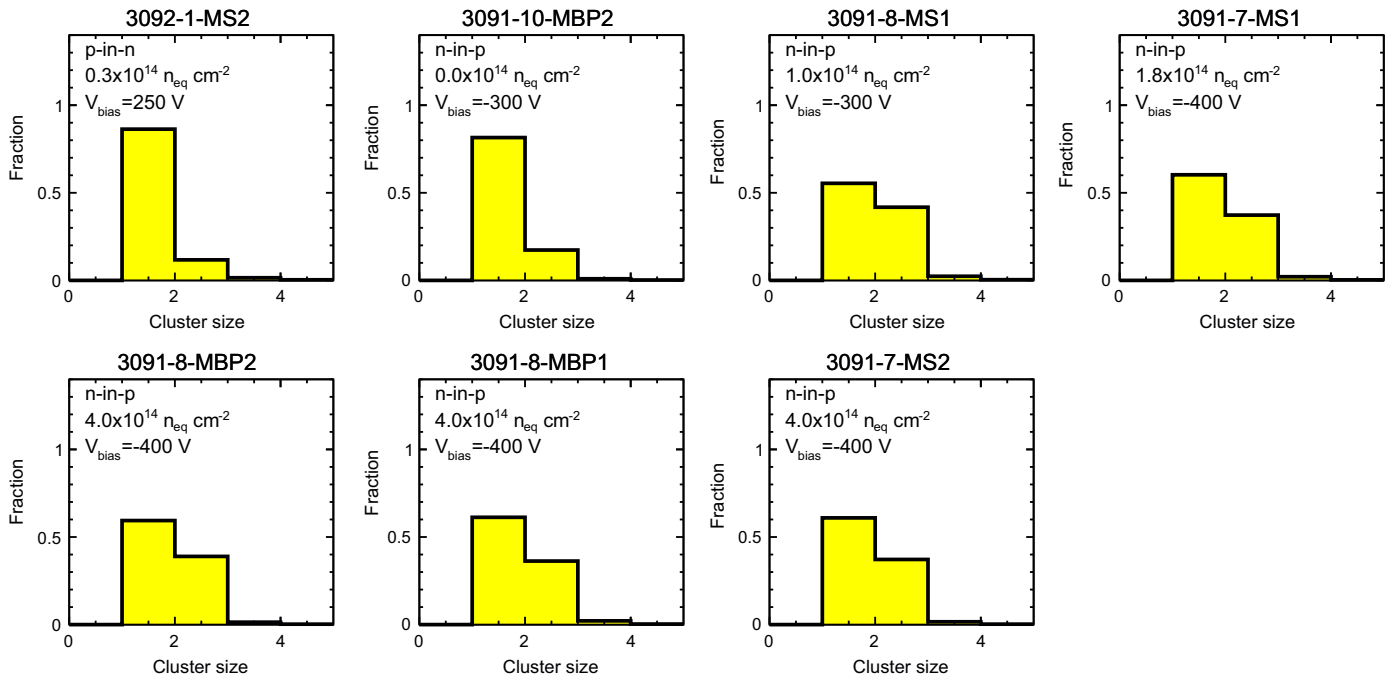


Fig. 13. Distributions of the cluster size for the DUTs, after all corrections for tracks at normal incidence. The runs are taken from the plateau region of each DUT.

