

A CENSUS OF STAR-FORMING GALAXIES IN THE $z \sim 9$ –10 UNIVERSE BASED ON *HST*+*SPITZER* OBSERVATIONS OVER 19 CLASH CLUSTERS: THREE CANDIDATE $z \sim 9$ –10 GALAXIES AND IMPROVED CONSTRAINTS ON THE STAR FORMATION RATE DENSITY AT $z \sim 9.2$ *

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ABSTRACT

We utilize a two-color Lyman-break selection criterion to search for $z \sim 9$ –10 galaxies over the first 19 clusters in the CLASH program. A systematic search yields three $z \sim 9$ –10 candidates. While we have already reported the most robust of these candidates, MACS1149-JD, two additional $z \sim 9$ candidates are also found and have H_{160} -band magnitudes of ~ 26.2 – 26.9 . A careful assessment of various sources of contamination suggests $\lesssim 1$ contaminants for our $z \sim 9$ –10 selection. To determine the implications of these search results for the luminosity function (LF) and star formation rate density at $z \sim 9$, we introduce a new differential approach to deriving these quantities in lensing fields. Our procedure is to derive the evolution by comparing the number of $z \sim 9$ –10 galaxy candidates found in CLASH with the number of galaxies in a slightly lower-redshift sample (after correcting for the differences in selection volumes), here taken to be $z \sim 8$. This procedure takes advantage of the fact that the relative volumes available for the $z \sim 8$ and $z \sim 9$ –10 selections behind lensing clusters are not greatly dependent on the details of the lensing models. We find that the normalization of the UV LF at $z \sim 9$ is just $0.28^{+0.39}_{-0.20} \times$ that at $z \sim 8$, which is $\sim 1.4^{+3.0}_{-0.8} \times$ lower than extrapolating $z \sim 4$ –8 LF results. While consistent with the evolution in the UV LF seen at $z \sim 4$ –8, these results marginally favor a more rapid evolution at $z > 8$. Compared to similar evolutionary findings from the HUDF, our result is less insensitive to large-scale structure uncertainties, given our many independent sightlines on the high-redshift universe.

Key words: galaxies: evolution – galaxies: high-redshift

Online-only material: color figures

1. INTRODUCTION

Since the discovery of large numbers of $z \sim 3$ galaxies with the Lyman-break selection technique 17 yr ago (Steidel et al. 1996), there has been a persistent effort to use the latest

facilities to identify galaxies at higher and higher redshifts through photometric selections and follow-up spectroscopy. These efforts allow us to probe galaxies during the epoch of reionization to ascertain what role they may have in driving this process. Progressively, the high-redshift frontier has been extended to $z \sim 4$ –5 (e.g., Madau et al. 1996; Steidel et al. 1999), $z \sim 6$ (e.g., Stanway et al. 2003; Bouwens et al. 2003; Dickinson et al. 2004), $z \sim 7$ (e.g., Bouwens et al. 2004; Yan & Windhorst 2004; Bouwens & Illingworth 2006b; Iye et al.

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2006; Fontana et al. 2010; Schenker et al. 2012), and $z \sim 8$ (e.g., Bouwens et al. 2010; McLure et al. 2010; Bunker et al. 2010; Yan et al. 2010).

The current frontier for identifying high-redshift galaxies now seems to lie firmly at $z \sim 10$, with three distinct $z \sim 10$ galaxy candidates having been reported.²⁴ Bouwens et al. (2011a) presented the discovery of a plausible $z \sim 10.3$ galaxy in the full two-year HUDF09 observations over the HUDF (see also Oesch et al. 2012a). More recently, Zheng et al. (2012, hereafter Z12) presented evidence for a highly magnified $z \sim 9.6$ galaxy within the 524 orbit CLASH program (Postman et al. 2012), and Coe et al. (2013; hereafter C13) reported the discovery of an even higher redshift triply lensed $z \sim 10.8$ galaxy.

Despite the very interesting nature of earlier exploratory work, the total number of $z \sim 9$ –11 galaxies is small, and hence it is still somewhat challenging to obtain accurate constraints on how rapidly the luminosity function (LF) or star formation rate (SFR) density evolved in the very early universe, at $z > 8$. Earlier $z \sim 10$ searches using the very deep HUDF09 data (Bouwens et al. 2011a; Oesch et al. 2012a) found tentative evidence for a deficit of $z \sim 10$ galaxies, relative to simple extrapolations from lower redshifts, pointing toward a very rapid evolution in the UV LF and SFR density at $z > 8$ (Oesch et al. 2012a). A rapid evolution of the UV LF at $z > 8$ is supported by several theoretical models (Trenti et al. 2010; Lacey et al. 2011), but may be in some tension with the discovery of one bright, multiply lensed $z \sim 10.8$ galaxy in the CLASH program (C13), because one might have expected such sources to be quite rare assuming a rapid evolution of the UV LF.

Fortunately, there is an ever-increasing quantity of observations now available to identify $z \sim 9$ –10 galaxies. One noteworthy near-term opportunity exists in the moderately deeper WFC3/IR observations acquired over the HUDF (GO 12498; Ellis et al. 2013). This program has made it possible to extend $z \sim 9$ –10 samples in the HUDF deeper by ~ 0.4 mag, while increasing the number of sources by a factor of 2–4 (McLure et al. 2013). However, another significant opportunity exists in ongoing observations over lensing clusters as part of the 524 orbit CLASH program (Postman et al. 2012). The initial discovery papers of Z12 and C13 only reported on the brightest and most robust $z \sim 10$ and $z \sim 11$ galaxy candidates from the CLASH program; however, it should be possible to extend these searches somewhat fainter by ~ 0.5 –1.0 mag to the magnitude limit of the survey (~ 27 AB mag). At such magnitudes we would expect to identify other plausible $z \sim 9$ –10 galaxies, potentially increasing the overall sample size to ~ 3 –5 sources in total.

The purpose of this paper is to capitalize on the opportunity that exists within lensing clusters from the CLASH program. A deeper search for $z \sim 9$ –10 galaxies can be performed in a reasonably reliable manner, taking full advantage of the substantial observations with the *Spitzer*/IRAC instrument over the CLASH program (Egami et al. 2008; Bouwens et al. 2011c), allowing us to distinguish potential star forming galaxy candidates at $z \sim 9$ –10 from lower-redshift interlopers. We also incorporate *Hubble Space Telescope* (HST) observations over two more clusters from the CLASH program (utilizing a total

²⁴ Following the initial submission of this paper, seven additional $z \sim 10$ candidates were identified: a $z \sim 9.5$ candidate from Ellis et al. (2013) over the HUDF, a $z \sim 9.8$ candidate from Oesch et al. (2014) over the HUDF, four $z \sim 9.5$ –10.2 candidates from Oesch et al. (2014) over CANDELS, and a triply lensed $z \sim 9.8$ candidate from Zitrin et al. (2014) behind A2744.

Table 1
The 19 Cluster Fields from the CLASH Program Considered
in the Present $z \sim 9$ Search

Cluster	Redshift	High Magnification ^a
A209	0.206	
A383	0.187	
A611	0.288	
A2261	0.224	
MACS0329.7–0211	0.450	
MACS0416.1–2403	0.42	Y
MACS0647.8+7015	0.584	Y
MACS0717.5+3745	0.548	Y
MACS0744.9+3927	0.686	
MACS1115.9+0129	0.352	
MACS1149.6+2223	0.544	Y
MACS1206.2–0847	0.440	
MACSJ1720.3+3536	0.391	
MACSJ1931.8–2635	0.352	
MACSJ2129.4–0741	0.570	Y
MS2137–2353	0.313	
RXJ1347.5–1145	0.451	
RXJ1532.9+3021	0.345	
RXJ2129.7+0005	0.234	

Notes. ^a Clusters in the CLASH program were selected based on either their X-ray or magnification properties (Postman et al. 2012). Clusters marked here with a “Y” were included because of their magnification properties.

of 19 clusters) to expand the total search area by 50% and 10% over what was considered in Z12 and C13, respectively.

The plan for this paper is as follows. In Section 2 we describe our observational data set. In Section 3 we discuss our procedure for catalog creation, the selection of $z \sim 9$ –10 galaxy candidates, quantifying their properties, and estimating the extent to which contamination may be a concern for our selection. In Section 4 we introduce a new differential approach to derive the evolution in the UV LF and SFR density at $z \gtrsim 9$ and then apply it to our search results at $z \sim 9$. Finally, in Section 5, we summarize the results from this paper and offer a prospective. Throughout this work, we quote results in terms of the luminosity $L_{z=3}^*$ Steidel et al. (1999) derived at $z \sim 3$: $M_{1700,AB} = -21.07$. For simplicity we refer to the HST F225W, F390W, F435W, F475W, F606W, F625W, F775W, F814W, F850LP, F105W, F110W, F125W, F140W, and F160W bands as UV_{225} , U_{390} , B_{435} , g_{475} , V_{606} , r_{625} , i_{775} , I_{814} , z_{850} , Y_{105} , J_{110} , J_{125} , JH_{140} , and H_{160} , respectively. Where necessary, we assume $\Omega_0 = 0.3$, $\Omega_\Lambda = 0.7$, and $H_0 = 70$ km/s/Mpc. All magnitudes are in the AB system (Oke & Gunn 1983).

2. OBSERVATIONAL DATA

Our primary data set for this study are the 20 orbit HST observations over the first 19 clusters with data from the 524 orbit CLASH multi-cycle treasury program (Postman et al. 2012; see Table 1). The HST observations over each of the CLASH clusters is typically distributed over 16 different bands using the WFC3/UVIS camera, the Advanced Camera for Surveys (ACS) wide field camera, and the WFC3/IR camera. These observations extend from $0.2 \mu\text{m}$ (UV_{225}) to $1.6 \mu\text{m}$ (H_{160}) and reach to depths of 26.4–27.7 AB mag (5σ : $0''.4$ diameter aperture) depending upon the passband.

Our reductions of these data were conducted using standard procedures, aligned, and then drizzled on the same frame ($0''.065$ pixel scale) with the multidrizzle software (Koekemoer et al. 2003). The FWHM for the point-spread function (PSF) is

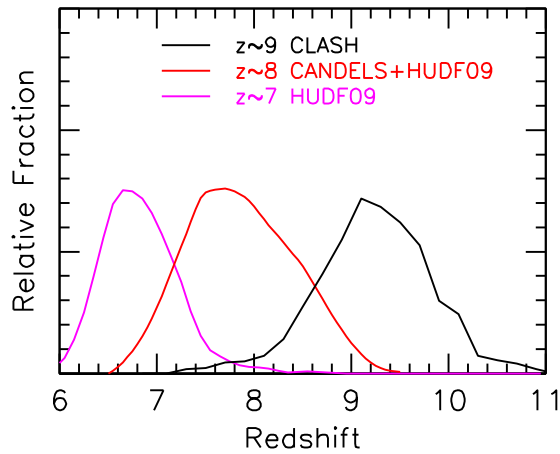


Figure 1. Redshift distribution we would expect for our present $z \sim 9$ selection based on the simulations we run in Section 4.3. These simulations allow us to assess the relative selection volume for our $z \sim 9$ selection and our comparison sample at $z \sim 8$. The mean redshift for our selection is 9.2. Our $z \sim 9$ selection cuts off at $z > 10$ due to our use of a $JH_{140} - H_{160} < 0.5$ criterion (Section 3.2; see also Figure 2). For context, we also show the expected redshift distributions for the $z \sim 7$ and $z \sim 8$ selections of Bouwens et al. (2011b) and Oesch et al. (2012b), respectively.

