



A search for an excited muon decaying to a muon and two jets in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

The ATLAS Collaboration

Abstract

A new search signature for excited leptons is explored. Excited muons are sought in the channel $pp \rightarrow \mu\mu^* \rightarrow \mu\mu \text{ jet jet}$, assuming both the production and decay occur via a contact interaction. The analysis is based on 20.3 fb^{-1} of pp collision data at a centre-of-mass energy of $\sqrt{s} = 8$ TeV taken with the ATLAS detector at the Large Hadron Collider. No evidence of excited muons is found, and limits are set at the 95% confidence level on the cross section times branching ratio as a function of the excited-muon mass m_{μ^*} . For m_{μ^*} between 1.3 TeV and 3.0 TeV, the upper limit on $\sigma B(\mu^* \rightarrow \mu q\bar{q})$ is between 0.6 and 1 fb. Limits on σB are converted to lower bounds on the compositeness scale Λ . In the limiting case $\Lambda = m_{\mu^*}$, excited muons with a mass below 2.8 TeV are excluded. With the same model assumptions, these limits at larger μ^* masses improve upon previous limits from traditional searches based on the gauge-mediated decay $\mu^* \rightarrow \mu\gamma$.

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1 Introduction

The Standard Model (SM) of particle physics successfully describes a wide range of phenomena but does not explain the generational structure and mass hierarchy of quarks and leptons. Composite models of fermions [1–7] aim to reduce the number of matter constituents by postulating that SM fermions are bound states of more fundamental particles. A direct consequence of substructure would be the existence of excited fermion states.

This paper reports on a search for an excited muon μ^* using 20.3 fb^{-1} of pp collision data at a centre-of-mass energy of $\sqrt{s} = 8 \text{ TeV}$ recorded in 2012 with the ATLAS detector at the Large Hadron Collider (LHC). The search is based on a benchmark model [7] that describes excited-fermion interactions with an effective Lagrangian containing four-fermion contact interactions and gauge-mediated interactions. A contact interaction decay signature not previously employed in excited leptons searches, $\mu^* \rightarrow \mu j j$ (j represents a jet), is used.

In this paper, as in Ref. [7], the model is assumed to be valid for μ^* masses up to the compositeness scale. The contact interaction terms are described by the Lagrangian

$$\mathcal{L}_{\text{contact}} = \frac{g_*^2}{2\Lambda^2} j^\mu j_\mu, \text{ with } j_\mu = \eta \bar{f}_L \gamma_\mu f_L + \eta' \bar{f}_L^* \gamma_\mu f_L^* + \eta'' \bar{f}_L^* \gamma_\mu f_L + \text{h.c.},$$

where Λ is the compositeness scale; j_μ is the fermion current for ground states (f) and excited states (f^*); g_* and the η 's are constants; “h.c.” stands for Hermitian conjugate; and only left-handed fermion interactions are assumed. As is done in Ref. [7], g_*^2 is set to 4π , and η , η' , and η'' are taken to be one for all fermions. To calculate branching ratios, the compositeness scale Λ is assumed to be the same for gauge-mediated interactions, and the parameters f and f' in Ref. [7] are taken to be one.

The search described here focuses on the predominant single- μ^* production via the contact interaction ($q\bar{q} \rightarrow \mu^*\mu$) followed by the decay of the excited muon via the contact interaction to $\mu q\bar{q}$, leading to a

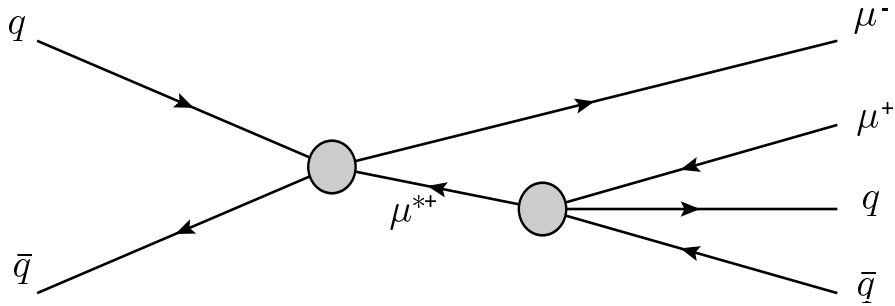


Figure 1: Feynman diagram for the process $q\bar{q} \rightarrow \mu^*\mu \rightarrow \mu\mu q\bar{q}$, where both the production and decay are via contact interactions.

final state with two muons and two jets (figure 1). Top quarks from excited muons with masses accessible in the 8-TeV LHC data would not have sufficient energy to form narrow jets and are excluded from the analysis in this paper. Previous searches at LEP [8–11], the Tevatron [12–15], and the LHC [16–20] looked for the gauge-mediated decay $\mu^* \rightarrow \mu\gamma$. The analysis reported in Ref. [20] also includes the gauge-mediated decay $\mu^* \rightarrow \mu Z$ followed by $Z \rightarrow \ell\ell$ or $q\bar{q}$. In the model of Ref. [7], this decay is dominant at low μ^* mass, but for $m_{\mu^*} \gtrsim 0.25\Lambda$, the $\mu q\bar{q}$ decay mode is expected to have the largest branching ratio, rising to 65% for $m_{\mu^*} = \Lambda$. The search reported here complements the search in the $\mu\gamma$ channel and increases the sensitivity of the search for excited muons at higher masses. The ATLAS Collaboration recently published [21] another new search signature for excited muons decaying via a contact interaction to $\mu\ell\ell$, where ℓ is an electron or a muon.

2 ATLAS detector

The ATLAS experiment [22] uses a multi-purpose particle detector with a forward-backward symmetric cylindrical geometry and a near 4π coverage in solid angle.¹ It consists of an inner tracking detector surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field, electromagnetic and hadron calorimeters, and a muon spectrometer. The inner tracking detector covers the pseudorapidity range $|\eta| < 2.5$. It consists of silicon pixel, silicon microstrip, and transition radiation tracking detectors. Lead/liquid-argon (LAr) sampling calorimeters provide electromagnetic (EM) energy measurements with high granularity. A hadronic steel/scintillator-tile calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The endcap and forward regions are instrumented with LAr calorimeters for EM and hadronic energy measurements up to $|\eta| = 4.9$. The muon spectrometer surrounding the calorimeters covers the pseudorapidity range $|\eta| < 2.7$ and is based on three large air-core toroid superconducting magnets with eight coils each. Their bending power is in the range from 2.0 to 7.5 T m. The muon spectrometer consists of three stations of precision tracking chambers and fast detectors for triggering. The majority of the precision tracking chambers are composed of drift tubes, while cathode-strip chambers provide coverage in the inner stations of the forward region for $2.0 < |\eta| < 2.7$. A three-level trigger system is

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in terms of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to at most 75 kHz. This is followed by two software-based trigger levels that together reduce the accepted event rate to 400 Hz on average, depending on the data-taking conditions during 2012.

3 Signal and background simulation

Simulation of the excited-muon signal is based on calculations from Ref. [7]. Signal samples are generated at leading order (LO) with CompHEP 4.5.1 [23] using MSTW2008lo [24] parton distribution functions (PDFs). CompHEP is interfaced with PYTHIA 8.170 [25, 26] with the AU2 parameters settings [27] for the simulation of parton showers and hadronization. Only the production of $\mu\mu^*$ followed by the decay $\mu^* \rightarrow \mu q\bar{q}$ is simulated. Signal samples are produced for $\Lambda = 5$ TeV and for the μ^* masses given in table 2. The distributions of kinematic variables should be independent of Λ , which was checked with generator-level studies. For a compositeness scale of $\Lambda = 5$ TeV, cross section times branching ratios are 10.4, 2.9, and 0.21 fb for μ^* masses of 500, 1500, and 2500 GeV, respectively. The intrinsic total width of the μ^* is expected to be less than 8% for $m_{\mu^*} < \Lambda$, which is smaller than the mass resolution of about 20% over the range of μ^* masses considered here.

The dominant background is from the process $Z/\gamma^* \rightarrow \mu\mu$ produced in association with jets ($Z/\gamma^* + \text{jets}$). The second most important background is $t\bar{t}$ production. Other processes, such as diboson (WW , WZ , and ZZ), single-top, $W + \text{jets}$, and multi-jet production, give small contributions to the background.

The $Z/\gamma^* + \text{jets}$ samples are produced by the multi-leg LO generator SHERPA 1.4.1 [28] using CT10 [29] PDFs. The cross section for $Z/\gamma^* \rightarrow \mu\mu$ ($m_{\mu\mu} > 70$ GeV) plus any number of jets is 1.24 nb, calculated at next-to-leading order (NLO), corrected by a K -factor [30, 31] to next-to-next-to-leading order (NNLO) in QCD couplings and NLO in electroweak couplings. The $t\bar{t}$ events are generated at the parton level at NLO with POWHEG 1.0 [32] and the Perugia 2011c parameter settings [33], and the parton showering is done with PYTHIA 6.426 [34]. At least one of the t or \bar{t} must have a semileptonic decay (e , μ , or τ), giving a cross section for this process of 137 pb, calculated at NNLO + next-to-next-to-leading-log (NNLL) accuracy [35]. The diboson background samples are produced at LO by HERWIG 6.52 [36] with the AUET2 parameter settings [37] using CTEQ6L1 PDFs, and it is required that at least one light lepton (e or μ) with transverse momentum (p_T) above 10 GeV be produced. The $W + \text{jets}$ samples are produced by the multi-leg LO generator ALPGEN 2.14 [38] with JIMMY 4.31 [39] and HERWIG 6.52 using the AUET2 parameter settings with CTEQ6L1 PDFs, and the cross section is calculated at NNLO [30, 31]. The multi-jet samples are generated at LO by PYTHIA 8.160 using the AU2 parameter settings with CT10 PDFs. The single-top t -channel samples are generated at LO corrected to NLO + NNLL by AcerMC 3.8 [40] using the AUET2B parameters settings [41] with the CTEQ6L1 PDFs, and the parton showering is done with PYTHIA 6.426. The single-top s - and Wt -channel samples are generated at NLO with MC@NLO 4.01 [42–44] using the AUET2 parameters settings with CT10 PDFs. The background predictions from the $Z/\gamma^* + \text{jets}$ and $t\bar{t}$ samples are normalized using control regions discussed in Section 5. Cross sections for diboson processes are evaluated at NLO [45] with an uncertainty of 5%. The $W + \text{jets}$ and multi-jet backgrounds are determined from the Monte Carlo (MC) samples but are verified using data-driven methods. A summary of the Standard Model samples used in this analysis is given in table 1.

The generated samples are processed using a detailed detector simulation [46] based on GEANT4 [47] to propagate the particles through the detector material and account for the detector response. Simulated minimum-bias events are overlaid on both the signal and background samples to reproduce the effect

Process	Generator	Parton showering/ hadronization	PDF
$Z/\gamma^* (\rightarrow \mu\mu) + \text{jets}$	SHERPA 1.4.1	SHERPA 1.4.1	CT10
$t\bar{t} (\geq 1\ell)$	POWHEG 1.0	PYTHIA 6.426	CT10
$WW, WZ, ZZ (\geq 1\ell)$	HERWIG 6.52	HERWIG 6.52	CTEQ6L1
Single top, t -channel	AcerMC 3.8	PYTHIA 6.426	CTEQ6L1
Single top, s -channel	MC@NLO 4.01	JIMMY 4.31 + HERWIG 6.52	CT10
Single top, Wt -channel	MC@NLO 4.01	JIMMY 4.31 + HERWIG 6.52	CT10
$W (\rightarrow \mu\nu) + \text{jets}$	ALPGEN 2.14	JIMMY 4.31 + HERWIG 6.52	CTEQ6L1
Multi-jet	PYTHIA 8.160	PYTHIA 8.160	CT10
Signal ($\mu\mu^* \rightarrow \mu\mu jj$)	CompHEP 4.5.1	PYTHIA 8.170	MSTW2008lo

Table 1: Summary of the background and signal MC sample generation used in this search. The columns give the process generated, the generator program, the parton shower program, and the PDF utilized.

of additional pp collisions. The simulated events are weighted to give a distribution of the number of interactions per bunch crossing that agrees with the data. The simulated background and signal events are processed with the same reconstruction programs as used for the data.

4 Data set and event selection

The data were collected in 2012 during stable-beam periods of $\sqrt{s} = 8$ TeV pp collisions. After selecting events where the relevant parts of the detector were functioning properly, the data correspond to an integrated luminosity of 20.3 fb^{-1} . The events are required to pass at least one of two single-muon triggers. The first has a nominal p_T threshold of 36 GeV, and the second has a lower nominal threshold of 24 GeV but also has an isolation requirement that the sum of the p_T of tracks with p_T above 1 GeV and within a distance of $\Delta R = 0.2$ of the muon, excluding the muon from the sum, divided by the p_T of the muon is less than 0.12.

A primary vertex with at least three tracks with $p_T > 0.4$ GeV within 200 mm of the centre of the detector along the beam direction is required. If there is more than one primary vertex in an event, the one with the highest sum of p_T^2 is selected, where the sum is over all tracks associated with the vertex.

Each muon candidate must be reconstructed independently in both the inner detector and the muon spectrometer. Its momentum is determined by a combination of the two measurements using their covariance matrices. Only muon candidates with p_T^μ above 30 GeV are considered. Muons must have a minimum number of hits in the inner detector and hits in each of the inner, middle, and outer layers of the muon spectrometer. These hit requirements, which restrict the muon acceptance to $|\eta| < 2.5$, guarantee a precise momentum measurement. To suppress background from cosmic rays, the muon tracks are required to have transverse and longitudinal impact parameters $|d_0| < 0.2$ mm and $|z_0| < 1$ mm with respect to the selected primary vertex. To reduce background from semileptonic decays of heavy-flavour hadrons, each muon is required to be isolated such that $\sum p_T/p_T^\mu < 0.05$, where the sum is over inner-detector tracks with $p_T > 1$ GeV within a distance $\Delta R = 0.3$ of the candidate muon, excluding the muon from the sum. The muon trigger and reconstruction efficiencies are evaluated using tag-and-probe techniques with

$Z \rightarrow \mu\mu$ events [48, 49], and p_T - and η -dependent corrections are applied to simulated events. Events are required to have exactly two muons of opposite charge that meet these selection requirements.

Although electrons are not part of the signal for this search, they are used to define one of the control regions (see Section 5). Each electron candidate is formed from the energy in a cluster of cells in the electromagnetic calorimeter associated with a charged-particle track in the inner detector. Each electron must have p_T above 30 GeV and have $|\eta| < 2.47$ but not be in the interval $1.37 < |\eta| < 1.52$ to avoid the transition region between the barrel and endcap calorimeters. The ATLAS tight electron identification criteria (based on the methodology described in [50] and updated for 2012 running conditions) for the transverse shower shape, longitudinal leakage into the hadronic calorimeter, the association with an inner-detector track, and hits in the transition radiation detector are applied to the cluster. An electron track is required to have transverse and longitudinal impact parameters $|d_0| < 1$ mm and $|z_0| < 5$ mm with respect to the selected primary vertex. Finally, the electrons must pass the isolation requirement $\sum E_T < 0.007E_T^e + 5$ GeV, where the sum is of transverse energies deposited in cells within a cone of $\Delta R = 0.2$ around the electron, excluding those cells associated with the electron, and E_T^e is the transverse energy of the electron.

Jets of hadrons are reconstructed using the anti- k_r algorithm [51] with a radius parameter of $R = 0.4$ applied to clusters of calorimeter cells that are topologically connected. The jets are calibrated using energy- and η -dependent correction factors derived from simulation and with residual corrections from in-situ measurements [52]. Jets are required to have $|\eta| < 2.8$ and $p_T > 30$ GeV. Jets that overlap ($\Delta R < 0.4$) any electron or muon candidate satisfying the selection criteria described above are removed. The two jets with the highest p_T are then selected.

The missing transverse momentum vector is the negative of the vector sum of the transverse momenta of muons, electrons, photons [53], jets, and clusters of calibrated calorimeter cells not associated with these objects. The missing transverse energy is the magnitude of the missing transverse momentum vector.

5 Background determination

Most of the SM background contributions are estimated from the MC samples. The expected yields from the $Z/\gamma^* +$ jets and $t\bar{t}$ production processes are normalized to the data using control regions. The $Z/\gamma^* +$ jets control region is defined by $70 < m_{\mu\mu} < 110$ GeV in addition to the selection criteria given in Section 4. The $t\bar{t}$ control region is defined as events that meet the selection requirements given in Section 4, except there is exactly one muon and one electron of opposite sign, so it should contain no signal events. The normalization scale factors are determined from simultaneous fits to data in the control and signal regions (see Section 8). The scale factors are primarily determined from the control regions, giving the same values in all cases. From the fits, the scale factor is $1.010^{+0.087}_{-0.066}$ for the $Z/\gamma^* +$ jets sample and is 1.050 ± 0.013 for the $t\bar{t}$ sample. The MC predictions agree well with the data in the control regions, as can be seen, for example, in figure 2(a).

A jet can produce a prompt muon candidate either from the semileptonic decay of a heavy quark or from misidentification of a charged hadron in the jet as a muon. The expected background from jets, primarily from $W +$ jets and multi-jet processes, is determined from MC samples, giving zero expected events. This prediction is checked by the data-driven matrix method [54], which uses isolated and non-isolated muons and their data-determined efficiencies and misidentification rates to determine the number of prompt muons. The matrix method predicts -0.07 ± 0.55 events from these backgrounds.

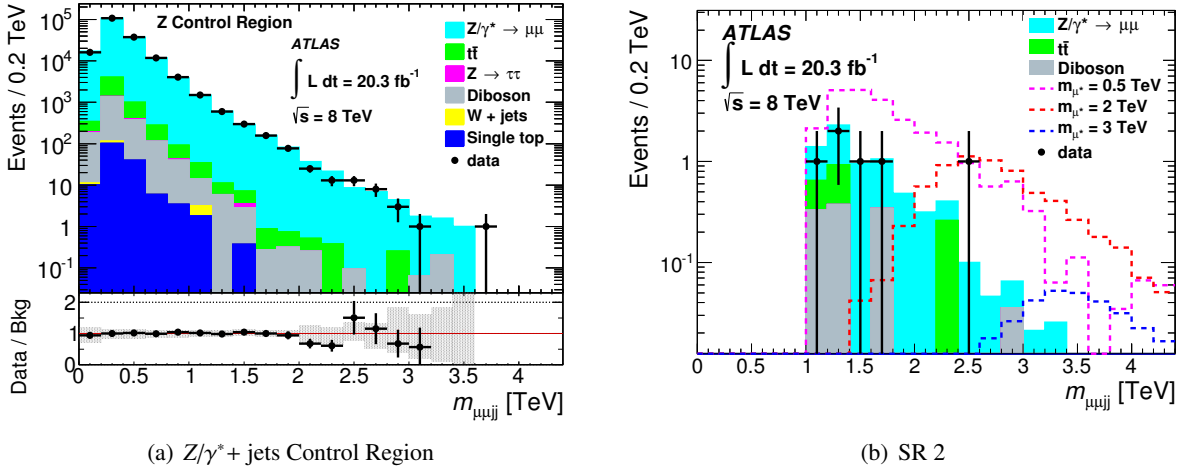


Figure 2: The $\mu\mu jj$ mass distribution for (a) the Z/γ^* + jets control region with the MC predictions and (b) for SR 2 ($m_{\mu\mu} > 550$ GeV, $S_T > 900$ GeV, and $m_{\mu\mu jj} > 1000$ GeV) with three representative signal distributions for μ^* masses of 500, 2000, and 3000 GeV and for $\Lambda = 5$ TeV. The background expectations are determined after the fit, and the grey band on the ratio plot in (a) gives the systematic uncertainty. For (b) there are no expected events from the single top, W + jets, and $Z \rightarrow \tau\tau$ processes. SR 2 is not the most sensitive signal region for the latter two m_{μ^*} masses. They are shown for comparison.

