

Existence and uniqueness for stochastic 2D Euler flows with bounded vorticity

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Abstract

The strong existence and the pathwise uniqueness of solutions with L^∞ -vorticity of the 2D stochastic Euler equations are proved. The noise is multiplicative and it involves the first derivatives. A Lagrangian approach is implemented, where a stochastic flow solving a nonlinear flow equation is constructed. The stability under regularizations is also proved.

1 Introduction

The aim of this paper is to prove the strong existence and the pathwise uniqueness of L^∞ solutions to the stochastic 2D Euler equation in vorticity form

$$d\xi + u^\xi \cdot \nabla \xi dt + \sum_{k=1}^{\infty} \sigma_k \cdot \nabla \xi \circ dW^k = 0, \quad \xi|_{t=0} = \xi_0, \quad (1.1)$$

where the initial vorticity ξ_0 also belongs to the L^∞ space. The equation above is subject to the periodic boundary conditions and thus can be reformulated as a problem on a 2-dimensional torus $\mathbb{T}^2 = (\mathbb{R}/\mathbb{Z})^2$, see for instance [41, chapter 2]. In other words the space variable is assumed to be an element of \mathbb{T}^2 and all fields are assumed to be 1-periodic (or simply defined on \mathbb{T}^2). The noise coefficients σ_k 's are bounded, regular enough, divergence-free vector fields, $(W^k)_{k=1}^\infty$ is a family of independent Brownian motions and the velocity field u^ξ is defined as

$$u_t^\xi(x) = K * \xi_t(x) = \int_{\mathbb{T}^2} K(x-y) \xi_t(y) dy, \quad x \in \mathbb{T}^2,$$

where $K = \nabla^\perp G = (-\partial_2 G, \partial_1 G)$ and G is the Green function of the Laplacian on the torus \mathbb{T}^2 with mean 0, i.e.,

$$u^\xi = -\nabla^\perp(-\Delta)^{-1}\xi.$$

We will also prove the stability of the solutions under regularization of the kernel K .

The Stratonovich form is the natural one for several reasons, including physical intuition related to the Wong-Zakai principle and the fact that an Itô term of the form $\sum_{k=1}^{\infty} \sigma_k \cdot \nabla \xi dW^k$ would require a compensating second order operator to hope for a well defined system, see [39]. Besides, the Stratonovich form preserves the L^2 norm of the solution and is the right one to deal with manifold-valued SPDEs, see [11]. However, for the opportunity of mathematical analysis, we will formally rewrite the equation in the Itô form

$$\begin{aligned} d\xi + u^\xi \cdot \nabla \xi dt &+ \sum_{k=1}^{\infty} \sigma_k \cdot \nabla \xi dW^k - \frac{1}{2} \sum_{k=1}^{\infty} (\sigma_k \cdot \nabla) \sigma_k \cdot \nabla \xi dt \\ &= \frac{1}{2} \sum_{k=1}^{\infty} \text{tr}[\sigma_k \sigma_k^* D^2 \xi] dt, \end{aligned} \tag{1.2}$$

and we will give a rigorous interpretation of the latter one (under some simplified assumptions). Nonetheless it is useful to think sometimes heuristically in form of the Stratonovich expression and it would be misleading to believe that the equation has a parabolic character due to the term $\text{tr}(aD^2\xi)$ in the Itô formulation.

The noise in equation (1.1) has a very special form, compared to general abstract models of Stochastic Partial Differential Equations (SPDEs). Our aim is not an abstract generality. We have chosen this noise for two reasons. Firstly, because it occupies a relevant position in the Mathematical Physics literature on fluids and secondly because it is of transport type, hence allowing us to use special tools related to the transport equations (flows, L^∞ -bounds). The applied and theoretical literature on SPDEs driven by this type of noise is rich, see for instance [5], [6], [14], [19], [22], [31], [33], [36], [30], in particular for its relation with turbulent transport of passive scalars and the so called Kraichnan model ([27], [28]), one of the most remarkable theories where stochastic models have been applied with success to explain phenomena in fluid mechanics. The transport structure of the nonlinear deterministic part of the equation (the vorticity in 2D is only transported) and

of the stochastic part (Stratonovich choice is important here), allow one to use stochastic the flows and to control the L^∞ -norm of solutions (the vorticity) by the L^∞ -norm of initial conditions. This control is ω -wise, uniform also in ω in Ω . Thus, having assumed that initial vorticity is bounded, the solution is uniformly bounded in all parameters (also ω), opposite to several other stochastic cases, like the additive noise. This property is an important tool of our approach and it cannot be readily extended to other stochastic perturbations of the Euler equations.

What concerns the theory of the deterministic Euler equations, the uniqueness for L^∞ -vorticity in the deterministic case is the celebrated result of Wolibner [42] and Yudovich ([44], [45]). In addition to an excellent recent monograph [34] where some additional information about the trajectory method can be found, one should also mention more recent publications as for instance a recent review paper [15] by Chemin and a new approach to the old non-uniqueness results of Schaeffer and Shnirelman by De Lellis and Székelyhidi in [16].

The literature on the stochastic Euler equations counts a number of works, including [7], [8], [9], [10], [12], [13], [21], [24], [25], [26], [37], [38], [43]. The differences are in the structure of the noise, the results and topologies involved and sometimes the domain and boundary conditions. A full discussion is not possible so we limit ourselves to few remarks. Some of the works deal with additive noise, some others with more general, namely multiplicative, noise but not of the form treated here which involves the derivatives of the solution, and one paper with noise with derivatives of the solution. When the noise is additive, the theory is more complete, also because the equation can be studied pathwise. First results were given in [10], where the existence is proved when the initial data belongs to the space V and the solution is an H -valued continuous and V -valued square integrable process (where V is the space of divergence free vector fields with finite enstrophy and H is the space of divergence free square integrable vector fields, periodic in an appropriate sense. However this solution, constructed pathwise on a given probability space, is not known to be progressively measurable. Moreover, if the vorticity of the initial data is bounded and the external forces (deterministic and random) satisfy certain assumptions, the solution is proved to be unique. These results in the additive noise case have been improved and generalized in the interesting paper [25], based on different techniques with respect to [10], which relaxes various regularity and boundary conditions on the noise for the result of existence and uniqueness of solutions with bounded vorticity

and proves very careful measurability properties in the case of solutions in V , those which are not necessarily unique. Let us also mention that in the additive noise case the more recent paper [23] gives delicate L^∞ -vorticity estimates on invariant measures for the stochastic Navier-Stokes equations with and their inviscid limit. Multiplicative noise, depending on the velocity field u (and not on the gradient) has been initially treated in the paper [13] by nonstandard analysis tools. That paper is devoted to the stochastic Euler equations on a 2-dimensional torus and the authors prove the existence of a solution on the Loeb space and the existence of a corresponding notion of statistical solution and it does not deal with the uniqueness. Then, in the paper [12] the authors prove the existence (but again not the uniqueness) of a solution to a problem with multiplicative noise as in [13] and possibly unbounded domains, but the state space is the space H intersected with the Sobolev space $H^{1,p}$ for $p > 2$. In this way they are able to prove the existence of solutions which are Hölder continuous with respect to the space variables. The question of uniqueness in the case of multiplicative noise, which was left open by these and other works (like [7]), has been recently investigated in the paper [24], however only for one-dimensional Brownian motion (so that the Doss-Sussmann transformation can be used). Finally, let us mention the recent paper [43] by Yokoyama, which is the closest to our model (the stochastic Euler equations in Stratonovitch form with the noise coefficients depending linearly on the gradient of the solution), where the author proves the existence of a martingale solution with the state space V ; the paper does not deal with the uniqueness.

We solve here the problem in the space L^∞ following the Lagrangian approach of [35]. It is based, in the stochastic case, on the investigation of the stochastic flow equation

$$\begin{aligned} \Phi_t(x) &= x + \int_0^t \int_{\mathbb{T}^2} K(\Phi_s(x) - \Phi_s(y)) \xi_0(y) dy \\ &+ \sum_k \int_0^t \sigma_k(\Phi_s(x)) dW_s^k, \quad t \in [0, T], \quad x \in \mathbb{T}^2 \end{aligned}$$

which is a problem of interest in itself, even when the kernel K is smooth. This equation is not trivial because of the global dependence of $\Phi_t(x)$ on $(\Phi_s(y))_{y \in \mathbb{T}^2}$ and the difficulty to develop stochastic calculus (for instance a fixed point argument) in the space of (measure preserving, continuous) maps $\psi : \mathbb{T}^2 \rightarrow \mathbb{T}^2$. The approach inspired by [35] allows us to study this equation

and apply the result to the existence and the uniqueness of equation (1.1) in L^∞ .

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2 The main results

Before stating the results, we list the hypotheses with some preliminary remarks.

Condition 2.1. *In the paper, we will always assume that ξ_0 , the initial vorticity, belongs to the space $L^\infty(\mathbb{T}^2)$.*

Condition 2.2. *The family of processes $W = (W^k)_{k=1}^\infty$ is a cylindrical Brownian motion (i.e. W^k 's are independent Brownian motions), defined on a probability space (Ω, \mathcal{A}, P) , with respect to the filtration $\mathbb{F} = (\mathcal{F}_t)_{t \in [0, T]}$.*

Condition 2.3. *The vector fields σ_k 's are divergence-free and belong to $C^{0,1}(\mathbb{T}^2)$ (Lipschitz periodic functions, hence a.e. differentiable); moreover the family $(\sigma_k)_{k=1}^\infty$ is in $W^{1,\infty}(\ell^2)$, that is*

$$L_\sigma^2 := \sup_{x \in \mathbb{T}^2} \sum_{k=1}^{\infty} |\sigma_k(x)|^2 + \left\| \sum_{k=1}^{\infty} |D\sigma_k|^2 \right\|_{L^\infty} < +\infty.$$

We call $a(x) := \sum_{k=1}^{\infty} \sigma_k(x) \sigma_k(x)^*$ (A^* denotes the transpose matrix of A). We assume also that $a \equiv cI_2$, where c is a non-negative constant (possibly equal to 0) and I_2 is the constant identity matrix.

Remark 2.4. *If $a(x) \equiv cI_2$, for all x in \mathbb{T}^2 , the Itô formulation (1.2) of the stochastic Euler Equations simplifies to*

$$d\xi + u^\xi \cdot \nabla \xi dt + \sum_{k=1}^{\infty} \sigma_k \cdot \nabla \xi dW^k = \frac{1}{2} c \Delta \xi dt. \quad (2.1)$$

Indeed, since the vector fields σ_k 's are divergence-free, the first order Itô correction term, namely $\frac{1}{2} \sum_k (\sigma_k \cdot \nabla) \sigma_k \cdot \nabla \xi dt$, disappears:

$$\sum_k \sum_i \sigma_{k,i}(x) \partial_i \sigma_{k,j}(x) = \sum_i \partial_i \left(\sum_k \sigma_{k,i}(x) \sigma_{k,j}(x) \right) = \sum_i \partial_i a_{ij}(x) = 0.$$

Remark 2.5. Condition $a(x) \equiv cI_2$, $x \in \mathbb{T}^2$, can be avoided at the price of requiring more regularity on the functions σ_k 's and of a few additional computations, which would obscure the main arguments. Indeed, the fact that a is constant implies the absence of the first order Itô correction term, which contains the derivatives of σ , and that the operator $\frac{1}{2} \text{tr}[aD^2] = \frac{1}{2} c \Delta$ commutes with the convolution with a given function; this will avoid the use of a second order commutator lemma (not difficult but boring and requiring maybe more regularity on the σ_k 's).

Remark 2.6. Let us briefly discuss examples of noise covered by the class above. The trivial example of a noise term of the form $\nabla \xi_t \cdot dW_t$ where W is a 2-dimensional Brownian motion is covered by taking $\sigma_k = e_k$ for $k = 1, 2$ (where (e_1, e_2) is the canonical basis of \mathbb{R}^2) and $\sigma_k = 0$ for $k \geq 3$. In this case it should be noticed that the stochastic Euler equations can be reduced to the classical deterministic ones by the simple transformation $\tilde{\xi}(t, x) = \xi(t, x + W_t)$. More than this one, we are mainly motivated by the examples described in the Mathematical Physics literature quoted in the Introduction, where the noise term has heuristically the form $\nabla \xi(t, x) \cdot \partial_t W(t, x)$, for a space-dependent random field $W(t, x)$, Brownian in time, with a given incremental covariance function $Q(x, y) = E[W(1, x) \otimes W(1, y)]$, sometimes prescribed through its Fourier spectrum like $Q(x - y) = \int_{\mathbb{Z}^2} e^{ik \cdot x} f(|k|) dk$, for suitable functions f . A rigorous and simple way to deal with such space-time noise (correlated in space) is the one adopted above, namely to prescribe a sequence of independent real valued Brownian motions W_t^k and a sequence of vector fields $\sigma_k(x)$. The space-dependent noise $W(t, x)$ is then given by $W(t, x) = \sum_{k=1}^{\infty} \sigma_k(x) W_t^k$ and the function $Q(x, y) = \sum_{k=1}^{\infty} \sigma_k(x) \otimes \sigma_k(y)$ is its incremental covariance. If one starts with a prescribed covariance function $Q(x, y)$ (with suitable properties), the σ_k 's are an orthonormal basis of a certain Hilbert space, thus their form is not explicitly given (though their existence is guaranteed, see [5] and [33]).