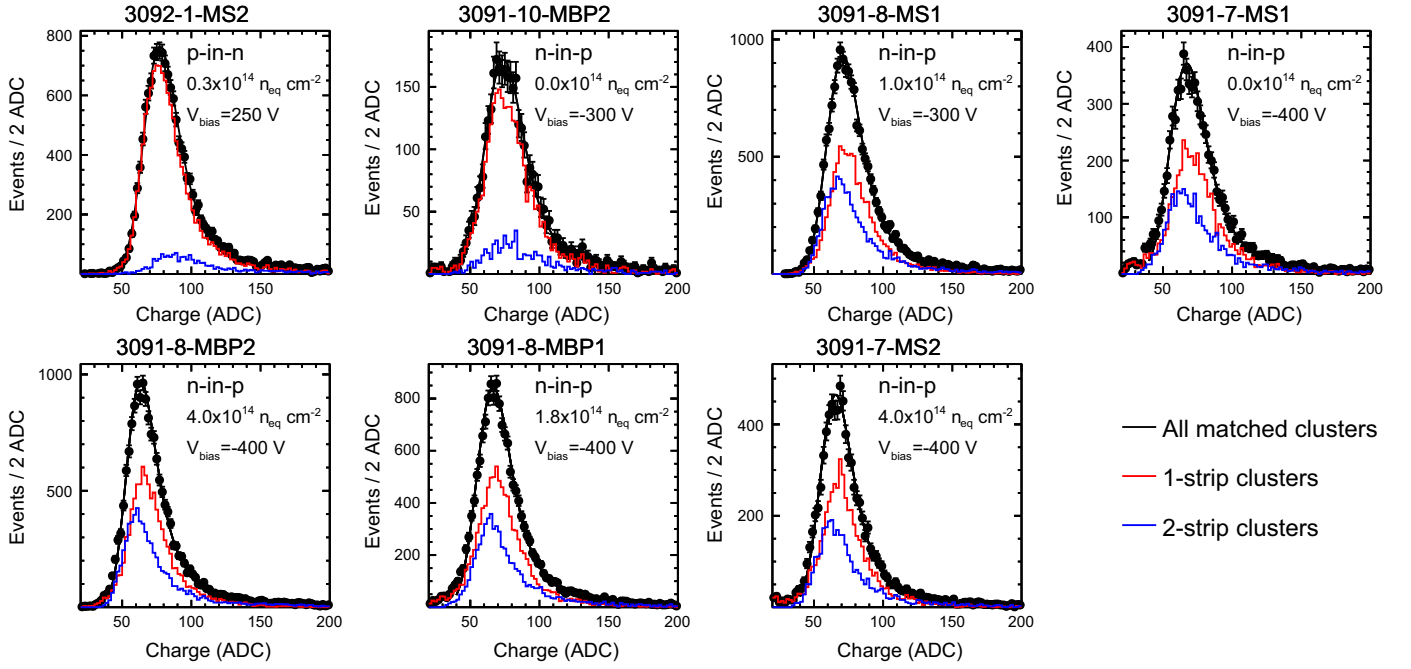


Fig. 14. Distributions of charge collected in each DUT in the plateau region for tracks at normal incidence. The black points show all clusters matched within 100 μm of a track, and the colored histograms show 1-strip and 2-strip cluster contributions. The data are fit to a Landau convoluted with a Gaussian resolution function, and the fit is shown. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