6 Signal regions

Signal regions (SR) are defined by three kinematic variables - the dimuon invariant mass $m_{\mu\mu}$, the invariant mass $m_{\mu\mu jj}$ of the two muons and two jets (j), and S_T , the scalar sum of transverse momenta of the four signal objects, that is, $S_T = p_T^{\mu 1} + p_T^{\mu 2} + p_T^{j 1} + p_T^{j 2}$.²

For all three of these variables, the signal tends to have higher values than the backgrounds, so all criteria are lower bounds in the selection. The values of these bounds are chosen to maximize the search sensitivity for each signal mass considered by scanning the three-dimensional parameter space for the values that minimize the expected 95% confidence level (CL) upper limit on the cross section times branching ratio. The selection criteria for the signal regions are shown in table 2. The $m_{\mu\mu jj}$ and S_T criteria increase with increasing signal mass, but the $m_{\mu\mu}$ criterion decreases. The latter is because the increase in the other parameters sufficiently reduces the expected background so that the signal efficiency may be increased by decreasing the $m_{\mu\mu}$ criterion.

The dominant background in all signal regions is from the Z/γ^* + jets process, which is 50% of the background in SR 1, rising to 90% or more in SR 5 through SR 10. The $t\bar{t}$ process contributes 40% of the background in SR 1, but this contribution falls quickly to 10% or less in SR 3 through SR 10. The contribution to the background from all other processes is between 10% and 20% in SR 1 through SR 5 and is less than 5% for SR 6 through SR 10.

² The μjj invariant mass was considered as a discriminating variable instead of one of the three selection variables. Several methods for selecting the correct μjj combination and the possibility of using both μjj combinations were considered. No method that improved the sensitivity was found.

m_{μ^*} [GeV]	Signal Region	$m_{\mu\mu}$ [GeV]	S_T [GeV]	$m_{\mu\mu jj}$ [GeV]	Acc \times Eff	Exp Signal	Exp BG (prefit)	Exp BG (postfit)	Obs Events
100	1	500	450	0	0.041	3.0 \pm 0.3	73 \pm 17	71.7 \pm 8.6	71
300	2	550	900	1000	0.088	12.5 \pm 0.9	10.1 \pm 3.5	7.8 \pm 2.2	6
500	2	550	900	1000	0.15	29.4 \pm 1.6	10.1 \pm 3.5	7.8 \pm 2.2	6
750	3	450	900	1300	0.23	43.7 \pm 2.2	6.6 \pm 2.5	5.8 \pm 1.9	5
1000	4	450	1050	1300	0.31	44.1 \pm 1.8	4.6 \pm 1.8	4.6 \pm 1.9	5
1250	5	450	1200	1500	0.38	29.8 \pm 1.3	2.1 \pm 0.9	2.1 \pm 1.0	3
1500	6	400	1200	1700	0.38	19.8 \pm 0.8	1.5 \pm 0.7	1.3 \pm 0.6	1
1750	7	300	1350	1900	0.41	11.8 \pm 0.5	0.9 \pm 0.7	0.9 \pm 0.5	2
2000	8	300	1350	2000	0.40	6.6 \pm 0.2	0.7 \pm 0.4	0.7 \pm 0.4	2
2250	9	300	1500	2100	0.37	3.4 \pm 0.1	0.4 \pm 0.3	0.4 \pm 0.3	2
2500	10	110	1650	2300	0.39	1.65 \pm 0.07	0.2 $^{+1.3}_{-0.2}$	0.2 $^{+1.0}_{-0.2}$	2
2750	10	110	1650	2300	0.45	0.72 \pm 0.02	0.2 $^{+1.3}_{-0.2}$	0.2 $^{+1.0}_{-0.2}$	2
2900	10	110	1650	2300	0.45	0.52 \pm 0.02	0.2 $^{+1.3}_{-0.2}$	0.2 $^{+1.0}_{-0.2}$	2
3000	10	110	1650	2300	0.46	0.38 \pm 0.02	0.2 $^{+1.3}_{-0.2}$	0.2 $^{+1.0}_{-0.2}$	2
3100	10	110	1650	2300	0.45	0.30 \pm 0.02	0.2 $^{+1.3}_{-0.2}$	0.2 $^{+1.0}_{-0.2}$	2

Table 2: The signal masses considered and the corresponding signal regions are listed. The $m_{\mu\mu}$, S_T , and $m_{\mu\mu jj}$ values giving the lower bound of each signal region are listed, along with the acceptance times efficiency, the expected number of signal events ($\Lambda = 5$ TeV), expected number of background events before and after the fit discussed in Section 8, and the number of events observed in the data. The uncertainties in the expected numbers of signal and background events are the systematic uncertainties. The numbers of events observed are discussed in Section 8.

7 Systematic uncertainties

Contributions to the systematic uncertainties in the background and signal yield predictions stem from both experimental and theoretical sources, as discussed below.

The luminosity is derived using the methodology in Ref. [55] and has an uncertainty of 2.8%. The luminosity uncertainty for the backgrounds is less than this because the largest backgrounds ($Z/\gamma^* +$ jets and $t\bar{t}$) are normalized using control regions.

Uncertainties in the MC modelling of the detector, particularly for muons and jets in this analysis, must be taken into account and are derived from detailed studies of data. One-standard-deviation variation of a given parameter is determined, and then the parameter is varied up and down in the simulation by this amount to determine the effect on the signal and background yields.

The uncertainty in the jet energy scale is the largest contribution to the systematic uncertainty in the signal yield and a significant contribution to the uncertainty in the backgrounds. The uncertainty in the jet energy resolution also makes a contribution. These uncertainties are determined from p_T balance in $\gamma +$ jet and $Z +$ jet events and in events with high- p_T jets recoiling against multiple, low- p_T jets [52, 56]. The uncertainty in contributions from additional energy deposited in the calorimeters from other pp interactions in the event is also included. The various effects are investigated separately and combined to give the values summarized in tables 3 and 4.

Muon performance is determined in $Z \rightarrow \mu\mu$ events. The most important parameters for this analysis are the muon efficiency and the muon spectrometer p_T resolution. The inner-detector resolution and the muon p_T scale are found to have negligible effect. The uncertainty in the trigger efficiency is less than 2% for the backgrounds and less than 1% for the signal yield.

The uncertainties in the signal and background yield predictions due to uncertainties in PDFs have two contributions. The first is from one-standard-deviation variation of the parameters of the relevant PDFs (Section 3). The second is a comparison with the alternative NNPDF2.1 PDF set [57]. These variations produce changes in the predicted cross section and in kinematical distributions, which in turn affect the acceptance times efficiency. For the background, both effects are included in the systematic uncertainty. For the signal yield, the uncertainty in the acceptance times efficiency is included, but the uncertainty in the cross section is considered part of the uncertainty in the theoretical prediction and is not included in the statistical analysis.

The uncertainty in the background modelling in the signal regions is estimated by examining how well the MC prediction agrees with the data in two validation regions selected to be similar in kinematics to the signal regions but containing no signal. Both validation regions require the same selection as the signal regions except that $m_{\mu\mu jj} < 500$ GeV and $m_{\mu\mu} > 200$ GeV with no selection on S_T . Requiring the missing transverse energy be greater (less) than 50 GeV (40 GeV) selects a validation region dominated by $t\bar{t}$ (Z/γ^*) events. For some of the kinematic variables, an extrapolation of the predicted yield from the validation regions to the signal regions is necessary to evaluate possible mismodelling effects. Of the several kinematic variables studied, only the modelling of the S_T variable is found to have a significant effect. A linear fit to the ratio of the number of data events to the MC expectation is extrapolated to higher values of S_T , and the deviation from unity symmetrized about zero gives the uncertainty, referred to as “ Z/γ^* + jets modelling” and “ $t\bar{t}$ modelling” in table 4. For both validation regions, the linear fit is consistent within the statistical uncertainties with a flat line at a ratio of one.

To produce sufficient numbers of events for high dimuon masses, the Z/γ^* MC samples were produced in bins of dimuon mass above the Z mass. For the S_T and $m_{\mu\mu jj}$ criteria in this analysis, this yields zero events in SR 7 through SR 10 for some ranges of the $m_{\mu\mu}$ distribution (for example, 110 to 400 GeV for SR 10). For these signal regions, an additional systematic uncertainty (referred to as “ Z/γ^* + jets extrapolation” in table 4) is estimated by loosening the S_T criteria and extrapolating into the signal region. The uncertainty introduced by this procedure is small except in SR 10, where the effect on the statistical analysis is still small because the predicted number of background events is only 0.2.

Additional sources of uncertainty in the acceptance times efficiency are initial-state radiation, final-state radiation, renormalization and factorization scales, and the beam energy. The effects of initial- and final-state radiation are determined in generator-level studies by varying the relevant PYTHIA parameters and are less than 1%. The effect of the beam energy uncertainty (0.65%) [58] is determined by varying the momentum fraction of the initial partons in the PDFs by this amount, giving a change of less than 1%. The renormalization and factorization scales are independently varied in the simulation by factors of 2 and 1/2, changing the expected signal acceptance times efficiency by about 2% at low mass and by less than 1% for masses above 750 GeV.

The uncertainties in the signal yield depend on the μ^* mass, and the largest contributions are summarized in table 3 for three representative masses. For the signal yield, uncertainties in jet energy scale, PDFs, and luminosity are the dominant sources. The uncertainties in the background depend on the signal region, and the largest contributions are shown in table 4 for three representative regions. The most significant contributions to the background uncertainty are from the modelling of the Z/γ^* + jets and $t\bar{t}$ processes.

m_{μ^*} [GeV]	500	1500	2500
Luminosity	2.8	2.8	2.8
Jet energy scale	2.7	1.5	1.1
Hadronization and factorization scales	2.0	0.5	0.1
PDFs	3.0	2.5	2.7
Muon efficiency	0.7	0.8	0.9
Jet energy resolution	0.9	0.2	0.5
Muon spectrometer resolution	0.3	<0.1	0.2
Total	5.1	4.1	4.2

Table 3: Largest contributions to the relative systematic uncertainty in the signal yield. The uncertainties for the hadronization and factorization scales and for the PDFs are only for the signal acceptance times efficiency. All uncertainties are given in percent and are determined after the fit discussed in Section 8.

Signal region	2	6	10
Z/γ^* + jets modelling	25	47	65
Jet energy scale	19	9.0	6.2
$t\bar{t}$ modelling	12	<0.1	<0.1
Muon spectrometer resolution	6.2	0.6	63
PDFs	4.2	8.8	17
Jet energy resolution	3.2	1.7	0.6
Muon efficiency	0.7	0.8	0.9
Luminosity	0.4	0.1	<0.1
Z/γ^* + jets extrapolation	–	–	500
Total	35	49	500

Table 4: Largest contributions to the relative systematic uncertainty in the expected background for three representative signal regions. All uncertainties are given in percent and are determined after the fit discussed in Section 8.

The jet energy scale and the parton distribution functions also make significant contributions. Any source of systematic uncertainty contributing less than 2% to the background for all signal regions and less than 1% to the signal yield for all μ^* masses would have negligible effect in the statistical analysis in Section 8 and is not included.

8 Results

For each μ^* mass considered, the numbers of events in the corresponding signal region and in the two control regions are simultaneously fit [59] using a profile likelihood method [60, 61]. The likelihood function models the number of events as a Poisson distribution and the systematic effects are modelled using nuisance parameters with lognormal constraints. The parameters of interest in the fit are the signal yield in the corresponding signal region and the normalizations of the Z/γ^* and $t\bar{t}$ backgrounds, with the latter two being primarily determined in the fit by the events in the control regions. The possible contribution of signal to the control regions is included in the fit and found to be negligible. Correlations of the systematic uncertainties are taken into account.

Event	SR	$m_{\mu\mu}$ [GeV]	$m_{\mu\mu jj}$ [GeV]	S_T [GeV]	$m_{\mu_1 jj}$ [GeV]	$m_{\mu_2 jj}$ [GeV]	$p_T^{\mu_1}$ [GeV]	$p_T^{\mu_2}$ [GeV]	$p_T^{j_1}$ [GeV]	$p_T^{j_2}$ [GeV]
A	all	1800	2410	1820	1200	1090	650	630	350	190
B	7–9	310	2250	2010	2200	630	840	46	990	130
C	10	113	2440	1760	2230	1850	150	35	890	690

Table 5: Values of $\mu\mu$ mass, $\mu\mu jj$ mass, S_T , μjj mass for each μjj combination, and p_T of each muon and jet for the three events in SR 9 or 10.

As an example of the result of the fit, the $m_{\mu\mu jj}$ distribution for signal region 2 is shown in figure 2(b) for the data, expected backgrounds, and three signal predictions for $\Lambda = 5$ TeV (the signal regions for the higher masses have fewer background events). The expected and observed numbers of events for each signal mass considered are shown in table 2 for $\Lambda = 5$ TeV. Due to correlations among the nuisance parameters, the uncertainties on the expected backgrounds are reduced after the fits. The data are consistent with the Standard Model expectations, and no significant excess is observed. Thus, limits on the cross section times branching ratio as a function of the μ^* mass are calculated.

A modified frequentist CL_s method [62, 63] is used to derive the 95% CL upper limits on the signal yield. The expected limit is the median limit for a large number of background-only pseudo-experiments. The one- and two-standard-deviations bands cover 68% and 95%, respectively, of the pseudo-experiment limits. The observed limit is the 95% CL limit for the observed number of events. The p -value is a measure of how well the background-only hypothesis models the data. For a signal region, it is the fraction of background-only pseudo-experiments where the fitted signal value is greater than that for the observed data.

The smallest p -values are for SR 9 and 10 with values of 0.034 and 0.099, respectively, corresponding to 1.8 and 1.3 standard deviations on one side of a Gaussian distribution. Some kinematic properties of the events in these signal regions are given in table 5. There is one event (event A) that is in all signal regions.

An upper limit on the cross section times branching ratio $\sigma(pp \rightarrow \mu\mu^*) \times B(\mu^* \rightarrow \mu q \bar{q})$ (figure 3) is determined for each signal mass from the limit on the signal yield at the 95% CL. The theoretical uncertainties are not included in either the σB or Λ limit determinations. For m_{μ^*} above 1.3 TeV, the limit is between 0.6 and 1 fb. The theoretical expectation for $\Lambda = m_{\mu^*}$ is also shown. The theoretical band represents uncertainties from PDFs and from renormalization and factorization scales.

The expected cross section and branching ratio depend on the μ^* mass and on Λ [7]. For each signal mass, the limit on σB is translated into a lower bound on the compositeness scale (figure 4). The bound is the value of Λ for which the theoretical prediction of $\sigma B(m_{\mu^*}, \Lambda)$ is equal to the upper limit on σB . The region with $m_{\mu^*} > \Lambda$ is unphysical. For the limiting case where $\Lambda = m_{\mu^*}$, excited-muon masses below 2.8 TeV are excluded. Previous limits set by ATLAS [17, 21] are also shown. The analysis presented here improves upon the limits from $\mu^* \rightarrow \mu\gamma$ for masses above 1100 GeV and upon those from $\mu^* \rightarrow \mu\ell\ell$ for masses from 700 to 2100 GeV.

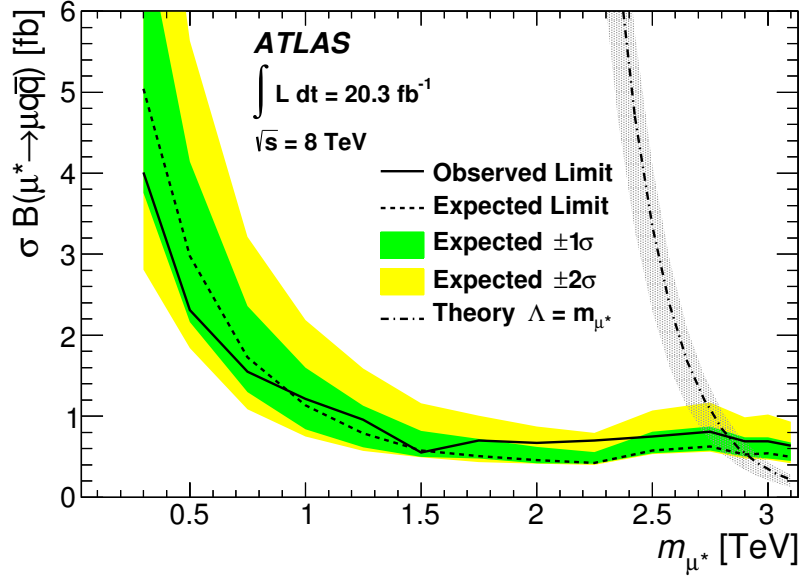


Figure 3: Limit at 95% CL on cross section times branching ratio $\sigma(pp \rightarrow \mu\mu^*)B(\mu^* \rightarrow \mu q\bar{q})$ as a function of the μ^* mass. The theory curve only includes contact-interaction decays and does not include the top quark. The solid line is the limit and the dotted line is the expected limit. The theoretical σB for the limiting case $\Lambda = m_{\mu^*}$ along with its uncertainties is also shown (dot-dashed curve).