However, we should notice that in comparison with the literature on the Kraichnan model of turbulent advection (related to the original Kraichnan's papers [27], [28]), we impose regularity properties on the vector fields σ_k

which forbid us from considering certain singular examples treated there. The covariance function $Q(x, y)$, corresponding to our case, is always relatively regular, while it scales with fractional powers of $|x - y|$ in Kraichnan model, see for example [22], [33].

Definition 2.7. Let ξ be an element of $L^\infty([0, T] \times \mathbb{T}^2 \times \Omega)$. We say that ξ is \mathbb{F} -weakly progressively measurable if, for every f in $L^1(\mathbb{T}^2)$, the process $t \rightarrow \langle \xi_t, f \rangle = \int_{\mathbb{T}^2} f \xi_t dx$ is \mathbb{F} -progressively measurable.

Given an element w in $L^\infty(\mathbb{T}^2)$, we will write

$$u = u^w = K * w.$$

If w is also time-dependent, we will write $u_t^w = u^{wt}$. It is well known, see Corollary 2.18, that $|u^w(x) - u^w(y)| \leq L_K \|w\|_{L^\infty} |x - y| (1 - \log |x - y|)$ for some constant L_K if $|x - y| \leq 1$.

Now we give a precise definition of a solution. We use the Itô formulation, having in mind Remark 2.4. In what follows, $\langle f, g \rangle := \int_{\mathbb{T}^2} fg dx$ denotes the scalar product in $L^2(\mathbb{T}^2)$.

Definition 2.8. Let ξ_0 be in $L^\infty(\mathbb{T}^2)$. A distributional L^∞ solution to the stochastic Euler vorticity equation (2.1) is an \mathbb{F} -weakly progressively measurable element ξ in $L^\infty([0, T] \times \mathbb{T}^2 \times \Omega)$, such that, for every φ in $C^\infty(\mathbb{T}^2)$, it holds P -a.s.

$$\begin{aligned} \langle \xi_t, \varphi \rangle &= \langle \xi_0, \varphi \rangle + \int_0^t \langle \xi_r, u_r^\xi \cdot \nabla \varphi \rangle dr + \sum_k \int_0^t \langle \xi_r, \sigma_k \cdot \nabla \varphi \rangle dW_r \\ &+ \frac{1}{2} \int_0^t \langle \xi_r, \text{tr}[aD^2\varphi] \rangle dr \quad \forall t \in [0, T]. \end{aligned} \quad (2.2)$$

It is implicit in the definition that the process $\langle \xi_t, \varphi \rangle$ has continuous trajectories.

Remark 2.9. If a process $\xi \in L^\infty([0, T] \times \mathbb{T}^2 \times \Omega)$ is weakly progressive measurable then so is the process $u^\xi \xi$. Indeed it implies that, for every h in $L^1(\mathbb{T}^2 \times \mathbb{T}^2)$, the process

$$t \mapsto \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} \xi(t, x) \xi(t, y) h(x, y) dx dy \quad (2.3)$$

is progressive measurable (this can be verified first for h of the form $h(x, y) = f(x)g(y)$, then approximating every h with sums of such separable functions). Now, for a test function φ , it is enough to write $\int u^\xi \xi \cdot \varphi dx$ as

$$\int_{\mathbb{T}^2} \int_{\mathbb{T}^2} K(x - y) \xi_t(y) \xi_t(x) \cdot \varphi(x) dx dy$$

and take $h(x, y) = K(x - y)\varphi(x)$.

The main result about the stochastic Euler vorticity equation is as follows.

Theorem 2.10. *Given ξ_0 in $L^\infty(\mathbb{T}^2)$ and the cylindrical Brownian motion W (with the associated filtration), under Conditions 2.2 and 2.3 on the coefficients of the noise, the stochastic Euler vorticity equation (2.1) admits a unique L^∞ distributional solution.*

Remark 2.11. *Notice that the filtration is given a-priori. Thus both the existence and the uniqueness are in the strong sense: there exists a solution ξ adapted to the (completed) Brownian filtration (the smallest possible filtration) and any solution, defined on a possibly larger filtered space, must coincide with ξ . The same kind of existence and uniqueness will hold for every equation we will meet.*

Theorem 2.10 will be proved by solving the associated non-local SDE:

$$\begin{aligned} \Phi_t(x) &= x + \int_0^t \int_{\mathbb{T}^2} K(\Phi_r(x) - \Phi_r(y)) \xi_0(y) dy dr \\ &+ \sum_k \int_0^t \sigma_k(\Phi_r(x)) dW_r^k. \end{aligned} \quad (2.4)$$

Notice that here the drift, namely

$$u^\Phi(t, x) = \int_{\mathbb{T}^2} K(x - \Phi_t(y)) \xi_0(y) dy, \quad (2.5)$$

depends on the whole flow.

Definition 2.12. • *A stochastic continuous flow is a measurable map $\Phi : [0, T] \times \mathbb{T}^2 \times \Omega \rightarrow \mathbb{T}^2$ such that, for a.e. ω in Ω , $\Phi(\omega) : [0, T] \times \mathbb{T}^2 \rightarrow \mathbb{T}^2$ is continuous and, for every $x \in \mathbb{T}^2$, the process $\Phi(x) : [0, T] \times \Omega \rightarrow \mathbb{T}^2$ is progressively measurable.*

- A stochastic continuous flow Φ is said to be measure-preserving iff¹, for a.e. ω , $\Phi_t(\omega) : \mathbb{T}^2 \rightarrow \mathbb{T}^2$ preserves the Lebesgue measure on \mathbb{T}^2 for every t .

Definition 2.13. We say that a stochastic continuous flow Φ is solution to the SDE (2.4) if, for every x , the process $X := \Phi(x)$ solves the SDE

$$dX = u^\Phi(X) dt + \sum_k \sigma_k(X) dW^k \quad (2.6)$$

with initial condition $X_0 = x$.

Theorem 2.14. Given ξ_0 in $L^\infty(\mathbb{T}^2)$ and the cylindrical Brownian motion W (with the associated filtration), there exists a unique measure-preserving stochastic flow solution to equation (2.4). This solution is a continuous flow Φ of class C^α in space and C^β in time, for some $\alpha > 0$ and for every $\beta < 1/2$.

Remark 2.15. Actually the uniqueness holds in a larger class of flows, namely the class SM defined at the beginning of Section 4.2, as it can be seen from the proof of Theorem 2.14.

2.1 The strategy

There are two ways to prove our results. We will develop mainly the one which requires the weakest regularity assumptions on the σ_k 's. This strategy will be as follows.

First we will prove that, for a log-Lipschitz random vector field u , the SDE

$$dX_t = u(X_t) dt + \sum_k \sigma_k(X_t) dW_t^k$$

admits a unique solution, given by a stochastic measure-preserving continuous flow (Lemma 4.5). This includes the case of a “linear” version of (2.4), where the drift is replaced by u^ψ for some fixed stochastic flow ψ , see also next paragraph for notation. Then, using an iteration scheme, we will build a unique solution to (2.4), reaching the assertion of Theorem 2.14.

In the subsequent section, we will use Theorem 2.14 to prove Theorem 2.10.

¹To avoid any ambiguity here and other similar situations we assume that there exists a measurable set $\tilde{\Omega}$ of full \mathbb{P} -measure, such that for all $\omega \in \tilde{\Omega}$ and every $t \in [0, T]$, the map $\Phi_t(\omega) : \mathbb{T}^2 \rightarrow \mathbb{T}^2$ preserves the Lebesgue measure on \mathbb{T}^2 .

In the last section, we will show the second method: a “trick” allows us to reduce the stochastic case to a modified deterministic case. This seems to be more rapid but requires the σ_k 's to be at least C^2 (at least if one wants to use classical results), while the first method requires only a Lipschitz-type (precisely $W^{1,\infty}(\ell^2)$) hypothesis on the diffusion coefficients. That is why we will not develop this second method in all the details.

2.2 Log-Lipschitz property of K and other useful facts

First we state the fundamental log-Lipschitz property for K and the drift u^ε . The key inequality (2.8) is stated in [35, section 1.2], and it follows from standard estimates of the Green function G , see e.g. [3, section 4.2]. For the completeness sake, we have recalled the proof in the appendix. For $r \geq 0$, call

$$\gamma(r) = r(1 - \log r)1_{]0,1/e[}(r) + (r + (1/e))1_{[1/e,+\infty[}(r).$$

Remark 2.16. *The following elementary properties of γ will be of use: the function γ is increasing, concave and for every $0 < \varepsilon < 1/e$, we have*

$$\gamma(r) \leq -r \log \varepsilon + \varepsilon, \quad \forall r \geq 0. \quad (2.7)$$

Lemma 2.17. *The map K , introduced before, is an $L^p(\mathbb{T}^2)$ divergence-free (in the distributional sense) vector field, for every $p < 2$, and verifies for certain constants $L_{0,K}$, L_K :*

$$\begin{aligned} \|K\|_{L^1(\mathbb{T}^2)} &\leq L_{0,K}, \\ \int_{\mathbb{T}^2} |K(x-y) - K(x'-y)| dy &\leq L_K \gamma(|x-x'|), \quad \forall x, x' \in \mathbb{T}^2. \end{aligned} \quad (2.8)$$

The divergence-free property is a consequence of the fact that K is orthogonal to a gradient of a scalar field.

Corollary 2.18. *For every w in $L^\infty(\mathbb{T}^2)$, $u^w = K * w$ is divergence-free and satisfies*

$$\begin{aligned} \|u^w\|_{L^\infty} &\leq L_{0,K} \|w\|_{L^\infty}, \\ |u^w(x) - u^w(x')| &\leq L_K \|w\|_{L^\infty} \gamma(|x-x'|), \quad \forall x, x' \in \mathbb{T}^2. \end{aligned} \quad (2.9)$$

We will use also the following elementary result. We recall that, for a finite signed measure μ on a space E and a measurable map $F : E \rightarrow E'$, $\nu = F_{\#}\mu$ denotes the image measure of μ on E' , namely $\nu(A) = \mu(F^{-1}(A))$ for every measurable set A in E' . Notice that ν is a finite signed measure and that $|\nu| \leq F_{\#}|\mu|$ (since $|\nu|(A) \leq F_{\#}|\mu|(A)$ for every A).

Lemma 2.19. *Let F be a measure preserving map on \mathbb{T}^2 and let w be in $L^\infty(\mathbb{T}^2)$. Let μ the (signed) measure on \mathbb{T}^2 with density w (with respect to the Lebesgue measure) and define $\nu = F_{\#}w$. Then ν has a density (denoted by v) with respect to Lebesgue measure and $\|v\|_{L^\infty} \leq \|w\|_{L^\infty}$.*

Proof. It is enough to prove the Lemma when w is nonnegative. Since F is measure-preserving, if A is a set of zero Lebesgue measure, then $\mathcal{L}^2\{F \in A\} = \mathcal{L}^2(A) = 0$, and so $\int_A d\nu = \int_{\mathbb{T}^2} 1_A(F)w dx = 0$. So ν admits a (nonnegative) density v . Now, taking $\varepsilon > 0$, $B = \{v > \|w\|_{L^\infty} + \varepsilon\}$, we have

$$\begin{aligned} (\|w\|_{L^\infty} + \varepsilon)\mathcal{L}^2(B) &\leq \int_{\mathbb{T}^2} 1_B v dx = \int_{\mathbb{T}^2} 1_B(F)w dx \leq \\ &\leq \|w\|_{L^\infty}\mathcal{L}^2\{F \in B\} = \|w\|_{L^\infty}\mathcal{L}^2(B), \end{aligned}$$

which implies that $\mathcal{L}^2(B) = 0$. By arbitrariness of ε , we get $\|v\|_{L^\infty} \leq \|w\|_{L^\infty}$. \square

Given Lemma 2.19, we will use often $v = F_{\#}w$ instead of $\nu = F_{\#}\mu$. Finally some other notation. Let ψ a measurable measure-preserving flow on \mathbb{T}^2 . With the notation in the previous section define $\xi_t^\psi = (\psi_t)_{\#}\xi_0$ (which is in $L^\infty([0, T] \times \mathbb{T}^2)$ by Lemma 2.19) and $u^\psi = u^{\xi^\psi}$, which also reads

$$u^\psi(t, x) = \int_{\mathbb{T}^2} K(x - \psi_t(y)) \xi_0(y) dy.$$

As already noticed, the SDE (2.4) reads as

$$\Phi_t(x) = x + \int_0^t u_r^\Phi(\Phi_r(x)) dr + \sum_k \int_0^t \sigma_k(\Phi_r(x)) dW_r^k.$$

Remark 2.20. *By the definition of u^ψ , Corollary 2.18 and Lemma 2.19, u^ψ enjoys $\|u^\psi\|_{L^\infty} \leq L_{0,K}\|\xi_0\|_{L^\infty}$ and the following log-Lipschitz property:*

$$|u^\psi(x) - u^\psi(x')| \leq L_K \|\xi_0\|_{L^\infty} \gamma(|x - x'|), \quad \forall x, x' \in \mathbb{T}^2. \quad (2.10)$$

Given λ a positive constant and z_0 in $[0, 1/e]$, we will also denote by $z^\lambda(t, z_0)$ (omitting the λ when not necessary) the solution to the ODE

$$z_t = z_0 + \int_0^t \lambda \gamma_r(z_r) dr$$

This z is unique and has the explicit formula

$$z(t, z_0) = z_0^{\exp[-\lambda t]} e^{1 - \exp[-\lambda t]} 1_{t < t_0} + (2e^{-1} \exp[\lambda(t - t_0)] - e^{-1}) 1_{t \geq t_0}, \quad (2.11)$$

where $t_0 = t_0(\lambda, z_0) = \frac{1}{\lambda} \log \frac{1 - \log z_0}{2}$ is the time such that $z(t_0) = 1/e$. Notice that, for z_0 in $[0, \exp[1 - 2e^{\lambda T}]]$, it holds $t_0 \geq T$ and so, for t in $[0, T]$,

$$z(t, z_0) \leq e z_0^{\exp[-\lambda t]}. \quad (2.12)$$

3 The deterministic case

We first treat the deterministic case, in order to show the basic ideas. The scheme of the proof, strongly inspired by [35], is a suitable rewriting of [35], convenient for generalization to the stochastic case.