(A color version of this figure is available in the online journal.)

$\sim 0'.1$ in the WFC3/UVIS or ACS observations and $\sim 0.16\text{--}0'.17$ for the WFC3/IR observations.

The typical area available over each cluster to search for $z \sim 9\text{--}10$ galaxies is $\sim 4 \text{ arcmin}^2$ and is dictated by the area available within the WFC3/IR field of view. In total, we make use of $\sim 77 \text{ arcmin}^2$ over the first 19 CLASH clusters to search for $z \sim 9\text{--}10$ galaxies. This corresponds to an approximate search volume of $\sim 7000 \text{ Mpc}^3$ (comoving) at $z \sim 9$ to probe faint, highly magnified $\mu > 5$ galaxies (assuming that $\sim 25\%$ of our WFC3/IR area is high magnification $\mu \gtrsim 5$ and a $\Delta z \sim 1$ width for our redshift selection window; see Figure 1). To ensure that we have the maximum depth and filter coverage available for the candidates uncovered in our search, we do not consider the small amount of data over each cluster with observations in only one of the two roll angles used for the CLASH program (see Figure 11 of Postman et al. 2012 for an illustration of the two roll-angle strategy).

Each of the CLASH clusters also has a substantial amount of observations with the *Spitzer*/IRAC instrument (Fazio et al. 2004). The typical integration times range from ~ 3.5 hr per IRAC band from the ICLASH program (GO #80168; Bouwens et al. 2011c) to ~ 5 hr per IRAC band from the *Spitzer* IRAC Lensing Survey program (GO #60034; PI Egami). Even deeper observations are available from the Surfs'Up program (Bradač et al. 2014: 30 hr), Frontier Field program (T. Soifer and P. Capak: 50 hr), and follow-up observations on MACS0647 and MACS1720 (PI Bouwens [90213]: 11/24 hr; Coe [10140]: 56 hr). These observations reach to 1σ depths of $\sim 26.2\text{--}27.4$ mag in both the $3.6 \mu\text{m}$ and $4.5 \mu\text{m}$ IRAC channels, allowing us to set useful constraints on the color of possible $z \sim 9\text{--}10$ candidates redward of the break. The FWHM for the IRAC PSF at $3.6 \mu\text{m}$ and $4.5 \mu\text{m}$ is $\sim 1''.8$. We reduced the *Spitzer*/IRAC observations using the public MOPEX software available from the *Spitzer* Science Center (Makovoz & Khan 2005), excluding roll angles where artifacts from bright stars had an impact on the photometry of the candidate $z \sim 9\text{--}10$ galaxies under study. The reductions were drizzled onto a common output frame ($0'.6$ pixel scale).

3. RESULTS

3.1. Catalog Construction

Our procedure for constructing catalogs is similar to that previously utilized by Bouwens et al. (2007, 2011b, 2012b). These catalogs are distinct from those distributed as part of the CLASH program, but overall the results are in very good agreement.

We provide a brief outline of the procedure we use here. More details are provided in several of our previous publications (e.g., Bouwens et al. 2007, 2011b, 2012b). SExtractor (Bertin & Arnouts 1996) is run in dual-image mode using the square root of the χ^2 image (Szalay et al. 1999) to detect sources and the PSF-matched images for photometry. The χ^2 image (similar to a coadded frame) is constructed from the imaging observations in the two passbands where we expect $z \sim 9$ candidates to show significant signal (i.e., the JH_{140} and H_{160} bands). For the photometry, PSF-matching is done to the WFC3/IR H_{160} band. Fluxes and colors of sources are measured in apertures that scale with the size of sources, as recommended by Kron (1980) and using a Kron factor of 1.2. The small-aperture fluxes are then corrected to total magnitudes in two steps. First the excess flux around the source in a larger scalable aperture (Kron factor 2.5) is derived based on the square root of χ^2 image and this correction is applied to the measured fluxes in all *HST* bands. Second, a correction is made for the expected light outside the larger scaled aperture and on the wings of the PSF using the tabulated encircled energy distribution (e.g., from Sirianni et al. 2005).

The measurement of IRAC fluxes is important for a more secure identification of $z \sim 9$ candidates in our fields, because it allows us to quantify the approximate spectral slope of the sources redward of the spectral break observed at $\sim 1.2 \mu\text{m}$ and therefore distinguish potential star forming galaxies at $z \sim 9\text{--}10$ from interlopers at $z \sim 1\text{--}2$. IRAC photometry can be challenging due to the significant overlap between nearby sources in existing data. Fortunately, there are well-established procedures to use the positions and spatial profiles of sources in available *HST* observations to model the IRAC image observations and extract fluxes (e.g., Shapley et al. 2005; Labbé et al. 2006; Grazian et al. 2006; Laidler et al. 2007).

Here we make use of the MOPHONGO software (Labbé et al. 2006, 2010a, 2010b, 2013) to do photometry on sources in our fields, given the confusion. Since this software has been presented more extensively in other places, we only include a brief description here. The most important step for doing photometry on faint sources with this software is to remove confusion from neighboring sources. This is accomplished by using the deep WFC3/IR observations as a template to model the positions and isolated flux profiles of the foreground sources. These flux profiles are convolved to match the IRAC PSF and then simultaneously fit to the IRAC imaging data, leaving only the fluxes of the sources as unknowns. The best-fit model is then used to subtract the flux from neighboring sources and normal aperture photometry is performed on sources in a $2''.5$ diameter aperture. The measured $3.6 \mu\text{m}$ and $4.6 \mu\text{m}$ fluxes are then corrected to account for the light on the wings of the IRAC PSF (typically the correction is a factor of ~ 2.2).

3.2. Source Selection

In this paper we adopt a two-color Lyman-break selection to search for promising $z \sim 9\text{--}10$ galaxy candidates in the CLASH program. This work takes advantage of the sharp break in the

spectrum of star forming galaxies due to absorption by neutral hydrogen. Many years of spectroscopic work have demonstrated that the Lyman-break selection technique provides us with a very efficient means of identifying high-redshift galaxies (Steidel et al. 1996, 2003; Bunker et al. 2003; Dow-Hygelund et al. 2007; Popesso et al. 2009; Vanzella et al. 2009; Stark et al. 2010), with generally minimal contamination, albeit with a few notable exceptions at brighter magnitudes (e.g., Steidel et al. 2003; Bowler et al. 2012; Hayes et al. 2012). In the latter case, deeper mid-IR data can be valuable for guarding against such contamination.

In analogy with lower-redshift Lyman-break selections (e.g., Giavalisco et al. 2004; Bouwens et al. 2007; Bouwens et al. 2011b), we devised the following two-color $z \sim 9$ –10 selection for the CLASH cluster fields:

$$((J_{110} + J_{125})/2 - H_{160} > 0.7) \wedge (JH_{140} - H_{160} < 0.5),$$

where \wedge represents the logical **AND** symbol. This is very similar to the criteria previously presented in Z12 (i.e., $(J_{110} - JH_{140} > 0.5) \wedge (JH_{140} - H_{160} < 0.5)$), but probe to slightly higher-redshift sources on average, also folding in information from the redder J_{125} -band filter and requiring a sharper break in the spectrum. In general, it makes sense to combine the flux information from both the J_{110} and J_{125} bands to search for $z \gtrsim 9$ candidates because of their similar red-side cutoffs at $1.4 \mu\text{m}$. In applying the above criteria, the magnitudes of sources not detected at 1σ are set to their 1σ upper limits.

It is also important that we detect sources at sufficient signal-to-noise ratio (S/N) that we can rely on the color information (and optical non-detections) to provide reliable redshift information on the sources and guarantee they are real. After some experimentation and extensive simulations (Section 3.5), we elected to require sources in our $z \sim 9$ selection be detected at $\geq 6\sigma$ in combined JH_{140} and H_{160} bands (using a fixed $0'.35$ diameter aperture). For significance thresholds less than 6σ , our simulations (Section 3.5) suggest that our $z \sim 9$ selection would be subject to significant contamination from lower-redshift interlopers.

To ensure that sources really have no flux in the spectrum blueward of the Lyman break, we also require sources to be undetected ($< 2.5\sigma$) in the Y_{105} band and any passband blueward of this.²⁵ Moreover, we combine the flux in all the bluer bands (U_{390} , B_{435} , g_{475} , V_{606} , r_{625} , i_{775} , I_{814} , z_{850} , and Y_{105}) to construct a χ^2 statistic for sources in our catalogs and exclude sources from our selection if the $\chi_{\text{opt+Y}}^2$ statistic is greater than 3.8. The particular threshold for $\chi_{\text{opt+Y}}^2$ (i.e., 3.8), was chosen to keep contamination in our $z \sim 9$ sample relatively low, while not overly affecting the completeness of our samples (see Figure 19 from Bouwens et al. 2011b for an illustration of how such a choice can be made). This criterion is very effective at guarding against contamination from sources that are consistently faint in all optical bands; it ensures that sources are not consistently detected at $> 1\sigma$ in more than three optical bands.

Here χ^2 is calculated as follows: $\chi_{\text{opt+Y}}^2 = \sum_i \text{SGN}(f_i)(f_i/\sigma_i)^2$ where f_i is the flux in band i in a consistent aperture, σ_i is the uncertainty in this flux, and $\text{SGN}(f_i)$ is equal to 1 if $f_i > 0$ and -1 if $f_i < 0$ (Bouwens et al. 2011b). As in Bouwens et al. (2011b), we calculate this χ^2 statistic in three

different apertures (scalable Kron apertures [Kron factor of 1.2], $0'.35$ diameter circular apertures, $0'.18$ diameter circular apertures) to ensure that there is absolutely no evidence for a significant excess of light blueward of the break, whether this light be tightly concentrated on the source itself or more diffuse. When computing the χ^2 statistic with $0'.18$ diameter apertures, we use the original unsmoothed ACS or WFC3/IR images (i.e., before PSF-matching to the WFC3/IR H_{160} -band data) to retain the maximum S/N for the purposes of rejecting low-redshift interlopers.