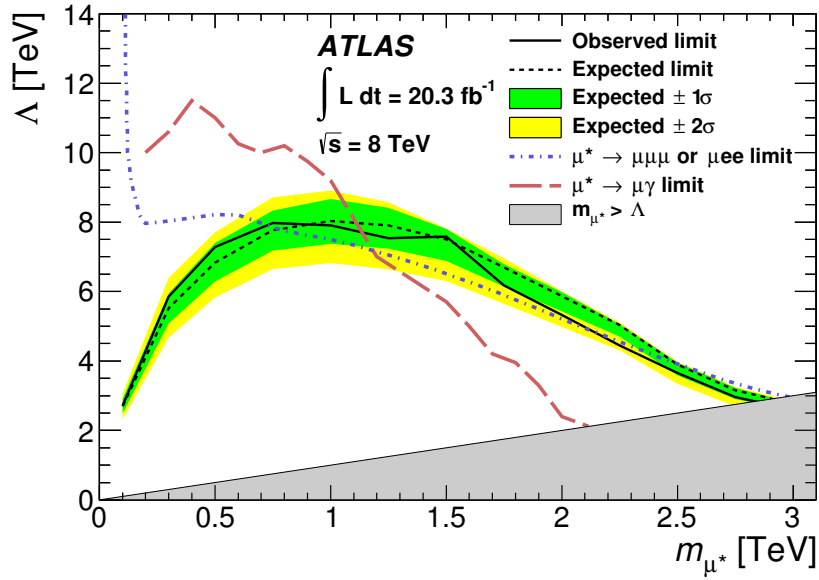


Figure 4: Limit at 95% CL on the compositeness scale Λ as a function of the μ^* mass. The solid black line is the limit and the short dashed black line is the expected limit. Also indicated are previous results from ATLAS based on $\mu^* \rightarrow \mu\gamma$ (long dashed red line) and $\mu^* \rightarrow \mu\ell\ell$ (dot-dashed blue line), where ℓ is an electron or muon.

9 Conclusion

The results of a search for excited muons decaying to μjj via a contact interaction are reported based on data from $\sqrt{s} = 8$ TeV pp collisions collected with the ATLAS detector at the LHC corresponding to an integrated luminosity of 20.3 fb^{-1} . The observed data are consistent with SM expectations. An upper limit is set at 95% CL on the cross section times branching ratio $\sigma B(\mu^* \rightarrow \mu q\bar{q})$ as a function of the excited-muon mass. For m_{μ^*} between 1.3 and 3.0 TeV, the limit on σB is between 0.6 and 1 fb.

The σB upper limits are converted to lower bounds on the compositeness scale Λ . In the limiting case where $\Lambda = m_{\mu^*}$, excited-muon masses below 2.8 TeV are excluded. At higher μ^* masses, the signature explored in this paper, $\mu^* \rightarrow \mu jj$, has better sensitivity than the traditional signature $\mu^* \rightarrow \mu \gamma$. For μ^* masses above 0.8 TeV, the sensitivity is up to 15% better than a previous search using the signature $\mu^* \rightarrow \mu \ell\ell$. In models other than the benchmark model used here, the branching ratios to these modes could be different, affecting their relative importance for limits on the compositeness scale.

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The ATLAS Collaboration

G. Aad⁸⁵, B. Abbott¹¹², J. Abdallah¹⁵⁰, O. Abdinov¹¹, B. Abeloos¹¹⁶, R. Aben¹⁰⁶, M. Abolins⁹⁰, O.S. AbouZeid¹⁵⁷, H. Abramowicz¹⁵², H. Abreu¹⁵¹, R. Abreu¹¹⁵, Y. Abulaiti^{145a,145b}, B.S. Acharya^{163a,163b,a}, L. Adamczyk^{38a}, D.L. Adams²⁵, J. Adelman¹⁰⁷, S. Adomeit⁹⁹, T. Adye¹³⁰, A.A. Affolder⁷⁴, T. Agatonovic-Jovin¹³, J. Agricola⁵⁴, J.A. Aguilar-Saavedra^{125a,125f}, S.P. Ahlen²², F. Ahmadov^{65,b}, G. Aielli^{132a,132b}, H. Akerstedt^{145a,145b}, T.P.A. Åkesson⁸¹, A.V. Akimov⁹⁵, G.L. Alberghi^{20a,20b}, J. Albert¹⁶⁸, S. Albrand⁵⁵, M.J. Alconada Verzini⁷¹, M. Aleksa³⁰, I.N. Aleksandrov⁶⁵, C. Alexa^{26b}, G. Alexander¹⁵², T. Alexopoulos¹⁰, M. Alhroob¹¹², G. Alimonti^{91a}, L. Alio⁸⁵, J. Alison³¹, S.P. Alkire³⁵, B.M.M. Allbrooke¹⁴⁸, B.W. Allen¹¹⁵, P.P. Allport¹⁸, A. Aloisio^{103a,103b}, A. Alonso³⁶, F. Alonso⁷¹, C. Alpigiani¹³⁷, B. Alvarez Gonzalez³⁰, D. Álvarez Piqueras¹⁶⁶, M.G. Alvigi^{103a,103b}, B.T. Amadio¹⁵, K. Amako⁶⁶, Y. Amaral Coutinho^{24a}, C. Amelung²³, D. Amidei⁸⁹, S.P. Amor Dos Santos^{125a,125c}, A. Amorim^{125a,125b}, S. Amoroso³⁰, N. Amram¹⁵², G. Amundsen²³, C. Anastopoulos¹³⁸, L.S. Ancu⁴⁹, N. Andari¹⁰⁷, T. Andeen³¹, C.F. Anders^{58b}, G. Anders³⁰, J.K. Anders⁷⁴, K.J. Anderson³¹, A. Andreazza^{91a,91b}, V. Andrei^{58a}, S. Angelidakis⁹, I. Angelozzi¹⁰⁶, P. Anger⁴⁴, A. Angerami³⁵, F. Anghinolfi³⁰, A.V. Anisenkov^{108,c}, N. Anjos¹², A. Annovi^{123a,123b}, M. Antonelli⁴⁷, A. Antonov⁹⁷, J. Antos^{143b}, F. Anulli^{131a}, M. Aoki⁶⁶, L. Aperio Bella¹⁸, G. Arabidze⁹⁰, Y. Arai⁶⁶, J.P. Araque^{125a}, A.T.H. Arce⁴⁵, F.A. Arduh⁷¹, J-F. Arguin⁹⁴, S. Argyropoulos⁶³, M. Arik^{19a}, A.J. Armbruster³⁰, O. Arnaez³⁰, H. Arnold⁴⁸, M. Arratia²⁸, O. Arslan²¹, A. Artamonov⁹⁶, G. Artoni¹¹⁹, S. Artz⁸³, S. Asai¹⁵⁴, N. Asbah⁴², A. Ashkenazi¹⁵², B. Åsman^{145a,145b}, L. Asquith¹⁴⁸, K. Assamagan²⁵, R. Astalos^{143a}, M. Atkinson¹⁶⁴, N.B. Atlay¹⁴⁰, K. Augsten¹²⁷, G. Avolio³⁰, B. Axen¹⁵, M.K. Ayoub¹¹⁶, G. Azuelos^{94,d}, M.A. Baak³⁰, A.E. Baas^{58a}, M.J. Baca¹⁸, H. Bachacou¹³⁵, K. Bachas¹⁵³, M. Backes³⁰, M. Backhaus³⁰, P. Bagiachi^{131a,131b}, P. Bagnaia^{131a,131b}, Y. Bai^{33a}, J.T. Baines¹³⁰, O.K. Baker¹⁷⁵, E.M. Baldwin^{108,c}, P. Balek¹²⁸, T. Balestri¹⁴⁷, F. Balli⁸⁴, W.K. Balunas¹²¹, E. Banas³⁹, Sw. Banerjee^{172,e}, A.A.E. Bannoura¹⁷⁴, L. Barak³⁰, E.L. Barberio⁸⁸, D. Barberis^{50a,50b}, M. Barbero⁸⁵, T. Barillari¹⁰⁰, M. Barisonzi^{163a,163b}, T. Barklow¹⁴², N. Barlow²⁸, S.L. Barnes⁸⁴, B.M. Barnett¹³⁰, R.M. Barnett¹⁵, Z. Barnovska⁵, A. Baroncelli^{133a}, G. Barone²³, A.J. Barr¹¹⁹, L. Barranco Navarro¹⁶⁶, F. Barreiro⁸², J. Barreiro Guimarães da Costa^{33a}, R. Bartoldus¹⁴², A.E. Barton⁷², P. Bartos^{143a}, A. Basalae¹²², A. Bassalat¹¹⁶, A. Basye¹⁶⁴, R.L. Bates⁵³, S.J. Batista¹⁵⁷, J.R. Batley²⁸, M. Battaglia¹³⁶, M. Bauce^{131a,131b}, F. Bauer¹³⁵, H.S. Bawa^{142,f}, J.B. Beacham¹¹⁰, M.D. Beattie⁷², T. Beau⁸⁰, P.H. Beauchemin¹⁶⁰, R. Beccherle^{123a,123b}, P. Bechtel²¹, H.P. Beck^{17,g}, K. Becker¹¹⁹, M. Becker⁸³, M. Beckingham¹⁶⁹, C. Becot¹¹⁶, A.J. Beddall^{19b}, A. Beddall^{19b}, V.A. Bednyakov⁶⁵, M. Bedognetti¹⁰⁶, C.P. Bee¹⁴⁷, L.J. Beemster¹⁰⁶, T.A. Beermann³⁰, M. Begel²⁵, J.K. Behr¹¹⁹, C. Belanger-Champagne⁸⁷, W.H. Bell⁴⁹, G. Bella¹⁵², L. Bellagamba^{20a}, A. Bellerive²⁹, M. Bellomo⁸⁶, K. Belotskiy⁹⁷, O. Beltramello³⁰, O. Benary¹⁵², D. Bencheekroun^{134a}, M. Bender⁹⁹, K. Bendtz^{145a,145b}, N. Benekos¹⁰, Y. Benhammou¹⁵², E. Benhar Nocchioli¹⁷⁵, J.A. Benitez Garcia^{158b}, D.P. Benjamin⁴⁵, J.R. Bensinger²³, S. Bentvelsen¹⁰⁶, L. Beresford¹¹⁹, M. Beretta⁴⁷, D. Berge¹⁰⁶, E. Bergeaas Kuutmann¹⁶⁵, N. Berger⁵, F. Berghaus¹⁶⁸, J. Beringer¹⁵, C. Bernard²², N.R. Bernard⁸⁶, C. Bernius¹⁰⁹, F.U. Bernlochner²¹, T. Berry⁷⁷, P. Berta¹²⁸, C. Bertella⁸³, G. Bertoli^{145a,145b}, F. Bertolucci^{123a,123b}, C. Bertsche¹¹², D. Bertsche¹¹², G.J. Besjes³⁶, O. Bessidskaia Bylund^{145a,145b}, M. Bessner⁴², N. Besson¹³⁵, C. Betancourt⁴⁸, S. Bethke¹⁰⁰, A.J. Bevan⁷⁶, W. Bhimji¹⁵, R.M. Bianchi¹²⁴, L. Bianchini²³, M. Bianco³⁰, O. Biebel⁹⁹, D. Biedermann¹⁶, N.V. Biesuz^{123a,123b}, M. Biglietti^{133a}, J. Bilbao De Mendizabal⁴⁹, H. Bilokon⁴⁷, M. Bindi⁵⁴, S. Binet¹¹⁶, A. Bingul^{19b}, C. Bini^{131a,131b}, S. Biondi^{20a,20b}, D.M. Bjergaard⁴⁵, C.W. Black¹⁴⁹, J.E. Black¹⁴², K.M. Black²², D. Blackburn¹³⁷, R.E. Blair⁶, J.-B. Blanchard¹³⁵, J.E. Blanco⁷⁷, T. Blazek^{143a}, I. Bloch⁴², C. Blocker²³, W. Blum^{83,*}, U. Blumenschein⁵⁴, S. Blunier^{32a}, G.J. Bobbink¹⁰⁶, V.S. Bobrovnikov^{108,c}, S.S. Bocchetta⁸¹, A. Bocci⁴⁵,

C. Bock⁹⁹, M. Boehler⁴⁸, D. Boerner¹⁷⁴, J.A. Bogaerts³⁰, D. Bogavac¹³, A.G. Bogdanchikov¹⁰⁸,
 C. Bohm^{145a}, V. Boisvert⁷⁷, T. Bold^{38a}, V. Boldea^{26b}, A.S. Boldyrev⁹⁸, M. Bomben⁸⁰, M. Bona⁷⁶,
 M. Boonekamp¹³⁵, A. Borisov¹²⁹, G. Borissov⁷², J. Bortfeldt⁹⁹, V. Bortolotto^{60a,60b,60c}, K. Bos¹⁰⁶,
 D. Boscherini^{20a}, M. Bosman¹², J.D. Bossio Sola²⁷, J. Boudreau¹²⁴, J. Bouffard²,
 E.V. Bouhova-Thacker⁷², D. Boumediene³⁴, C. Bourdarios¹¹⁶, N. Bousson¹¹³, S.K. Boutle⁵³,
 A. Boveia³⁰, J. Boyd³⁰, I.R. Boyko⁶⁵, J. Bracinik¹⁸, A. Brandt⁸, G. Brandt⁵⁴, O. Brandt^{58a},
 U. Bratzler¹⁵⁵, B. Brau⁸⁶, J.E. Brau¹¹⁵, H.M. Braun^{174,*}, W.D. Breaden Madden⁵³, K. Brendlinger¹²¹,
 A.J. Brennan⁸⁸, L. Brenner¹⁰⁶, R. Brenner¹⁶⁵, S. Bressler¹⁷¹, T.M. Bristow⁴⁶, D. Britton⁵³, D. Britzger⁴²,
 F.M. Brochu²⁸, I. Brock²¹, R. Brock⁹⁰, G. Brooijmans³⁵, T. Brooks⁷⁷, W.K. Brooks^{32b}, J. Brosamer¹⁵,
 E. Brost¹¹⁵, P.A. Bruckman de Renstrom³⁹, D. Bruncko^{143b}, R. Bruneliere⁴⁸, A. Bruni^{20a}, G. Bruni^{20a},
 BH Brunt²⁸, M. Bruschi^{20a}, N. Bruscino²¹, P. Bryant³¹, L. Bryngemark⁸¹, T. Buanes¹⁴, Q. Buat¹⁴¹,
 P. Buchholz¹⁴⁰, A.G. Buckley⁵³, I.A. Budagov⁶⁵, F. Buehrer⁴⁸, L. Bugge¹¹⁸, M.K. Bugge¹¹⁸,
 O. Bulekov⁹⁷, D. Bullock⁸, H. Burckhart³⁰, S. Burdin⁷⁴, C.D. Burgard⁴⁸, B. Burghgrave¹⁰⁷, S. Burke¹³⁰,
 I. Burmeister⁴³, E. Busato³⁴, D. Büscher⁴⁸, V. Büscher⁸³, P. Bussey⁵³, J.M. Butler²², A.I. Butt³,
 C.M. Buttar⁵³, J.M. Butterworth⁷⁸, P. Butti¹⁰⁶, W. Buttinger²⁵, A. Buzatu⁵³, A.R. Buzykaev^{108,c},
 S. Cabrera Urbán¹⁶⁶, D. Caforio¹²⁷, V.M. Cairo^{37a,37b}, O. Cakir^{4a}, N. Calace⁴⁹, P. Calafiura¹⁵,
 A. Calandri⁸⁵, G. Calderini⁸⁰, P. Calfayan⁹⁹, L.P. Caloba^{24a}, D. Calvet³⁴, S. Calvet³⁴, T.P. Calvet⁸⁵,
 R. Camacho Toro³¹, S. Camarda⁴², P. Camarri^{132a,132b}, D. Cameron¹¹⁸, R. Caminal Armadans¹⁶⁴,
 C. Camincher⁵⁵, S. Campana³⁰, M. Campanelli⁷⁸, A. Campoverde¹⁴⁷, V. Canale^{103a,103b}, A. Canepa^{158a},
 M. Cano Bret^{33e}, J. Cantero⁸², R. Cantrill^{125a}, T. Cao⁴⁰, M.D.M. Capeans Garrido³⁰, I. Caprini^{26b},
 M. Caprini^{26b}, M. Capua^{37a,37b}, R. Caputo⁸³, R.M. Carbone³⁵, R. Cardarelli^{132a}, F. Cardillo⁴⁸, T. Carli³⁰,
 G. Carlino^{103a}, L. Carminati^{91a,91b}, S. Caron¹⁰⁵, E. Carquin^{32a}, G.D. Carrillo-Montoya³⁰, J.R. Carter²⁸,
 J. Carvalho^{125a,125c}, D. Casadei⁷⁸, M.P. Casado^{12,h}, M. Casolino¹², D.W. Casper¹⁶²,
 E. Castaneda-Miranda^{144a}, A. Castelli¹⁰⁶, V. Castillo Gimenez¹⁶⁶, N.F. Castro^{125a,i}, A. Catinaccio³⁰,
 J.R. Catmore¹¹⁸, A. Cattai³⁰, J. Caudron⁸³, V. Cavaliere¹⁶⁴, D. Cavalli^{91a}, M. Cavalli-Sforza¹²,
 V. Cavasinni^{123a,123b}, F. Ceradini^{133a,133b}, L. Cerda Alberich¹⁶⁶, B.C. Cerio⁴⁵, A.S. Cerqueira^{24b},
 A. Cerri¹⁴⁸, L. Cerrito⁷⁶, F. Cerutti¹⁵, M. Cerv³⁰, A. Cervelli¹⁷, S.A. Cetin^{19c}, A. Chafaq^{134a},
 D. Chakraborty¹⁰⁷, I. Chalupkova¹²⁸, Y.L. Chan^{60a}, P. Chang¹⁶⁴, J.D. Chapman²⁸, D.G. Charlton¹⁸,
 C.C. Chau¹⁵⁷, C.A. Chavez Barajas¹⁴⁸, S. Che¹¹⁰, S. Cheatham⁷², A. Chegwidden⁹⁰, S. Chekanov⁶,
 S.V. Chekulaev^{158a}, G.A. Chelkov^{65,j}, M.A. Chelstowska⁸⁹, C. Chen⁶⁴, H. Chen²⁵, K. Chen¹⁴⁷,
 S. Chen^{33c}, S. Chen¹⁵⁴, X. Chen^{33f}, Y. Chen⁶⁷, H.C. Cheng⁸⁹, Y. Cheng³¹, A. Cheplakov⁶⁵,
 E. Cheremushkina¹²⁹, R. Cherkaoui El Moursli^{134e}, V. Chernyatin^{25,*}, E. Cheu⁷, L. Chevalier¹³⁵,
 V. Chiarella⁴⁷, G. Chiarelli^{123a,123b}, G. Chiodini^{73a}, A.S. Chisholm¹⁸, R.T. Chislett⁷⁸, A. Chitan^{26b},
 M.V. Chizhov⁶⁵, K. Choi⁶¹, S. Chouridou⁹, B.K.B. Chow⁹⁹, V. Christodoulou⁷⁸,
 D. Chromek-Burckhart³⁰, J. Chudoba¹²⁶, A.J. Chuinard⁸⁷, J.J. Chwastowski³⁹, L. Chytka¹¹⁴,
 G. Ciapetti^{131a,131b}, A.K. Ciftci^{4a}, D. Cinca⁵³, V. Cindro⁷⁵, I.A. Cioara²¹, A. Ciocio¹⁵, F. Ciotto^{103a,103b},
 Z.H. Citron¹⁷¹, M. Ciubancan^{26b}, A. Clark⁴⁹, B.L. Clark⁵⁷, P.J. Clark⁴⁶, R.N. Clarke¹⁵,
 C. Clement^{145a,145b}, Y. Coadou⁸⁵, M. Cokal^{163a,163c}, A. Coccaro⁴⁹, J. Cochran⁶⁴, L. Coffey²³,
 L. Colasurdo¹⁰⁵, B. Cole³⁵, S. Cole¹⁰⁷, A.P. Colijn¹⁰⁶, J. Collot⁵⁵, T. Colombo^{58c}, G. Compostella¹⁰⁰,
 P. Conde Muiño^{125a,125b}, E. Coniavitis⁴⁸, S.H. Connell^{144b}, I.A. Connelly⁷⁷, V. Consorti⁴⁸,
 S. Constantinescu^{26b}, C. Conta^{120a,120b}, G. Conti³⁰, F. Conventi^{103a,k}, M. Cooke¹⁵, B.D. Cooper⁷⁸,
 A.M. Cooper-Sarkar¹¹⁹, T. Cornelissen¹⁷⁴, M. Corradi^{131a,131b}, F. Corriveau^{87,l}, A. Corso-Radu¹⁶²,
 A. Cortes-Gonzalez¹², G. Cortiana¹⁰⁰, G. Costa^{91a}, M.J. Costa¹⁶⁶, D. Costanzo¹³⁸, G. Cottin²⁸,
 G. Cowan⁷⁷, B.E. Cox⁸⁴, K. Cranmer¹⁰⁹, S.J. Crawley⁵³, G. Cree²⁹, S. Crépe-Renaudin⁵⁵, F. Crescioli⁸⁰,
 W.A. Cribbs^{145a,145b}, M. Crispin Ortuzar¹¹⁹, M. Cristinziani²¹, V. Croft¹⁰⁵, G. Crosetti^{37a,37b},
 T. Cuhadar Donszelmann¹³⁸, J. Cummings¹⁷⁵, M. Curatolo⁴⁷, J. Cúth⁸³, C. Cuthbert¹⁴⁹, H. Cziri¹⁴⁰,
 P. Czodrowski³, S. D'Auria⁵³, M. D'Onofrio⁷⁴, M.J. Da Cunha Sargedas De Sousa^{125a,125b}, C. Da Via⁸⁴,