Euler flows in 2D (on the torus \mathbb{T}^2) are described by the following non-local ODE:

$$\Phi_t(x) = x + \int_0^t \int_{\mathbb{T}^2} K(\Phi_s(x) - \Phi_s(y)) \xi_0(y) dy. \quad (3.1)$$

Equation (3.1) reads as $\dot{\Phi} = u^\Phi(\Phi)$ (with initial condition $\Phi_0 = id$), notice that the drift is log-Lipschitz. That is why we consider the auxiliary equation (linear problem):

$$X_t^x = x + \int_0^t u(s, X_s^x) ds, \quad (3.2)$$

where u is a fixed measurable vector field with the following property: for every t, x, y ,

$$|u(t, x) - u(t, y)| \leq L_u \gamma(|x - y|) \quad (3.3)$$

for some L_u independent of t, x, y .

Lemma 3.1. *For every initial datum x , equation (3.2) has a unique solution. This solution is described by a (unique) flow ψ of measure-preserving homeomorphisms of class C^α in space and Lipschitz in time, with $\alpha = \exp[-L_u T]$.*

Proof. The existence of a global solution (in \mathbb{R}^2) to (3.2) follows from the Peano Theorem, since u is continuous bounded. The uniqueness holds by the Osgood criterion (since $\int_0^\varepsilon \gamma(r)^{-1} dr = +\infty$) or even by the Hölder estimate below (simply take $x = y$).

The Lipschitz continuity in time follows by boundedness of u . As for the Hölder continuity, property (3.3) implies that, for every x and x' ,

$$|\psi_t(x) - \psi_t(x')| \leq |x - x'| + L_u \int_0^t \gamma(|\psi_s(x) - \psi_s(x')|) ds.$$

By a comparison result, $|\psi_t(x) - \psi_t(x')| \leq z^{L_u}(t, |x - x'|)$ (recall that z^λ is the unique solution to $z_t = z_0 + \int_0^t \lambda \gamma(z_s) ds$). The bound (2.12) for z gives the desired regularity. The invertibility and the continuity of the inverse map are due to the classical cocycle law, so that the inverse flow of ψ_t is ψ_{-t} . The measure-preserving property follows by a simple approximation argument, see the proof of Lemma 4.5 in the stochastic case. \square

Now we use the Picard iteration scheme to prove the existence and the uniqueness of solutions to (3.1). Consider the set

$$M_T = \left\{ \psi : [0, T] \times \mathbb{T}^2 \rightarrow \mathbb{T}^2 \mid \begin{array}{l} \psi \text{ measurable, } \sup_{[0, T]} \int_{\mathbb{T}^2} |\psi_t(x)| dx < +\infty, \\ \psi_t \text{ measure-preserving for a.e. } t \end{array} \right\}.$$

It is a complete metric space, endowed with the distance $dist(\psi^1, \psi^2) = \sup_{[0, T]} \int_{\mathbb{T}^2} |\psi_t^1(x) - \psi_t^2(x)| dx$. For any ψ in M_T , define $G(\psi)$ as the unique flow solution to (3.2) with $u = u^\psi$, i.e.

$$\frac{d}{dt} G(\psi) = u^\psi(G(\psi))$$

(with initial condition $G(\psi)(0, x) = x$). Recall that

$$u^\psi(t, x) = \int_{\mathbb{T}^2} K(x - \psi_t(y)) \xi_0(y) dy.$$

enjoys the log-Lipschitz property (2.10), so that by the previous Lemma G takes values in M_T .

Lemma 3.2. *For every $\varepsilon > 0$, for every two flows ψ^1, ψ^2 in M_T , we have:*

$$\begin{aligned}
& \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\
& \leq L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx \right) ds \\
& + L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} |G(\psi^1)_s(x) - G(\psi^2)_s(x)| dx \right) ds \quad (3.4)
\end{aligned}$$

and also

$$\begin{aligned}
& \int_{\mathbb{T}^2} |G(\psi^1)_t - G(\psi^2)_t| dx \\
& \leq L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} |G(\psi^1)_s - G(\psi^2)_s| dx ds \\
& + L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} |\psi_s^1 - \psi_s^2| dx ds + 2L_K \|\xi_0\|_{L^\infty} t \varepsilon. \quad (3.5)
\end{aligned}$$

Proof. We have

$$\begin{aligned}
& \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \leq \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} \\
& |K(G(\psi^1)_s(x) - \psi_s^1(y)) - K(G(\psi^2)_s(x) - \psi_s^2(y))| dx dy ds.
\end{aligned}$$

In order to use (2.8), we add and subtract $K(G(\psi^1)_s(x) - \psi_s^2(y))$ to the integrand of the right-hand side. Thus we get

$$\begin{aligned}
& \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \leq \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} \\
& \left[|K(G(\psi^1)_s(x) - \psi_s^1(y)) - K(G(\psi^1)_s(x) - \psi_s^2(y))| \right. \\
& \left. + |K(G(\psi^1)_s(x) - \psi_s^2(y)) - K(G(\psi^2)_s(x) - \psi_s^2(y))| \right] dx dy ds \\
& \leq \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} \left[|K(x - \psi_s^1(y)) - K(x - \psi_s^2(y))| \right. \\
& \left. |K(G(\psi^1)_s(x) - y) - K(G(\psi^2)_s(x) - y)| \right] dx dy ds \\
& \leq L_K \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \gamma(|\psi_s^1(y) - \psi_s^2(y)|) dy ds \\
& + L_K \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \gamma(|G(\psi^1)_s(x) - G(\psi^2)_s(x)|) dx ds,
\end{aligned}$$

where in the second passage we used the measure-preserving property. Finally, by the Jensen inequality applied to the concave function γ , we have

$$\begin{aligned} & \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx \right) ds \\ & + L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} |G(\psi^1)_s(x) - G(\psi^2)_s(x)| dx \right) ds, \end{aligned}$$

that is the first estimate (3.4). Now we apply property (2.7):

$$\begin{aligned} & \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx ds + \\ & + L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} |G(\psi^1)_s(x) - G(\psi^2)_s(x)| dx ds + \\ & + 2L_K \|\xi_0\|_{L^\infty} t\varepsilon, \end{aligned}$$

i.e. the second estimate (3.5). \square

The following continuity result is a consequence of the previous Lemma.

Corollary 3.3. *The map $G : M_T \rightarrow M_T$ is continuous. In fact, it is locally Hölder continuous.*

Proof. Let us denote $w_t = \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx$. Then the estimate (3.4) in Lemma 3.2, together with monotonicity of γ , gives

$$\begin{aligned} w_t & \leq L_K \|\xi_0\|_{L^\infty} T \gamma \left(\sup_{s \in [0, T]} \int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx \right) \\ & + L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma(w_s) ds. \end{aligned}$$

Again by a comparison theorem (recall the definition of z in (2.11)), we get that

$$\begin{aligned} & \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq z^{L_K \|\xi_0\|_{L^\infty}} \left(t, L_K \|\xi_0\|_{L^\infty} T \gamma \left(\sup_{s \in [0, T]} \int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx \right) \right) \end{aligned}$$

When $L_K \|\xi_0\|_{L^\infty} T \gamma(\text{dist}(\psi^1, \psi^2)) \leq \exp[1 - 2e^{L_K \|\xi_0\|_{L^\infty} T}]$ (a condition which is verified for $\text{dist}(\psi^1, \psi^2)$ small enough, for fixed $\|\xi_0\|$ and T), the estimate (2.12) gives:

$$\begin{aligned} & \int_{\mathbb{T}^2} |G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq e \left(L_K \|\xi_0\|_{L^\infty} T \gamma \left(\sup_{s \in [0, T]} \int_{\mathbb{T}^2} |\psi_s^1(x) - \psi_s^2(x)| dx \right) \right)^{\exp[-L_K \|\xi_0\|_{L^\infty} t]} \end{aligned}$$

From this and the continuity of γ , we see that G is continuous on M_T . The Hölder continuity of G follows from the fact that γ is Hölder continuous. \square

We are ready to prove:

Theorem 3.4. *There exists a unique solution in M_T to equation (3.1), which is a flow Φ of measure-preserving homeomorphisms of class C^α in space and Lipschitz in time.*

Proof. First step. First we prove the existence and the uniqueness on an interval $[0, T_1]$, for T_1 small enough. For the existence, we define the approximating sequence for the solution to problem (3.1). Choose $\psi_t^0 = I$. For any n , put $\psi^{n+1} = G(\psi^n)$ (G being defined on M_{T_1}) and denote $\rho_t^n = \sup_{k \geq n} \int_{\mathbb{T}^2} |\psi_t^{k+1}(x) - \psi_t^k(x)| dx$. [The reason for the supremum in $k \geq n$ is to have, in the formula (3.6) below, ρ^n on the left hand side and ρ^{n-1} on the right hand side: otherwise it seems difficult to have good estimates.] The estimate (3.5) in Lemma 3.2 gives immediately that for all $n \in \mathbb{N}$

$$\rho_t^n \leq 2L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \rho_s^{n-1} ds + 2L_K \|\xi_0\|_{L^\infty} t \varepsilon. \quad (3.6)$$

By Lemma A.1 we infer that

$$\sup_{[0, T_1]} \rho_t^n \leq \frac{(2eL_K \|\xi_0\|_{L^\infty} T_1)^n}{\sqrt{2\pi n}} \sup_{[0, T_1]} \rho_t^0 + 2L_K \|\xi_0\|_{L^\infty} T_1 \exp[n(2L_K \|\xi_0\|_{L^\infty} T_1 - 1)] \quad (3.7)$$

and so, provided $\alpha := 2eL_K \|\xi_0\|_{L^\infty} T_1 < 1$, there exists a unique $\psi \in M_{T_1}$ such that the sequence $(\psi^n)_n$ converges in M_{T_1} to ψ . By Corollary 3.3 it follows that $G(\psi) = \psi$.

The uniqueness follows by applying the previous iterative scheme to two solutions Φ^1, Φ^2 . More precisely we take $\Phi^{i,0} = \Phi^i$ and $\Phi^{i,n+1} = G(\Phi^{i,n})$,

$i = 1, 2$, and we define $\bar{\rho}_t^n = \sup_{k \geq n} \int_{\mathbb{T}^2} |\Phi_t^{1,k}(x) - \Phi_t^{2,k}(x)| dx$. Then (3.6) and so (3.7) hold for the sequence $\bar{\rho}^n$ (in place of ρ^n). But, since Φ^1, Φ^2 are solutions and hence fixed points of G , $\Phi^{i,n} = \Phi^i$ and $\bar{\rho}_t^n = \bar{\rho}_t^0$ for every n , $i = 1, 2$ and so we get

$$\text{dist}(\Phi^1, \Phi^2) = \sup_{[0, T_1]} \bar{\rho}_t^n \leq \alpha^n \sup_{[0, T_1]} \bar{\rho}_t^0 + e^{-1} \alpha e^{-n(1-\alpha)} = \alpha^n \text{dist}(\Phi^1, \Phi^2) + \alpha e^{-n(1-\alpha)},$$

for any integer n . Since $\alpha < 1$, taking n large, we get $\text{dist}(\Phi^1, \Phi^2) = 0$.