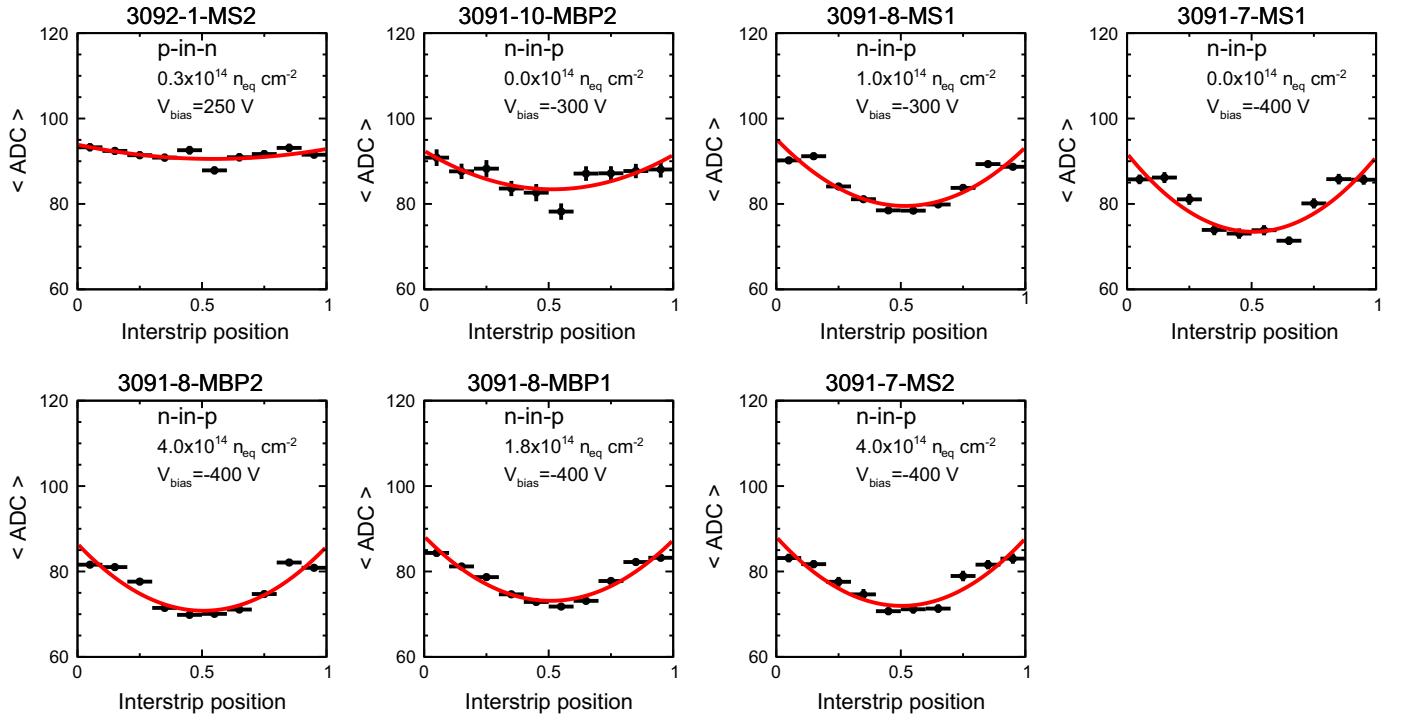


Fig. 15. Distributions of the average cluster charge as a function of the relative position within a strip, for tracks at normal incidence. A second order polynomial fit is overlaid just to guide the eye.

charge in the mid-region between two strips in irradiated detectors has been reported in test beams of the ATLAS SCT [20].

From this, it is evident that a sizeable part of the drop in collected charge of the irradiated sensors is due to loss of charge when the track has an interstrip position between 0.2 and 0.8, where a 10–15% loss in charge is seen. When the track has an interstrip position less

than 0.2 or greater than 0.8, the loss in collected charge is not as pronounced. Nevertheless, it is the average over the strip which is relevant for the detector performance. The efficiency vs. the interstrip position has also been investigated, and in all cases, it is at least 99%. Despite the loss in collected charge, there is no indication of any loss of efficiency.

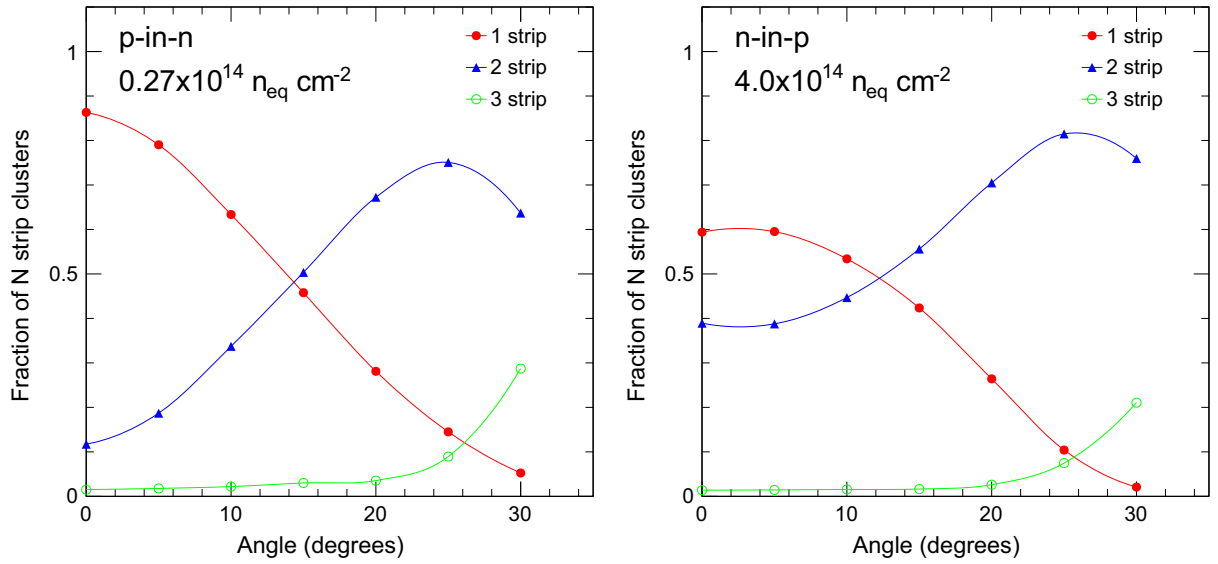


Fig. 16. Fraction of 1-strip, 2-strip, and 3-strip clusters as a function of angle for (left) 3092-1-MS2 ($0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ p-in-n) at $V_{\text{bias}} = 250 \text{ V}$, and (right) 3091-8-MBP2 ($4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ n-in-p) at $V_{\text{bias}} = 350 \text{ V}$.