W. Dabrowski^{38a}, A. Dafinca¹¹⁹, T. Dai⁸⁹, O. Dale¹⁴, F. Dallaire⁹⁴, C. Dallapiccola⁸⁶, M. Dam³⁶,
 J.R. Dandoy³¹, N.P. Dang⁴⁸, A.C. Daniells¹⁸, M. Danninger¹⁶⁷, M. Dano Hoffmann¹³⁵, V. Dao⁴⁸,
 G. Darbo^{50a}, S. Darmora⁸, J. Dassoulas³, A. Dattagupta⁶¹, W. Davey²¹, C. David¹⁶⁸, T. Davidek¹²⁸,
 E. Davies^{119,m}, M. Davies¹⁵², P. Davison⁷⁸, Y. Davygora^{58a}, E. Dawe⁸⁸, I. Dawson¹³⁸,
 R.K. Daya-Ishmukhametova⁸⁶, K. De⁸, R. de Asmundis^{103a}, A. De Benedetti¹¹², S. De Castro^{20a,20b},
 S. De Cecco⁸⁰, N. De Groot¹⁰⁵, P. de Jong¹⁰⁶, H. De la Torre⁸², F. De Lorenzi⁶⁴, D. De Pedis^{131a},
 A. De Salvo^{131a}, U. De Sanctis¹⁴⁸, A. De Santo¹⁴⁸, J.B. De Vivie De Regie¹¹⁶, W.J. Dearnaley⁷²,
 R. Debbe²⁵, C. Debenedetti¹³⁶, D.V. Dedovich⁶⁵, I. Deigaard¹⁰⁶, J. Del Peso⁸², T. Del Prete^{123a,123b},
 D. Delgove¹¹⁶, F. Deliot¹³⁵, C.M. Delitzsch⁴⁹, M. Deliyergiyev⁷⁵, A. Dell'Acqua³⁰, L. Dell'Asta²²,
 M. Dell'Orso^{123a,123b}, M. Della Pietra^{103a,k}, D. della Volpe⁴⁹, M. Delmastro⁵, P.A. Delsart⁵⁵,
 C. Deluca¹⁰⁶, D.A. DeMarco¹⁵⁷, S. Demers¹⁷⁵, M. Demichev⁶⁵, A. Demilly⁸⁰, S.P. Denisov¹²⁹,
 D. Denysiuk¹³⁵, D. Derendarz³⁹, J.E. Derkaoui^{134d}, F. Derue⁸⁰, P. Dervan⁷⁴, K. Desch²¹, C. Deterre⁴²,
 K. Dette⁴³, P.O. Deviveiros³⁰, A. Dewhurst¹³⁰, S. Dhaliwal²³, A. Di Ciaccio^{132a,132b}, L. Di Ciaccio⁵,
 A. Di Domenico^{131a,131b}, C. Di Donato^{131a,131b}, A. Di Girolamo³⁰, B. Di Girolamo³⁰, A. Di Mattia¹⁵¹,
 B. Di Micco^{133a,133b}, R. Di Nardo⁴⁷, A. Di Simone⁴⁸, R. Di Sipio¹⁵⁷, D. Di Valentino²⁹, C. Diaconu⁸⁵,
 M. Diamond¹⁵⁷, F.A. Dias⁴⁶, M.A. Diaz^{32a}, E.B. Diehl⁸⁹, J. Dietrich¹⁶, S. Diglio⁸⁵, A. Dimitrievska¹³,
 J. Dingfelder²¹, P. Dita^{26b}, S. Dita^{26b}, F. Dittus³⁰, F. Djama⁸⁵, T. Djobava^{51b}, J.I. Djuvsland^{58a},
 M.A.B. do Vale^{24c}, D. Dobos³⁰, M. Dobre^{26b}, C. Doglioni⁸¹, T. Dohmae¹⁵⁴, J. Dolejsi¹²⁸, Z. Dolezal¹²⁸,
 B.A. Dolgoshein^{97,*}, M. Donadelli^{24d}, S. Donati^{123a,123b}, P. Dondero^{120a,120b}, J. Donini³⁴, J. Dopke¹³⁰,
 A. Doria^{103a}, M.T. Dova⁷¹, A.T. Doyle⁵³, E. Drechsler⁵⁴, M. Dris¹⁰, Y. Du^{33d}, J. Duarte-Campderros¹⁵²,
 E. Dubreuil³⁴, E. Duchovni¹⁷¹, G. Duckeck⁹⁹, O.A. Ducu^{26b,85}, D. Duda¹⁰⁶, A. Dudarev³⁰, L. Duflot¹¹⁶,
 L. Duguid⁷⁷, M. Dührssen³⁰, M. Dunford^{58a}, H. Duran Yildiz^{4a}, M. Düren⁵², A. Durglishvili^{51b},
 D. Duschinger⁴⁴, B. Dutta⁴², M. Dyndal^{38a}, C. Eckardt⁴², K.M. Ecker¹⁰⁰, R.C. Edgar⁸⁹, W. Edson²,
 N.C. Edwards⁴⁶, T. Eifert³⁰, G. Eigen¹⁴, K. Einsweiler¹⁵, T. Ekelof¹⁶⁵, M. El Kacimi^{134c},
 V. Ellajosyula⁸⁵, M. Ellert¹⁶⁵, S. Elles⁵, F. Ellinghaus¹⁷⁴, A.A. Elliot¹⁶⁸, N. Ellis³⁰, J. Elmsheuser⁹⁹,
 M. Elsing³⁰, D. Emelianov¹³⁰, Y. Enari¹⁵⁴, O.C. Endner⁸³, M. Endo¹¹⁷, J. Erdmann⁴³, A. Ereditato¹⁷,
 G. Ernis¹⁷⁴, J. Ernst², M. Ernst²⁵, S. Errede¹⁶⁴, E. Ertel⁸³, M. Escalier¹¹⁶, H. Esch⁴³, C. Escobar¹²⁴,
 B. Esposito⁴⁷, A.I. Etienne¹³⁵, E. Etzion¹⁵², H. Evans⁶¹, A. Ezhilov¹²², L. Fabbri^{20a,20b}, G. Facini³¹,
 R.M. Fakhruddinov¹²⁹, S. Falciano^{131a}, R.J. Falla⁷⁸, J. Faltova¹²⁸, Y. Fang^{33a}, M. Fanti^{91a,91b}, A. Farbin⁸,
 A. Farilla^{133a}, C. Farina¹²⁴, T. Farooque¹², S. Farrell¹⁵, S.M. Farrington¹⁶⁹, P. Farthouat³⁰, F. Fassi^{134e},
 P. Fassnacht³⁰, D. Fassouliotis⁹, M. Faucci Giannelli⁷⁷, A. Favareto^{50a,50b}, L. Fayard¹¹⁶, O.L. Fedin^{122,n},
 W. Fedorko¹⁶⁷, S. Feigl¹¹⁸, L. Feligioni⁸⁵, C. Feng^{33d}, E.J. Feng³⁰, H. Feng⁸⁹, A.B. Fenyuk¹²⁹,
 L. Feremenga⁸, P. Fernandez Martinez¹⁶⁶, S. Fernandez Perez¹², J. Ferrando⁵³, A. Ferrari¹⁶⁵,
 P. Ferrari¹⁰⁶, R. Ferrari^{120a}, D.E. Ferreira de Lima⁵³, A. Ferrer¹⁶⁶, D. Ferrere⁴⁹, C. Ferretti⁸⁹,
 A. Ferretto Parodi^{50a,50b}, F. Fiedler⁸³, A. Filipčič⁷⁵, M. Filipuzzi⁴², F. Filthaut¹⁰⁵, M. Fincke-Keeler¹⁶⁸,
 K.D. Finelli¹⁴⁹, M.C.N. Fiolhais^{125a,125c}, L. Fiorini¹⁶⁶, A. Firan⁴⁰, A. Fischer², C. Fischer¹²,
 J. Fischer¹⁷⁴, W.C. Fisher⁹⁰, N. Flaschel⁴², I. Fleck¹⁴⁰, P. Fleischmann⁸⁹, G.T. Fletcher¹³⁸, G. Fletcher⁷⁶,
 R.R.M. Fletcher¹²¹, T. Flick¹⁷⁴, A. Floderus⁸¹, L.R. Flores Castillo^{60a}, M.J. Flowerdew¹⁰⁰,
 G.T. Forcolin⁸⁴, A. Formica¹³⁵, A. Forti⁸⁴, D. Fournier¹¹⁶, H. Fox⁷², S. Fracchia¹², P. Francavilla⁸⁰,
 M. Franchini^{20a,20b}, D. Francis³⁰, L. Franconi¹¹⁸, M. Franklin⁵⁷, M. Frate¹⁶², M. Fraternali^{120a,120b},
 D. Freeborn⁷⁸, S.M. Fressard-Batraneanu³⁰, F. Friedrich⁴⁴, D. Froidevaux³⁰, J.A. Frost¹¹⁹,
 C. Fukunaga¹⁵⁵, E. Fullana Torregrosa⁸³, T. Fusayasu¹⁰¹, J. Fuster¹⁶⁶, C. Gabaldon⁵⁵, O. Gabizon¹⁷⁴,
 A. Gabrielli^{20a,20b}, A. Gabrielli¹⁵, G.P. Gach^{38a}, S. Gadatsch³⁰, S. Gadomski⁴⁹, G. Gagliardi^{50a,50b},
 P. Gagnon⁶¹, C. Galea¹⁰⁵, B. Galhardo^{125a,125c}, E.J. Gallas¹¹⁹, B.J. Gallop¹³⁰, P. Gallus¹²⁷, G. Galster³⁶,
 K.K. Gan¹¹⁰, J. Gao^{33b,85}, Y. Gao⁴⁶, Y.S. Gao^{142,f}, F.M. Garay Walls⁴⁶, C. García¹⁶⁶,
 J.E. García Navarro¹⁶⁶, M. Garcia-Sciveres¹⁵, R.W. Gardner³¹, N. Garelli¹⁴², V. Garonne¹¹⁸,
 A. Gascon Bravo⁴², C. Gatti⁴⁷, A. Gaudiello^{50a,50b}, G. Gaudio^{120a}, B. Gaur¹⁴⁰, L. Gauthier⁹⁴,

I.L. Gavrilenko⁹⁵, C. Gay¹⁶⁷, G. Gaycken²¹, E.N. Gazis¹⁰, Z. Gecse¹⁶⁷, C.N.P. Gee¹³⁰,
 Ch. Geich-Gimbel²¹, M.P. Geisler^{58a}, C. Gemme^{50a}, M.H. Genest⁵⁵, C. Geng^{33b,o}, S. Gentile^{131a,131b},
 S. George⁷⁷, D. Gerbaudo¹⁶², A. Gershon¹⁵², S. Ghasemi¹⁴⁰, H. Ghazlane^{134b}, B. Giacobbe^{20a},
 S. Giagu^{131a,131b}, P. Giannetti^{123a,123b}, B. Gibbard²⁵, S.M. Gibson⁷⁷, M. Gignac¹⁶⁷, M. Gilchriese¹⁵,
 T.P.S. Gillam²⁸, D. Gillberg²⁹, G. Gilles³⁴, D.M. Gingrich^{3,d}, N. Giokaris⁹, M.P. Giordani^{163a,163c},
 F.M. Giorgi^{20a}, F.M. Giorgi¹⁶, P.F. Giraud¹³⁵, P. Giromini⁵⁷, D. Giugni^{91a}, C. Giuliani¹⁰⁰, M. Giuliani^{58b},
 B.K. Gjelsten¹¹⁸, S. Gkaitatzis¹⁵³, I. Gkialas¹⁵³, E.L. Gkoukousis¹¹⁶, L.K. Gladilin⁹⁸, C. Glasman⁸²,
 J. Glatzer³⁰, P.C.F. Glaysher⁴⁶, A. Glazov⁴², M. Goblirsch-Kolb¹⁰⁰, J. Godlewski³⁹, S. Goldfarb⁸⁹,
 T. Golling⁴⁹, D. Golubkov¹²⁹, A. Gomes^{125a,125b,125d}, R. Gonçalo^{125a},
 J. Goncalves Pinto Firmino Da Costa¹³⁵, L. Gonella²¹, S. González de la Hoz¹⁶⁶, G. Gonzalez Parra¹²,
 S. Gonzalez-Sevilla⁴⁹, L. Goossens³⁰, P.A. Gorbounov⁹⁶, H.A. Gordon²⁵, I. Gorelov¹⁰⁴, B. Gorini³⁰,
 E. Gorini^{73a,73b}, A. Gorišek⁷⁵, E. Gornicki³⁹, A.T. Goshaw⁴⁵, C. Gössling⁴³, M.I. Gostkin⁶⁵,
 C.R. Goudet¹¹⁶, D. Goujdami^{134c}, A.G. Goussiou¹³⁷, N. Govender^{144b}, E. Gozani¹⁵¹, L. Graber⁵⁴,
 I. Grabowska-Bold^{38a}, P.O.J. Gradin¹⁶⁵, P. Grafström^{20a,20b}, J. Gramling⁴⁹, E. Gramstad¹¹⁸,
 S. Grancagnolo¹⁶, V. Gratchev¹²², H.M. Gray³⁰, E. Graziani^{133a}, Z.D. Greenwood^{79,p}, C. Greife²¹,
 K. Gregersen⁷⁸, I.M. Gregor⁴², P. Grenier¹⁴², K. Grevtsov⁵, J. Griffiths⁸, A.A. Grillo¹³⁶, K. Grimm⁷²,
 S. Grinstein^{12,q}, Ph. Gris³⁴, J.-F. Grivaz¹¹⁶, S. Groh⁸³, J.P. Grohs⁴⁴, E. Gross¹⁷¹, J. Grosse-Knetter⁵⁴,
 G.C. Grossi⁷⁹, Z.J. Grout¹⁴⁸, L. Guan⁸⁹, J. Guenther¹²⁷, F. Guescini⁴⁹, D. Guest¹⁶², O. Gueta¹⁵²,
 E. Guido^{50a,50b}, T. Guillemin⁵, S. Guindon², U. Gul⁵³, C. Gumpert³⁰, J. Guo^{33e}, Y. Guo^{33b,o},
 S. Gupta¹¹⁹, G. Gustavino^{131a,131b}, P. Gutierrez¹¹², N.G. Gutierrez Ortiz⁷⁸, C. Gutschow⁴⁴, C. Guyot¹³⁵,
 C. Gwenlan¹¹⁹, C.B. Gwilliam⁷⁴, A. Haas¹⁰⁹, C. Haber¹⁵, H.K. Hadavand⁸, N. Haddad^{134e}, A. Hadeef⁸⁵,
 P. Haefner²¹, S. Hageböck²¹, Z. Hajduk³⁹, H. Hakobyan¹⁷⁶, M. Haleem⁴², J. Haley¹¹³, D. Hall¹¹⁹,
 G. Halladjian⁹⁰, G.D. Hallewell⁸⁵, K. Hamacher¹⁷⁴, P. Hamal¹¹⁴, K. Hamano¹⁶⁸, A. Hamilton^{144a},
 G.N. Hamity¹³⁸, P.G. Hamnett⁴², L. Han^{33b}, K. Hanagaki^{66,r}, K. Hanawa¹⁵⁴, M. Hance¹³⁶, B. Haney¹²¹,
 P. Hanke^{58a}, R. Hanna¹³⁵, J.B. Hansen³⁶, J.D. Hansen³⁶, M.C. Hansen²¹, P.H. Hansen³⁶, K. Hara¹⁵⁹,
 A.S. Hard¹⁷², T. Harenberg¹⁷⁴, F. Hariri¹¹⁶, S. Harkusha⁹², R.D. Harrington⁴⁶, P.F. Harrison¹⁶⁹,
 F. Hartjes¹⁰⁶, M. Hasegawa⁶⁷, Y. Hasegawa¹³⁹, A. Hasib¹¹², S. Hassani¹³⁵, S. Haug¹⁷, R. Hauser⁹⁰,
 L. Hauswald⁴⁴, M. Havranek¹²⁶, C.M. Hawkes¹⁸, R.J. Hawkins³⁰, A.D. Hawkins⁸¹, T. Hayashi¹⁵⁹,
 D. Hayden⁹⁰, C.P. Hays¹¹⁹, J.M. Hays⁷⁶, H.S. Hayward⁷⁴, S.J. Haywood¹³⁰, S.J. Head¹⁸, T. Heck⁸³,
 V. Hedberg⁸¹, L. Heelan⁸, S. Heim¹²¹, T. Heim¹⁵, B. Heinemann¹⁵, L. Heinrich¹⁰⁹, J. Hejbal¹²⁶,
 L. Helary²², S. Hellman^{145a,145b}, C. Helsens³⁰, J. Henderson¹¹⁹, R.C.W. Henderson⁷², Y. Heng¹⁷²,
 S. Henkelmann¹⁶⁷, A.M. Henriques Correia³⁰, S. Henrot-Versille¹¹⁶, G.H. Herbert¹⁶,
 Y. Hernández Jiménez¹⁶⁶, G. Herten⁴⁸, R. Hertenberger⁹⁹, L. Hervas³⁰, G.G. Hesketh⁷⁸, N.P. Hessey¹⁰⁶,
 J.W. Hetherly⁴⁰, R. Hickling⁷⁶, E. Higón-Rodríguez¹⁶⁶, E. Hill¹⁶⁸, J.C. Hill²⁸, K.H. Hiller⁴²,
 S.J. Hillier¹⁸, I. Hinchliffe¹⁵, E. Hines¹²¹, R.R. Hinman¹⁵, M. Hirose¹⁵⁶, D. Hirschbuehl¹⁷⁴, J. Hobbs¹⁴⁷,
 N. Hod¹⁰⁶, M.C. Hodgkinson¹³⁸, P. Hodgson¹³⁸, A. Hoecker³⁰, M.R. Hoefkamp¹⁰⁴, F. Hoenic⁹⁹,
 M. Hohlfield⁸³, D. Hohn²¹, T.R. Holmes¹⁵, M. Homann⁴³, T.M. Hong¹²⁴, B.H. Hooberman¹⁶⁴,
 W.H. Hopkins¹¹⁵, Y. Horii¹⁰², A.J. Horton¹⁴¹, J.-Y. Hostachy⁵⁵, S. Hou¹⁵⁰, A. Hoummada^{134a},
 J. Howard¹¹⁹, J. Howarth⁴², M. Hrabovsky¹¹⁴, I. Hristova¹⁶, J. Hrivnac¹¹⁶, T. Hryn'ova⁵,
 A. Hrynevich⁹³, C. Hsu^{144c}, P.J. Hsu^{150,s}, S.-C. Hsu¹³⁷, D. Hu³⁵, Q. Hu^{33b}, Y. Huang⁴², Z. Hubacek¹²⁷,
 F. Hubaut⁸⁵, F. Huegging²¹, T.B. Huffman¹¹⁹, E.W. Hughes³⁵, G. Hughes⁷², M. Huhtinen³⁰,
 T.A. Hülsing⁸³, N. Huseynov^{65,b}, J. Huston⁹⁰, J. Huth⁵⁷, G. Iacobucci⁴⁹, G. Iakovidis²⁵, I. Ibragimov¹⁴⁰,
 L. Iconomidou-Fayard¹¹⁶, E. Ideal¹⁷⁵, Z. Idrissi^{134e}, P. Iengo³⁰, O. Igonkina¹⁰⁶, T. Iizawa¹⁷⁰,
 Y. Ikegami⁶⁶, M. Ikeno⁶⁶, Y. Ilchenko^{31,t}, D. Iliadis¹⁵³, N. Ilic¹⁴², T. Ince¹⁰⁰, G. Introzzi^{120a,120b},
 P. Ioannou⁹, M. Iodice^{133a}, K. Iordanidou³⁵, V. Ippolito⁵⁷, A. Irlles Quiles¹⁶⁶, C. Isaksson¹⁶⁵,
 M. Ishino⁶⁸, M. Ishitsuka¹⁵⁶, R. Ishmukhametov¹¹⁰, C. Issever¹¹⁹, S. Istin^{19a}, J.M. Iturbe Ponce⁸⁴,
 R. Iuppa^{132a,132b}, J. Ivarsson⁸¹, W. Iwanski³⁹, H. Iwasaki⁶⁶, J.M. Izen⁴¹, V. Izzo^{103a}, S. Jabbar³,