As one final step to ensure that our $z \sim 9$ candidates show no evidence for flux blueward of the break, we construct a second χ^2 statistic for each source utilizing only the information in the three bands immediately blueward of the break (i.e., the I_{814} , z_{850} , and Y_{105} bands). We then exclude any source that has a χ_{I+z+Y}^2 value greater than 3. This criterion provides us with better discrimination against dusty lower-redshift interlopers (which we would not expect to be detected in the bluer bands) and serves as an effective complement to our other χ^2 criterion (which is better at discriminating against sources that are consistently faint in all optical bands). Sources detected at $> 2\sigma$ in the Y_{105} band are also excluded to minimize the contribution of $z \sim 8$ galaxies to our selection.

In Figure 1, we show the approximate redshift selection window for our current selection. Details on how it is calculated will be presented in Section 4.3, but approximately involve adding artificial sources to the real data with realistic colors, sizes, and magnitudes, and then attempting to reselect them with the criteria given above. The mean redshift we derive for our $z \sim 9$ –10 selection from the simulations is 9.2. For context, we also present the redshift selection windows for samples at $z \sim 7$ and $z \sim 8$, as selected by Bouwens et al. (2011b) and Oesch et al. (2012b), respectively.

3.3. Resulting $z \sim 9$ Sample

We applied the selection criteria given in the previous section to the *HST* WFC3/UVIS+ACS+WFC3/IR observations from all 19 clusters in the current data set. We identified three sources that satisfy these selection criteria (see Figure 2, which shows these sources relative to our CLASH photometric sample as a whole). The sources are found behind three different clusters: MACSJ1149.6+2223, MACSJ1115.9+0129, and MACSJ1720.3+3536. The brightest of our three candidates, MACS1149-JD, was already presented in Z12. These sources were also flagged as the most interesting $z > 8$ sources using an independent, purely photometric redshift selection (Bradley et al. 2014).

In narrowing our selection down to our three highest quality $z \sim 9$ –10 candidates, we found that our $\chi_{\text{opt+Y}}^2$ optical non-detection and $JH_{140} - H_{160} < 0.5$ criteria were particularly important. From the small sample of 29 sources that satisfied our $(J_{125} + JH_{140})/2 - H_{160} > 0.7$ Lyman-break criterion and optical non-detection criteria in individual bands, we found that our $z \sim 9$ –10 reduced down to nine sources if we applied our $\chi_{\text{opt+Y}}^2$ criterion and finally down to our three candidates if we applied the $JH_{140} - H_{160} < 0.5$ color criterion. Use of either of the two J -band filters in constructing a Lyman-break sample resulted in a similar selection of sources, modulo one or two sources. For example, application of $J_{110} - H_{160} > 0.7$ color selection instead of the $(J_{110} + J_{125})/2 - H_{160} > 0.7$ selection resulted in the same $z \sim 9$ –10 candidates as are featured in our paper, plus a source at 04:16:11.53, $-24:04:53.2$, which

²⁵ Since we combine the optical flux measurements into several χ^2 statistics that we use to test the plausibility of specific sources as $z \sim 9$ candidates, we only adopt a weaker 2.5σ threshold here to avoid unnecessarily excluding many plausible $z \sim 9$ candidates.

Table 2
Coordinates, Estimated Redshifts and Magnification Factors, and Photometry for the Present $z \sim 9$ Sample^a

	MACS1149-JD ^b	MACSJ1115-JD1	MACSJ1720-JD1	Stack ^c
R.A.	11:49:33.58	11:15:54.50	17:20:12.76	...
Decl.	22:24:45.7	01:29:47.9	35:36:17.5	...
$z_{\text{photo}}^{\text{d}}$	$9.7^{+0.1}_{-0.1}$ ^e	$9.2^{+0.4}_{-0.8}$	$8.9^{+0.3}_{-0.5}$...
Magnification	$14.5^{+4.2}_{-1.0}$	$9.3^{+5.8}_{-3.6}$	$5.0^{+4.7}_{-0.7}$...
$S/N (JH_{140} + H_{160})^{\text{f}}$	15.4	7.8	6.9	...
U_{390}	-8 ± 25	-14 ± 37	16 ± 32	1 ± 18
B_{435}	-1 ± 26	-117 ± 39	4 ± 32	-35 ± 19
g_{475}	-3 ± 19	-23 ± 25	-10 ± 20	-12 ± 12
V_{606}	-9 ± 14	-0 ± 35	-11 ± 28	-7 ± 16
r_{625}	-27 ± 22	10 ± 24	-9 ± 17	-6 ± 11
i_{775}	0 ± 27	49 ± 47	-35 ± 38	0 ± 22
J_{814}	-3 ± 11	-13 ± 20	-27 ± 17	-16 ± 10
z_{850}	-38 ± 34	-32 ± 55	4 ± 39	-15 ± 25
Y_{105}	-3 ± 17	-39 ± 23	-20 ± 20	-21 ± 12
J_{110}	27 ± 13	37 ± 19	22 ± 13	26 ± 8
J_{125}	56 ± 16	63 ± 21	44 ± 16	49 ± 10
JH_{140}	146 ± 15	80 ± 22	80 ± 15	86 ± 10
	(=26.0 \pm 0.1)	(=26.6 \pm 0.3)	(=26.6 \pm 0.2)	(=26.6 \pm 0.1)
H_{160}	193 ± 15	115 ± 19	66 ± 16	100 ± 10
	(=25.7 \pm 0.1)	(=26.2 \pm 0.2)	(=26.9 \pm 0.3)	(=26.4 \pm 0.1)
[3.6]	$175 \pm 44^{\text{g}}$	380 ± 118	-42 ± 132	142 ± 87
[4.5]	$366 \pm 71^{\text{g}}$	-56 ± 122	209 ± 133	102 ± 91
$\frac{1}{2}([3.6] + [4.5])^{\text{h}}$	271 ± 41	163 ± 85	83 ± 93	122 ± 62
	(=25.3 \pm 0.2)	(=25.9 \pm 0.5)	(>26.8)	(=26.2 \pm 0.5)

Notes.

^a The fluxes in this table are in units of nJy.

^b The same candidate as presented in Z12. The fluxes presented in this table were derived independently from those presented in Z12, but are very similar in general.

^c This column gives the average fluxes in all *HST*+IRAC bands blueward of $0.4 \mu\text{m}$ for the three $z \sim 9$ candidates in our selection. The fluxes of each source are rescaled such that its average $JH_{140}+H_{160}$ flux matches the average $JH_{140} + H_{160}$ flux of the sample (prior to rescaling).

^d These photometric redshift estimates are based on the EAZY photometric redshift software (Brammer et al. 2008; see Section 3.4). In Section 3.4, we also provide photometric redshift estimates for sources using BPZ and Le PHARE.

^e Z12 prefer a slightly lower redshift of 9.6 for this source based on the photometry, but within the uncertainties, the present estimate is fully consistent with that given in Z12.

^f S/N of our $z \sim 9$ candidates in the JH_{140} and H_{160} bands added in quadrature ($0''.35$ diameter circular aperture). The S/N limit for our $z \sim 9$ selection was 6.0. Our highest S/N candidates are much less likely to correspond to lower-redshift contaminants (see Section 3.5, Figures 5 and 6).

^g The fluxes we measure in the *Spitzer*/IRAC $3.6 \mu\text{m}$ and $4.5 \mu\text{m}$ channels are very similar to that measured by Bradač et al. (2014), that is, $25.70 \pm 0.17 \pm 0.49$ (196 ± 32 nJy) in $3.6 \mu\text{m}$ channel and $25.01 \pm 0.08 \pm 0.21$ (370 ± 30 nJy) in the $4.5 \mu\text{m}$ channel.

^h We also presented an average of the *Spitzer*/IRAC $3.6 \mu\text{m}$ and $4.5 \mu\text{m}$ fluxes due to the limited S/N of each of these measurements.

appears quite likely to be at $z \sim 8.4$ (i.e., just outside our redshift selection window).

We performed *Spitzer*/IRAC photometry on all three $z \sim 9$ –10 candidates using the software described in Section 3.1. None of the three sources is near a bright foreground source, and so all of our IRAC flux measurements should be reliable. Two of our three $z \sim 9$ –10 candidates (MACS1149-JD and MACS1720-JD1) are detected ($>2\sigma$) in the moderately deep *Spitzer*/IRAC observations now available for MACS1149 and MACS1720.

The coordinates and photometry of these candidates are provided in Table 2, while postage stamp images of the candidates are shown in Figure 3. In Table 2, we also present a mean spectral energy distribution for galaxies at $z \sim 9$, which we computed on the basis of our *HST*+*Spitzer* photometry for the three $z \sim 9$ candidates. In computing this mean spectral energy distribution (SED), the fluxes of each source are rescaled such that its aver-

age $JH_{140}+H_{160}$ flux matches the average $JH_{140} + H_{160}$ flux for the sample (prior to rescaling).

As shown in Figure 3, MACS1149-JD is clearly resolved (see the Supplementary Information to Z12). MACS1149-JD also shows distinct elongation along the shear axis (Figure 1 from Z12) predicted from our gravitational lensing model for MACSJ1149.6+2223 (Z12). The other two plausible $z \sim 9$ candidates in our selection are quite small and show no clear evidence for gravitational shearing in the expected directions. However, since we would expect faint $z \geq 9$ galaxies to be small and the predicted magnification to be only modest (magnifications of ~ 5 – $9\times$ in total), it is not clear that the structural properties of the sources teach us anything definitive.