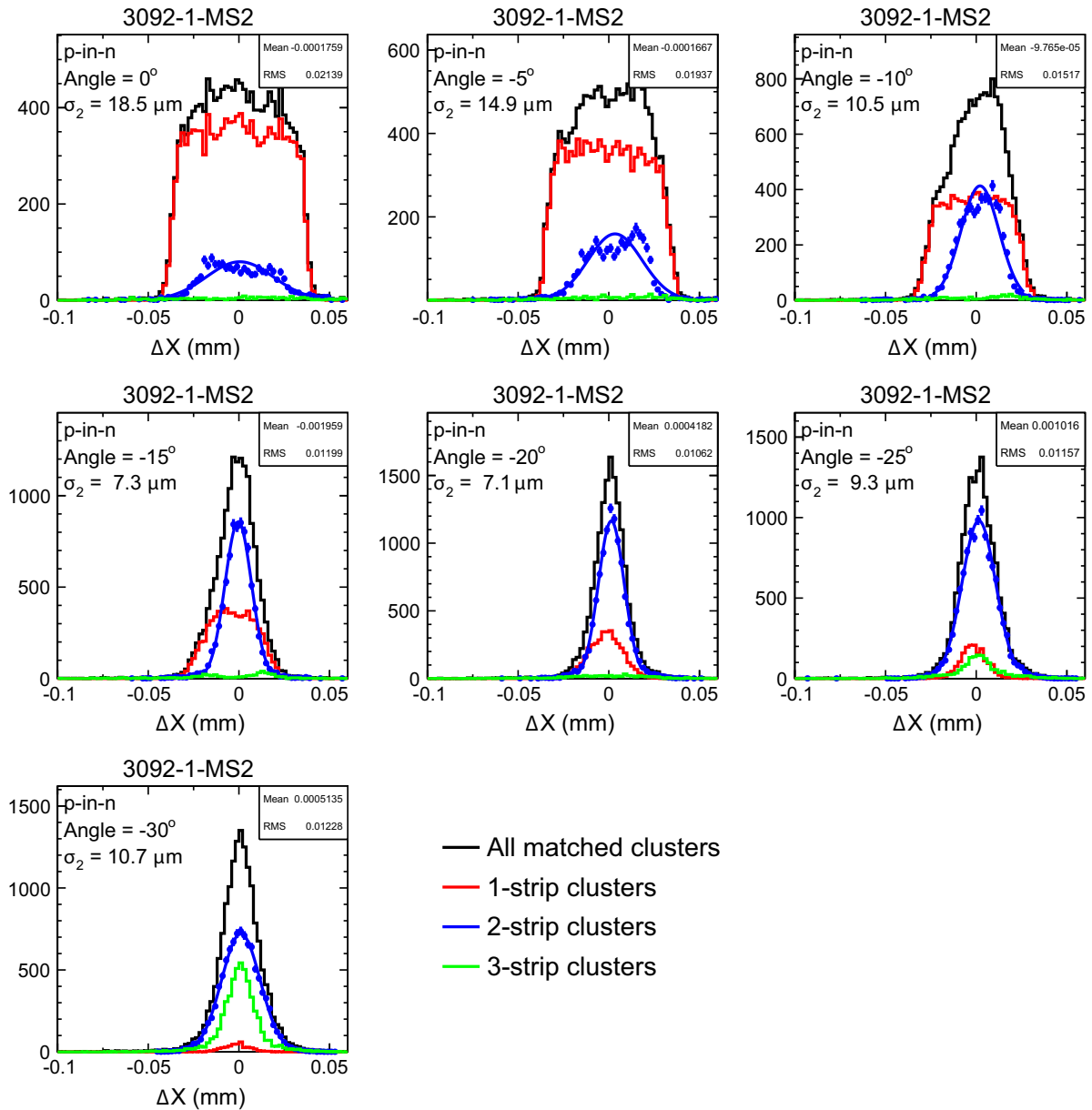


Fig. 17. Residual distributions for angles ranging from 0° to -30° for DUT 3092-1-MS2 ($0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ p-in-n) at $V_{\text{bias}} = 250 \text{ V}$. The contributions from 1-strip, 2-strip, and 3-strip clusters are also shown. The 2-strip cluster distributions are each fit to a single Gaussian function, and the widths, σ_2 , are also indicated.