Second step. We prove the global existence and uniqueness. They follow essentially by iteration in time, but we prefer to make this argument explicit, since the non-locality of the drift could create some confusion. The main point is to notice that, for fixed $0 < T' < T$, a flow Φ solves the non-local ODE (3.1) on $[0, T]$ if and only if it solves the non-local ODE on $[0, T']$ and it satisfies, for t in $[T', T]$,

$$\Phi_t(x) = \Phi_{T'}(x) + \int_{T'}^t \int_{\mathbb{T}^2} K(\Phi_s(x) - \Phi_s(y)) \xi_0(y) dy. \quad (3.8)$$

Hence we will prove the global result by showing the existence and the uniqueness for equation (3.8) on $[T_1, 2T_1]$, and then iterating the idea. As before, we define the approximating sequence $(\psi^n)_n$ of maps on $[T_1, 2T_1] \times \mathbb{T}^2$ by imposing

$$\psi_t^n(x) = \Phi_{T_1}(x) + \int_{T_1}^t \int_{\mathbb{T}^2} K(\psi_s^n(x) - \psi_s^{n-1}(y)) \xi_0(y) dy. \quad (3.9)$$

Here a small technical clarification is needed for the existence, the continuity and the measure-preserving property of ψ^n : they cannot be inferred directly from Lemma 3.1, since the initial datum is no more x (we could repeat the argument starting from $\Phi_{T_1}(x)$: this can be done, but at the price of introducing a flow map $\Phi_{T_1, t}$ which we avoid for simplicity). So we prove ψ^n exists continuous and is measure-preserving, by defining ψ^n on the whole interval $[0, 2T_1]$ as $\psi_t^0 = \Phi_t 1_{[0, T_1]} + \Phi_{T_1} 1_{[T_1, 2T_1]}$ and $\psi^n = G(\psi^{n-1})$, the map G relative to the interval $[0, 2T_1]$. In this way ψ^n coincides with Φ on $[0, T_1]$ (in particular it satisfies the condition $\psi_{T_1}^n = \Phi_{T_1}$) and it verifies equation (3.9) on $[T_1, 2T_1]$. The definition of ψ^n (with continuity and measure-preserving property) is now done.

Having the existence and the measure-preserving property, we can repeat the estimates in Lemma 3.2, starting from Φ_{T_1} , with no difference in the

proof; in particular the estimates hold with the same constant and with final time T which is replaced by $T - T_1$. In this way we get the existence on $T_1 \leq t \leq T_1 + T_1 = 2T_1$.

The uniqueness follows again applying the iterative scheme above to two solutions and concluding as in step 1.

Step 3. The regularity and homeomorphism properties hold by Lemma 3.1, since $\Phi = G(\Phi)$ is in the image of G (G now being defined on the whole $[0, T]$). \square

Remark 3.5. *In case ξ_0 is more smooth, more regularity of Φ can be obtained, using the usual iterative scheme: if Φ has some regularity, then u^Φ has more regularity, which implies again an improvement of regularity of Φ , and so on.*

4 The stochastic case

Now we prove the existence and the uniqueness of a stochastic continuous flow solving equation (2.4). Notice that, differently from the classical (linear) case, the drift depends on the whole flow, so Kunita's theory ([29], [30]) is not (at least easily) applicable.

We try to mimic the previous reasoning in the deterministic case. The last part, the iterative procedure from the proof of Theorem 3.4, works in this simple way. First we get a generalized Lemma 3.1 (with Itô formula to treat the modulus of the difference of two flows), then we repeat the scheme and obtain a measurable flow solution to (2.4).

The main difficulty is in the first part, precisely in the generalization of Lemma 3.1 to stochastic continuous flows (remember that we need a continuity property for ω fixed). In order to get rid of the first difficulty, we will apply Kolmogorov test, in the spirit of Kunita's results (see [29], [30]). For this we need some estimates on the linear equation.

4.1 The linear stochastic equation

Consider the following SDE ("linear" problem):

$$dX_t = u_t(X_t) dt + \sum_k \sigma_k(X_t) dW_t^k, \quad (4.1)$$

where u is a random vector field with the following properties: for every x , $(t, \omega) \rightarrow u(t, x, \omega)$ is a progressively measurable process and, for every t, x, y, ω ,

$$u(t, x, \omega) \leq L_{0,u}, \quad (4.2)$$

$$|u(t, x, \omega) - u(t, y, \omega)| \leq L_u \gamma(|x - y|) \quad (4.3)$$

for some $L_{0,u}, L_u$ independent of t, x, y, ω . These properties imply that, if X is a progressively measurable process with values in \mathbb{T}^2 , then $u(t, X_t)$ is progressively measurable too.

Lemma 4.1. *Let X, Y be two solutions of (4.1) starting from x, x' resp.. Then, for any $p \geq 2$, there exists $\delta = \delta(T, L_u, L_\sigma, p)$ such that, if $|x - x'| < \delta$, it holds for some constant $C_{p,T}$, depending only on p and T*

$$E[|X_t - Y_{t'}|^p] \leq e|x - x'|^{p \exp[-(2pL_u + L_\sigma)T]} + C_{p,T}(L_{0,u}^p + L_\sigma^p)|t - t'|^{p/2}. \quad (4.4)$$

Proof. It is enough to prove the formula in the two particular cases $t = t'$ and $x = x'$. Fix $t = t'$. By the Itô formula (applied to $f(x) = |x|^p$), calling $Z = X - Y$, we have

$$\begin{aligned} d[|Z|^p] &= p|Z|^{p-2}Z \cdot (u(X) - u(Y)) dt \\ &+ \left[\sum_k p|Z|^{p-2}|\sigma_k(X) - \sigma_k(Y)|^2 \right] dt \\ &+ \left[\sum_k p(p-2)|Z|^{p-4}|Z \cdot (\sigma_k(X) - \sigma_k(Y))|^2 \right] dt + \\ &+ \sum_k p|Z|^{p-2}Z \cdot (\sigma_k(X) - \sigma_k(Y))dW^k. \end{aligned}$$

We take the expectation and use the Lipschitz continuity of σ_k 's and the log-Lipschitz property of u :

$$\begin{aligned} &p|Z|^{p-1}|u(X) - u(Y)| \\ &\leq pL_u|Z|^p(1 - \log |Z|)1_{|Z| < 1/e} + pL_u|Z|^{p-1}(|Z| + 1/e)1_{|Z| \geq 1/e} \\ &\leq pL_u|Z|^p(1 - \log |Z|^p)1_{|Z| < 1/e} + 2pL_u|Z|^p + 1/e1_{|Z| \geq 1/e}pL_u\gamma(|Z|^p) \\ &\leq 2pL_u\gamma(|Z|^p). \end{aligned}$$

Then

$$E[|Z_t^p|] \leq |x - x'|^p + 2pL_u \int_0^t E[\gamma(|Z_s^p|)]ds + L_\sigma \int_0^t E[|Z_s^p|]ds,$$

from which, using Jensen inequality for the concave function γ and the fact that $r \leq \gamma(r)$, we obtain

$$\begin{aligned} E[|Z_t^p|] &\leq |x - x'|^p + 2pL_u \int_0^t \gamma(E[|Z_s^p|])ds + L_\sigma \int_0^t E[|Z_s^p|]ds \\ &\leq |x - x'|^p + (2pL_u + L_\sigma) \int_0^t \gamma(E[|Z_s^p|])ds. \end{aligned}$$

By a comparison principle, $E[|Z_t^p|] \leq z^{2pL_u + L_\sigma}(t, |x - x'|^p)$ (recall the definition of z in (2.11)). When $|x - x'|$ is small enough (precisely, $< \delta$ for some $\delta(T, L_u, L_\sigma, p)$), we can apply the estimate (2.12) and we get the thesis for $t = t'$.

Now put $x = x'$, $t' < t$. By the boundedness of u and σ_k 's, using the Hölder and the Burkholder inequalities, we get

$$\begin{aligned} E[|X_t - X_{t'}|^p] &\leq 2^{p-1} E\left[\left| \int_{t'}^t u(X_r) dr \right|^p + \left| \sum_k \int_{t'}^t \sigma_k(X_r) dW_r^k \right|^p \right] \\ &\leq C_p(L_{0,u}^p + L_\sigma^{p,T})(|t - t'|^p + |t - t'|^{p/2}). \end{aligned}$$

The proof is complete. \square

This will be enough to get the uniqueness and the continuity, but we still need the existence. For this, we will use a generalization of the previous lemma, exhibiting a Cauchy sequence of solutions of approximating equations. Let ρ be a $C_c^\infty(\mathbb{R}^2)$ function, define $\rho_\varepsilon(x) = \varepsilon^{-2}\rho(\varepsilon^{-1}x)$; consider the standard mollification of u : $u^\varepsilon(t, x, \omega) = u(t, \cdot, \omega) * \rho_\varepsilon(x)$, for $x \in \mathbb{T}^2$ (the convolution must be understood on the whole \mathbb{R}^d , where u is extended by periodicity). Notice that, since by (4.3) the field u is continuous and bounded in x , uniformly with respect to t and ω , we get that $(u^\varepsilon)_\varepsilon$ converges to u uniformly in (t, x, ω) : that is, we can find a continuous function $\theta : [0, 1] \rightarrow [0, +\infty[$, with $\theta(0) = 0$, such that, for every $\varepsilon > 0$, $\delta > 0$,

$$\sup_{[0,T] \times \mathbb{T}^2 \times \Omega} |u^\varepsilon - u^\delta| \leq \theta(|\varepsilon - \delta|). \quad (4.5)$$

Moreover, Corollary 3.3 holds uniformly in ε :

$$\sup_{\varepsilon > 0} |u^\varepsilon(t, x) - u^\varepsilon(t, x')| \leq L_u \gamma(|x - x'|). \quad (4.6)$$

Similarly, we define $\sigma_k^\varepsilon(t, x) := \sigma_k(t, \cdot) * \rho_\varepsilon(x)$; since the σ_k 's are Lipschitz-continuous (more precisely, by Condition 2.3), we get (possibly for another θ , with the same properties as above)

$$\sup_{[0, T] \times \mathbb{T}^2} \sum_k |\sigma_k^\varepsilon - \sigma_k^\delta|^2 \leq \theta(|\varepsilon - \delta|), \quad (4.7)$$

$$\sup_{\varepsilon > 0} \sum_k |\sigma_k^\varepsilon(t, x) - \sigma_k^\varepsilon(t, x')|^2 \leq L_\sigma^2 |x - x'|^2. \quad (4.8)$$

Lemma 4.2. *For any $\varepsilon > 0$, let ψ^ε be the stochastic continuous flow solution to*

$$dX_t^\varepsilon = u_t^\varepsilon(X_t^\varepsilon) dt + \sum_k \sigma_k^\varepsilon(X_t^\varepsilon) dW_t^k. \quad (4.9)$$

Then, for any $p \geq 2$, for every ε, δ close enough to 0, for every x, x' in \mathbb{T}^2 with $|x - x'|$ small enough, it holds

$$\sup_{[0, T]} E[|\psi_t^\varepsilon(x) - \psi_t^\delta(x')|^p] \leq C(|x - x'|^p + C\theta(\varepsilon - \delta))^{\exp[-Ct]}$$

for some $C > 0$ (independent of $\varepsilon, \delta, x, x'$). In particular, $(\psi^\varepsilon)_\varepsilon$ is a Cauchy sequence in $C([0, T] \times \mathbb{T}^2; L^p(\Omega))$.

For the sake of simplicity, we do not specify, in the result above and in the proof below, the constants involved (using the letter C for all of them), since the estimates will not be used in the proof of the main result.

Remark 4.3. *For every $\varepsilon > 0$, for every initial datum, equation (4.9) has a unique solution, which can be represented by a stochastic continuous flow ψ^ε of C^1 maps. Indeed, by the boundedness of u , the C^1 norm of u^ε is uniformly bounded, and Kunita's theory applies. Notice that here we need Kunita's result with a stochastic drift, namely [30], Theorem 4.6.5.*

Remark 4.4. *Again for $\varepsilon > 0$, since the stochastic integral is of Stratonovich type (which we have written in Itô form), usual calculus rules give the standard equation for the Jacobian, which depends only on the divergence of the vector fields. Since u^ε and σ_k^ε 's are divergence free, the Jacobian turns out to be constant and so the stochastic flow is measure-preserving.*

Proof. First we notice that, for $p \geq 2$, $E[|\psi_t^\varepsilon(x)|^p]$ is bounded by a constant independent of ε, t and x (simply estimate $|u^\varepsilon(X^\varepsilon)|$ and $|\sigma_k(X^\varepsilon)|$ with the

sup-norms of u and σ_k and use Hölder and Burkholder inequalities). Similarly, one sees that ψ^ε is in $C([0, T] \times \mathbb{T}^2; L^p(\Omega))$ for every $\varepsilon > 0$. By Itô formula (applied to $f(x) = |x|^p$), calling $Z = \psi_t^\varepsilon(y) - \psi_t^\delta(x)$, we have

$$\begin{aligned} d[|Z|^p] &= p|Z|^{p-2}Z \cdot (u^\varepsilon(\psi^\varepsilon(x)) - u^\delta(\psi^\delta(x'))) dt \\ &+ \left[\sum_k p|Z|^{p-2}|\sigma_k^\varepsilon(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\delta(x'))|^2 \right. \\ &+ \left. \sum_k p(p-2)|Z|^{p-4}|Z \cdot (\sigma_k^\varepsilon(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\delta(x')))|^2 \right] dt \\ &+ \sum_k p|Z|^{p-2}Z \cdot (\sigma_k^\varepsilon(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\delta(x')))dW^k. \end{aligned}$$