In Figure 4, we indicate the position of these candidates within the field of view of our MACSJ1149.6+2223, MACSJ1115.9+0129, and MACSJ1720.3+3536 observations (magenta circles). In Figure 4, we have also overplotted the

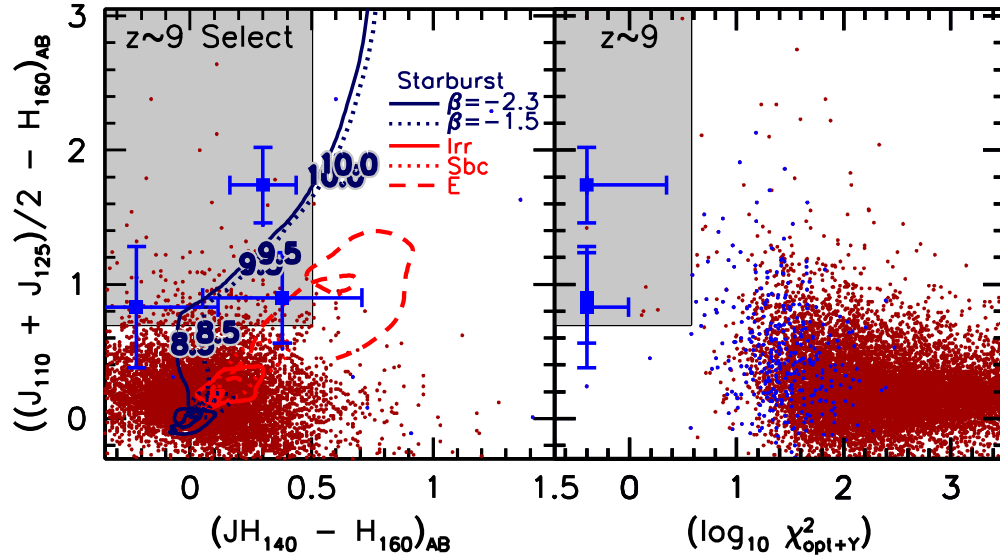


Figure 2. Selection criteria used to identify $z \sim 9$ –10 galaxies over the CLASH program. Left: the $((J_{110} + J_{125})/2 - H_{160})_{AB}$ vs. $(JH_{140} - H_{160})_{AB}$ diagram shows the first of the two primary criteria we use to identify $z \sim 9$ –10 galaxies from the CLASH program. Selected sources must fall in the gray region defined by two LBG-like color criteria, with a $(J_{110} + J_{125})/2 - H_{160} > 0.7$ criterion defining the Lyman break and a $JH_{140} - H_{160} < 0.5$ criterion providing a constraint on the spectral slope redward of the break. The large blue squares show the sources that made it into our $z \sim 9$ –10 sample. The error bars on these points are the 1σ uncertainties. The blue lines show the expected colors for star forming galaxies with varying UV -continuum slopes as a function of redshift, while the red lines show the expected colors for different SED templates at lower redshift (Coleman et al. 1980). The small dark red points show the colors of sources in our photometric sample where the χ^2_{opt+Y} statistic is > 3.8 . The blue points show these colors for sources where the χ^2_{opt+Y} statistic is < 3.8 . See Section 3.2 (and Bouwens et al. 2011b) for a definition of the χ^2_{opt+Y} statistic, but it roughly includes a stack of all the flux information in the Y_{105} band and bluer bands. Right: the $((J_{110} + J_{125})/2 - H_{160})_{AB}$ vs. χ^2_{opt+Y} diagram shows the second of the two primary criteria we use to identify $z \sim 9$ –10 galaxies from the CLASH program. The selected sources must fall in the gray region and therefore must show no flux in the optical or Y_{105} bands (i.e., $\chi^2_{opt+Y} < 3.8$). The three selected $z \sim 9$ candidates are the blue squares. The dark red points indicate sources in our photometric sample that are either detected in the Y_{105} band ($> 2\sigma$) or where the $JH_{140} - H_{160}$ colors are greater than 0.5. The blue points are those sources where neither condition is satisfied. This figure is similar to Figure 2 of Oesch et al. (2012b). Using both the two-color criteria and our χ^2_{opt+Y} criteria, we observe a clear separation between our $z \sim 9$ –10 candidates and the bulk of our photometric sample. While we cannot completely rule out certain classes of lower-redshift galaxies contaminating our selection (note that the color–color track for early-type galaxies overlaps our selection window in the left panel), the volume density of such contaminants would seem to be lower than that of bright $z \sim 9$ –10 galaxies (see Appendix A).

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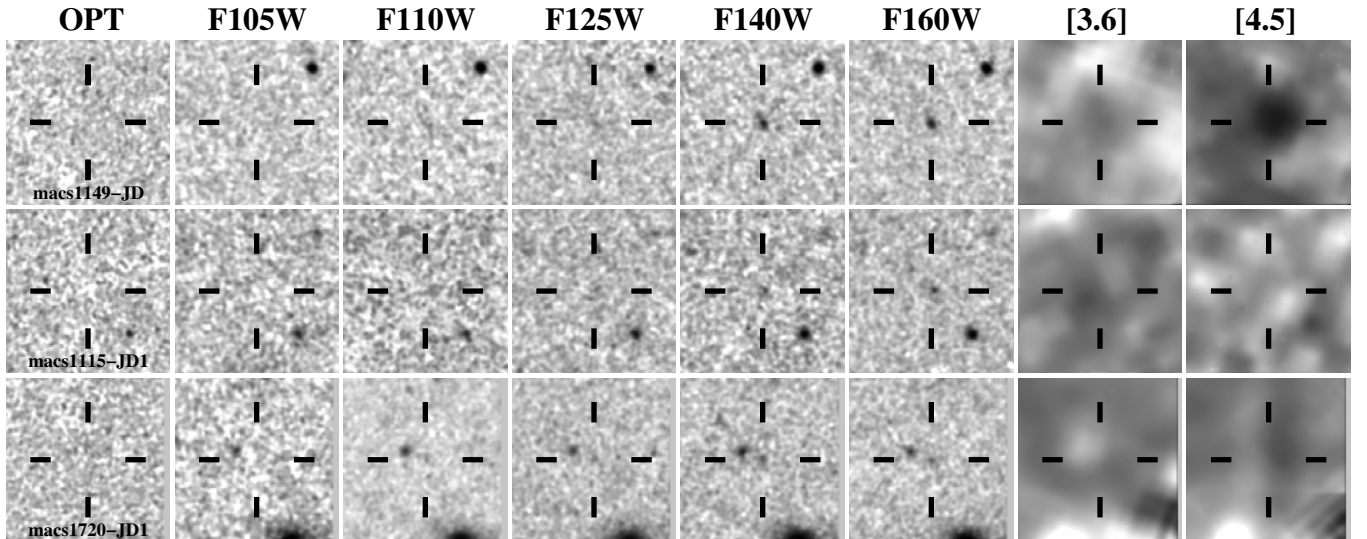


Figure 3. Postage stamp images ($6''.6 \times 6''.6$) of the three $z \sim 9$ galaxy candidates we identified in the current 19 cluster CLASH observations. The source in the uppermost row is the same $z \sim 9.6$ candidate we reported in Z12 (although our redshift estimate for this source is a very consistent $z \sim 9.7$; see Section 3.4). The leftmost postage stamp shows a stack of the deep ACS $B_{435} + g_{475} + V_{606} + r_{625} + i_{775} + I_{814} + z_{850}$ optical observations, while the other stamps show the observations in specific *HST* WFC3/IR and *Spitzer*/IRAC bands. On the IRAC postage stamps, flux from neighboring sources (as derived by MOPHONGO) has been subtracted off. All three of our $z \sim 9$ candidates are detected at $> 6.8\sigma$ in a coadded $JH_{140} + H_{160}$ image ($0''.35$ diameter aperture; see Table 2). The *Spitzer* fluxes we measure for the sources are sufficiently faint, so as to substantially prefer a $z > 6$ solution for the sources rather than a low-redshift solution. None of the sources show any significant detections in the optical ACS observations.

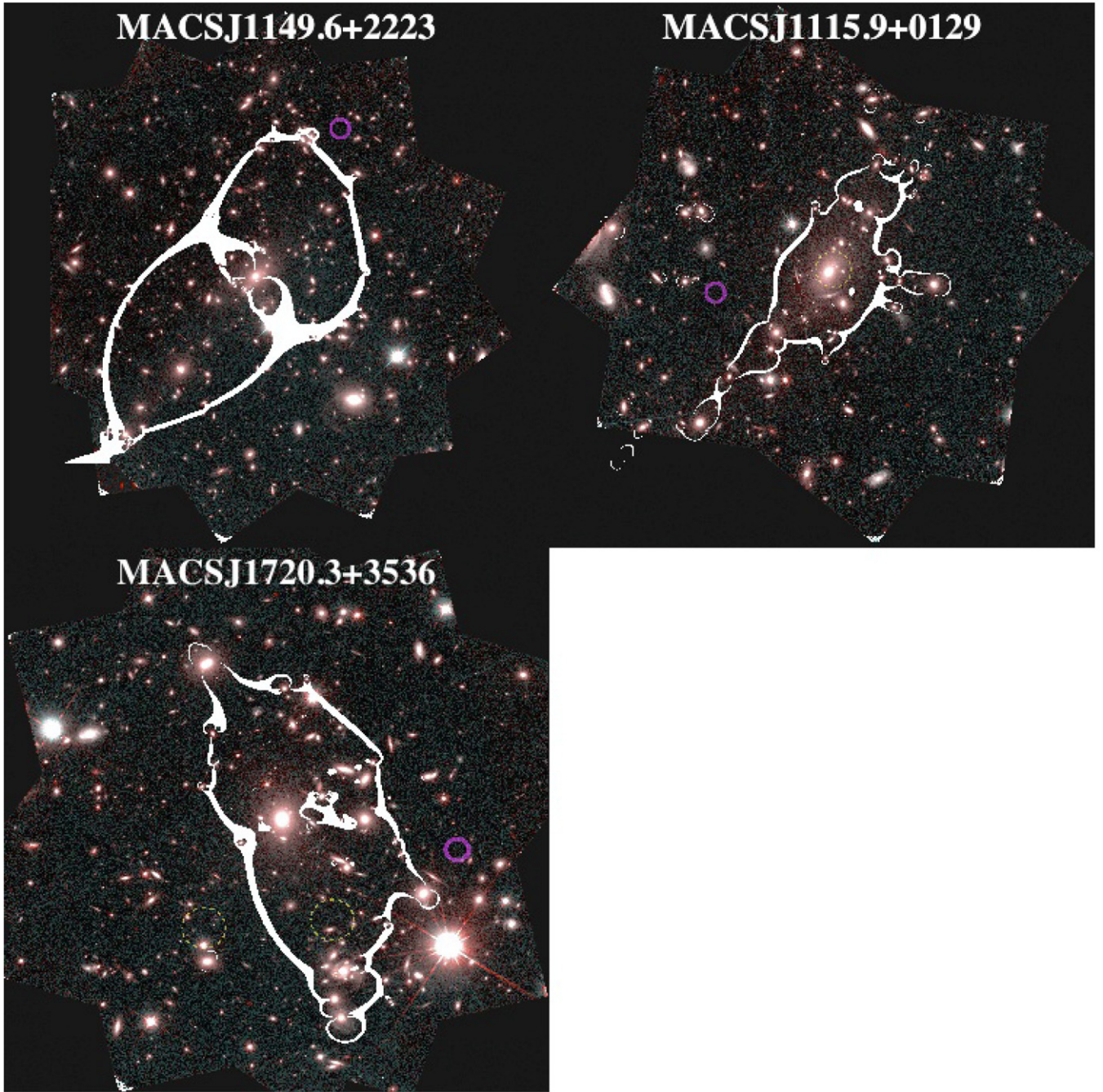


Figure 4. Position of the three $z \sim 9$ galaxy candidates that we identify over MACSJ1149.6+2223, MACSJ1115.9+0129, and MACSJ1720.3+3536. The color images shown are based on the *HST* $I_{814} + H_{160}$ observations of these clusters with CLASH and are shown over those regions with deep WFC3/IR observations. Overlaid on these images are the expected ultra high-magnification regions ($\mu > 100$) for a source at $z = 9.2$ based on the gravitational lensing models we have for the three clusters (Z12; A. Zitrin et al. 2014, in preparation; M. Carrasco et al. 2014, in preparation). Our lensing models for MACSJ1115.9+0129 and MACSJ1720.3+3536 are still preliminary and have not yet been finalized, constructed merely with the assumption that mass traces light, with typically only one lower-redshift system for normalization. The position of our three candidates is indicated by the large magenta circles. The dashed yellow circles indicate the position of possible counterimages as predicted by our preliminary lensing models.