4.8. Cluster size and resolution vs. angle

Studies of the detector performance were also carried out at angles ranging from normal incidence to 30° (with respect to the normal to the sensor). The fraction of 1, 2 and 3-strip clusters as a function of the angle for (left) the p-in-n sensor and (right) one of the 4.0×10^{14} n_{eq}/cm² n-in-p sensors is shown in Fig. 16. There is a clear increase in the fraction of 2-strip clusters in both cases, which reaches a maximum at about 22.5°, beyond which the 3-strip clusters start to become sizeable. The main difference between the two sensors is the cluster size at zero angle. For the heavily irradiated detectors, at normal incidence, the 2-strip cluster fraction is significantly larger (see cluster size distributions for normal incidence in Fig. 13).

The residual distributions between the UT hit and the track are shown in Fig. 17 for the 0.27×10^{14} n_{eq}/cm² p-in-n sensor. The angles range from 0° to 30° in 5° steps. The contributions from 1, 2, and 3-strip clusters are also shown. The 3-strip clusters make up

only a very small fraction, except at the two largest angles. As the angle increases the resolution improves, as indicated by the RMS of the distributions. The best resolution should be achieved when $\tan^{-1}(\text{pitch}/\text{thickness}) \sim 18^\circ$. The RMS is minimum for the 20° angle (RMS = 10.6 μm). The 2-strip residual distribution is fit to a Gaussian function, and the width returned by the fit, σ_2 , is shown. Two-strip clusters provide a resolution of about 7 μm for tracks with angles in the 15–20° range, which is about 3 times better than binary resolution ($80 \mu\text{m}/\sqrt{12}$). At small angles, most of the clusters are single strips, and the residual distribution is roughly flat, as one would expect. At very large angles, the 2-strip resolution worsens. This is likely due to a third strip that should have been included in the cluster but was below the minimum ADC side strip threshold to be included. Improved resolution, primarily for tracks at low incidence angle, was demonstrated by applying a non-linear charge weighting to two-strip clusters, but it is not presented here.