The difficult term is $u^\varepsilon(\psi^\varepsilon(x)) - u^\delta(\psi^\delta(x'))$. For this, by (4.5) and (4.6), we have

$$\begin{aligned} &|u^\varepsilon(\psi^\varepsilon(x)) - u^\delta(\psi^\delta(x'))| \\ &\leq |u^\varepsilon(\psi^\varepsilon(x)) - u^\delta(\psi^\varepsilon(x))| + |u^\delta(\psi^\varepsilon(x)) - u^\delta(\psi^\delta(x'))| \\ &\leq \theta(\varepsilon - \delta) + C\gamma(|Z|). \end{aligned}$$

The terms with σ_k^ε are easier: by (4.7) and (4.8), we have

$$\begin{aligned} &\sum_k |\sigma_k^\varepsilon(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\delta(x'))|^2 \\ &\leq 2 \sum_k [|\sigma_k^\varepsilon(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\varepsilon(x))|^2 + |\sigma_k^\delta(\psi^\varepsilon(x)) - \sigma_k^\delta(\psi^\delta(x'))|^2] \\ &\leq \theta(\varepsilon - \delta) + C|Z|^2. \end{aligned}$$

So, proceeding as before, using concavity of γ and uniform boundedness of $E[|Z|^{p-1}]$ and $E[|Z|^{p-2}]$, we get

$$E[|Z|_t^p] \leq |x - x'|^p + C\theta(|\varepsilon - \delta|) + \int_0^t \gamma(E[|Z|_s^p])ds.$$

We conclude that, if $|x - x'|^p + C\theta(|\varepsilon - \delta|)$ is small enough (precisely, smaller than a constant depending on T , C and p),

$$\sup_{[0, T]} E[|Z|_t^p] \leq C \left(|x - x'|^p + C\theta(|\varepsilon - \delta|) \right)^{\exp[-Ct]},$$

which implies that, if $x = x'$, the sequence $(\psi^\varepsilon)_\varepsilon$ is Cauchy in the space $C([0, T] \times \mathbb{T}^2; L^p(\Omega))$. \square

Lemma 4.5. *Equation (4.1) has a unique solution, for every deterministic initial datum. This solution is described by a (unique) stochastic measure-preserving continuous flow ψ of class C^α in space, for some $\alpha > 0$, and C^β in time, for every $\beta < 1/2$.*

Proof. By the previous Lemma, for every x , there exists the limit, in $C([0, T]; L^p(\Omega))$, X of the approximating processes $X^\varepsilon = \psi^\varepsilon(x)$'s. Then we can pass to the limit in equation (4.9), because the coefficients are continuous bounded. Hence we infer that the process X is progressively measurable and solves (4.1). The uniqueness follows from Lemma 4.1, with $x = y$. The Hölder continuity property is a consequence of the Kolmogorov criterion, applied again to (4.4). Indeed we get that ψ is α -Hölder continuous in space, for every $\alpha < e^{-CT} - 2/p$, and β -Hölder continuous in time, for every $\beta < 1/2 - 1/p$, so for every $\beta < 1/2$.

As for the measure-preserving property, we will prove that, for every bounded measurable $F : \Omega \rightarrow \mathbb{R}$, every bounded measurable $h : [0, T] \rightarrow \mathbb{R}$ and every continuous bounded $g : \mathbb{T}^2 \rightarrow \mathbb{R}$,

$$\int_0^T h(t) E \left[F \int_{\mathbb{T}^2} g(\psi_t(x)) dx \right] dt = \int_0^T h(t) E \left[F \int_{\mathbb{T}^2} g(x) dx \right] dt. \quad (4.10)$$

This will prove that, for a.e. (t, ω) , $\psi_t(\omega)$ is measure-preserving. By continuity in (t, x) at ω fixed, this implies easily that, for a.e. ω , $\psi_t(\omega)$ is measure-preserving for every t . Since the approximating flows ψ^ε 's are measure-preserving (remember Remark 4.4), equality (4.10) holds for the ψ^ε 's. By the convergence in L^p , we can find a subsequence ψ^{ε_n} such that $(\psi^{\varepsilon_n})_n$ converges to ψ for a.e. (t, x, ω) . Passing to the limit along this subsequence (using dominated convergence theorem), we get (4.10) for ψ . The proof is complete. \square

Remark 4.6. *With a small effort, one could also show the injectivity of $\psi_t(\omega)$ for all t , for a.e. ω (essentially, one has to extend Lemma 4.1 to negative p and use Kolmogorov criterion for $|\psi_t(x) - \psi_t(y)|^{-1}$). Surjectivity and continuity of the inverse map follow from the continuity and the measure-preserving property. The range of a measure-preserving continuous map is a compact set, whose complement (an open set) is Lebesgue-negligible. Thus this range must be the whole \mathbb{T}^2 . Thus the flow is actually a flow of homeomorphisms.*

Corollary 4.7. *Let ξ be an element of $L^\infty([0, T] \times \mathbb{T}^2 \times \Omega)$. Then equation (4.1) with $u = u^\xi$ has a unique solution, for every deterministic initial datum, which enjoys the properties in Lemma 4.5.*

4.2 Stochastic Euler flows

The rest of the section goes on in analogy with the deterministic case. We define a space

$$SM_T = \left\{ \psi : [0, T] \times \mathbb{T}^2 \times \Omega \rightarrow \mathbb{T}^2 : \psi \text{ measurable w.r.t. } \mathcal{P} \times (B)(\mathbb{T}^2), \right. \\ \left. \sup_{[0, T]} \int_{\mathbb{T}^2} E[|\psi_t(x)|] dx < +\infty, \psi_t \text{ meas.-pres. for a.e. } (t, \omega) \right\}.$$

Here \mathcal{P} is the predictable σ -algebra associated with the filtration $(\mathcal{F}_t)_t$. It is a complete metric space, endowed with the distance $dist(\psi^1, \psi^2) = \sup_{[0, T]} \int_{\mathbb{T}^2} E|\psi_t^1(x) - \psi_t^2(x)| dx$. For a given measure-preserving stochastic flow ψ in SM_T , we call $G(\psi)$ the unique solution to the SDE (4.1) with $u = u^\psi$. Recall again that

$$u^\psi(t, x) = \int_{\mathbb{T}^2} K(x - \psi_t(y)) \xi_0(y) dy.$$

enjoys the log-Lipschitz property (2.10) and it is also progressively measurable as required in the previous section, so that G takes values in SM_T .

Remark 4.8. *One may ask at this point why, in the definition on SM_T , we have the supremum in time outside the expectation and not inside (while Burkholder inequality allows supremum inside, in some cases). The reason is that the argument works with the supremum outside and putting the supremum inside could create additional difficulties. A posteriori, since the flow solution Φ to (2.4) is in the image of G , it is continuous and also Hölder continuous.*

Lemma 4.9. *For every $\varepsilon > 0$ (small enough), for every ψ^1, ψ^2 flows in*

SM_T , the following estimates hold:

$$\begin{aligned}
& \int_{\mathbb{T}^2} E|G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\
& \leq L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} E|\psi_s^1(x) - \psi_s^2(x)| dx \right) ds \\
& + L_K \|\xi_0\|_{L^\infty} \int_0^t \gamma \left(\int_{\mathbb{T}^2} E|G(\psi^1)_s(x) - G(\psi^2)_s(x)| dx \right) ds, \\
& + 2L_\sigma^2 \int_0^t \int_{\mathbb{T}^2} E|G(\psi^1)_s(x) - G(\psi^2)_s(x)| dx ds, \\
& \int_{\mathbb{T}^2} E|G(\psi^1)_t - G(\psi^2)_t| dx \\
& \leq (L_K \|\xi_0\|_{L^\infty} + 2L_\sigma^2)(-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} E|G(\psi^1)_s - G(\psi^2)_s| dx ds \\
& + L_K \|\xi_0\|_{L^\infty} (-\log \varepsilon) \int_0^t \int_{\mathbb{T}^2} E|\psi_s^1 - \psi_s^2| dx ds + 2L_K \|\xi_0\|_{L^\infty} t \varepsilon.
\end{aligned}$$

Proof. We would like to apply Itô formula to the modulus function and get an estimate for $|G(\psi^1)_t(x) - G(\psi^2)_t(x)|$. Since the modulus is not C^2 , we use the approximate functions $f_\delta(x) = (|x|^2 + \delta)^{1/2}$, for $\delta > 0$. Calling $Z = G(\psi^1)_t(x) - G(\psi^2)_t(x)$, we have

$$\begin{aligned}
d[f_\delta(Z)] &= f_\delta(Z)^{-1} Z \cdot [u^{\psi^1}(G(\psi^1)) - u^{\psi^2}(G(\psi^2))] dt \\
& + \sum_k f_\delta(Z)^{-1} |\sigma_k(G(\psi^1)) - \sigma_k(G(\psi^2))|^2 dt \\
& + \sum_k f_\delta(Z)^{-3} [(G(\psi^1) - G(\psi^2)) \cdot (\sigma_k(G(\psi^1)) - \sigma_k(G(\psi^2)))]^2 dt \\
& + \sum_k f_\delta(Z)^{-1} Z \cdot [\sigma_k(G(\psi^1)) - \sigma_k(G(\psi^2))] dW.
\end{aligned}$$

Taking the expectation and using the Lipschitz property of σ , since $f_\delta(x) \geq |x|$, we get

$$\begin{aligned}
E[|Z_t|] &\leq \int_0^t E[|u_s^{\psi^1}(G(\psi^1)_s(x)) - u_s^{\psi^2}(G(\psi^2)_s(x))|] ds \\
& + 2L_\sigma^2 \int_0^t E[|Z_s|] ds.
\end{aligned}$$

The rest of the proof follows the lines of Lemma 3.2: we estimate $\int_{\mathbb{T}^2} |u_s^{\psi^1}(G(\psi^1)_s(x)) - u_s^{\psi^2}(G(\psi^2)_s(x))| dx$ and use Jensen inequality to pass γ outside the integral in x and outside the expectation. The second inequality is a consequence of the first one. \square

Proof of Theorem 2.14. Similar to the proof of Theorem 3.4, we only recall the main passages.

First step. We prove the existence and the uniqueness on an interval $[0, T_1]$, with T_1 small enough (but deterministic). The iteration scheme is completely similar to the one in the deterministic case: we consider $\psi_t^0(x) = x$, $\psi^{n+1} = G(\psi^n)$, $\rho_t^n = \sup_{k \geq n} \int_{\mathbb{T}^2} E|\psi_t^{k+1}(x) - \psi_t^k(x)| dx$ and proceed as in the deterministic case, getting a limit flow Φ in SM_{T_1} , for T_1 such that $\alpha := 2e(L_K \|\xi_0\|_{L^\infty} + L_\sigma^2)T_1 < 1$ (notice that T_1 is independent of ω , since all the estimates are in expectation). Such a flow solves (2.4), because G is continuous in SM_T : indeed, from Lemma 4.9 again by comparison with z

$$\begin{aligned} & \int_{\mathbb{T}^2} E|G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq z^{L_K \|\xi_0\|_{L^\infty} + 2L_\sigma^2} \left(t, L_K \|\xi_0\|_{L^\infty} T \gamma \left(\sup_{s \in [0, T]} \int_{\mathbb{T}^2} E|\psi_s^1(x) - \psi_s^2(x)| dx \right) \right) \end{aligned}$$

and, if $dist(\psi^1, \psi^2)$ is small enough,

$$\begin{aligned} & \int_{\mathbb{T}^2} E|G(\psi^1)_t(x) - G(\psi^2)_t(x)| dx \\ & \leq e \left(L_K \|\xi_0\|_{L^\infty} T \gamma \left(\sup_{s \in [0, T]} \int_{\mathbb{T}^2} E|\psi_s^1(x) - \psi_s^2(x)| dx \right) \right)^{\exp[-(L_K \|\xi_0\|_{L^\infty} + 2L_\sigma^2)t]}. \end{aligned}$$

The uniqueness on $[0, T_1]$ is also proved in the same way of the deterministic case.

Second step. We prove the global existence and uniqueness. For this, as in the deterministic case, we solve the equation on $[T_1, 2T_1]$

$$\Phi_t(x) = \Phi_{T_1}(x) + \int_{T_1}^t \int_{\mathbb{T}^2} K(\Phi_s(x) - \Phi_s(y)) \xi_0(y) dy + \sum_k \int_{T_1}^t \sigma_k(\Phi_s(x)) dW_s^k.$$

To get the existence for this equation, we define the approximating sequence $(\psi^n)_n$ of maps on $[T_1, 2T_1] \times \mathbb{T}^2$ by imposing

$$\psi_t^n(x) = \Phi_{T_1}(x) + \int_{T_1}^t \int_{\mathbb{T}^2} K(\psi_s^n(x) - \psi_s^{n-1}(y)) \xi_0(y) dy + \sum_k \int_{T_1}^t \sigma_k(\psi_s^{n-1}(x)) dW_s^k \quad (4.11)$$

The existence, the continuity and the measure-preserving property for equation (4.11) are again not a direct consequence of Lemmata 4.1 and 4.2, since here we start from Φ_{T_1} and not from the identity; here we also have the problem of the randomness of Φ_{T_1} , which brings us to consider the strategy in the deterministic case. Following that strategy of the deterministic case, we can build ψ^n and prove the continuity and the measure-preserving property. Then we apply the previous estimates, again with no change in the constants and with the final time T replaced by $T - T_1$ (again deterministic). This allows to conclude the existence on $[T_1, 2T_1]$. The uniqueness on this interval is as in the step 1.