B. Jackson¹²¹, M. Jackson⁷⁴, P. Jackson¹, V. Jain², K.B. Jakobi⁸³, K. Jakobs⁴⁸, S. Jakobsen³⁰,
 T. Jakoubek¹²⁶, D.O. Jamin¹¹³, D.K. Jana⁷⁹, E. Jansen⁷⁸, R. Jansky⁶², J. Janssen²¹, M. Janus⁵⁴,
 G. Jarlskog⁸¹, N. Javadov^{65,b}, T. Javůrek⁴⁸, F. Jeanneau¹³⁵, L. Jeanty¹⁵, J. Jejelava^{51a,u}, G.-Y. Jeng¹⁴⁹,
 D. Jennens⁸⁸, P. Jenni^{48,v}, J. Jentzsch⁴³, C. Jeske¹⁶⁹, S. Jézéquel⁵, H. Ji¹⁷², J. Jia¹⁴⁷, H. Jiang⁶⁴,
 Y. Jiang^{33b}, S. Jiggins⁷⁸, J. Jimenez Pena¹⁶⁶, S. Jin^{33a}, A. Jinaru^{26b}, O. Jinnouchi¹⁵⁶, P. Johansson¹³⁸,
 K.A. Johns⁷, W.J. Johnson¹³⁷, K. Jon-And^{145a,145b}, G. Jones¹⁶⁹, R.W.L. Jones⁷², S. Jones⁷, T.J. Jones⁷⁴,
 J. Jongmanns^{58a}, P.M. Jorge^{125a,125b}, J. Jovicevic^{158a}, X. Ju¹⁷², A. Juste Rozas^{12,q}, M.K. Köhler¹⁷¹,
 M. Kaci¹⁶⁶, A. Kaczmarska³⁹, M. Kado¹¹⁶, H. Kagan¹¹⁰, M. Kagan¹⁴², S.J. Kahn⁸⁵, E. Kajomovitz⁴⁵,
 C.W. Kalderon¹¹⁹, A. Kaluza⁸³, S. Kama⁴⁰, A. Kamenshchikov¹²⁹, N. Kanaya¹⁵⁴, S. Kaneti²⁸,
 V.A. Kantserov⁹⁷, J. Kanzaki⁶⁶, B. Kaplan¹⁰⁹, L.S. Kaplan¹⁷², A. Kapliy³¹, D. Kar^{144c}, K. Karakostas¹⁰,
 A. Karamaoun³, N. Karastathis^{10,106}, M.J. Kareem⁵⁴, E. Karentzos¹⁰, M. Karnevskiy⁸³, S.N. Karpov⁶⁵,
 Z.M. Karpova⁶⁵, K. Karthik¹⁰⁹, V. Kartvelishvili⁷², A.N. Karyukhin¹²⁹, K. Kasahara¹⁵⁹, L. Kashif¹⁷²,
 R.D. Kass¹¹⁰, A. Kastanas¹⁴, Y. Kataoka¹⁵⁴, C. Kato¹⁵⁴, A. Katre⁴⁹, J. Katzy⁴², K. Kawade¹⁰²,
 K. Kawagoe⁷⁰, T. Kawamoto¹⁵⁴, G. Kawamura⁵⁴, S. Kazama¹⁵⁴, V.F. Kazanin^{108,c}, R. Keeler¹⁶⁸,
 R. Kehoe⁴⁰, J.S. Keller⁴², J.J. Kempster⁷⁷, H. Keoshkerian⁸⁴, O. Kepka¹²⁶, B.P. Kerševan⁷⁵,
 S. Kersten¹⁷⁴, R.A. Keyes⁸⁷, F. Khalil-zada¹¹, H. Khandanyan^{145a,145b}, A. Khanov¹¹³,
 A.G. Kharlamov^{108,c}, T.J. Khoo²⁸, V. Khovanskii⁹⁶, E. Khramov⁶⁵, J. Khubua^{51b,w}, S. Kido⁶⁷,
 H.Y. Kim⁸, S.H. Kim¹⁵⁹, Y.K. Kim³¹, N. Kimura¹⁵³, O.M. Kind¹⁶, B.T. King⁷⁴, M. King¹⁶⁶,
 S.B. King¹⁶⁷, J. Kirk¹³⁰, A.E. Kiryunin¹⁰⁰, T. Kishimoto⁶⁷, D. Kisielewska^{38a}, F. Kiss⁴⁸, K. Kiuchi¹⁵⁹,
 O. Kivernyk¹³⁵, E. Kladiva^{143b}, M.H. Klein³⁵, M. Klein⁷⁴, U. Klein⁷⁴, K. Kleinknecht⁸³,
 P. Klimek^{145a,145b}, A. Klimentov²⁵, R. Klingenberg⁴³, J.A. Klinger¹³⁸, T. Klioutchnikova³⁰,
 E.-E. Kluge^{58a}, P. Kluit¹⁰⁶, S. Kluth¹⁰⁰, J. Knapik³⁹, E. Kneringer⁶², E.B.F.G. Knoop⁸⁵, A. Knue⁵³,
 A. Kobayashi¹⁵⁴, D. Kobayashi¹⁵⁶, T. Kobayashi¹⁵⁴, M. Kobel⁴⁴, M. Kocian¹⁴², P. Kodys¹²⁸, T. Koffas²⁹,
 E. Koffeman¹⁰⁶, L.A. Kogan¹¹⁹, S. Kohlmann¹⁷⁴, T. Kohriki⁶⁶, T. Koi¹⁴², H. Kolanoski¹⁶, M. Kolb^{58b},
 I. Koletsou⁵, A.A. Komar^{95,*}, Y. Komori¹⁵⁴, T. Kondo⁶⁶, N. Kondrashova⁴², K. Köneke⁴⁸,
 A.C. König¹⁰⁵, T. Kono^{66,x}, R. Konoplich^{109,y}, N. Konstantinidis⁷⁸, R. Kopeliansky⁶¹, S. Koperny^{38a},
 L. Köpke⁸³, A.K. Kopp⁴⁸, K. Korcyl³⁹, K. Kordas¹⁵³, A. Korn⁷⁸, A.A. Korol^{108,c}, I. Korolkov¹²,
 E.V. Korolkova¹³⁸, O. Kortner¹⁰⁰, S. Kortner¹⁰⁰, T. Kosek¹²⁸, V.V. Kostyukhin²¹, V.M. Kotov⁶⁵,
 A. Kotwal⁴⁵, A. Kourkoumeli-Charalampidi¹⁵³, C. Kourkoumelis⁹, V. Kouskoura²⁵, A. Koutsman^{158a},
 R. Kowalewski¹⁶⁸, T.Z. Kowalski^{38a}, W. Kozanecki¹³⁵, A.S. Kozhin¹²⁹, V.A. Kramarenko⁹⁸,
 G. Kramberger⁷⁵, D. Krasnopevtsev⁹⁷, M.W. Krasny⁸⁰, A. Krasznahorkay³⁰, J.K. Kraus²¹,
 A. Kravchenko²⁵, M. Kretz^{58c}, J. Kretzschmar⁷⁴, K. Kreutzfeldt⁵², P. Krieger¹⁵⁷, K. Krizka³¹,
 K. Kroeninger⁴³, H. Kroha¹⁰⁰, J. Kroll¹²¹, J. Kroseberg²¹, J. Krstic¹³, U. Kruchonak⁶⁵, H. Krüger²¹,
 N. Krumnack⁶⁴, A. Kruse¹⁷², M.C. Kruse⁴⁵, M. Kruskal²², T. Kubota⁸⁸, H. Kucuk⁷⁸, S. Kудay^{4b},
 J.T. Kuechler¹⁷⁴, S. Kuehn⁴⁸, A. Kugel^{58c}, F. Kuger¹⁷³, A. Kuhl¹³⁶, T. Kuhl⁴², V. Kukhtin⁶⁵,
 R. Kukla¹³⁵, Y. Kulchitsky⁹², S. Kuleshov^{32b}, M. Kuna^{131a,131b}, T. Kunigo⁶⁸, A. Kupco¹²⁶,
 H. Kurashige⁶⁷, Y.A. Kurochkin⁹², V. Kus¹²⁶, E.S. Kuwertz¹⁶⁸, M. Kuze¹⁵⁶, J. Kvita¹¹⁴, T. Kwan¹⁶⁸,
 D. Kyriazopoulos¹³⁸, A. La Rosa¹⁰⁰, J.L. La Rosa Navarro^{24d}, L. La Rotonda^{37a,37b}, C. Lacasta¹⁶⁶,
 F. Lacava^{131a,131b}, J. Lacey²⁹, H. Lacker¹⁶, D. Lacour⁸⁰, V.R. Lacuesta¹⁶⁶, E. Ladygin⁶⁵, R. Lafaye⁵,
 B. Laforge⁸⁰, T. Lagouri¹⁷⁵, S. Lai⁵⁴, L. Lambourne⁷⁸, S. Lammers⁶¹, C.L. Lampen⁷, W. Lampl⁷,
 E. Lançon¹³⁵, U. Landgraf⁴⁸, M.P.J. Landon⁷⁶, V.S. Lang^{58a}, J.C. Lange¹², A.J. Lankford¹⁶², F. Lanni²⁵,
 K. Lantzsich²¹, A. Lanza^{120a}, S. Laplace⁸⁰, C. Lapoire³⁰, J.F. Laporte¹³⁵, T. Lari^{91a},
 F. Lasagni Manghi^{20a,20b}, M. Lassnig³⁰, P. Laurelli⁴⁷, W. Lavrijsen¹⁵, A.T. Law¹³⁶, P. Laycock⁷⁴,
 T. Lazovich⁵⁷, O. Le Dortz⁸⁰, E. Le Guirriec⁸⁵, E. Le Menedeu¹², M. LeBlanc¹⁶⁸, T. LeCompte⁶,
 F. Ledroit-Guillon⁵⁵, C.A. Lee²⁵, S.C. Lee¹⁵⁰, L. Lee¹, G. Lefebvre⁸⁰, M. Lefebvre¹⁶⁸, F. Legger⁹⁹,
 C. Leggett¹⁵, A. Lehan⁷⁴, G. Lehmann Miotto³⁰, X. Lei⁷, W.A. Leight²⁹, A. Leisos^{153,z}, A.G. Leister¹⁷⁵,
 M.A.L. Leite^{24d}, R. Leitner¹²⁸, D. Lellouch¹⁷¹, B. Lemmer⁵⁴, K.J.C. Leney⁷⁸, T. Lenz²¹, B. Lenzi³⁰,