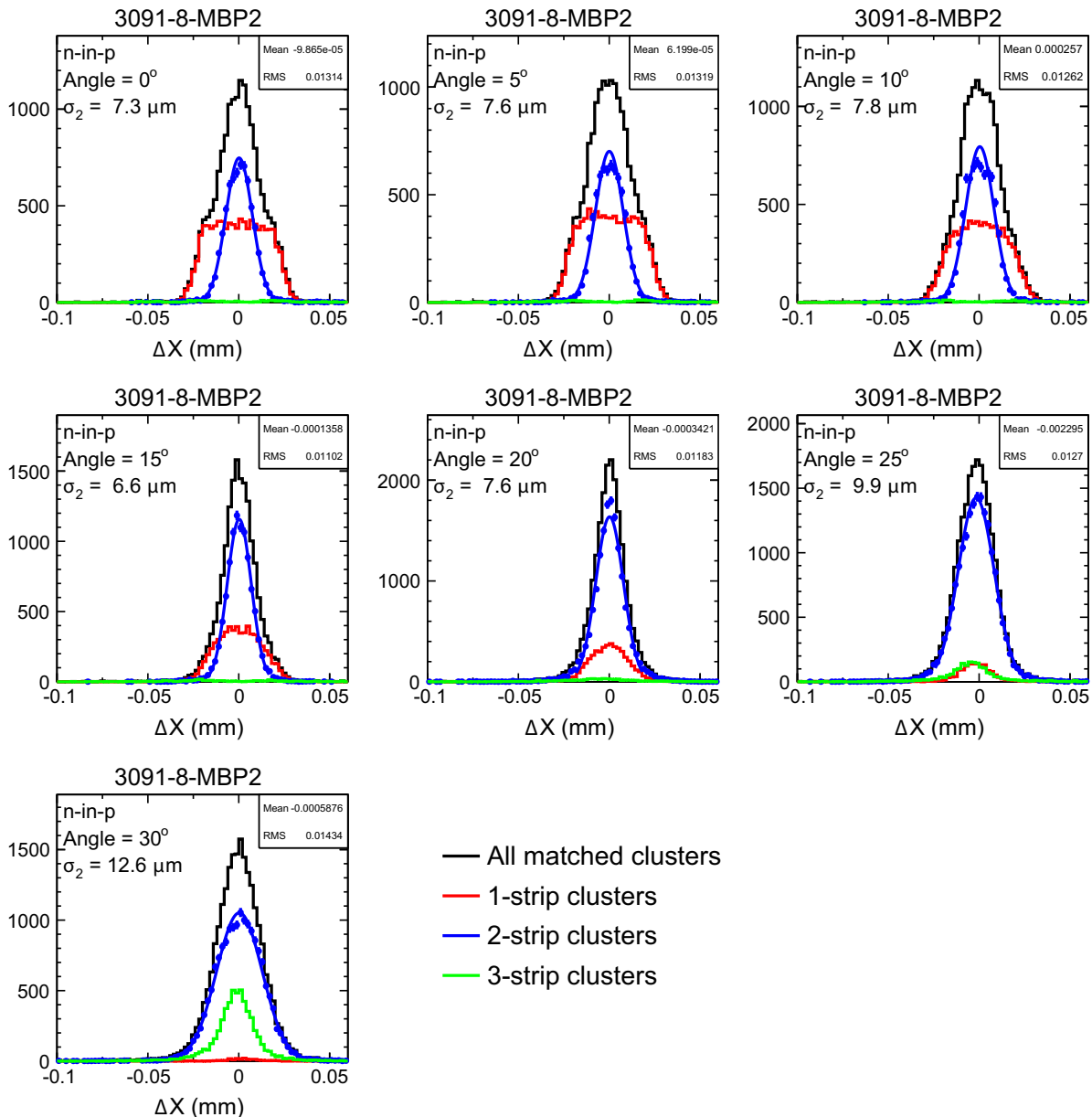


Fig. 18. Residual distributions for angles ranging from 0° to 30° for DUT 3091-8-MBP2 (4.0×10^{14} n_{eq}/cm² n-in-p at $V_{\text{bias}} = -350$ V). The contributions from 1, 2, and 3-strip clusters are also shown. The 2-strip cluster distributions are each fit to a single Gaussian function, and the widths, σ_2 , are also indicated.

Fig. 18 shows the corresponding residual distributions for one of the $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ n-in-p sensors. A similar resolution is obtained in the 15–20 ° angular range, as for the p-in-n sensor. However, the notable difference (with respect to the p-in-n sensor) is that at low angles, the 2-strip contribution is sizeable, and shows the improved resolution that is generally accompanied by a 2-strip cluster. This clearly illustrates that the larger cluster sizes

associated with the irradiated n-in-p sensors are not an anomaly, since they are also accompanied by better position resolution.

The position resolution as a function of the angle is presented for 3 of the sensors in Fig. 19, the $0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ p-in-n, the $1.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ n-in-p and one of the $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ n-in-p sensors. The resolution of the p-in-n sensor increases as one approaches normal incidence, reaching a maximum of about 23 μm , which is approximately binary, e.g. $80 \mu\text{m}/\sqrt{12}$. For the n-in-p sensors at normal incidence, the position resolution is much better, about 13 μm , owing to the larger amount of charge sharing between the strips (see Fig. 16). The $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ irradiated n-in-p sensor shows slightly worse position resolution than the $1.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ sensor, but it is only at the 2 μm level in the worst case. This difference is not of great concern for the UT detector.