Third step. The regularity properties hold by Lemma 4.1, since $\Phi = G(\Phi)$ is in the image of G (G now being defined on the whole $[0, T]$). \square

5 The stochastic Euler vorticity equation

In this section we will prove Theorem 2.10. First we need the existence of solutions to the stochastic Euler vorticity equation (2.1).

Proposition 5.1. *Let Φ be a solution to (2.4). For $t \geq 0$, define $\xi_t = (\Phi_t)_\# \xi_0$. Then ξ has a density (still denoted by ξ) in $L^\infty([0, T] \times \mathbb{T}^2 \times \Omega)$, which is a distributional L^∞ solution to the stochastic Euler equation (2.1).*

Proof. Fix $t > 0$ and the probabilistic datum ω (omitted in the sequel). By Lemma 2.19, since Φ_t is measure preserving, ξ_t is absolutely continuous with respect to the Lebesgue measure on \mathbb{T}^2 and $\|\xi_t\|_{L^\infty} \leq \|\xi_0\|_{L^\infty}$.

Let φ be a test function, Itô formula applied to $\varphi(\Phi_t)$ gives

$$\begin{aligned} d[\varphi(\Phi_t)] &= u_t^\Phi(\Phi_t) \cdot \nabla \varphi(\Phi_t) dt + \sum_k \sigma_k(\Phi_t) \cdot \nabla \varphi(\Phi_t) dW_r^k \\ &\quad + \frac{1}{2} \text{tr}[a(\Phi_t) D^2 \varphi(\Phi_t)] dt. \end{aligned}$$

Now notice that, by definition of ξ_t , $u_t^\Phi = K * \xi_t$; so, integrating in $\xi_0 dx$, we get (2.2). \square

For the proof of the uniqueness, we will adapt a classical argument for the transport equation. We first recall the idea in the case $\sigma_k \equiv 0$ for simplicity. A formal application of the chain rule gives

$$\frac{d}{dt} \xi_t(\Phi_t) = \partial_t \xi_t(\Phi_t) + D \xi_t(\Phi_t) \frac{d\Phi_t}{dt} = (\partial_t \xi_t + u_t \cdot \nabla \xi_t)(\Phi_t) = 0.$$

This implies that $\xi_t(\Phi_t) = \xi_0$, so that $\xi_t = \xi_0(\Phi_t^{-1})$ is completely determined by the flow. But we have used the chain rule for an object (ξ_t) which is not regular in general (and in fact there are counterexamples for irregular drifts). Thus we need to regularize ξ . This regularization ξ^ε solves a transport-type equation with an additional term, a commutator, which we need to control to conclude the argument. We use for this the argument in [17], [1], [2], where the commutator is an essential tool for the uniqueness of the transport equation.

First we need approximate identities. For this, let ρ be a $C^\infty(\mathbb{R}^2)$ nonnegative even function, with support in $[-1/2, 1/2]^2$ and $\int_{\mathbb{R}^2} \rho dx = 1$. For $\varepsilon > 0$, define $\rho_\varepsilon(x) = \varepsilon^{-2} \rho(x/\varepsilon)$. If f is an integrable function on \mathbb{T}^2 , f can be extended periodically to a locally integrable function on the whole \mathbb{R}^2 , so that the convolution $\rho_\varepsilon * f$ makes sense and is still a C^∞ periodic function.

For a vector field v and a function w on the torus, we define formally the commutator as

$$[v \cdot \nabla, \rho_\varepsilon *]w := v \cdot \nabla(\rho_\varepsilon * w) - \rho_\varepsilon * (v \cdot \nabla w). \quad (5.1)$$

Suppose that v and w are integrable and v is divergence free. Then the expression above defines a measurable function on \mathbb{T}^2 . Indeed, the following equalities hold in distribution (the functions being thought as extended to the whole \mathbb{R}^2):

$$\rho_\varepsilon * (v \cdot \nabla w) = \rho_\varepsilon * \operatorname{div}(vw) = - \int_{\mathbb{R}^2} \nabla \rho_\varepsilon(z) \cdot v(\cdot - z) w(\cdot - z) dz. \quad (5.2)$$

Besides, by (5.1) and (5.2), the commutator reads

$$[v \cdot \nabla, \rho_\varepsilon *]w(x) = \int_{\mathbb{R}^2} (v(x) - v(x - z)) \cdot \nabla \rho_\varepsilon(z) w(x - z) dz.$$

With the change of variable $y = z/\varepsilon$, $x' = x'_\varepsilon = x - \varepsilon y$ we get

$$[v \cdot \nabla, \rho_\varepsilon *]w(x) = \int_{\mathbb{R}^2} \frac{v(x' + \varepsilon y) - v(x')}{\varepsilon} \cdot \nabla \rho(y) w(x') dy.$$

If v is in $W^{1,1}(\mathbb{T}^2)$, then, for every y in \mathbb{R}^2 , for a.e. x' in \mathbb{T}^2 , $v(x' + \varepsilon y) - v(x') = \varepsilon \int_0^1 Dv(x' + \xi \varepsilon y) y d\xi$. Indeed, this is true for $v^\delta = \rho_\delta * v$ and, for fixed y , $v^\delta(x' + \varepsilon y) - v^\delta(x') - \varepsilon \int_0^1 Dv^\delta(x' + \xi \varepsilon y) y d\xi$, as function of x' , converges to

0 a.e. as $\delta \rightarrow 0$ (possibly passing to a subsequence). So, in this case, the commutator has the following expression:

$$[v \cdot \nabla, \rho_\varepsilon *]w(x') = \int_{\mathbb{R}^2} \int_0^1 Dv(x' + \xi \varepsilon y) y d\xi \cdot \nabla \rho(y) w(x') dy. \quad (5.3)$$

Lemma 5.2 (Commutator lemma). *Let p be in $[1, +\infty[$, let v be in $W^{1,p}(\mathbb{T}^2)$ with zero divergence, let w be in $L^\infty(\mathbb{T}^2)$. Then*

$$\lim_{\varepsilon \rightarrow 0} [v \cdot \nabla, \rho_\varepsilon *]w = 0 \quad \text{in } L^p(\mathbb{T}^2)$$

and we have the inequality

$$\|[v \cdot \nabla, \rho_\varepsilon *]w\|_{L^p(\mathbb{T}^2)} \leq C \|Dv\|_{L^p(\mathbb{T}^2)} \|w\|_{L^\infty(\mathbb{T}^2)}.$$

Proof. The inequality follows integrating in x the p -power of the expression on the LHS of (5.3). Precisely, since ρ is supported on $[-1/2, 1/2]^2$, we have by Hölder inequality (remember $x' = x + \varepsilon y$)

$$\begin{aligned} & \int_{\mathbb{T}^2} |[v \cdot \nabla, \rho_\varepsilon *]w|^p dx \\ & \leq \int_{\mathbb{R}^2} \int_{\mathbb{T}^2} \int_0^1 |Dv(x' + \xi \varepsilon y)|^p d\xi |w(x')|^p dx' |y|^p |\nabla \rho(y)|^p dy \\ & \leq \|Dv\|_{L^p(\mathbb{T}^2)}^p \|w\|_{L^\infty(\mathbb{T}^2)}^p \int_{\mathbb{R}^2} |y|^p |\nabla \rho(y)|^p dy \end{aligned}$$

(the integral in x' should be on $\mathbb{T}^2 - \varepsilon y$, but by periodicity we can integrate on \mathbb{T}^2 as well).

For the limit, it is enough to show that

$$L^p(\mathbb{T}^2)\text{-}\lim_{\varepsilon \rightarrow 0} [v \cdot \nabla, \rho_\varepsilon *]w = w(\cdot) \left(\int_{\mathbb{R}^2} Dv(\cdot) y \cdot \nabla \rho(y) dy \right).$$

Indeed, by the symmetry property of ρ , $\int_{\mathbb{R}^2} y_i \partial_j \rho(y) dy = -C \delta_{ij}$ (where C is independent of i) and so $\int_{\mathbb{R}^2} Dv(x) y \cdot \nabla \rho(y) dy = -C \operatorname{div} w = 0$. By (5.3) we have

$$\begin{aligned} & \int_{\mathbb{T}^2} \left| [v \cdot \nabla, \rho_\varepsilon *]w(x) - w(x) \left(\int_{\mathbb{R}^2} Dv(x) y \cdot \nabla \rho(y) dy \right) \right|^p dx \leq \int_{\mathbb{R}^2} \int_{\mathbb{T}^2} \int_0^1 \\ & |w(x') Dv(x' + \xi \varepsilon y) - w(x' + \varepsilon y) Dv(x' + \varepsilon y)|^p d\xi dx' |y|^p |\nabla \rho(y)|^p dy, \end{aligned}$$

hence it is enough to prove that

$$\int_{\mathbb{T}^2} \int_0^1 |w(x') Dv(x' + \xi \varepsilon y) - w(x' + \varepsilon y) Dv(x' + \varepsilon y)|^p d\xi dx' \rightarrow 0$$

uniformly in y . Using the continuity of translations in L^p for the function wDv , we need only to show that

$$\int_{\mathbb{T}^2} \int_0^1 |w(x') Dv(x' + \xi \varepsilon y) - w(x') Dv(x')|^p d\xi dx' \rightarrow 0.$$

Since w is in L^∞ , this follows from $\int_{\mathbb{T}^2} \int_0^1 |Dv(x' + \xi \varepsilon y) - Dv(x')|^p d\xi dx' \rightarrow 0$, which is again a consequence of continuity of translation in L^p applied to Dv . \square

Proposition 5.3. *Let ξ be a (distributional) L^∞ solution to the stochastic Euler vorticity equation. Let Φ be a measure-preserving stochastic flow, which solves (4.1) with $u = u^\xi$ (it exists by Corollary 4.7). Then $\xi_t = (\Phi_t)_\# \xi_0$.*

Proof. We will prove that $\xi_t(\Phi_t) = \xi_0$ Lebesgue-a.e.. Having this, then, for every measurable bounded function φ on \mathbb{T}^2 , $\langle \xi_t, \varphi \rangle = \langle \xi_t(\Phi_t), \varphi(\Phi_t) \rangle = \langle \xi_0, \varphi(\Phi_t) \rangle$ (in the first equality we used the measure-preserving property) and so $\xi_t = (\Phi_t)_\# \xi_0$.

As mentioned before, we need to consider $\xi_t^\varepsilon = \xi_t * \rho_\varepsilon$ instead of ξ_t . Notice that, for every x , $\xi_t^\varepsilon(x) = \langle \xi_t, \rho_\varepsilon(x - \cdot) \rangle$. So $\xi^\varepsilon(x)$ is a progressively measurable process, with continuous trajectories, and the stochastic Euler vorticity equation, applied to the test function $\rho_\varepsilon(x - \cdot)$, gives the following equality:

$$d\xi^\varepsilon + (u \cdot \nabla \xi) * \rho_\varepsilon dt + \sum_k (\sigma_k \cdot \nabla \xi) * \rho_\varepsilon dW^k - \frac{1}{2} \text{tr}[aD^2 \xi^\varepsilon] dt = 0, \quad (5.4)$$

which also reads

$$\begin{aligned} d\xi^\varepsilon &+ u \cdot \nabla \xi^\varepsilon dt + \sum_k \sigma_k \cdot \nabla \xi^\varepsilon dW^k - \frac{1}{2} \text{tr}[aD^2 \xi^\varepsilon] dt = [u \cdot \nabla, \rho_\varepsilon *] \xi dt \\ &+ \sum_k [\sigma_k \cdot \nabla, \rho_\varepsilon *] \xi dW^k. \end{aligned}$$

Now, by (5.4), since ξ^ε is adapted regular (together with $(u \cdot \nabla \xi) * \rho_\varepsilon$, $\sigma_k \cdot \nabla \xi) * \rho_\varepsilon$, $aD^2 \xi^\varepsilon$), we can apply Itô-Kunita-Wentzell formula (see e.g. Theorem 8.3,

page 188 of [29], with easy modifications for the case of an infinite number of k 's), obtaining for $\xi_t^\varepsilon(\Phi_t)$

$$d\xi_t^\varepsilon(\Phi_t) = [u_t \cdot \nabla, \rho_\varepsilon^*] \xi_t(\Phi_t) dt + \sum_k [\sigma_k \cdot \nabla, \rho_\varepsilon^*] \xi_t(\Phi_t) dW^k.$$

Since Φ is measure-preserving, integrating in space we get

$$\begin{aligned} E\left[\int_{\mathbb{T}^2} |\xi_t^\varepsilon(\Phi_t) - \xi_0| dx\right] &\leq \int_0^t \int_{\mathbb{T}^2} E[|[u_r \cdot \nabla, \rho_\varepsilon^*] \xi_r|] dx dr \\ + \sum_k \int_0^t \int_{\mathbb{T}^2} E[|[\sigma_k \cdot \nabla, \rho_\varepsilon^*] \xi_r|^2]^{1/2} dx dr. \end{aligned}$$