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approximate critical lines for these clusters based on the lens models we have for these clusters (white contours: Z12; A. Zitrin et al. 2014, in preparation; M. Carrasco et al. 2014, in preparation). We caution that the lens models we have for MACSJ1115.9+0129 and MACSJ1720.3+3536 are still somewhat preliminary and are not totally finalized. The models are constructed based on the assumption that mass traces light, with typically only one lower-redshift system for normalization.

We can use these magnification models to estimate the approximate magnification factors for our candidate $z \sim 9$ galaxies. The approximate magnification factors are $14.5^{+4.2}_{-1.0}$, $9.3^{+5.8}_{-3.6}$, and $5.0^{+4.7}_{-0.7}$ and suggest intrinsic delensed $H_{160,AB}$ magnitudes for the sources of 28.5, 28.6, and 28.6 mag, respectively, for MACSJ1149-JD, MACSJ1115-JD1, and MACSJ1720-JD1. The intrinsic magnitudes inferred for the first three $z \sim 9$ galaxy candidates in the CLASH sample are only slightly brighter than

was found for the Bouwens et al. (2011a) $z \sim 10$ candidate (i.e., $H_{160,AB} \sim 28.7$ mag), and seem consistent with expectations.

The predicted positions of any possible counterimages to our $z \sim 9$ candidates are also shown in Figure 4 (dashed yellow circles). The only case where the counterimages are expected to be bright enough to detect is for MACSJ1720-JD1. Unfortunately, we were unable to locate the counterimages to MACSJ1720-JD1 at the predicted positions, which could mean that our lensing model may require further refinements, the counterimages are blended with foreground sources, or that the redshift identification is incorrect. For MACSJ1115-JD1, the counterimage is expected to be too faint to detect.

3.4. Best-fit Photometric Redshifts

The three candidate $z \sim 9$ galaxies presented in the previous section were selected using a two-color Lyman-break selection, and therefore their photometry is likely a reasonable fit to a model star forming galaxy SED at $z \sim 9$. However, as one can often fit the same photometry with SED templates at different redshifts, it is worthwhile for us to examine these candidates using standard photometric redshift procedures to look for possible degeneracies. Our use of photometric redshift procedures also allows us to naturally fold in the IRAC flux information we have for our $z \sim 9$ candidates.

To this end, we used the EAZY photometric redshift software (Brammer et al. 2008) to estimate photometric redshifts for the sources based on the observed photometry and to calculate the relative probability that sources in sample are more likely star-forming galaxies at $z \sim 9$ or galaxies at lower redshift (i.e., $z < 3$). The photometric redshift fitting is conducted using the EAZY_v1.0 template set supplemented by SED templates from the Galaxy Evolutionary Synthesis Models (Kotulla et al. 2009), which include nebular continuum and emission lines as described in Anders & Fritze-v. Alvensleben (2003). The EAZY_v1.0 template set consists of five SED templates from the PEGASE library (Fioc & Rocca-Volmerange 1997), which was derived based on the Blanton & Roweis (2007) algorithm and one young, dusty template (50 Myr, $A_V = 2.75$).

We consider three different priors in looking at the redshift likelihood distribution of our three $z \sim 9$ candidates: (1) a flat prior, (2) a prior calibrated to published LFs or LF trends, and (3) a prior tuned to reproduce the contamination rate estimated in the next section (Section 3.5). Our second prior is based on the LF results of Giallongo et al. (2005) and R. Quadri et al. (2012, private communication) for red $z \sim 1.3-2$ galaxies, while at $z > 7$ we utilize the LF-fitting formula of Bouwens et al. (2011b). The third prior accounts for the effect of noise on the photometry of lower-redshift galaxies in our search fields and the fact that in some rare events, noise could cause $\sim 1-2$ sources from our fields to seem like highly probable $z \sim 9$ galaxies (Section 3.5). Our third prior is calibrated to reproduce the results from our photometric scattering experiments. For simplicity (and because of the similar luminosities and magnitudes of all three $z \sim 9$ candidates), these priors are only a function of redshift; no luminosity dependence is considered. A more detailed description of these priors is provided in Appendix A.

The results are shown in Figure 5. The left panels show a comparison of the observed photometry with the best-fit $z \sim 9-10$ galaxy (blue line) and best-fit $z < 3$ galaxy (red line), while the right panels show the probability that a given source in our sample has a particular redshift. The best-fit redshifts for MACS1149-JD, MACSJ1115-JD1, and MACSJ1720-JD1 using the flat priors were 9.7, 9.2, and 8.9, respectively. The 68%

confidence intervals on the derived redshifts based on these priors are [9.57, 9.78], [8.38, 9.57], and [8.38, 9.26], respectively.

No substantial changes in these results are seen using our other two priors, except for the integrated probability within the $z \sim 1-2$ peak. For our second LF-calibrated prior (red line), the lower-redshift peak is actually smaller than in the case of the flat prior. This simply reflects the extreme rarity of faint red (old and/or dusty) galaxies at $z \sim 1.3-2$ as found in the Giallongo et al. (2005) and R. Quadri et al. (2012, private communication) probes (see also Stutz et al. 2008 and Figure 11 from Oesch et al. 2012a). For our third prior (dotted blue lines), the lower-redshift peak is larger, particularly for MACSJ1115-JD1 and MACS1720-JD1. Indeed, we might expect the lower-redshift peak to be higher than we would estimate from the photometry (and a flat prior), due to the impact that the selection process has on the observed SEDs of sources that satisfy our selection criteria. The selection process picks out those particular noise realizations for individual sources that are most consistent with $z \sim 9-10$ galaxies (even if that is not actually the case).

For the likelihood distributions given for the third prior, we should emphasize that the likelihood distributions were tuned so as to reproduce the expected contamination level for our $z \sim 9$ selection over the first 19 CLASH clusters (suggesting some possible contamination of our selection by lower-redshift interlopers) and we have no evidence that one particular source from our selection (e.g., MACSJ1115-JD1 or MACSJ1720-JD1) is in fact a contaminant.

As many filters in the CLASH program have overlapping wavelength coverage, we can further test the robustness of our best-fit photometric redshifts by using either the J_{110} or J_{125} flux measurements and either the JH_{140} or H_{160} flux measurement. The best-fit redshifts we find for MACS1149-JD, MACS1115-JD1, and MACS1720-JD1 range from 9.6 to 9.8, 9.1 to 9.1, and 9.0 to 9.4, respectively. If we exclude the Y_{105} flux measurement in deriving the best-fit photometric redshift, we find similar photometric redshifts for MACS1149-JD and MACS1720-JD1, but find a best-fit photometric redshift of 1.2 for MACS1115-JD1.

We also derived redshift likelihood confidence intervals using the Le PHARE photometric redshift package (Arnouts et al. 1999; Ilbert et al. 2006, 2009) for our three candidates. The SED templates we used with Le PHARE were the same as those optimized for the COSMOS survey (Scoville et al. 2007) and included three elliptical and six spiral SEDs as generated by Polletta et al. (2007) using the GRASIL code (Silva et al. 1998), as well as 12 starburst galaxies ranging in age from 30 Myr to 3 Gyr using the Bruzual & Charlot (2003) GALAXEV library. We supplemented these with four additional elliptical templates for a total of seven elliptical templates. Dust extinction was added in 10 steps up to $E(B - V) = 0.6$.²⁶ With these templates, we used Le PHARE to derive the following 68% confidence intervals for the candidates: [9.49, 9.85] for MACS1149-JD, [8.77, 9.57] for MACSJ1115-JD1, and [8.65, 9.31] for MACSJ1720-JD1. The best-fit redshifts for these three candidates are 9.68, 9.17, and 8.93, respectively. These results are for a flat prior and are quite similar to what we derived using EAZY. Use of the two other priors resulted in similar changes to the redshift likelihood distributions as shown in Figure 5.

²⁶ Of course, allowing for an even larger range of reddenings would be useful for more fully considering the possibility that these candidates might correspond to ULIRGs. However, the moderately blue color of our three candidates likely rules out this possibility.

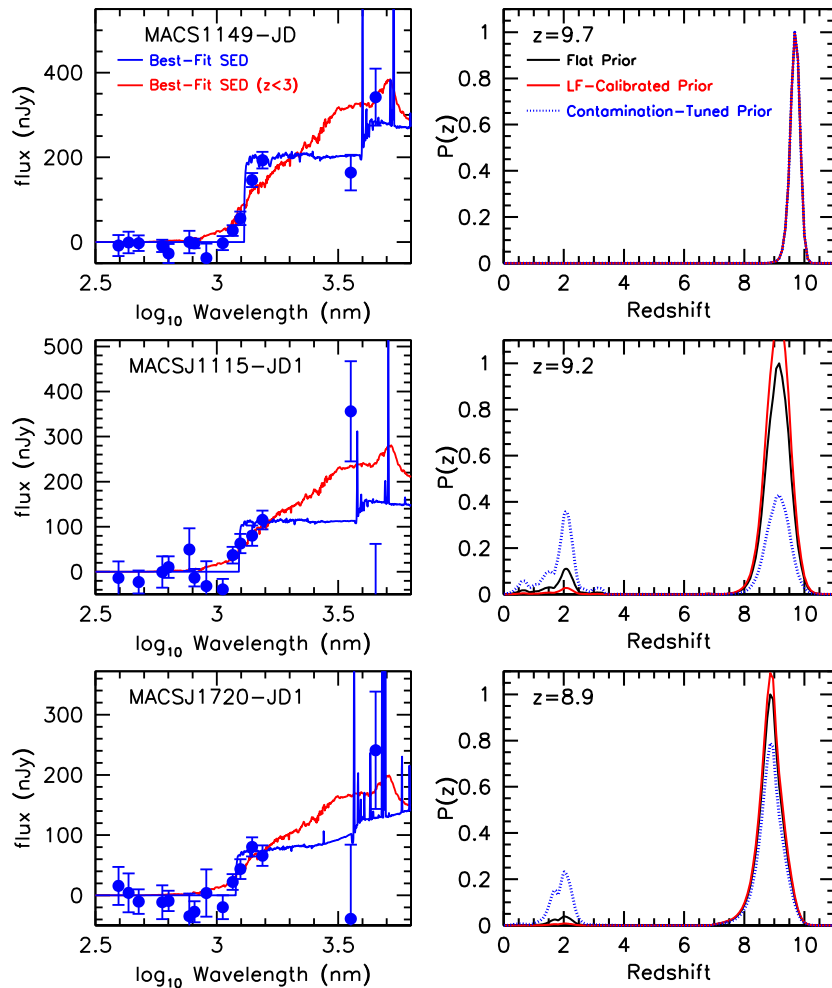


Figure 5. Left: observed spectral energy distributions (solid blue circles) for three $z \sim 9$ galaxy candidates in our selection. The blue line shows the SED template that best fits our observed photometry (using the EAZY photometric redshift code), while the red line shows the best-fit $z \sim 0-3$ SED template. The candidate in the uppermost row was previously presented in Z12. Right: the redshift likelihood distribution computed for our three $z \sim 9$ candidates using the EAZY photometric redshift software (see Section 3.4). We consider three different priors in computing the redshift likelihood distributions: (1) a flat prior (black line), (2) a prior calibrated to reproduce published LFs or LF trends (red line; Giallongo et al. 2005; Bouwens et al. 2011b; R. Quadri et al. 2012, private communication), and (3) a prior tuned to reproduce the results from our photometric scattering simulations (dotted blue line; Section 3.5). Appendix A provides a more detailed description of these priors. Results from our third prior account for the fact that ~ 1 faint source from our selection might be expected to resemble plausible $z \sim 9$ galaxies, due to the effects of noise (see Section 3.5). However, even though we might expect a source to possibly scatter into our $z \sim 9$ selection, we have no evidence that any particular source in our sample actually corresponds to such a low-redshift interloper.