4.9. Charge collection near the quarter-circle region

Three of the sensors studied during the test beam have a quarter-circle region where there are no strips (see Fig. 3). One of the sensors with this quarter-circle is the unirradiated n-in-p sensor, and the other two are board 4 sensors 0 and 1. The latter two have different guard-ring structures, as shown in Fig. 3. We investigate whether there is any drop in charge collected as one approaches the quarter-circle radially. Using the data, the center of the circle is determined using the 3D information from the TP3 tracks that have matched DUT clusters. The average ADC value as a function of the radial distance from the center of the circle is then analyzed. The results for the 3 sensors are shown in Fig. 20, where each row shows the results for a single sensor. The left column shows the XY profile of tracks that have matched DUT hits, along with an arc that shows the estimated circle to match the quarter-circle region void of strips. The right hand plots show an overlay of the radial distribution of tracks with matched DUT hits, with the

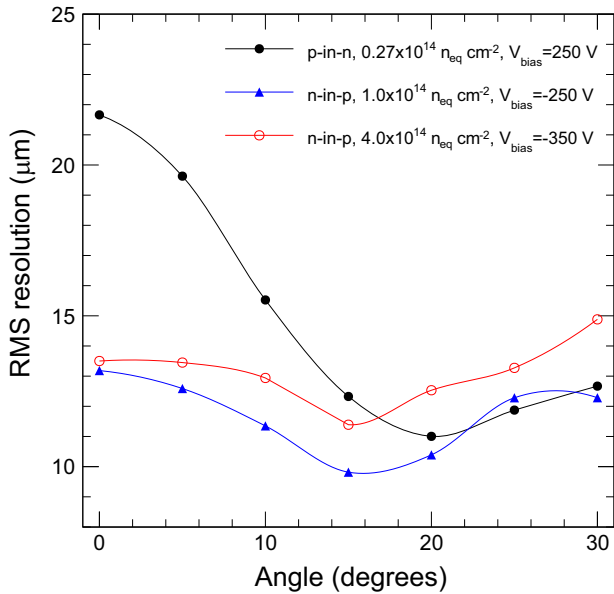


Fig. 19. RMS of the residual distribution as a function of angle (with respect to the normal to the sensor) for 3 different sensors. The bias voltages for the 3 sensors are 250 V.