By the Commutator Lemma, for a.e. r and ω in Ω , $\int_{\mathbb{T}^2} |[u_r \cdot \nabla, \rho_\varepsilon^*] \xi_r| dx$ tends to 0 as $\varepsilon \rightarrow 0$. Besides, this term is dominated by

$$C \|Du_r\|_{L^1(\mathbb{T}^2)} \|\xi_r\|_{L^\infty(\mathbb{T}^2)} \leq C' \|\xi\|_{L^\infty([0, T] \times \mathbb{T}^2 \times \xi)}^2.$$

Indeed, for every v in $L^\infty(\mathbb{T}^2)$ and every finite $p \geq 1$, $\|D(K * v)\|_{L^p(\mathbb{T}^2)} \leq C \|D^2(-\Delta)^{-1}v\|_{L^p(\mathbb{T}^2)} \leq C' \|v\|_{L^\infty(\mathbb{T}^2)}$. So dominated convergence theorem gives that

$$\lim_{\varepsilon \rightarrow 0} \int_0^t \int_{\mathbb{T}^2} E[|[u_r \cdot \nabla, \rho_\varepsilon^*] \xi_r|] dx dr = 0.$$

Similarly, for every k , for a.e. r and ω in Ω , $\int_{\mathbb{T}^2} |[\sigma_k r \cdot \nabla, \rho_\varepsilon^*] \xi_r|^2 dx$ tends to 0 as $\varepsilon \rightarrow 0$ and is dominated by

$$C \|D\sigma_k\|_{L^2(\mathbb{T}^2)}^2 \|\xi_r\|_{L^\infty(\mathbb{T}^2)}^2.$$

Since $\sum_k \|D\sigma_k\|_{L^2(\mathbb{T}^2)}^2 \leq \|\sum_k |D\sigma_k|^2\|_{L^\infty(\mathbb{T}^2)} < +\infty$ by hypothesis, then we have (again by dominated convergence theorem)

$$\lim_{\varepsilon \rightarrow 0} \sum_k \int_0^t \int_{\mathbb{T}^2} E[|[\sigma_k \cdot \nabla, \rho_\varepsilon^*] \xi_r|^2] dx dr = 0.$$

Thus, for any fixed $t > 0$, $\xi_t^\varepsilon(\Phi_t)$ tends to ξ_0 in $L^1(\mathbb{T}^2 \times \Omega)$ as $\varepsilon \rightarrow 0$. Since ξ_t^ε converges to ξ_t in $L^1(\mathbb{T}^2 \times \Omega)$ (the convergence in $L^1(\mathbb{T}^2)$ being dominated by $\|\xi\|_{L^\infty}$) and Φ_t is measure-preserving, $\xi_t^\varepsilon(\Phi_t)$ converges to $\xi_t(\Phi_t)$ in $L^1(\mathbb{T}^2 \times \Omega)$ and thus $\xi_t(\Phi_t) = \xi_0$, which is our thesis. \square

Corollary 5.4. *The uniqueness for the stochastic Euler vorticity equation (in the class of L^∞ solutions) holds.*

Proof. The above Proposition 5.3 tells that a solution ξ to the stochastic Euler vorticity equation is completely determined by the associated flow Φ which solves (4.1) with $u = u^\xi$; again for the proposition, $u = u^\Phi$ and so Φ solves (2.4). Thus the uniqueness for (2.4) implies the uniqueness for the stochastic Euler vorticity equation. \square

This concludes the proof of Theorem 2.10.

6 Stability

In this section we want to prove a stability result, both at Lagrangian and Eulerian points of view, when the kernel K is regularized.

Precisely, take a family $(\rho_\varepsilon)_\varepsilon$ of even compactly supported resolutions of identity and define $K^\varepsilon := K * \rho_\varepsilon$. Consider the approximated non-local ODE

$$\Phi_t^\varepsilon(x) = x + \int_0^t \int_{\mathbb{T}^2} K^\varepsilon(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) \xi_0(y) dy + \sum_{k=1}^{\infty} \int_0^t \int_{\mathbb{T}^2} \sigma_k(\Phi_r^\varepsilon(x)) dW_r^k \quad (6.1)$$

and the approximated stochastic Euler vorticity equation

$$d\xi^\varepsilon + u^{\varepsilon, \xi^\varepsilon} \cdot \nabla \xi^\varepsilon dt + \sum_k \sigma_k \cdot \nabla \xi^\varepsilon dW^k = \frac{1}{2} C \Delta \xi^\varepsilon, \quad (6.2)$$

where $u^{\varepsilon, \xi^\varepsilon} := K^\varepsilon * \xi^\varepsilon$.

One can repeat all the previous definitions and arguments with K^ε in place of K , to get the analogues of Theorem 2.14 and Theorem 2.10: there exists a unique measure-preserving stochastic continuous flow Φ solving (6.1), which is also C^α in space, for every $\alpha < 1$ and C^β in time, for every $\beta < 1/2$; there exists a unique L^∞ distributional solution ξ^ε for (6.2). Moreover it holds

$$\xi_t^\varepsilon = (\Phi_t^\varepsilon)_\# \xi_0. \quad (6.3)$$

The first stability result is for flows:

Proposition 6.1. *The family $(\Phi^\varepsilon)_\varepsilon$ converges to Φ (as $\varepsilon \rightarrow 0$) in $C([0, T]; L^1(\mathbb{T}^2 \times \Omega))$.*

Proof. The fact that Φ^ε and Φ belong to $C([0, T]; L^1(\mathbb{T}^2 \times \Omega))$ can be proved easily, using similar techniques to those below. For the convergence, call $Z_t^\varepsilon(x) = \Phi_t^\varepsilon(x) - \Phi_t(x)$. As in the proof of Lemma 4.9, we would like to apply Itô formula for $|Z^\varepsilon|$. Proceeding as in that proof (applying Itô formula to $f_\delta(x) = (|x|^2 + \delta)^{1/2}$), we get

$$\begin{aligned} E|Z_t^\varepsilon(x)| &\leq \int_0^t \int_{\mathbb{T}^2} E |K^\varepsilon(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r(x) - \Phi_r(y))| |\xi_0(y)| dy dr \\ &\quad + 2L_\sigma^2 \int_0^t E|Z_r^\varepsilon(x)| dr. \end{aligned}$$

Integrating this inequality in x , since ξ_0 is bounded, we obtain

$$\begin{aligned} &\int_{\mathbb{T}^2} E|Z_t^\varepsilon(x)| dx \\ &\leq \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K^\varepsilon(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r(x) - \Phi_r(y))| dx dy dr \\ &\quad + 2L_\sigma^2 \int_0^t \int_{\mathbb{T}^2} E|Z_r^\varepsilon(x)| dx dr \\ &\leq \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K^\varepsilon(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y))| dx dy dr \\ &\quad + \|\xi_0\|_{L^\infty} \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r(x) - \Phi_r(y))| dx dy dr \\ &\quad + 2L_\sigma^2 \int_0^t \int_{\mathbb{T}^2} E|Z_r^\varepsilon(x)| dx dr. \end{aligned} \tag{6.4}$$

For the first integral of (6.4), we exploit the fact that Φ^ε is measure-preserving, for every ε ; so we have

$$\begin{aligned} &\int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K^\varepsilon(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y))| dx dy dr \\ &= \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K^\varepsilon(x - y) - K(x - y)| dx dy dr \\ &\leq T \int_{\mathbb{T}^2} |K^\varepsilon(x') - K(x')| dx', \end{aligned}$$

where we have used, in the last passage, the change of variable $x - y = x'$, $x + y = y'$ (this implies a change of domain, but the L^1 norm of $K^\varepsilon(x') - K(x')$

on the new domain is comparable with that on the torus). For the second integral of (6.4), we exploit the log-Lipschitz property of K (estimate (2.8)) and get

$$\begin{aligned}
& \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E |K(\Phi_r^\varepsilon(x) - \Phi_r^\varepsilon(y)) - K(\Phi_r(x) - \Phi_r(y))| \, dx dy dr \\
& \leq L_K \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E \gamma(|Z_r^\varepsilon(x) - Z_r^\varepsilon(y)|) \, dx dy dr \\
& \leq L_K \int_0^t \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} E [\gamma(|Z_r^\varepsilon(x)|) + \gamma(|Z_r^\varepsilon(y)|)] \, dx dy dr \\
& \leq 2L_K \int_0^t \int_{\mathbb{T}^2} \gamma(E|Z_r^\varepsilon(x)|) \, dx dr,
\end{aligned}$$

where we have used the sub-additivity of γ ($\gamma(|x+y|) \leq \gamma(|x|) + \gamma(|y|)$, as it can be easily checked) and Jensen inequality. Putting all together, we have

$$\int_{\mathbb{T}^2} E|Z_t^\varepsilon(x)| \, dx \leq T \|K^\varepsilon - K\|_{L^1(\mathbb{T}^2)} + (2L_K \|\xi_0\|_{L^\infty} + 2L_\sigma^2) \int_0^t \int_{\mathbb{T}^2} \gamma(E|Z_r^\varepsilon(x)|) \, dx dr.$$

Again by comparison, we get $\int_{\mathbb{T}^2} E|Z_t^\varepsilon(x)| \, dx \leq z^{2L_K \|\xi_0\|_{L^\infty} + 2L_\sigma^2}(t, T \|K^\varepsilon - K\|_{L^1(\mathbb{T}^2)})$, where z is defined as in (2.11). Since K is in $L^1(\mathbb{T}^2)$, $\|K^\varepsilon - K\|_{L^1(\mathbb{T}^2)}$ tends to 0 (as $\varepsilon \rightarrow 0$), so

$$\sup_{t \in [0, T]} \int_{\mathbb{T}^2} E|Z_t^\varepsilon(x)| \, dx \leq \sup_{t \in [0, T]} z(t, T \|K^\varepsilon - K\|_{L^1(\mathbb{T}^2)}) \rightarrow 0.$$

The proof is complete. □

Here is the result for the vorticity:

Proposition 6.2. *The family $(\xi^\varepsilon)_\varepsilon$ converges weakly to ξ (as $\varepsilon \rightarrow 0$), in the following sense. For every φ in $C_b(\mathbb{T}^2)$,*

$$E \left| \int_{\mathbb{T}^2} \varphi \xi_t^\varepsilon \, dx - \int_{\mathbb{T}^2} \varphi \xi_t \, dx \right| \rightarrow 0$$

for every t and in $L^p([0, T])$, for any $p \in [1, \infty)$.

Proof. First, notice that, by (6.3),

$$\int_{\mathbb{T}^2} \varphi \xi_t^\varepsilon dx = \int_{\mathbb{T}^2} \varphi(\Phi_t^\varepsilon) \xi_0 dx$$

and the same without ε . In particular, $\varphi(\Phi_t^\varepsilon) \xi_0$ is dominated a.e. by a constant. Now fix the time t . We use here a classical argument in measure theory. Suppose by contradiction that there exist $\delta > 0$ and a sequence $\varepsilon_n \rightarrow 0$ such that

$$E \left| \int_{\mathbb{T}^2} \varphi(\Phi_t^{\varepsilon_n}) \xi_0 dx - \int_{\mathbb{T}^2} \varphi(\Phi_t) \xi_0 dx \right| \geq \delta. \quad (6.5)$$

The previous proposition gives that $\Phi_t^{\varepsilon_n}$ converges to Φ_t in $L^1(\mathbb{T}^2 \times \Omega)$. So we have for a subsequence ε_{n_k} that $\Phi_t^{\varepsilon_{n_k}}$ tends to Φ_t for a.e. (x, ω) and similarly for $\varphi(\Phi_t^{\varepsilon_{n_k}})$, since φ is continuous. Hence, by dominated convergence theorem, we get that

$$E \left| \int_{\mathbb{T}^2} \varphi(\Phi_t^{\varepsilon_{n_k}}) \xi_0 dx - \int_{\mathbb{T}^2} \varphi(\Phi_t) \xi_0 dx \right| \rightarrow 0,$$

which contradicts (6.5). We have proved convergence at t fixed. Convergence in $L^p([0, T])$, for any finite p , follows from this result and the Lebesgue Dominated Convergence Theorem. \square

7 An alternative way: reduction to the deterministic case

In this section we will see how to deduce the results in the stochastic case by a suitable transformation, assuming the deterministic case and more regularity for the σ_k 's. As we already said, we will not develop this method in all the details.