R. Leone⁷, S. Leone^{123a,123b}, C. Leonidopoulos⁴⁶, S. Leontsinis¹⁰, C. Leroy⁹⁴, C.G. Lester²⁸, M. Levchenko¹²², J. Levêque⁵, D. Levin⁸⁹, L.J. Levinson¹⁷¹, M. Levy¹⁸, A.M. Leyko²¹, M. Leyton⁴¹, B. Li^{33b,aa}, H. Li¹⁴⁷, H.L. Li³¹, L. Li⁴⁵, L. Li^{33e}, S. Li⁴⁵, X. Li⁸⁴, Y. Li¹⁴⁰, Z. Liang¹³⁶, H. Liao³⁴, B. Liberti^{132a}, A. Liblong¹⁵⁷, P. Lichard³⁰, K. Lie¹⁶⁴, J. Liebal²¹, W. Liebig¹⁴, C. Limbach²¹, A. Limosani¹⁴⁹, S.C. Lin^{150,ab}, T.H. Lin⁸³, B.E. Lindquist¹⁴⁷, E. Lipeles¹²¹, A. Lipniacka¹⁴, M. Lisovyi^{58b}, T.M. Liss¹⁶⁴, D. Lissauer²⁵, A. Lister¹⁶⁷, A.M. Litke¹³⁶, B. Liu^{150,ac}, D. Liu¹⁵⁰, H. Liu⁸⁹, H. Liu²⁵, J. Liu⁸⁵, J.B. Liu^{33b}, K. Liu⁸⁵, L. Liu¹⁶⁴, M. Liu⁴⁵, M. Liu^{33b}, Y.L. Liu^{33b}, Y. Liu^{33b}, M. Livan^{120a,120b}, A. Lleres⁵⁵, J. Llorente Merino⁸², S.L. Lloyd⁷⁶, F. Lo Sterzo¹⁵⁰, E. Lobodzinska⁴², P. Loch⁷, W.S. Lockman¹³⁶, F.K. Loebinger⁸⁴, A.E. Loevschall-Jensen³⁶, K.M. Loew²³, A. Loginov¹⁷⁵, T. Lohse¹⁶, K. Lohwasser⁴², M. Lokajicek¹²⁶, B.A. Long²², J.D. Long¹⁶⁴, R.E. Long⁷², K.A.Looper¹¹⁰, L. Lopes^{125a}, D. Lopez Mateos⁵⁷, B. Lopez Paredes¹³⁸, I. Lopez Paz¹², A. Lopez Solis⁸⁰, J. Lorenz⁹⁹, N. Lorenzo Martinez⁶¹, M. Losada¹⁶¹, P.J. Lösel⁹⁹, X. Lou^{33a}, A. Lounis¹¹⁶, J. Love⁶, P.A. Love⁷², H. Lu^{60a}, N. Lu⁸⁹, H.J. Lubatti¹³⁷, C. Luci^{131a,131b}, A. Lucotte⁵⁵, C. Luedtke⁴⁸, F. Luehring⁶¹, W. Lukas⁶², L. Luminari^{131a}, O. Lundberg^{145a,145b}, B. Lund-Jensen¹⁴⁶, D. Lynn²⁵, R. Lysak¹²⁶, E. Lytken⁸¹, H. Ma²⁵, L.L. Ma^{33d}, G. Maccarrone⁴⁷, A. Macchiolo¹⁰⁰, C.M. Macdonald¹³⁸, B. Maček⁷⁵, J. Machado Miguens^{121,125b}, D. Madaffari⁸⁵, R. Madar³⁴, H.J. Maddocks¹⁶⁵, W.F. Mader⁴⁴, A. Madsen⁴², J. Maeda⁶⁷, S. Maeland¹⁴, T. Maeno²⁵, A. Maevskiy⁹⁸, E. Magradze⁵⁴, J. Mahlstedt¹⁰⁶, C. Maiani¹¹⁶, C. Maidantchik^{24a}, A.A. Maier¹⁰⁰, T. Maier⁹⁹, A. Maio^{125a,125b,125d}, S. Majewski¹¹⁵, Y. Makida⁶⁶, N. Makovec¹¹⁶, B. Malaescu⁸⁰, Pa. Malecki³⁹, V.P. Maleev¹²², F. Malek⁵⁵, U. Mallik⁶³, D. Malon⁶, C. Malone¹⁴², S. Maltezos¹⁰, V.M. Malyshev¹⁰⁸, S. Malyukov³⁰, J. Mamuzic⁴², G. Mancini⁴⁷, B. Mandelli³⁰, L. Mandelli^{91a}, I. Mandić⁷⁵, J. Maneira^{125a,125b}, L. Manhaes de Andrade Filho^{24b}, J. Manjarres Ramos^{158b}, A. Mann⁹⁹, B. Mansoulie¹³⁵, R. Mantifel⁸⁷, M. Mantoani⁵⁴, S. Manzoni^{91a,91b}, L. Mapelli³⁰, L. March⁴⁹, G. Marchiori⁸⁰, M. Marcisovsky¹²⁶, M. Marjanovic¹³, D.E. Marley⁸⁹, F. Marroquim^{24a}, S.P. Marsden⁸⁴, Z. Marshall¹⁵, L.F. Marti¹⁷, S. Marti-Garcia¹⁶⁶, B. Martin⁹⁰, T.A. Martin¹⁶⁹, V.J. Martin⁴⁶, B. Martin dit Latour¹⁴, M. Martinez^{12,q}, S. Martin-Haugh¹³⁰, V.S. Martoiu^{26b}, A.C. Martyniuk⁷⁸, M. Marx¹³⁷, F. Marzano^{131a}, A. Marzin³⁰, L. Masetti⁸³, T. Mashimo¹⁵⁴, R. Mashinistov⁹⁵, J. Masik⁸⁴, A.L. Maslennikov^{108,c}, I. Massa^{20a,20b}, L. Massa^{20a,20b}, P. Mastrandrea⁵, A. Mastroberardino^{37a,37b}, T. Masubuchi¹⁵⁴, P. Mättig¹⁷⁴, J. Mattmann⁸³, J. Maurer^{26b}, S.J. Maxfield⁷⁴, D.A. Maximov^{108,c}, R. Mazini¹⁵⁰, S.M. Mazza^{91a,91b}, N.C. Mc Fadden¹⁰⁴, G. Mc Goldrick¹⁵⁷, S.P. Mc Kee⁸⁹, A. McCarn⁸⁹, R.L. McCarthy¹⁴⁷, T.G. McCarthy²⁹, K.W. McFarlane^{56,*}, J.A. Mcfayden⁷⁸, G. Mchedlidze⁵⁴, S.J. McMahon¹³⁰, R.A. McPherson^{168,l}, M. Medinnis⁴², S. Meehan¹³⁷, S. Mehlhase⁹⁹, A. Mehta⁷⁴, K. Meier^{58a}, C. Meineck⁹⁹, B. Meirose⁴¹, B.R. Mellado Garcia^{144c}, F. Meloni¹⁷, A. Mengarelli^{20a,20b}, S. Menke¹⁰⁰, E. Meoni¹⁶⁰, K.M. Mercurio⁵⁷, S. Mergelmeyer¹⁶, P. Mermod⁴⁹, L. Merola^{103a,103b}, C. Meroni^{91a}, F.S. Merritt³¹, A. Messina^{131a,131b}, J. Metcalfe⁶, A.S. Mete¹⁶², C. Meyer⁸³, C. Meyer¹²¹, J-P. Meyer¹³⁵, J. Meyer¹⁰⁶, H. Meyer Zu Theenhausen^{58a}, R.P. Middleton¹³⁰, S. Miglioranzi^{163a,163c}, L. Mijović²¹, G. Mikenberg¹⁷¹, M. Mikestikova¹²⁶, M. Mikuž⁷⁵, M. Milesi⁸⁸, A. Milic³⁰, D.W. Miller³¹, C. Mills⁴⁶, A. Milov¹⁷¹, D.A. Milstead^{145a,145b}, A.A. Minaenko¹²⁹, Y. Minami¹⁵⁴, I.A. Minashvili⁶⁵, A.I. Mincer¹⁰⁹, B. Mindur^{38a}, M. Mineev⁶⁵, Y. Ming¹⁷², L.M. Mir¹², K.P. Mistry¹²¹, T. Mitani¹⁷⁰, J. Mitrevski⁹⁹, V.A. Mitsou¹⁶⁶, A. Miucci⁴⁹, P.S. Miyagawa¹³⁸, J.U. Mjörnmark⁸¹, T. Moa^{145a,145b}, K. Mochizuki⁸⁵, S. Mohapatra³⁵, W. Mohr⁴⁸, S. Molander^{145a,145b}, R. Moles-Valls²¹, R. Monden⁶⁸, M.C. Mondragon⁹⁰, K. Mönig⁴², J. Monk³⁶, E. Monnier⁸⁵, A. Montalbano¹⁴⁷, J. Montejo Berlingen³⁰, F. Monticelli⁷¹, S. Monzani^{91a,91b}, R.W. Moore³, N. Morange¹¹⁶, D. Moreno¹⁶¹, M. Moreno Llácer⁵⁴, P. Morettini^{50a}, D. Mori¹⁴¹, T. Mori¹⁵⁴, M. Morii⁵⁷, M. Morinaga¹⁵⁴, V. Morisbak¹¹⁸, S. Moritz⁸³, A.K. Morley¹⁴⁹, G. Mornacchi³⁰, J.D. Morris⁷⁶, S.S. Mortensen³⁶, L. Morvaj¹⁴⁷, M. Mosidze^{51b}, J. Moss¹⁴², K. Motohashi¹⁵⁶, R. Mount¹⁴², E. Mountricha²⁵, S.V. Mouraviev^{95,*}, E.J.W. Moyse⁸⁶, S. Muanza⁸⁵, R.D. Mudd¹⁸, F. Mueller¹⁰⁰, J. Mueller¹²⁴, R.S.P. Mueller⁹⁹, T. Mueller²⁸, D. Muenstermann⁷²,

P. Mullen⁵³, G.A. Mullier¹⁷, F.J. Munoz Sanchez⁸⁴, J.A. Murillo Quijada¹⁸, W.J. Murray^{169,130},
 H. Musheghyan⁵⁴, A.G. Myagkov^{129,ad}, M. Myska¹²⁷, B.P. Nachman¹⁴², O. Nackenhorst⁴⁹, J. Nadal⁵⁴,
 K. Nagai¹¹⁹, R. Nagai⁶⁶, Y. Nagai⁸⁵, K. Nagano⁶⁶, Y. Nagasaka⁵⁹, K. Nagata¹⁵⁹, M. Nagel¹⁰⁰,
 E. Nagy⁸⁵, A.M. Nairz³⁰, Y. Nakahama³⁰, K. Nakamura⁶⁶, T. Nakamura¹⁵⁴, I. Nakano¹¹¹,
 H. Namasivayam⁴¹, R.F. Naranjo Garcia⁴², R. Narayan³¹, D.I. Narrias Villar^{58a}, I. Naryshkin¹²²,
 T. Naumann⁴², G. Navarro¹⁶¹, R. Nayyar⁷, H.A. Neal⁸⁹, P.Yu. Nechaeva⁹⁵, T.J. Neep⁸⁴, P.D. Nef¹⁴²,
 A. Negri^{120a,120b}, M. Negrini^{20a}, S. Nektarijevic¹⁰⁵, C. Nellist¹¹⁶, A. Nelson¹⁶², S. Nemecek¹²⁶,
 P. Nemethy¹⁰⁹, A.A. Nepomuceno^{24a}, M. Nessi^{30,ae}, M.S. Neubauer¹⁶⁴, M. Neumann¹⁷⁴,
 R.M. Neves¹⁰⁹, P. Nevski²⁵, P.R. Newman¹⁸, D.H. Nguyen⁶, R.B. Nickerson¹¹⁹, R. Nicolaidou¹³⁵,
 B. Niquevert³⁰, J. Nielsen¹³⁶, A. Nikiforov¹⁶, V. Nikolaenko^{129,ad}, I. Nikolic-Audit⁸⁰,
 K. Nikolopoulos¹⁸, J.K. Nilsen¹¹⁸, P. Nilsson²⁵, Y. Ninomiya¹⁵⁴, A. Nisati^{131a}, R. Nisius¹⁰⁰, T. Nobe¹⁵⁴,
 L. Nodulman⁶, M. Nomachi¹¹⁷, I. Nomidis²⁹, T. Nooney⁷⁶, S. Norberg¹¹², M. Nordberg³⁰,
 O. Novgorodova⁴⁴, S. Nowak¹⁰⁰, M. Nozaki⁶⁶, L. Nozka¹¹⁴, K. Ntekas¹⁰, E. Nurse⁷⁸, F. Nuti⁸⁸,
 F. O'grady⁷, D.C. O'Neil¹⁴¹, V. O'Shea⁵³, F.G. Oakham^{29,d}, H. Oberlack¹⁰⁰, T. Obermann²¹,
 J. Ocariz⁸⁰, A. Ochi⁶⁷, I. Ochoa³⁵, J.P. Ochoa-Ricoux^{32a}, S. Oda⁷⁰, S. Odaka⁶⁶, H. Ogren⁶¹, A. Oh⁸⁴,
 S.H. Oh⁴⁵, C.C. Ohm¹⁵, H. Ohman¹⁶⁵, H. Oide³⁰, H. Okawa¹⁵⁹, Y. Okumura³¹, T. Okuyama⁶⁶,
 A. Olariu^{26b}, L.F. Oleiro Seabra^{125a}, S.A. Olivares Pino⁴⁶, D. Oliveira Damazio²⁵, A. Olszewski³⁹,
 J. Olszowska³⁹, A. Onofre^{125a,125e}, K. Onogi¹⁰², P.U.E. Onyisi^{31,t}, C.J. Oram^{158a}, M.J. Oreglia³¹,
 Y. Oren¹⁵², D. Orestano^{133a,133b}, N. Orlando¹⁵³, R.S. Orr¹⁵⁷, B. Osculati^{50a,50b}, R. Ospanov⁸⁴,
 G. Otero y Garzon²⁷, H. Otono⁷⁰, M. Ouchrif^{134d}, F. Ould-Saada¹¹⁸, A. Ouraou¹³⁵, K.P. Oussoren¹⁰⁶,
 Q. Ouyang^{33a}, A. Ovcharova¹⁵, M. Owen⁵³, R.E. Owen¹⁸, V.E. Ozcan^{19a}, N. Ozturk⁸, K. Pachal¹⁴¹,
 A. Pacheco Pages¹², C. Padilla Aranda¹², M. Pagáčová⁴⁸, S. Pagan Griso¹⁵, F. Paige²⁵, P. Pais⁸⁶,
 K. Pajchel¹¹⁸, G. Palacino^{158b}, S. Palestini³⁰, M. Palka^{38b}, D. Pallin³⁴, A. Palma^{125a,125b},
 E.St. Panagiotopoulou¹⁰, C.E. Pandini⁸⁰, J.G. Panduro Vazquez⁷⁷, P. Pani^{145a,145b}, S. Panitkin²⁵,
 D. Pantea^{26b}, L. Paolozzi⁴⁹, Th.D. Papadopoulou¹⁰, K. Papageorgiou¹⁵³, A. Paramonov⁶,
 D. Paredes Hernandez¹⁷⁵, M.A. Parker²⁸, K.A. Parker¹³⁸, F. Parodi^{50a,50b}, J.A. Parsons³⁵, U. Parzefall⁴⁸,
 V. Pascuzzi¹⁵⁷, E. Pasqualucci^{131a}, S. Passaggio^{50a}, F. Pastore^{133a,133b,*}, Fr. Pastore⁷⁷, G. Pásztor²⁹,
 S. Pataria¹⁷⁴, N.D. Patel¹⁴⁹, J.R. Pater⁸⁴, T. Pauly³⁰, J. Pearce¹⁶⁸, B. Pearson¹¹², L.E. Pedersen³⁶,
 M. Pedersen¹¹⁸, S. Pedraza Lopez¹⁶⁶, R. Pedro^{125a,125b}, S.V. Peleganchuk^{108,c}, D. Pelikan¹⁶⁵, O. Penc¹²⁶,
 C. Peng^{33a}, H. Peng^{33b}, B. Penning³¹, J. Penwell⁶¹, D.V. Perepelitsa²⁵, E. Perez Codina^{158a},
 L. Perini^{91a,91b}, H. Pernegger³⁰, S. Perrella^{103a,103b}, R. Peschke⁴², V.D. Peshekhonov⁶⁵, K. Peters³⁰,
 R.F.Y. Peters⁸⁴, B.A. Petersen³⁰, T.C. Petersen³⁶, E. Petit⁵⁵, A. Petridis¹, C. Petridou¹⁵³, P. Petroff¹¹⁶,
 E. Petrolo^{131a}, F. Petrucci^{133a,133b}, N.E. Pettersson¹⁵⁶, A. Peyaud¹³⁵, R. Pezoa^{32b}, P.W. Phillips¹³⁰,
 G. Piacquadio¹⁴², E. Pianori¹⁶⁹, A. Picazio⁸⁶, E. Piccaro⁷⁶, M. Piccinini^{20a,20b}, M.A. Pickering¹¹⁹,
 R. Piegaia²⁷, J.E. Pilcher³¹, A.D. Pilkington⁸⁴, A.W.J. Pin⁸⁴, J. Pina^{125a,125b,125d},
 M. Pinamonti^{163a,163c,af}, J.L. Pinfold³, A. Pingel³⁶, S. Pires⁸⁰, H. Pirumov⁴², M. Pitt¹⁷¹, L. Plazak^{143a},
 M.-A. Pleier²⁵, V. Pleskot⁸³, E. Plotnikova⁶⁵, P. Plucinski^{145a,145b}, D. Pluth⁶⁴, R. Poettgen^{145a,145b},
 L. Poggioli¹¹⁶, D. Pohl²¹, G. Polesello^{120a}, A. Poley⁴², A. Policicchio^{37a,37b}, R. Polifka¹⁵⁷, A. Polini^{20a},
 C.S. Pollard⁵³, V. Polychronakos²⁵, K. Pommès³⁰, L. Pontecorvo^{131a}, B.G. Pope⁹⁰, G.A. Popeneciu^{26c},
 D.S. Popovic¹³, A. Poppleton³⁰, S. Pospisil¹²⁷, K. Potamianos¹⁵, I.N. Potrap⁶⁵, C.J. Potter²⁸,
 C.T. Potter¹¹⁵, G. Poulard³⁰, J. Poveda³⁰, V. Pozdnyakov⁶⁵, M.E. Pozo Astigarraga³⁰, P. Pralavorio⁸⁵,
 A. Pranko¹⁵, S. Prell⁶⁴, D. Price⁸⁴, L.E. Price⁶, M. Primavera^{73a}, S. Prince⁸⁷, M. Proissl⁴⁶,
 K. Prokofiev^{60c}, F. Prokoshin^{32b}, E. Protopapadaki¹³⁵, S. Protopopescu²⁵, J. Proudfoot⁶,
 M. Przybycien^{38a}, D. Puddu^{133a,133b}, D. Poldon¹⁴⁷, M. Purohit^{25,ag}, P. Puzo¹¹⁶, J. Qian⁸⁹, G. Qin⁵³,
 Y. Qin⁸⁴, A. Quadt⁵⁴, D.R. Quarrie¹⁵, W.B. Quayle^{163a,163b}, M. Queitsch-Maitland⁸⁴, D. Quilty⁵³,
 S. Raddum¹¹⁸, V. Radeka²⁵, V. Radescu⁴², S.K. Radhakrishnan¹⁴⁷, P. Radloff¹¹⁵, P. Rados⁸⁸,
 F. Ragusa^{91a,91b}, G. Rahal¹⁷⁷, S. Rajagopalan²⁵, M. Rammensee³⁰, C. Rangel-Smith¹⁶⁵, F. Rauscher⁹⁹,