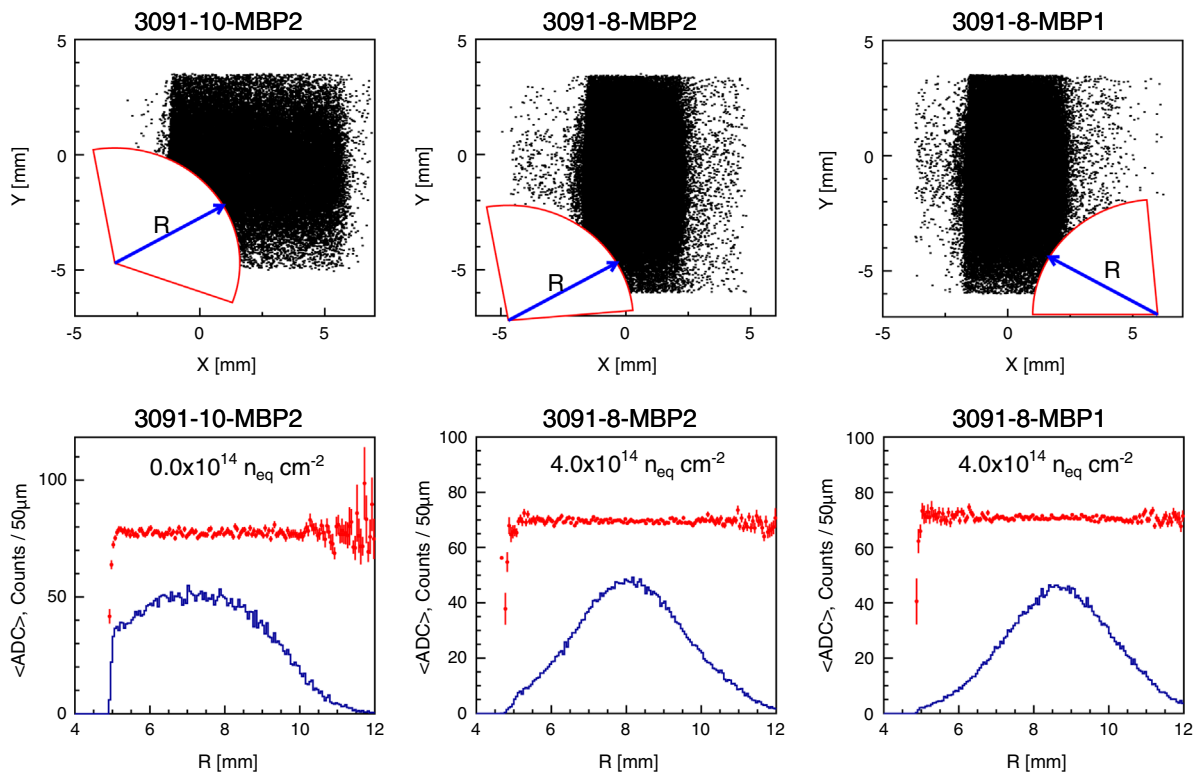


Fig. 20. (Top row) XY distributions of tracks matched to clusters for each of the 3 mini beam-pipe sensors for a bias voltage of -350 V . (Bottom row) Overlay of average ADC value vs. radial coordinate (red with error bars) and the radial distribution of tracks at the DUT (blue histogram). The radius is computed using the center of the red circles shown in the left column. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

average ADC value of clusters as a function of the radius. In this figure, several runs with bias voltage above 350 V are combined, and the average ADC is the raw value with no TDC time requirement. This is done to increase the sample size, as the main interest is to uncover a trend, while the absolute value is of less concern.

There is no clear indication of large loss of charge collection near the quarter-circle edge. A small loss in average charge in the final 50 μm or so from the edge cannot be ruled out, but the drop is most likely just the precision on which the center of the circle is determined, which is about 50 μm . No significant difference between the two guard-ring structure designs (MBP1 vs. MPB2 in bottom set of plots) is observed. Moreover, no significant differences between the unirradiated and irradiated detectors are seen.

5. Summary

Results from a test beam at the CERN SPS of 7 pre-prototype UT sensors, one p-in-n and 6 n-in-p, have been presented. The n-in-p sensors were irradiated at fluences of $1.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$, $1.8 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ and $4.0 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$, and showed a gradual loss in total charge collected with increased fluence. The maximum charge loss is about 10% for the detectors with the largest accumulated fluence. All detectors plateaued in the 300–400 V region, yielding a S/N of at least 15. With the full size UT sensors (10 cm \times 10 cm, or 10 cm \times 5 cm), there will be a larger input capacitance, and hence it is likely that the S/N will be lower. However, the studies presented indicate that the anticipated 500 V maximum will be sufficient for operation of these detectors for the full 50 fb^{-1} data set expected to be collected in Run III of LHCb.

We also find that the irradiated n-in-p DUTs have substantially larger cluster sizes at normal incidence, compared to either the unirradiated n-in-p DUT, or the $0.27 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ p-in-n DUT. These larger cluster sizes at normal incidence are accompanied by improved position resolution, as one expects when there is charge sharing between strips.

The dependence of the charge collection on the interstrip position has also been studied. For the irradiated n-in-p DUTs, it is found that about 15% less charge is collected when a track goes through the middle of two strips, relative to the amount collected when it passes through the center of a given strip. This is consistent with findings from other experiments. The efficiency as function of interstrip position has also been studied, and in all cases it is found to exceed about 99%.

The position resolution in the various sensors has also been investigated. At normal incidence, the irradiated detectors show substantially better position resolution than the unirradiated detector. This stems from the larger fraction of 2-strip clusters in the irradiated detectors. This goes against the conventional thinking that in irradiated detectors, less charge is collected, resulting in a reduction of the fraction of multi-strip clusters. At large angles, all detectors studied show improved spatial resolution, and the best resolution achieved is about 8 μm at an incidence angle in the 15–20° range. A resolution of about 6.5 μm is found for the subset of 2-strip clusters.

The charge collection close to the circular portion of the DUTs was studied, and there is no indication of any significant degradation of charge collection in the region close to this region.

Acknowledgments

We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the SPS, INFN (Italy); MNiSW and NCN (Poland); SNSF and SER (Switzerland); and NSF (USA). We also acknowledge the efforts of the VELO TimePix3 group, most notably Heinrich Schindler, Martin van Beuzekom, Hella Snoek, Paula Collins, Raphael Dumps, Kazu Akiba, Xabier Cid Vidal, Alvaro Dosil, Bas van der Heijden Sophie Richards, Dan Saunders, Panos Tsopelas, Jaap Velthuis and Mark Williams, for providing a high precision telescope to enable many of the studies presented in this paper. We also thank Ethan Cascio of the MGH proton irradiation facility for his support in carrying out the sensor irradiations.

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