At a Lagrangian level (trajectories), consider the SDE with only the stochastic integral, namely

$$d\psi = \sum_k \sigma_k(\psi) \circ dW^k. \quad (7.1)$$

It is well known that, if the fields σ_k 's are regular enough (C^3 should be sufficient, C^2 is assumed in every “classical” result) and divergence-free, then there

exists a stochastic flows ψ of $C^{1,1}(\mathbb{T}^2)$ measure-preserving diffeomorphisms solving (7.1) (a $C^{1,1}$ diffeomorphism is a C^1 map with Lipschitz-continuous derivatives, together with its inverse). The inverse flow ψ_t^{-1} satisfies

$$d\psi_t^{-1}(x) = - \sum_k \sigma_k(x) \cdot \nabla \psi_t^{-1}(x) \circ dW^k.$$

Now let Φ be the Euler stochastic flow (solving (2.4)) and make a change of variable, composing with ψ_t^{-1} : call

$$\tilde{\Phi}(t, x, \omega) = \psi_{t,\omega}^{-1}(\Phi_{t,\omega}(x)). \quad (7.2)$$

Using the Itô-Kunita-Wentzell formula, we obtain the following random ODE for $\tilde{\Phi}$:

$$d\tilde{\Phi}_t = (D\psi_t)^{-1} u_t^\Phi(\psi_t(\tilde{\Phi}_t)) dt,$$

where u^Φ is as in (2.5). This equation reads also as

$$d\tilde{\Phi}_t = \tilde{u}_t^{\tilde{\Phi}}(\tilde{\Phi}_t) dt, \quad (7.3)$$

where

$$\tilde{u}^{\tilde{\Phi}}(t, x, \omega) = (D\psi_{t,\omega}(x))^{-1} \int_{\mathbb{T}^2} K(\psi_{t,\omega}(x) - \psi_{t,\omega}(\tilde{\Phi}_{t,\omega}(y))) \xi_0(y) dy.$$

The equation (7.3) is not (3.1), but the drift $\tilde{u}^{\tilde{\Phi}}$ has the same regularity properties of the drift u^Φ of (3.1), provided ψ is a flow of $C^{1,1}(\mathbb{T}^2)$ diffeomorphisms, since the term $D\psi_t$ appears; here we need σ to be at least C^2 . Thus, one could proceed as follows:

1. first we can repeat the argument in the deterministic part, to get the existence and the uniqueness for $\tilde{\Phi}$ satisfying (7.3); since ψ is a regular flow adapted to the Brownian filtration, this implies the strong existence and the strong uniqueness for Φ itself (plus the homeomorphism property), i.e. Theorem 2.14;
2. then Section 5 applies and we deduce Theorem 2.10.

This can be seen also at an Eulerian level (velocity field). Heuristically, with the change of variable (7.2), we should consider, as new vorticity, $\tilde{\xi}_t =$

$\xi_0(\tilde{\Phi}_t^{-1}) = \xi_t(\psi_t)$. Indeed, let ξ be a solution to (1.1) and let ψ be as above, call

$$\tilde{\xi}(t, x, \omega) = \xi(t, \psi(t, x, \omega), \omega).$$

Applying, this time formally, the Itô-Kunita-Wentzell formula, we obtain the following random PDE for $\tilde{\xi}$:

$$\partial_t \tilde{\xi} + \tilde{u}^{\tilde{\xi}} \cdot \nabla \tilde{\xi} = 0, \quad (7.4)$$

where

$$\tilde{u}^{\tilde{\xi}} = (D\psi_{t,\omega}(x))^{-1} \int_{\mathbb{T}^2} K(\psi_{t,\omega}(x) - \psi_{t,\omega}(y)) \tilde{\xi}_t(y) dy.$$

This fact, as well as its converse (the passage from $\tilde{\xi}$ to ξ), can be made rigorous in the following way. First, we take $\xi^\varepsilon = \xi * \rho_\varepsilon$ (where ρ_ε are even compactly supported mollifiers) and write the equation for ξ^ε (using commutators only for the σ_k 's):

$$\partial_t \xi^\varepsilon + (u^\xi \cdot \nabla \xi)^\varepsilon + \sum_k \sigma_k \cdot \nabla \xi^\varepsilon \circ \dot{W}^k - \sum_k [\sigma_k \cdot \nabla, \rho_\varepsilon *] \xi \circ \dot{W}^k = 0.$$

Then we multiply this equation by $\varphi(\psi^{-1})$, where φ is any regular test function on \mathbb{T}^2 . In this way we obtain (7.4) for $\xi^\varepsilon(\psi)$, with $(u^\xi \cdot \nabla \xi)^\varepsilon(\psi)$ in place of $\tilde{u}^{\tilde{\xi}} \cdot \nabla \tilde{\xi}$ and with the additional commutator term $\sum_k [\sigma_k \cdot \nabla, \rho_\varepsilon *] \xi(\psi) \circ \dot{W}^k$. Finally we let ε go to 0, getting (7.4).

Again (7.4) is not the deterministic Euler vorticity equation ((1.1) with $\sigma = 0$), but its drift $\tilde{u}^{\tilde{\xi}}$ has the same regularity properties of the drift u^ξ of (3.1), provided ψ is a flow $C^{1,1}(\mathbb{T}^2)$ diffeomorphisms. So one can repeat the arguments in the deterministic case (flows and commutator lemma), to get the existence and the uniqueness for the random PDE (7.4), then the strong existence and the strong uniqueness for 1.1 follow immediately.

Finally we mention that the passage between ξ and $\tilde{\xi}$ can be seen at a more abstract level; this is a classical remark, due at least to Lamperti, Doss and Sussmann ([32], [18], [40]). Suppose to have an SPDE of the form

$$d\xi + A(\xi) \xi dt + \sum_k B_k \xi \circ dW^k = 0,$$

where $A(x)$ and B_k are linear operators (for simplicity assume B_k time-independent); in our case, $A(\xi) = u^\xi \cdot \nabla$ and $B_k = \sigma_k \cdot \nabla$. Consider formally

$$\tilde{\xi}_t = e^{\sum_k B_k W_t^k} \xi_t;$$

in our case, this corresponds to the composition $\xi(\psi)$. Then formally $\tilde{\xi}$ satisfies the following random PDE:

$$\partial_t \tilde{\xi} + e^{\sum_k B_k W_t^k} A(e^{-\sum_k B_k W_t^k} \tilde{\xi}) e^{-\sum_k B_k W_t^k} \tilde{\xi} = 0.$$

Thus we have reduced an SPDE to a random PDE, which can be treated through deterministic techniques.

A A useful inequality

This section contains a proof of an auxiliary inequality used in a crucial way twice in our paper.

Lemma A.1. *Assume that $A, B > 0$ and $T > 0$. Suppose that $(\rho_n)_{n=0}^\infty$ is a sequence of continuous nonnegative functions defined on the interval $[0, T]$ such that for every $\varepsilon \in (0, 1)$ and every n ,*

$$\rho_t^n \leq A \log \frac{1}{\varepsilon} \int_0^t \rho_s^{n-1} ds + \varepsilon B t, \quad t \in [0, T]. \quad (\text{A.1})$$

Then

$$\rho_t^n \leq \frac{(At)^n}{\sqrt{2\pi n}} \sup_{s \in [0, t]} |\rho_s^0| + Bt(e^{At-1})^n, \quad t \in [0, T].$$

Proof of Lemma A.1. By Induction one can show that for every $n \in \mathbb{N}^*$ and every $\varepsilon \in (0, 1)$

$$\begin{aligned} \rho_t^n &\leq (A(-\log \varepsilon))^n \int_0^t \dots \int_0^{s_2} \rho_{s_1}^0 ds_1 \dots ds_n \\ &+ B\varepsilon t \sum_{k=0}^{n-1} (A(-\log \varepsilon))^k \int_0^t \dots \int_0^{s_2} ds_1 \dots ds_k \\ &\leq \frac{(A(-\log \varepsilon)t)^n}{n!} \sup_{s \in [0, t]} |\rho_s^0| + B\varepsilon t \sum_{k=0}^{n-1} \frac{(A(-\log \varepsilon)t)^k}{k!}, \quad t \in [0, T] \end{aligned} \quad (\text{A.2})$$

Let us take $n \in \mathbb{N}$. Choose $\varepsilon = e^{-n}$. Then by the above inequality and

Stirling's inequality,

$$\begin{aligned}
\rho_t^n &\leq \frac{(Ant)^n}{n!} \sup_{s \in [0,t]} |\rho_s^0| + Be^{-nt} \sum_{k=0}^{n-1} \frac{(Ant)^k}{k!} \\
&\leq \frac{(eAnt)^n}{n!} \sup_{s \in [0,t]} |\rho_s^0| + Bt(e^{At-1})^n \\
&\leq \frac{(eAt)^n}{\sqrt{2\pi n}} \sup_{s \in [0,t]} |\rho_s^0| + Bt(e^{At-1})^n, \quad t \in [0, T]. \tag{A.3}
\end{aligned}$$

This concludes the proof. \square

Corollary A.2. *In the framework of the above Lemma, if $eAT^* < 1$, then $\sup_{t \in [0, T^*]} \rho_t^n \rightarrow 0$.*

Proof. If $eAT^* < 1$, then $\sup_{t \in [0, T^*]} \rho_t^n$ is bounded from above by a sum of the n -th terms of two convergent geometrical series. \square

B Proof of inequality (2.8)

We give a sketch of the proof of inequality (2.8). Call G is the Green function of the Laplace operator $-\Delta$ on the torus $\mathbb{T}^2 = [-1/2, 1/2]^2$ (with periodic boundary condition). We will prove:

Proposition B.1. *The function G is in $C^\infty(\mathbb{T}^2 \setminus \{0\})$. Its behaviour in 0 is given by*

$$|G(x)| \leq C(-\log|x| + 1)$$

and that of its derivative $D^{(n)}$, n positive integer, by

$$|D^n G(x)| \leq C_n(|x|^{-n} + 1).$$

Assuming this result, we get that $|K(x)| \leq C_1(|x|^{-1} + 1)$. This implies the estimate (2.8) by an elementary argument (see [35], Appendix 2.3).

Proposition B.1 is a special case (at least for $n \leq 2$) of a general fact, valid for compact C^∞ Riemannian manifolds of finite dimensions, see [3, section 4.2], for the statement and a proof. We give here a different proof, taken in spirit from [4] (which studies the 3D case).

Sketch of the proof. It is easy to see that the Fourier expansion of G is

$$G(x) = -\frac{1}{4\pi^2} \sum_{k \in \mathbb{Z}^2, k \neq 0} \frac{1}{|k|^2} e^{2\pi i k \cdot x}$$

Since this expression seems not helpful in the analysis of regularity around 0, we will use the solution v , in $L^2([0, T] \times \mathbb{T}^2)$, of the heat equation

$$\partial_t v = \Delta v,$$

with initial condition $v_0 = \delta_0 - 1$ (more precisely, $v_t \rightarrow \delta_0 - 1$ as $t \rightarrow 0$). It is easy to see that this unique solution can be expressed in two ways: one with its Fourier expansion, which is

$$v(t, x) = \sum_{k \in \mathbb{Z}^2, k \neq 0} e^{-4\pi^2 |k|^2 t} e^{2\pi i k \cdot x}, \quad (\text{B.1})$$

the other with Gaussian densities, that is

$$v(t, x) = -1 + \frac{1}{4\pi t} \sum_{l \in \mathbb{Z}^2} \exp \frac{-|x - l|^2}{4t}. \quad (\text{B.2})$$

One verifies, e.g. using (B.1), that

$$\begin{aligned} G(x) &= -\int_0^{+\infty} v(t, x) dt = -\int_1^{+\infty} v(t, x) dt - \int_0^1 v(t, x) dt \\ &=: -G_1(x) - G_2(x). \end{aligned}$$

Now G_1 is in $C^\infty(\mathbb{T}^2)$, as one can see from its Fourier expansion, again from (B.1). For G_2 we exploit (B.2):

$$\begin{aligned} G_2(x) &= \left(-1 + \int_0^1 \frac{1}{4\pi t} \sum_{l \in \mathbb{Z}^2, l \neq 0} \exp \frac{-|x - l|^2}{4t} dt \right) \\ &+ \int_0^1 \frac{1}{4\pi t} \exp \frac{-|x|^2}{4t} dt =: G_3(x) + G_4(x), \end{aligned}$$

the sum being between functions on \mathbb{R}^2 (though x is still in $[-1/2, 1/2]^2$). The first addend G_3 is C^∞ on an open neighborhood of $[-1/2, 1/2]^2$ (e.g.

$] - 3/4, 3/4[$): indeed, for any n nonnegative integer, we have

$$\begin{aligned}
& \int_0^1 \left| D^{(n)} \frac{1}{4\pi t} \sum_{l \in \mathbb{Z}^2, l \neq 0} \exp \frac{-|x-l|^2}{4t} \right| dt \\
& \lesssim \int_0^1 t^{-(2n+1)} \sum_{l \neq 0} \exp \frac{-|x-l|^2}{4t} dt \\
& \lesssim \int_0^1 t^{-(2n+1)} \sum_{h=1}^{\infty} \exp \frac{-h}{ct} dt \\
& \sim \int_0^1 t^{-(2n+1)} e^{-1/(ct)} dt < +\infty,
\end{aligned}$$

for some $c > 0$ independent of x , when x is in $] - 3/4, 3/4[$. The second addend G_4 is in $C^\infty(] - 3/4, 3/4[\setminus \{0\})$. So G is in $C^\infty(\mathbb{T}^2 \setminus \{0\})$. For the behaviour in 0, this is given by the behaviour of G_4 , which is computed by standard techniques. We have, with the change of variable $s = |x|^{-1/2}t$,

$$G_4(x) \sim \int_0^{|x|^{-1/2}} s^{-1} e^{-1/(4s)} ds \sim -\log |x|$$

and, for $n \geq 1$,

$$|D^{(n)} G_4(x)| \sim |x|^{-n} \int_0^{|x|^{-1/2}} s^{-(2n+1)} e^{-1/(4s)} ds \sim |x|^{-n}.$$

The proof is complete. □

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