S. Rave⁸³, T. Ravenscroft⁵³, M. Raymond³⁰, A.L. Read¹¹⁸, N.P. Readioff⁷⁴, D.M. Rebuzzi^{120a,120b},
 A. Redelbach¹⁷³, G. Redlinger²⁵, R. Reece¹³⁶, K. Reeves⁴¹, L. Rehnisch¹⁶, J. Reichert¹²¹, H. Reisin²⁷,
 C. Rembser³⁰, H. Ren^{33a}, M. Rescigno^{131a}, S. Resconi^{91a}, O.L. Rezanova^{108.c}, P. Reznicek¹²⁸,
 R. Rezvani⁹⁴, R. Richter¹⁰⁰, S. Richter⁷⁸, E. Richter-Was^{38b}, O. Ricken²¹, M. Ridel⁸⁰, P. Rieck¹⁶,
 C.J. Riegel¹⁷⁴, J. Rieger⁵⁴, O. Rifki¹¹², M. Rijssenbeek¹⁴⁷, A. Rimoldi^{120a,120b}, L. Rinaldi^{20a}, B. Ristic⁴⁹,
 E. Ritsch³⁰, I. Riu¹², F. Rizatdinova¹¹³, E. Rizvi⁷⁶, S.H. Robertson^{87,l}, A. Robichaud-Veronneau⁸⁷,
 D. Robinson²⁸, J.E.M. Robinson⁴², A. Robson⁵³, C. Roda^{123a,123b}, Y. Rodina⁸⁵, A. Rodriguez Perez¹²,
 S. Roe³⁰, C.S. Rogan⁵⁷, O. Røhne¹¹⁸, A. Romaniouk⁹⁷, M. Romano^{20a,20b}, S.M. Romano Saez³⁴,
 E. Romero Adam¹⁶⁶, N. Rompotis¹³⁷, M. Ronzani⁴⁸, L. Roos⁸⁰, E. Ros¹⁶⁶, S. Rosati^{131a}, K. Rosbach⁴⁸,
 P. Rose¹³⁶, O. Rosenthal¹⁴⁰, V. Rossetti^{145a,145b}, E. Rossi^{103a,103b}, L.P. Rossi^{50a}, J.H.N. Rosten²⁸,
 R. Rosten¹³⁷, M. Rotaru^{26b}, I. Roth¹⁷¹, J. Rothberg¹³⁷, D. Rousseau¹¹⁶, C.R. Royon¹³⁵, A. Rozanov⁸⁵,
 Y. Rozen¹⁵¹, X. Ruan^{144c}, F. Rubbo¹⁴², I. Rubinskiy⁴², V.I. Rud⁹⁸, M.S. Rudolph¹⁵⁷, F. Rühr⁴⁸,
 A. Ruiz-Martinez³⁰, Z. Rurikova⁴⁸, N.A. Rusakovich⁶⁵, A. Ruschke⁹⁹, H.L. Russell¹³⁷,
 J.P. Rutherford⁷, N. Ruthmann³⁰, Y.F. Ryabov¹²², M. Rybar¹⁶⁴, G. Rybkin¹¹⁶, N.C. Ryder¹¹⁹,
 A. Ryzhov¹²⁹, A.F. Saavedra¹⁴⁹, G. Sabato¹⁰⁶, S. Sacerdoti²⁷, H.F-W. Sadrozinski¹³⁶, R. Sadykov⁶⁵,
 F. Safai Tehrani^{131a}, P. Saha¹⁰⁷, M. Sahinsoy^{58a}, M. Saimpert¹³⁵, T. Saito¹⁵⁴, H. Sakamoto¹⁵⁴,
 Y. Sakurai¹⁷⁰, G. Salamanna^{133a,133b}, A. Salamon^{132a}, J.E. Salazar Loyola^{32b}, D. Salek¹⁰⁶,
 P.H. Sales De Bruin¹³⁷, D. Salihagic¹⁰⁰, A. Salnikov¹⁴², J. Salt¹⁶⁶, D. Salvatore^{37a,37b}, F. Salvatore¹⁴⁸,
 A. Salvucci^{60a}, A. Salzburger³⁰, D. Sammel⁴⁸, D. Sampsonidis¹⁵³, A. Sanchez^{103a,103b}, J. Sánchez¹⁶⁶,
 V. Sanchez Martinez¹⁶⁶, H. Sandaker¹¹⁸, R.L. Sandbach⁷⁶, H.G. Sander⁸³, M.P. Sanders⁹⁹,
 M. Sandhoff¹⁷⁴, C. Sandoval¹⁶¹, R. Sandstroem¹⁰⁰, D.P.C. Sankey¹³⁰, M. Sannino^{50a,50b}, A. Sansoni⁴⁷,
 C. Santoni³⁴, R. Santonico^{132a,132b}, H. Santos^{125a}, I. Santoyo Castillo¹⁴⁸, K. Sapp¹²⁴, A. Saponov⁶⁵,
 J.G. Saraiva^{125a,125d}, B. Sarrazin²¹, O. Sasaki⁶⁶, Y. Sasaki¹⁵⁴, K. Sato¹⁵⁹, G. Sauvage^{5,*}, E. Sauvan⁵,
 G. Savage⁷⁷, P. Savard^{157,d}, C. Sawyer¹³⁰, L. Sawyer^{79,p}, J. Saxon³¹, C. Sbarra^{20a}, A. Sbrizzi^{20a,20b},
 T. Scanlon⁷⁸, D.A. Scannicchio¹⁶², M. Scarcella¹⁴⁹, V. Scarfone^{37a,37b}, J. Schaarschmidt¹⁷¹,
 P. Schacht¹⁰⁰, D. Schaefer³⁰, R. Schaefer⁴², J. Schaeffer⁸³, S. Schaepe²¹, S. Schaezel^{58b}, U. Schäfer⁸³,
 A.C. Schaffer¹¹⁶, D. Schaile⁹⁹, R.D. Schamberger¹⁴⁷, V. Scharf^{58a}, V.A. Schegelsky¹²², D. Scheirich¹²⁸,
 M. Schernau¹⁶², C. Schiavi^{50a,50b}, C. Schillo⁴⁸, M. Schioppa^{37a,37b}, S. Schlenker³⁰, K. Schmieden³⁰,
 C. Schmitt⁸³, S. Schmitt^{58b}, S. Schmitt⁴², S. Schmitz⁸³, B. Schneider^{158a}, Y.J. Schnellbach⁷⁴,
 U. Schnoor⁴⁸, L. Schoeffel¹³⁵, A. Schoening^{58b}, B.D. Schoenrock⁹⁰, E. Schopf²¹,
 A.L.S. Schorlemmer⁵⁴, M. Schott⁸³, D. Schouten^{158a}, J. Schovancova⁸, S. Schramm⁴⁹, M. Schreyer¹⁷³,
 N. Schuh⁸³, M.J. Schultens²¹, H.-C. Schultz-Coulon^{58a}, H. Schulz¹⁶, M. Schumacher⁴⁸,
 B.A. Schumm¹³⁶, Ph. Schune¹³⁵, C. Schwanenberger⁸⁴, A. Schwartzman¹⁴², T.A. Schwarz⁸⁹,
 Ph. Schwegler¹⁰⁰, H. Schweiger⁸⁴, Ph. Schwemling¹³⁵, R. Schwienhorst⁹⁰, J. Schwindling¹³⁵,
 T. Schwindt²¹, G. Sciolla²³, F. Scuri^{123a,123b}, F. Scutti⁸⁸, J. Searcy⁸⁹, P. Seema²¹, S.C. Seidel¹⁰⁴,
 A. Seiden¹³⁶, F. Seifert¹²⁷, J.M. Seixas^{24a}, G. Sekhniaidze^{103a}, K. Sekhon⁸⁹, S.J. Sekula⁴⁰,
 D.M. Seliverstov^{122,*}, N. Semprini-Cesari^{20a,20b}, C. Serfon³⁰, L. Serin¹¹⁶, L. Serkin^{163a,163b},
 M. Sessa^{133a,133b}, R. Seuster^{158a}, H. Severini¹¹², T. Sfiligoj⁷⁵, F. Sforza³⁰, A. Sfyrla⁴⁹, E. Shabalina⁵⁴,
 N.W. Shaikh^{145a,145b}, L.Y. Shan^{33a}, R. Shang¹⁶⁴, J.T. Shank²², M. Shapiro¹⁵, P.B. Shatalov⁹⁶,
 K. Shaw^{163a,163b}, S.M. Shaw⁸⁴, A. Shcherbakova^{145a,145b}, C.Y. Shehu¹⁴⁸, P. Sherwood⁷⁸, L. Shi^{150,ah},
 S. Shimizu⁶⁷, C.O. Shimmin¹⁶², M. Shimojima¹⁰¹, M. Shiyakova⁶⁵, A. Shmeleva⁹⁵, D. Shoaleh Saadi⁹⁴,
 M.J. Shochet³¹, S. Shojaii^{91a,91b}, S. Shrestha¹¹⁰, E. Shulga⁹⁷, M.A. Shupe⁷, P. Sicho¹²⁶, P.E. Sidebo¹⁴⁶,
 O. Sidiropoulou¹⁷³, D. Sidorov¹¹³, A. Sidoti^{20a,20b}, F. Siegert⁴⁴, Dj. Sijacki¹³, J. Silva^{125a,125d},
 S.B. Silverstein^{145a}, V. Simak¹²⁷, O. Simard⁵, Lj. Simic¹³, S. Simion¹¹⁶, E. Simioni⁸³, B. Simmons⁷⁸,
 D. Simon³⁴, M. Simon⁸³, P. Sinervo¹⁵⁷, N.B. Sinev¹¹⁵, M. Sioli^{20a,20b}, G. Siragusa¹⁷³,
 S.Yu. Sivoklov⁹⁸, J. Sjölin^{145a,145b}, T.B. Sjursen¹⁴, M.B. Skinner⁷², H.P. Skottowe⁵⁷, P. Skubic¹¹²,
 M. Slater¹⁸, T. Slavicek¹²⁷, M. Slawinska¹⁰⁶, K. Sliwa¹⁶⁰, V. Smakhtin¹⁷¹, B.H. Smart⁴⁶, L. Smestad¹⁴,

S. Yu. Smirnov⁹⁷, Y. Smirnov⁹⁷, L.N. Smirnova^{98.ai}, O. Smirnova⁸¹, M.N.K. Smith³⁵, R.W. Smith³⁵,
 M. Smizanska⁷², K. Smolek¹²⁷, A.A. Snesarev⁹⁵, G. Snidero⁷⁶, S. Snyder²⁵, R. Sobie^{168.l}, F. Socher⁴⁴,
 A. Soffer¹⁵², D.A. Soh^{150.ah}, G. Sokhrannyi⁷⁵, C.A. Solans Sanchez³⁰, M. Solar¹²⁷, E. Yu. Soldatov⁹⁷,
 U. Soldevila¹⁶⁶, A.A. Solodkov¹²⁹, A. Soloshenko⁶⁵, O.V. Solovyanov¹²⁹, V. Solovyev¹²², P. Sommer⁴⁸,
 H.Y. Song^{33b.aa}, N. Soni¹, A. Sood¹⁵, A. Sopczak¹²⁷, V. Sopko¹²⁷, V. Sorin¹², D. Sosa^{58b},
 C.L. Sotiropoulou^{123a,123b}, R. Soualah^{163a,163c}, A.M. Soukharev^{108.c}, D. South⁴², B.C. Sowden⁷⁷,
 S. Spagnolo^{73a,73b}, M. Spalla^{123a,123b}, M. Spangenberg¹⁶⁹, F. Spanò⁷⁷, D. Sperlich¹⁶, F. Spettel¹⁰⁰,
 R. Spighi^{20a}, G. Spigo³⁰, L.A. Spiller⁸⁸, M. Spousta¹²⁸, R.D. St. Denis^{53,*}, A. Stabile^{91a}, S. Staerz³⁰,
 J. Stahlman¹²¹, R. Stamen^{58a}, S. Stamm¹⁶, E. Stanecka³⁹, R.W. Stanek⁶, C. Stanescu^{133a},
 M. Stanescu-Bellu⁴², M.M. Stanitzki⁴², S. Stapnes¹¹⁸, E.A. Starchenko¹²⁹, G.H. Stark³¹, J. Stark⁵⁵,
 P. Staroba¹²⁶, P. Starovoitov^{58a}, R. Staszewski³⁹, P. Steinberg²⁵, B. Stelzer¹⁴¹, H.J. Stelzer³⁰,
 O. Stelzer-Chilton^{158a}, H. Stenzel⁵², G.A. Stewart⁵³, J.A. Stillings²¹, M.C. Stockton⁸⁷, M. Stoebe⁸⁷,
 G. Stoicea^{26b}, P. Stolte⁵⁴, S. Stonjek¹⁰⁰, A.R. Stradling⁸, A. Straessner⁴⁴, M.E. Stramaglia¹⁷,
 J. Strandberg¹⁴⁶, S. Strandberg^{145a,145b}, A. Strandlie¹¹⁸, M. Strauss¹¹², P. Strizenc^{143b}, R. Ströhmer¹⁷³,
 D.M. Strom¹¹⁵, R. Stroynowski⁴⁰, A. Strubig¹⁰⁵, S.A. Stucci¹⁷, B. Stugu¹⁴, N.A. Styles⁴², D. Su¹⁴²,
 J. Su¹²⁴, R. Subramaniam⁷⁹, S. Suche^{58a}, Y. Sugaya¹¹⁷, M. Suk¹²⁷, V.V. Sulin⁹⁵, S. Sultansoy^{4c},
 T. Sumida⁶⁸, S. Sun⁵⁷, X. Sun^{33a}, J.E. Sundermann⁴⁸, K. Suruliz¹⁴⁸, G. Susinno^{37a,37b}, M.R. Sutton¹⁴⁸,
 S. Suzuki⁶⁶, M. Svatos¹²⁶, M. Swiatlowski³¹, I. Sykora^{143a}, T. Sykora¹²⁸, D. Ta⁴⁸, C. Taccini^{133a,133b},
 K. Tackmann⁴², J. Taenzer¹⁵⁷, A. Taffard¹⁶², R. Tafirout^{158a}, N. Taiblum¹⁵², H. Takai²⁵, R. Takashima⁶⁹,
 H. Takeda⁶⁷, T. Takeshita¹³⁹, Y. Takubo⁶⁶, M. Talby⁸⁵, A.A. Talyshev^{108.c}, J.Y.C. Tam¹⁷³, K.G. Tan⁸⁸,
 J. Tanaka¹⁵⁴, R. Tanaka¹¹⁶, S. Tanaka⁶⁶, B.B. Tannenwald¹¹⁰, S. Tapia Araya^{32b}, S. Tapprogge⁸³,
 S. Tarem¹⁵¹, G.F. Tartarelli^{91a}, P. Tas¹²⁸, M. Tasevsky¹²⁶, T. Tashiro⁶⁸, E. Tassi^{37a,37b},
 A. Tavares Delgado^{125a,125b}, Y. Tayalati^{134d}, A.C. Taylor¹⁰⁴, G.N. Taylor⁸⁸, P.T.E. Taylor⁸⁸,
 W. Taylor^{158b}, F.A. Teischinger³⁰, P. Teixeira-Dias⁷⁷, K.K. Temming⁴⁸, D. Temple¹⁴¹, H. Ten Kate³⁰,
 P.K. Teng¹⁵⁰, J.J. Teoh¹¹⁷, F. Tepel¹⁷⁴, S. Terada⁶⁶, K. Terashi¹⁵⁴, J. Terron⁸², S. Terzo¹⁰⁰, M. Testa⁴⁷,
 R.J. Teuscher^{157.l}, T. Theveneaux-Pelzer⁸⁵, J.P. Thomas¹⁸, J. Thomas-Wilsker⁷⁷, E.N. Thompson³⁵,
 P.D. Thompson¹⁸, R.J. Thompson⁸⁴, A.S. Thompson⁵³, L.A. Thomsen¹⁷⁵, E. Thomson¹²¹,
 M. Thomson²⁸, M.J. Tibbetts¹⁵, R.E. Ticse Torres⁸⁵, V.O. Tikhomirov^{95.aj}, Yu.A. Tikhonov^{108.c},
 S. Timoshenko⁹⁷, E. Tiouchichine⁸⁵, P. Tipton¹⁷⁵, S. Tisserant⁸⁵, K. Todome¹⁵⁶, T. Todorov^{5,*},
 S. Todorova-Nova¹²⁸, J. Tojo⁷⁰, S. Tokár^{143a}, K. Tokushuku⁶⁶, E. Tolley⁵⁷, L. Tomlinson⁸⁴,
 M. Tomoto¹⁰², L. Tompkins^{142.ak}, K. Toms¹⁰⁴, B. Tong⁵⁷, E. Torrence¹¹⁵, H. Torres¹⁴¹,
 E. Torrón Pastor¹³⁷, J. Toth^{85.al}, F. Touchard⁸⁵, D.R. Tovey¹³⁸, T. Trefzger¹⁷³, L. Tremblet³⁰, A. Tricoli³⁰,
 I.M. Trigger^{158a}, S. Trincaz-Duvoid⁸⁰, M.F. Tripiana¹², W. Trischuk¹⁵⁷, B. Trocmé⁵⁵, A. Trofymov⁴²,
 C. Troncon^{91a}, M. Trotter-McDonald¹⁵, M. Trovatelli¹⁶⁸, L. Truong^{163a,163c}, M. Trzebinski³⁹,
 A. Trzupek³⁹, J.C.-L. Tseng¹¹⁹, P.V. Tsiarshka⁹², G. Tsipolitis¹⁰, N. Tsirintanis⁹, S. Tsiskaridze¹²,
 V. Tsiskaridze⁴⁸, E.G. Tskhadadze^{51a}, K.M. Tsui^{60a}, I.I. Tsukerman⁹⁶, V. Tsulaia¹⁵, S. Tsuno⁶⁶,
 D. Tsybychev¹⁴⁷, A. Tudorache^{26b}, V. Tudorache^{26b}, A.N. Tuna⁵⁷, S.A. Tuppiti^{20a,20b},
 S. Turchikhin^{98.ai}, D. Turecek¹²⁷, D. Turgeman¹⁷¹, R. Turra^{91a,91b}, A.J. Turvey⁴⁰, P.M. Tuts³⁵,
 M. Tylmad^{145a,145b}, M. Tyndel¹³⁰, I. Ueda¹⁵⁴, R. Ueno²⁹, M. Ughetto^{145a,145b}, F. Ukegawa¹⁵⁹, G. Unal³⁰,
 A. Undrus²⁵, G. Unel¹⁶², F.C. Ungaro⁸⁸, Y. Unno⁶⁶, C. Unverdorben⁹⁹, J. Urban^{143b}, P. Urquijo⁸⁸,
 P. Urrejola⁸³, G. Usai⁸, A. Usanova⁶², L. Vacavant⁸⁵, V. Vacek¹²⁷, B. Vachon⁸⁷, C. Valderanis⁸³,
 N. Valencic¹⁰⁶, S. Valentinetti^{20a,20b}, A. Valero¹⁶⁶, L. Valery¹², S. Valkar¹²⁸, S. Vallecorsa⁴⁹,
 J.A. Valls Ferrer¹⁶⁶, W. Van Den Wollenberg¹⁰⁶, P.C. Van Der Deijl¹⁰⁶, R. van der Geer¹⁰⁶,
 H. van der Graaf¹⁰⁶, N. van Eldik¹⁵¹, P. van Gemmeren⁶, J. Van Nieuwkoop¹⁴¹, I. van Vulpen¹⁰⁶,
 M.C. van Woerden³⁰, M. Vanadia^{131a,131b}, W. Vandelli³⁰, R. Vanguri¹²¹, A. Vaniachine⁶,
 G. Vardanyan¹⁷⁶, R. Vari^{131a}, E.W. Varnes⁷, T. Varol⁴⁰, D. Varouchas⁸⁰, A. Vartapetian⁸, K.E. Varvell¹⁴⁹,
 F. Vazeille³⁴, T. Vazquez Schroeder⁸⁷, J. Veatch⁷, L.M. Veloce¹⁵⁷, F. Veloso^{125a,125c}, S. Veneziano^{131a},