(A color version of this figure is available in the online journal.)

Finally, we also estimated the photometric redshifts of our three candidates with BPZ (Bayesian Photometric Redshift Code; Benítez 2000; Coe et al. 2006). Similar to the analyses in C13 and Z12, we modeled the photometry using SEDs from PEGASE (Fioc & Rocca-Volmerange 1997) adjusted and recalibrated to match the observed photometry of galaxies with known spectroscopic redshifts in the FIREWORKS catalog (Wuyts et al. 2008). This FIREWORKS catalog includes photometry to 24.3 AB mag (5σ) in K_s , for galaxies with $z \sim 3.7$. The best-fit photometric redshifts we derive with BPZ are 9.7, 9.2, and 8.9 for MACS1149-JD, MACSJ1115-JD1, and MACSJ1720-JD1, respectively. For MACS1149-JD, the redshift likelihood distribution is predominantly uni-modal, although in the other two cases the distribution is more bimodal, with modest peaks at lower redshift. 29% and 5% of the total probability for MACS1115-JD1 and MACS1720-JD1, respectively, is at $z \sim 1.0-2.0$. Focusing on the dominant $z \sim 9$ peaks (excluding all $z < 5$ solutions), the 68% confidence intervals on the redshifts for our three candidates are [9.56, 9.87], [8.45, 9.55], and [8.30, 9.26], respectively. These results are for a flat

prior and are somewhat similar to what we derived using the other two photometric redshift codes, although the low-redshift peaks are slightly more significant with BPZ. We opted not to make use of the BPZ prior to computing the redshift distribution for our sources, due to the relative weight it assigns to faint red galaxies at $z \sim 1.3-2$ and blue galaxies at $z \sim 9$ (which differs by more than a factor of 30 from what we compute based on published LFs or LF trends: see Appendix A). We find that a flat prior comes much closer to accurately representing the relative surface densities of these two populations.

3.5. Possible Contamination

While the sources in our current selection are consistent with being $z \sim 9$ galaxies, these sources are faint enough that they could easily have a very different nature. Important sources of contamination for $z > 8$ selections include low-mass stars, supernovae (SNe), emission line galaxies (van der Wel et al. 2011), and the photometric scatter of various low-redshift galaxies. Readers are referred to Bouwens et al. (2011a), Z12, and C13 for extended discussions of these issues.

In general, the most important source of contamination for high-redshift samples results from faint Balmer-break galaxies entering these samples (and hence satisfying their selection criteria) due to the effects of noise (see discussion in Wilkins et al. 2011, Bouwens et al. 2011, Bouwens et al. 2014b). Noise can cause such galaxies (with faint optical flux and not especially red) to look bluer and disappear entirely at optical wavelengths.

Here we test for contamination from faint lower-redshift source scattering into our high-redshift selection through the effects of noise by using all intermediate magnitude sources in the CLASH cluster fields that are detected at $>2\sigma$ in the I_{814} band and Y_{105} bands (and therefore likely at redshifts $z < 6$) to implicitly define the color distribution for potential interlopers to our high-redshift samples. Then, we take all the faint sources in all the CLASH cluster fields (with their $H_{160,AB}$ magnitudes and errors), randomly match them up with a source from the sample that defines our color distribution, give this faint source the same colors as the intermediate-magnitude source, add noise to the photometry of the sources in bluer bands (assuming a normal distribution), and finally see if this source satisfies our $z \sim 9$ selection criteria. Our procedure is essentially identical to what we performed in many previous analyses (e.g., Bouwens et al. 2011a; Bouwens et al. 2011b). In modeling possible contamination of our selection, we only allow for three contaminants at maximum and the n th contaminant per CLASH data set must have a higher S/N than the n th lowest S/N source.

Applying this procedure to all the sources in the CLASH fields $100\times$, we find that only 0.7 lower-redshift ($z \lesssim 6$) sources enter our $z \sim 9$ selection by chance (per Monte-Carlo simulation for the entire CLASH program). The magnitude distribution of these contaminants is shown in Figure 6. The small number of contaminants we find from these simulations demonstrate that the overall level of contamination for the present probe is likely only modest ($\sim 23\%$), becoming important faintward of 26.5 mag.

We also considered the implications for contamination if we had restricted our selection to sources with a $JH_{140} + H_{160}$ detection significance of $>8\sigma$ and $>5.5\sigma$ (weaker than our $z \sim 9$ selection criteria). The results are shown in Figure 6 with the magenta and dotted red lines, respectively. Only ~ 0.2 contaminants are expected for a 8σ detection threshold, while for a $\sim 5.5\sigma$ threshold the expected number of contaminants is ~ 7.5 sources, and hence might be a significant concern if we had considered a lower detection threshold for our selection.

However, it is worth noting that our estimate of the total contamination may be a little high (perhaps by a factor of ~ 2 – 3), due to our use of an intermediate magnitude (~ 24.0 – 25.5 mag) population of galaxies to model the colors of somewhat fainter galaxies (i.e., 26.0–27.0 mag). Because the intermediate magnitude population is somewhat redder in general than the ~ 26.0 – 27.0 mag population (e.g., see Figure 11 of Oesch et al. 2012a), it is more likely to scatter into $z \sim 9$ selections via noise than is the actual situation for ~ 26.0 – 27.0 mag galaxies.

In any case, the results of these simulations strongly suggest that the two most significantly detected sources in our sample (i.e., MACSJ1115-JD1 and especially MACSJ1149-JD), are unlikely to correspond to such contaminants. For sources with lower S/N, we must remain concerned about contamination, even though we cannot establish the exact rate. The issue will contribute to the overall errors in our SFR density estimates at $z \sim 9$ (Section 4).

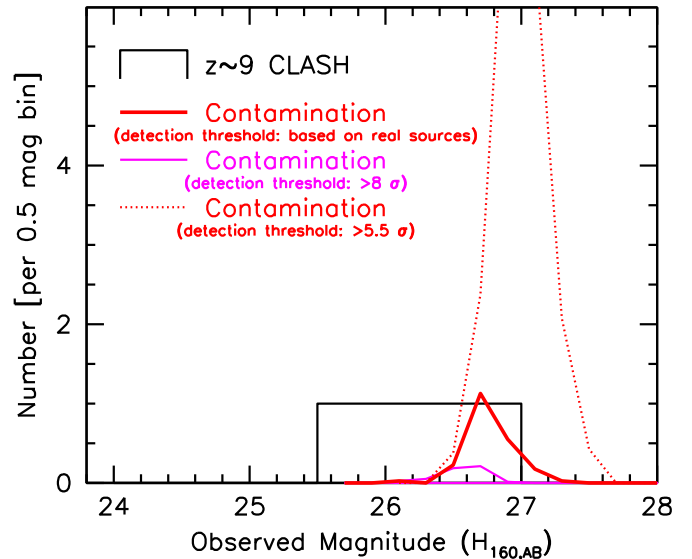


Figure 6. Number of $z \sim 9$ galaxy candidates we find in our CLASH cluster search as a function of the $H_{160,AB}$ -band magnitude. Also plotted (red line) is the number of contaminants we would expect to select in our search fields for sources due to the effects of noise on the photometry of other lower-redshift sources in search fields (see Section 3.5 for details). The total number that we estimate for our search fields is 0.7 (vs. the 3 $z \sim 9$ candidates in our selection). In modeling the possible contamination of our selection, we only allow for three contaminants at maximum and the n th contaminant must have a higher signal to noise than the n th lowest signal-to-noise source. For context, we also show the contamination expected for a $>8\sigma$ selection and for a $>5.5\sigma$ selection. Clearly, contamination from lower-redshift sources (due to photometric scatter) is only especially significant for sources with $H_{160,AB}$ -band magnitudes faintward of 26.5 AB mag. For sources detected at 5.5σ in the $JH_{140} + H_{160}$ bands (with magnitudes ~ 27 AB mag), contamination from lower redshift becomes very important.

(A color version of this figure is available in the online journal.)

Of course, faint moderately blue low-redshift galaxies are not the only galaxies that can contaminate high-redshift samples. Dust reddened galaxies can also occasionally contaminate high-redshift selections (Bowler et al. 2012; Laporte et al. 2011), as well as lower-redshift galaxies with somewhat unusual SEDs (Hayes et al. 2012; Boone et al. 2011). While it is difficult to be sure that sources in our sample do not correspond to such galaxies at lower redshift, all three sources in our samples generally have bluer colors than those lower-redshift contaminants, and so we suspect that such sources do not pose a problem for our selection. The moderately red color of MACS1720-JD1 in the $3.6\mu\text{m}$ and $4.5\mu\text{m}$ bands is similar to the colors seen in other $z \sim 8$ candidates (e.g., Ono et al. 2012; Labbé et al. 2013; Finkelstein et al. 2013; Laporte et al. 2014) likely showing strong [O III] emission.

Another possible source of contamination is from extreme emission-line galaxies (EELGs) with strong [O III] + $H\beta$ emission, such as that recently discovered by van der Wel et al. (2011) and Atek et al. (2011) in the CANDELS (Grogin et al. 2011; Koekemoer et al. 2011) or WISP (Atek et al. 2010) programs. Perhaps the most well-known high-redshift candidate thought to be such an EELG is the Bouwens et al. (2011) $z \sim 2/z \sim 12$ candidate UDFj-39546284 (Ellis et al. 2013; Brammer et al. 2013; Bouwens et al. 2013; Capak et al. 2013). However, it seems unlikely that any of our sources correspond to such candidates given their detection in multiple non-overlapping bands and blue UV -continuum slope of most EELGs (van der Wel et al. 2011).

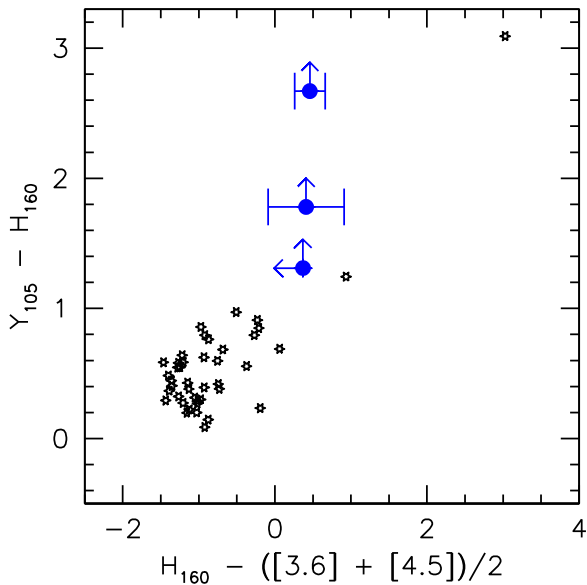


Figure 7. Comparison of the $Y_{105} - H_{160}$ vs. $H_{160} - ([3.6] + [4.5])/2$ colors of the three $z \sim 9$ candidates in our sample (solid blue circles, 1σ error bars, and arrows indicating 1σ limits) with the observed colors of various stars. The black star-like symbols are the colors derived from the substantial library of stellar spectra observed with IRTF (Cushing et al. 2005; Rayner et al. 2009), with sources ranging from very low mass stars to higher mass Mira-type variable stars (the black star-like symbol in the upper right of this figure). Two of our $z \sim 9$ candidate galaxies have $Y_{105} - H_{160}$ colors that are clearly too red to match those colors observed by the broad set of stars encompassed by this library.

(A color version of this figure is available in the online journal.)

Finally, there is the possibility that candidates from our selection could correspond to stars or SNe. Both possibilities would require that sources in our selection are unresolved. Comparing the coadded $JH_{140} + H_{160}$ profile of our candidates with the WFC3/IR PSF, it is clear that two out of our three candidates are resolved (see also discussion in Z12, which clearly demonstrate that MACS1149-JD is resolved). Only MACSJ1115-JD1 does not show any spatial extension. In any case, as Figure 7 demonstrates, the colors of the candidates do not clearly support a stellar origin. The redder Mira-variable stars would appear to give the best match, but their intrinsic luminosities are such that we would need to observe them well outside our own Milky Way galaxy (Whitelock et al. 1995; Dickinson et al. 2000). We can also safely exclude the possibility of a SNe, given that deep optical observations of our cluster fields were obtained over the same two month time window as our deep near-IR observations (Postman et al. 2012; see also Z12 and C13).

One final possibility is that some candidates may correspond to more local solar system or Oort cloud objects. To be consistent with the constraints we can set on the proper motion of our candidates based on our ~ 2 month observational baseline (see Figure 6 of C13 for an illustration of the constraints we can set), such a source would need to be at 50,000 AU. However, at such distances Oort cloud objects would be extremely faint (e.g., faintward of 40 mag), even if these sources were as large as the moon (see also discussion in Z12 and C13).

4. A NEW DIFFERENTIAL DETERMINATION OF THE UV LUMINOSITY FUNCTION AT $Z \sim 9$

The present $z \sim 9$ sample is the largest such sample available to date and should allow us to substantially improve

our constraints on the $z \sim 9$ LF. However, before providing a detailed discussion of the specific constraints we are able to set, we must first include a few words on the procedure we adopt.

4.1. UV LF Evolution from Lensing Cluster Searches: Rationale for Using a Differential Approach

Normally, we would derive the LF for $z \sim 9$ galaxies using the same approach that has been followed in the field, i.e., (1) distribute the sources in one's samples into different magnitude intervals, (2) count the total number of sources in a magnitude interval (after correcting for contamination), and (3) divide these numbers by the effective volume where such sources could be found. This procedure has been followed in a number of previous works (e.g., Santos et al. 2004; Richard et al. 2008).

However, even the simple process of placing sources into different intrinsic magnitude bins can be quite uncertain due to its dependence on a particular magnification model. Calculations of the selection volumes are just as equally model dependent. While in many cases these model dependencies may not result in large overall uncertainties in one's results, the uncertainties clearly become large (~ 0.3 – 0.4 mag or larger) near the critical curves of the lensing models, where the magnification factors become nominally infinite (e.g., see Figure 2 of Maizy et al. 2010). These issues can potentially have a huge effect on the LFs derived in the context of lensing clusters.²⁷

Ideally, we would like to determine the UV LF at $z \sim 9$ in a way that avoids these uncertainties. One possible way for us to do this is (1) to leverage existing well-determined LFs that already exist at $z \sim 7$ – 8 from blank field studies (e.g., Bouwens et al. 2011b; Oesch et al. 2012b; Bradley et al. 2012) not subject to potentially large selection volume uncertainties, and then (2) use our searches for $z \sim 7$ – 10 galaxies behind lensing clusters to derive the differential evolution in the LF from $z \sim 9$ to $z \sim 7$ – 8 . This provides us with a somewhat indirect approach to deriving the LF at $z \sim 9$ and takes advantage of the very similar effect gravitational lensing from low-redshift clusters has on light from the high-redshift universe, regardless of the exact redshift of the source. Fundamentally, this is due to the fact that the D_{LS}/D_S factor is very insensitive to redshift when the lensed source is at $z > 5$ (i.e., very distant) and the lensing cluster is relatively close (i.e., $z \sim 0.1$ – 0.5). For example, for a $z \sim 0.4$ lensing cluster the computed D_{LS}/D_S factor for $z \sim 9$ background sources is only $\sim 1\%$ higher than the D_{LS}/D_S factor for $z \sim 8$ sources. D_{LS} and D_S are the angular-diameter distances from the cluster lens to source and from observer to source, respectively (e.g., Narayan & Bartelmann 1996).

As a result, for sources seen behind a given lensing cluster, the $z \sim 8$ universe is magnified in almost exactly the same way as the $z \sim 9$ universe. This can be illustrated using the lensing models we have available for three of the CLASH clusters (Figure 8). The total area available behind a given cluster to magnify the background light by more than a factor of three is almost exactly the same for the $z \sim 8$ universe as for the $z \sim 9$

²⁷ Of course, for realistic LFs these uncertainties may not be especially problematic. Indeed, for LFs with an effective faint-end slope close to -2 , uncertainties in the magnification factor tradeoff almost perfectly with uncertainties in the search volume, so as to have no large effect on the inferred LFs. As a result, one potentially very effective approach for minimizing the impact of these uncertainties on the derived LFs is by marginalizing over the magnification factor in performing the comparisons with the observed numbers (C13). The excellent agreement between the present estimate of the SFR density at $z \sim 9$ and that obtained by C13 based on the Z12 search results would seem to support this conclusion.

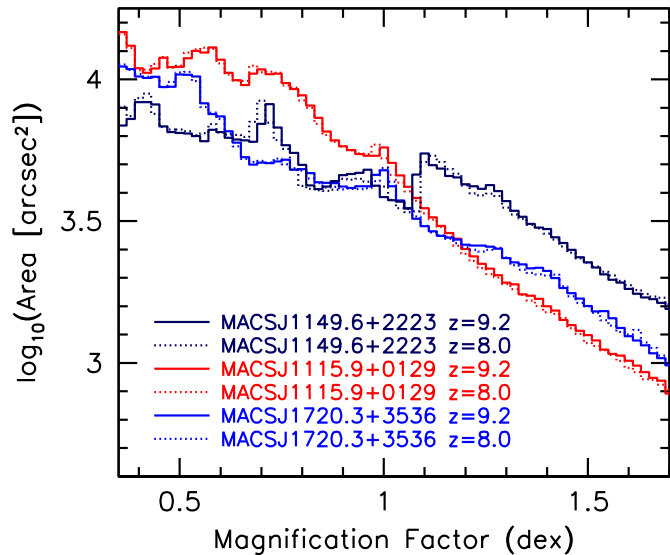


Figure 8. Search area (per unit dex) behind select galaxy clusters subject to varying levels of magnification by gravitational lensing. Results are shown for sources at $z = 8$ and $z = 9$ based on the lens models for MACSJ1149.6+2223, MACSJ1115.3+0129, and MACSJ1720.3+3536 (Z12; A. Zitrin et al. 2014, in preparation; M. Carrasco et al. 2014, in preparation). It is obvious from these results that the total search volume behind a cluster (given the area magnified to various levels) can show a huge variation from one cluster to another. However, if one utilizes the same cluster to search for sources at similar but slightly different redshifts (compare the dotted and solid lines representing $z \sim 8$ and $z \sim 9$ selections), almost exactly the same selection area is available for selecting sources at a given magnification factor (although we note that the selection area is slightly larger ($\sim 1\%$ – 3%) at $z \sim 9$ than at $z \sim 8$). As a result, we would expect the relative selection volumes for a $z \sim 9$ search behind lensing clusters and a $z \sim 8$ search behind lensing clusters to be very well defined, if the same clusters are utilized for the two searches.

(A color version of this figure is available in the online journal.)

universe. Note that this is true even if the precise position of the critical curves at $z \sim 9$ lies in a slightly different position from the critical curves at, for example, $z \sim 8$.

Because of the very similar way a given set of clusters magnifies galaxies at $z \sim 9$ and at other similar redshifts (e.g., $z \sim 8$), one might expect it to have the same effect on the total surface densities of these galaxies that one finds on the sky. Therefore, if one starts with the same LF of galaxies at both $z \sim 8$ and $z \sim 9$, one would expect to find roughly the same surface density of these galaxies on the sky, modulo two slight differences. The $z \sim 9$ galaxy distribution would be shifted to slightly fainter magnitudes (e.g., by ~ 0.3 mag versus $z \sim 8$) to reflect their slightly larger luminosity distances and would be present at slightly lower surface densities (by $\sim 10\%$ versus $z \sim 8$) reflecting the smaller cosmic volume available at $z \sim 9$.

Even multiple imaging of the same high-redshift sources would not appreciably affect the ratio of sources seen at different redshifts, as one would expect galaxies at $z \sim 9$ and similar redshifts to give rise to lensed multiples to approximately the same degree, and therefore the ratio of surface densities should be preserved. However, since multiple images of a single background source are not independent events, not accounting for this effect could have a slight effect on the uncertainties we estimate for the relative surface densities of galaxies at different redshifts.

Given this situation, it seems clear we should be able to use the relative surface densities of galaxies in different redshift samples to make reasonably reliable inferences about the relative

volume densities of the galaxy population at different epochs (after making small adjustments to the numbers to account for the factors discussed above).