A. Ventura^{73a,73b}, M. Venturi¹⁶⁸, N. Venturi¹⁵⁷, A. Venturini²³, V. Vercesi^{120a}, M. Verducci^{131a,131b}, W. Verkerke¹⁰⁶, J.C. Vermeulen¹⁰⁶, A. Vest^{44,am}, M.C. Vetterli^{141,d}, O. Viazlo⁸¹, I. Vichou¹⁶⁴, T. Vickey¹³⁸, O.E. Vickey Boeriu¹³⁸, G.H.A. Viehhauser¹¹⁹, S. Viel¹⁵, R. Vigne⁶², M. Villa^{20a,20b}, M. Villaplana Perez^{91a,91b}, E. Vilucchi⁴⁷, M.G. Vinciter²⁹, V.B. Vinogradov⁶⁵, I. Vivarelli¹⁴⁸, S. Vlachos¹⁰, M. Vlasak¹²⁷, M. Vogel^{32a}, P. Vokac¹²⁷, G. Volpi^{123a,123b}, M. Volpi⁸⁸, H. von der Schmitt¹⁰⁰, E. von Toerne²¹, V. Vorobel¹²⁸, K. Vorobev⁹⁷, M. Vos¹⁶⁶, R. Voss³⁰, J.H. Vossebeld⁷⁴, N. Vranjes¹³, M. Vranjes Milosavljevic¹³, V. Vrba¹²⁶, M. Vreeswijk¹⁰⁶, R. Vuillermet³⁰, I. Vukotic³¹, Z. Vykydal¹²⁷, P. Wagner²¹, W. Wagner¹⁷⁴, H. Wahlberg⁷¹, S. Wahrenmund⁴⁴, J. Wakabayashi¹⁰², J. Walder⁷², R. Walker⁹⁹, W. Walkowiak¹⁴⁰, V. Wallangen^{145a,145b}, C. Wang¹⁵⁰, C. Wang^{33d,85}, F. Wang¹⁷², H. Wang¹⁵, H. Wang⁴⁰, J. Wang⁴², J. Wang¹⁴⁹, K. Wang⁸⁷, R. Wang⁶, S.M. Wang¹⁵⁰, T. Wang²¹, T. Wang³⁵, X. Wang¹⁷⁵, C. Wanotayaroj¹¹⁵, A. Warburton⁸⁷, C.P. Ward²⁸, D.R. Wardrope⁷⁸, A. Washbrook⁴⁶, P.M. Watkins¹⁸, A.T. Watson¹⁸, I.J. Watson¹⁴⁹, M.F. Watson¹⁸, G. Watts¹³⁷, S. Watts⁸⁴, B.M. Waugh⁷⁸, S. Webb⁸⁴, M.S. Weber¹⁷, S.W. Weber¹⁷³, J.S. Webster⁶, A.R. Weidberg¹¹⁹, B. Weinert⁶¹, J. Weingarten⁵⁴, C. Weiser⁴⁸, H. Weits¹⁰⁶, P.S. Wells³⁰, T. Wenaus²⁵, T. Wengler³⁰, S. Wenig³⁰, N. Vermes²¹, M. Werner⁴⁸, P. Werner³⁰, M. Wessels^{58a}, J. Wetter¹⁶⁰, K. Whalen¹¹⁵, A.M. Wharton⁷², A. White⁸, M.J. White¹, R. White^{32b}, S. White^{123a,123b}, D. Whiteson¹⁶², F.J. Wickens¹³⁰, W. Wiedenmann¹⁷², M. Wielers¹³⁰, P. Wienemann²¹, C. Wiglesworth³⁶, L.A.M. Wiik-Fuchs²¹, A. Wildauer¹⁰⁰, H.G. Wilkens³⁰, H.H. Williams¹²¹, S. Williams¹⁰⁶, C. Willis⁹⁰, S. Willocq⁸⁶, J.A. Wilson¹⁸, I. Wingerter-Seez⁵, F. Winklmeier¹¹⁵, B.T. Winter²¹, M. Wittgen¹⁴², J. Wittkowski⁹⁹, S.J. Wollstadt⁸³, M.W. Wolter³⁹, H. Wolters^{125a,125c}, B.K. Wosiek³⁹, J. Wotschack³⁰, M.J. Woudstra⁸⁴, K.W. Wozniak³⁹, M. Wu⁵⁵, M. Wu³¹, S.L. Wu¹⁷², X. Wu⁴⁹, Y. Wu⁸⁹, T.R. Wyatt⁸⁴, B.M. Wynne⁴⁶, S. Xella³⁶, D. Xu^{33a}, L. Xu²⁵, B. Yabsley¹⁴⁹, S. Yacoob^{144a}, R. Yakabe⁶⁷, D. Yamaguchi¹⁵⁶, Y. Yamaguchi¹¹⁷, A. Yamamoto⁶⁶, S. Yamamoto¹⁵⁴, T. Yamanaka¹⁵⁴, K. Yamauchi¹⁰², Y. Yamazaki⁶⁷, Z. Yan²², H. Yang^{33e}, H. Yang¹⁷², Y. Yang¹⁵⁰, Z. Yang¹⁴, W-M. Yao¹⁵, Y.C. Yap⁸⁰, Y. Yasu⁶⁶, E. Yatsenko⁵, K.H. Yau Wong²¹, J. Ye⁴⁰, S. Ye²⁵, I. Yeletsikh⁶⁵, A.L. Yen⁵⁷, E. Yildirim⁴², K. Yorita¹⁷⁰, R. Yoshida⁶, K. Yoshihara¹²¹, C. Young¹⁴², C.J.S. Young³⁰, S. Youssef²², D.R. Yu¹⁵, J. Yu⁸, J.M. Yu⁸⁹, J. Yu⁶⁴, L. Yuan⁶⁷, S.P.Y. Yuen²¹, I. Yusuff^{28,an}, B. Zabinski³⁹, R. Zaidan^{33d}, A.M. Zaitsev^{129,ad}, N. Zakharchuk⁴², J. Zalieckas¹⁴, A. Zaman¹⁴⁷, S. Zambito⁵⁷, L. Zanello^{131a,131b}, D. Zanzi⁸⁸, C. Zeitnitz¹⁷⁴, M. Zeman¹²⁷, A. Zemla^{38a}, J.C. Zeng¹⁶⁴, Q. Zeng¹⁴², K. Zengel²³, O. Zenin¹²⁹, T. Ženiš^{143a}, D. Zerwas¹¹⁶, D. Zhang⁸⁹, F. Zhang¹⁷², G. Zhang^{33b,aa}, H. Zhang^{33c}, J. Zhang⁶, L. Zhang⁴⁸, R. Zhang²¹, R. Zhang^{33b,aa}, X. Zhang^{33d}, Z. Zhang¹¹⁶, X. Zhao⁴⁰, Y. Zhao^{33d,116}, Z. Zhao^{33b}, A. Zhemchugov⁶⁵, J. Zhong¹¹⁹, B. Zhou⁸⁹, C. Zhou⁴⁵, L. Zhou³⁵, L. Zhou⁴⁰, M. Zhou¹⁴⁷, N. Zhou^{33f}, C.G. Zhu^{33d}, H. Zhu^{33a}, J. Zhu⁸⁹, Y. Zhu^{33b}, X. Zhuang^{33a}, K. Zhukov⁹⁵, A. Zibell¹⁷³, D. Ziemska⁶¹, N.I. Zimine⁶⁵, C. Zimmermann⁸³, S. Zimmermann⁴⁸, Z. Zinonos⁵⁴, M. Zinser⁸³, M. Ziolkowski¹⁴⁰, L. Živković¹³, G. Zobernig¹⁷², A. Zoccoli^{20a,20b}, M. zur Nedden¹⁶, G. Zurzolo^{103a,103b}, L. Zwalinski³⁰.

¹ Department of Physics, University of Adelaide, Adelaide, Australia

² Physics Department, SUNY Albany, Albany NY, United States of America

³ Department of Physics, University of Alberta, Edmonton AB, Canada

⁴ (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c)

Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

⁵ LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France

⁶ High Energy Physics Division, Argonne National Laboratory, Argonne IL, United States of America

⁷ Department of Physics, University of Arizona, Tucson AZ, United States of America

⁸ Department of Physics, The University of Texas at Arlington, Arlington TX, United States of America

⁹ Physics Department, University of Athens, Athens, Greece

- ¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece
- ¹¹ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
- ¹² Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain, Spain
- ¹³ Institute of Physics, University of Belgrade, Belgrade, Serbia
- ¹⁴ Department for Physics and Technology, University of Bergen, Bergen, Norway
- ¹⁵ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA, United States of America
- ¹⁶ Department of Physics, Humboldt University, Berlin, Germany
- ¹⁷ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
- ¹⁸ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- ¹⁹ ^(a) Department of Physics, Bogazici University, Istanbul; ^(b) Department of Physics Engineering, Gaziantep University, Gaziantep; ^(c) Department of Physics, Dogus University, Istanbul, Turkey
- ²⁰ ^(a) INFN Sezione di Bologna; ^(b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
- ²¹ Physikalisches Institut, University of Bonn, Bonn, Germany
- ²² Department of Physics, Boston University, Boston MA, United States of America
- ²³ Department of Physics, Brandeis University, Waltham MA, United States of America
- ²⁴ ^(a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; ^(b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora; ^(c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; ^(d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
- ²⁵ Physics Department, Brookhaven National Laboratory, Upton NY, United States of America
- ²⁶ ^(a) Transilvania University of Brasov, Brasov, Romania; ^(b) National Institute of Physics and Nuclear Engineering, Bucharest; ^(c) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca; ^(d) University Politehnica Bucharest, Bucharest; ^(e) West University in Timisoara, Timisoara, Romania
- ²⁷ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- ²⁸ Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- ²⁹ Department of Physics, Carleton University, Ottawa ON, Canada
- ³⁰ CERN, Geneva, Switzerland
- ³¹ Enrico Fermi Institute, University of Chicago, Chicago IL, United States of America
- ³² ^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; ^(b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- ³³ ^(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; ^(b) Department of Modern Physics, University of Science and Technology of China, Anhui; ^(c) Department of Physics, Nanjing University, Jiangsu; ^(d) School of Physics, Shandong University, Shandong; ^(e) Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai; ^(f) Physics Department, Tsinghua University, Beijing 100084, China
- ³⁴ Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
- ³⁵ Nevis Laboratory, Columbia University, Irvington NY, United States of America
- ³⁶ Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
- ³⁷ ^(a) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; ^(b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
- ³⁸ ^(a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; ^(b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland

- 39 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- 40 Physics Department, Southern Methodist University, Dallas TX, United States of America
- 41 Physics Department, University of Texas at Dallas, Richardson TX, United States of America
- 42 DESY, Hamburg and Zeuthen, Germany
- 43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- 44 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- 45 Department of Physics, Duke University, Durham NC, United States of America
- 46 SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- 47 INFN Laboratori Nazionali di Frascati, Frascati, Italy
- 48 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
- 49 Section de Physique, Université de Genève, Geneva, Switzerland
- 50 ^(a) INFN Sezione di Genova; ^(b) Dipartimento di Fisica, Università di Genova, Genova, Italy
- 51 ^(a) E. Andronikashvili Institute of Physics, Iv. Javakishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- 52 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- 53 SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- 54 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
- 55 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
- 56 Department of Physics, Hampton University, Hampton VA, United States of America
- 57 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA, United States of America
- 58 ^(a) Kirchoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(c) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
- 59 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- 60 ^(a) Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, The University of Hong Kong, Hong Kong; ^(c) Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- 61 Department of Physics, Indiana University, Bloomington IN, United States of America
- 62 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- 63 University of Iowa, Iowa City IA, United States of America
- 64 Department of Physics and Astronomy, Iowa State University, Ames IA, United States of America
- 65 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
- 66 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- 67 Graduate School of Science, Kobe University, Kobe, Japan
- 68 Faculty of Science, Kyoto University, Kyoto, Japan
- 69 Kyoto University of Education, Kyoto, Japan
- 70 Department of Physics, Kyushu University, Fukuoka, Japan
- 71 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- 72 Physics Department, Lancaster University, Lancaster, United Kingdom
- 73 ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- 74 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- 75 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
- 76 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- 77 Department of Physics, Royal Holloway University of London, Surrey, United Kingdom

- ⁷⁸ Department of Physics and Astronomy, University College London, London, United Kingdom
- ⁷⁹ Louisiana Tech University, Ruston LA, United States of America
- ⁸⁰ Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
- ⁸¹ Fysiska institutionen, Lunds universitet, Lund, Sweden
- ⁸² Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain
- ⁸³ Institut für Physik, Universität Mainz, Mainz, Germany
- ⁸⁴ School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- ⁸⁵ CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
- ⁸⁶ Department of Physics, University of Massachusetts, Amherst MA, United States of America
- ⁸⁷ Department of Physics, McGill University, Montreal QC, Canada
- ⁸⁸ School of Physics, University of Melbourne, Victoria, Australia
- ⁸⁹ Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
- ⁹⁰ Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
- ⁹¹ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁹² B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
- ⁹³ National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
- ⁹⁴ Group of Particle Physics, University of Montreal, Montreal QC, Canada
- ⁹⁵ P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
- ⁹⁶ Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- ⁹⁷ National Research Nuclear University MEPhI, Moscow, Russia
- ⁹⁸ D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- ⁹⁹ Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- ¹⁰⁰ Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- ¹⁰¹ Nagasaki Institute of Applied Science, Nagasaki, Japan
- ¹⁰² Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
- ¹⁰³ ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- ¹⁰⁴ Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
- ¹⁰⁵ Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- ¹⁰⁶ Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- ¹⁰⁷ Department of Physics, Northern Illinois University, DeKalb IL, United States of America
- ¹⁰⁸ Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
- ¹⁰⁹ Department of Physics, New York University, New York NY, United States of America
- ¹¹⁰ Ohio State University, Columbus OH, United States of America
- ¹¹¹ Faculty of Science, Okayama University, Okayama, Japan
- ¹¹² Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
- ¹¹³ Department of Physics, Oklahoma State University, Stillwater OK, United States of America
- ¹¹⁴ Palacký University, RCPTM, Olomouc, Czech Republic
- ¹¹⁵ Center for High Energy Physics, University of Oregon, Eugene OR, United States of America

- ¹¹⁶ LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
- ¹¹⁷ Graduate School of Science, Osaka University, Osaka, Japan
- ¹¹⁸ Department of Physics, University of Oslo, Oslo, Norway
- ¹¹⁹ Department of Physics, Oxford University, Oxford, United Kingdom
- ¹²⁰ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ¹²¹ Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
- ¹²² National Research Centre "Kurchatov Institute" B.P.Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia
- ¹²³ ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ¹²⁴ Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
- ¹²⁵ ^(a) Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa; ^(b) Faculdade de Ciências, Universidade de Lisboa, Lisboa; ^(c) Department of Physics, University of Coimbra, Coimbra; ^(d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; ^(e) Departamento de Física, Universidade do Minho, Braga; ^(f) Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Granada (Spain); ^(g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
- ¹²⁶ Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
- ¹²⁷ Czech Technical University in Prague, Praha, Czech Republic
- ¹²⁸ Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
- ¹²⁹ State Research Center Institute for High Energy Physics (Protvino), NRC KI, Russia
- ¹³⁰ Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- ¹³¹ ^(a) INFN Sezione di Roma; ^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- ¹³² ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ¹³³ ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- ¹³⁴ ^(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; ^(b) Centre National de l'Energie des Sciences Techniques Nucleaires, Rabat; ^(c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; ^(d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; ^(e) Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ¹³⁵ DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
- ¹³⁶ Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America
- ¹³⁷ Department of Physics, University of Washington, Seattle WA, United States of America
- ¹³⁸ Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- ¹³⁹ Department of Physics, Shinshu University, Nagano, Japan
- ¹⁴⁰ Fachbereich Physik, Universität Siegen, Siegen, Germany
- ¹⁴¹ Department of Physics, Simon Fraser University, Burnaby BC, Canada
- ¹⁴² SLAC National Accelerator Laboratory, Stanford CA, United States of America
- ¹⁴³ ^(a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; ^(b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
- ¹⁴⁴ ^(a) Department of Physics, University of Cape Town, Cape Town; ^(b) Department of Physics, University of Johannesburg, Johannesburg; ^(c) School of Physics, University of the Witwatersrand,

Johannesburg, South Africa

¹⁴⁵ ^(a) Department of Physics, Stockholm University; ^(b) The Oskar Klein Centre, Stockholm, Sweden

¹⁴⁶ Physics Department, Royal Institute of Technology, Stockholm, Sweden

¹⁴⁷ Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America

¹⁴⁸ Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

¹⁴⁹ School of Physics, University of Sydney, Sydney, Australia

¹⁵⁰ Institute of Physics, Academia Sinica, Taipei, Taiwan

¹⁵¹ Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

¹⁵² Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

¹⁵³ Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

¹⁵⁴ International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

¹⁵⁵ Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

¹⁵⁶ Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

¹⁵⁷ Department of Physics, University of Toronto, Toronto ON, Canada

¹⁵⁸ ^(a) TRIUMF, Vancouver BC; ^(b) Department of Physics and Astronomy, York University, Toronto ON, Canada

¹⁵⁹ Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan

¹⁶⁰ Department of Physics and Astronomy, Tufts University, Medford MA, United States of America

¹⁶¹ Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

¹⁶² Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America

¹⁶³ ^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

¹⁶⁴ Department of Physics, University of Illinois, Urbana IL, United States of America

¹⁶⁵ Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

¹⁶⁶ Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

¹⁶⁷ Department of Physics, University of British Columbia, Vancouver BC, Canada

¹⁶⁸ Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada

¹⁶⁹ Department of Physics, University of Warwick, Coventry, United Kingdom

¹⁷⁰ Waseda University, Tokyo, Japan

¹⁷¹ Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

¹⁷² Department of Physics, University of Wisconsin, Madison WI, United States of America

¹⁷³ Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

¹⁷⁴ Fakult[ä]t für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany

¹⁷⁵ Department of Physics, Yale University, New Haven CT, United States of America

¹⁷⁶ Yerevan Physics Institute, Yerevan, Armenia

¹⁷⁷ Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

^a Also at Department of Physics, King's College London, London, United Kingdom

^b Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

- ^c Also at Novosibirsk State University, Novosibirsk, Russia
- ^d Also at TRIUMF, Vancouver BC, Canada
- ^e Also at Department of Physics & Astronomy, University of Louisville, Louisville, KY, United States of America
- ^f Also at Department of Physics, California State University, Fresno CA, United States of America
- ^g Also at Department of Physics, University of Fribourg, Fribourg, Switzerland
- ^h Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
- ⁱ Also at Departamento de Física e Astronomia, Faculdade de Ciências, Universidade do Porto, Portugal
- ^j Also at Tomsk State University, Tomsk, Russia
- ^k Also at Università di Napoli Parthenope, Napoli, Italy
- ^l Also at Institute of Particle Physics (IPP), Canada
- ^m Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- ⁿ Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia
- ^o Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
- ^p Also at Louisiana Tech University, Ruston LA, United States of America
- ^q Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain
- ^r Also at Graduate School of Science, Osaka University, Osaka, Japan
- ^s Also at Department of Physics, National Tsing Hua University, Taiwan
- ^t Also at Department of Physics, The University of Texas at Austin, Austin TX, United States of America
- ^u Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia
- ^v Also at CERN, Geneva, Switzerland
- ^w Also at Georgian Technical University (GTU), Tbilisi, Georgia
- ^x Also at O Chadai Academic Production, Ochanomizu University, Tokyo, Japan
- ^y Also at Manhattan College, New York NY, United States of America
- ^z Also at Hellenic Open University, Patras, Greece
- ^{aa} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
- ^{ab} Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
- ^{ac} Also at School of Physics, Shandong University, Shandong, China
- ^{ad} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
- ^{ae} Also at Section de Physique, Université de Genève, Geneva, Switzerland
- ^{af} Also at International School for Advanced Studies (SISSA), Trieste, Italy
- ^{ag} Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America
- ^{ah} Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China
- ^{ai} Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia
- ^{aj} Also at National Research Nuclear University MEPhI, Moscow, Russia
- ^{ak} Also at Department of Physics, Stanford University, Stanford CA, United States of America
- ^{al} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
- ^{am} Also at Flensburg University of Applied Sciences, Flensburg, Germany
- ^{an} Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia
- ^{ao} Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
- * Deceased