4.2. $z \sim 8$ Comparison Sample

Redshift $z \sim 8$ selections serve as the perfect comparison sample for our $z \sim 9$ studies. Not only is the $z \sim 8$ universe close enough to $z \sim 9$ to make differences in the lensing effects quite small overall, but the ~ 70 – 80 $z \sim 8$ galaxies available in current WFC3/IR surveys allow the LF there to be robustly established from field studies (e.g., Bouwens et al. 2011b; Lorenzoni et al. 2011; Oesch et al. 2012b; Bradley et al. 2012). This allows us to put together the new information we have on the differential evolution of the LF from $z \sim 9$ to $z \sim 8$ with previous $z \sim 8$ LF determinations to estimate the approximate UV LF at $z \sim 9$.

Finally, given the observed rate of evolution in M^* and α (e.g., using the fitting formula for Schechter (1976) parameterization given in Bouwens et al. 2014b), we would expect the shape of the LF at $z \sim 8$ to be similar to the shape of the LF at $z \sim 9$, that is, $\Delta(M^*(z = 8) - M^*(z = 9)) \lesssim 0.2$ and $\Delta\alpha(z = 8) - \alpha(z = 9) \lesssim 0.12$, so we can model any evolution in the LF very simply assuming a change in the normalization ϕ^* (although modeling the evolution in terms of the characteristic luminosity M^* is only slightly more involved).

For our $z \sim 8$ comparison sample, we use the same selection criteria as previously utilized in Bouwens et al. (2011b) and Oesch et al. (2012b), i.e.,

$$(Y_{105} - J_{125} > 0.45) \wedge (J_{125} - H_{160} < 0.5).$$

As in these two previous works (and as performed for our $z \sim 9$ – 10 selection), we also require sources to be undetected in the I_{814} band and blueward both in individual bands at $< 2\sigma$ and using the χ_{opt}^2 statistic discussed earlier (Section 3.2). We also demand that sources be detected at $> 6\sigma$ in a combined JH_{140} and H_{160} image ($0''.35$ diameter aperture), as performed for our primary $z \sim 9$ – 10 selection. Because these color criteria and selection criteria are very similar to that used by Bouwens et al. (2011b) and Oesch et al. (2012b) in identifying $z \sim 8$ galaxies, the redshift distribution for the present $z \sim 8$ selection should be approximately the same as shown in Figure 1 (red line).

Applying this selection criteria to the 19 cluster CLASH data set, we find a total of 19 sources that satisfy our $z \sim 8$ criteria. After excluding one candidate from the sample (19:31:48.7, $-26:34:03.0$), which is completely unresolved in the *HST* data²⁸ and has colors very similar to that of low-mass stars, we are left with a total sample of 18 $z \sim 8$ candidates. These sources have $H_{160,AB}$ magnitudes ranging from 25.0 to 27.3 mag. Coordinates of these candidates and their H_{160} -band magnitudes are provided in Table 5 from Appendix C. We allow for a potential contamination of ~ 1.5 source in our $z \sim 8$ sample, which is consistent with the contamination level found by Bouwens et al. (2011b) for their $z \sim 8$ sample and also allows for some possible contamination by low-mass stars in our search fields.

The current $z \sim 8$ selection includes more $z \sim 8$ candidates per cluster than the $z \sim 8$ selection from Bradley et al. (2014) using photometric redshifts over the same magnitude range. This is due to our current color criteria being effective at selecting galaxies at $z \gtrsim 7.2$, while the Bradley et al. (2014) photometric

²⁸ Median SExtractor stellarity parameter for this candidate is 0.94 in the $Y_{105}J_{110}J_{125}JH_{140}H_{160}$ data (where 1 and 0 correspond to a point and extended source, respectively).

redshift criteria only identify galaxies at $z \gtrsim 7.5$. Our choice of selection criteria should have little impact on our LF results, as the larger selection volumes we compute for our $z \sim 8$ sample (Section 4.3) will generally offset any changes in sample size.

4.3. Relative Selection Volumes at $z \sim 8$ and $z \sim 9$: Expected Sample Sizes Assuming no Evolution

In order to utilize the relative surface density of $z \sim 8$ and $z \sim 9$ galaxy candidates we observe to make inferences about the evolution of the LF, we must have an estimate for how many galaxies we would expect in the two samples if the UV LF did not evolve at all between the two epochs. Then, based on the relative number of sources expected in the two samples assuming no evolution, we can determine the approximate evolution in the LF from $z \sim 9$ to $z \sim 8$. With this step, we effectively account for the approximate difference in selection volume for our $z \sim 8$ and $z \sim 9$ samples.

The simplest way for us to account for any evolution in the UV LF is through the normalization ϕ^* , because it simply requires us to compare the number of sources we find in our $z \sim 8$ and $z \sim 9$ samples with that found in our simulations (see below) to derive the approximate evolution,

$$\phi^*(z=9) = \phi^*(z=8) \frac{n_{\text{obs},z=9} n_{\text{no-evol-sim},z=8}}{n_{\text{obs},z=8} n_{\text{no-evol-sim},z=9}}, \quad (1)$$

where $n_{\text{obs},z=9}$ is the number of sources in our $z \sim 9$ selection after correction for contamination (i.e., ~ 2.3), $n_{\text{obs},z=8}$ is the number of sources in our $z \sim 8$ selection after correction for contamination (i.e., ~ 18), $n_{\text{no-evol-sim},z=8}$ is the number of $z \sim 8$ candidates we find in our simulations for our $z \sim 8$ selection based on a fiducial lensed LF, and $n_{\text{no-evol-sim},z=9}$ is the number of $z \sim 9$ candidates we find in our simulations for our $z \sim 9$ selection based on this same LF.

As in our previous papers, we estimate the relative numbers of sources we would expect at both redshifts from simulations. To perform these simulations, we insert artificial sources with a variety of redshifts and luminosities into the real observations and then attempt to select these objects using our $z \sim 8$ and $z \sim 9$ –10 selection criteria. We generate artificial images for each source in these simulations using our well-tested cloning software (Bouwens et al. 1998; Bouwens et al. 2003; Bouwens et al. 2007), which we use to artificially redshift similar luminosity $z \sim 4$ galaxies from the HUDF to higher redshift. We scale the size of galaxies at fixed luminosity as $(1+z)^{-1}$ to match the observed size-redshift scaling at $z > 3$ (e.g., Bouwens et al. 2004, 2006a; Oesch et al. 2010; Mosleh et al. 2012; but see also Ferguson et al. 2004). We take the UV -continuum slope β of galaxies in our simulations to have a mean value and 1σ scatter of -2.3 and 0.45 , respectively, to match the observed trends extrapolated to $z \sim 8$ –9 (Bouwens et al. 2012b; Finkelstein et al. 2012; Bouwens et al. 2014a).

For simplicity, we estimate the relative numbers of sources we would expect in both samples without making use of the deflection maps estimated for all 19 CLASH clusters used in the present search. As we demonstrate in Appendix B, we can approximately ignore the effect of lensing in estimating the selectability of sources if the quantity we are interested in calculating is the relative number expected for sources in two adjacent redshift samples (e.g., $z \sim 8$ and $z \sim 9.2$). Lensing does not have a big influence on the relative number of sources seen in two adjacent samples due to the similar impact it has on

the selection efficiencies, volumes, and luminosities of galaxies in both samples. Nevertheless, for the simulations we run, the relative numbers of lensed $z \sim 9.2$ galaxies to lensed $z \sim 8$ galaxies is slightly lower (16%) in our simulations than if we ignore the impact of lensing in calculating its selectability (and only consider a boost to the LF from some fiducial magnification factor).

The UV LFs we input into the simulations have the following parameters: $M_{UV}^* = -22.4$, $\alpha = -2.0$, and $\phi^* = 5.5 \times 10^{-5} \text{ Mpc}^{-3}$. These luminosity parameters were chosen to implicitly include a factor of ~ 9 magnification from gravitational lensing—which is the median magnification estimated for sources in our selection—so the effective M^* at $z \sim 8$ is chosen to be ~ 2.4 mag brighter than seen in blank field studies (e.g., Oesch et al. 2012b). The faint-end slope we assume approximately matches what we would expect based on the UV LF results at $z \sim 7$ –8 (Bouwens et al. 2011b; Oesch et al. 2012b; Bradley et al. 2012), which point to faint-end slopes α of -2 . No change is required in the faint-end slope α of the LF, due to the perfect tradeoff between magnification and source dilution effects for slopes of -2 (e.g., Broadhurst et al. 1995). The normalization ϕ^* we choose has no effect on our final results (due to the differential nature of this calculation). While the LFs we adopt for these simulations could, in principle, affect our evolutionary results, the overall size of such effects will be small due to the differential nature of the comparison we are making. We also verified that the surface density of $z \sim 8$ sources predicted by this model LF showed a very similar magnitude dependence, as seen for our $z \sim 8$ sample.

Using the aforementioned simulation procedure and LF, we repeatedly added artificial sources to the real CLASH observations for all 19 CLASH clusters, created catalogs, and repeated our $z \sim 8$ and $z \sim 9$ selections. In total, we repeated the described simulations 20 times for each cluster field to obtain an accurate estimate of the total number of sources (selection volume and redshift distribution) we would expect to find in each sample, given the described LF.

In total, we find 657 sources that satisfy our $z \sim 8$ selection criteria and 383 sources that satisfy our $z \sim 9$ selection criteria, based on the same luminosity and simulation area (so $n_{\text{no-evol-sim},z=8} = 657$ and $n_{\text{no-evol-sim},z=9} = 383$ in Equation (1) above). These results imply that without the impact of gravitational lensing, we would expect to find 58% as many $z \sim 9$ galaxies as $z \sim 8$ galaxies behind our CLASH cluster sample. We adopt the Oesch et al. (2012b) LF in performing this estimate (keeping with our original treatment), however we would have obtained essentially the same result using other recent $z \sim 8$ LF results (e.g., Bouwens et al. 2014b; McLure et al. 2013; Schenker et al. 2013).

As we demonstrate in Appendix B, we would expect to find a very similar ratio of galaxies in these two samples, even including lensing in the simulations. In the case of lensing, we expect just 49% as many $z \sim 9$ galaxies as $z \sim 8$ candidates over our search fields (although we note that the precise factor depends slightly [i.e., $\lesssim 20\%$] on the lensing model and LF adopted in performing the calculation). Therefore, to make a fair comparison between our $z \sim 8$ and $z \sim 9$ samples we need to multiply the surface densities in our $z \sim 8$ sample by 0.49 (Equation (1) above). In Figure 9, we show the comparison of the surface densities of $z \sim 8$ galaxies found over the first 19 CLASH clusters (corrected for the difference in selection volume) with the surface densities of $z \sim 9$ galaxies found over these clusters.

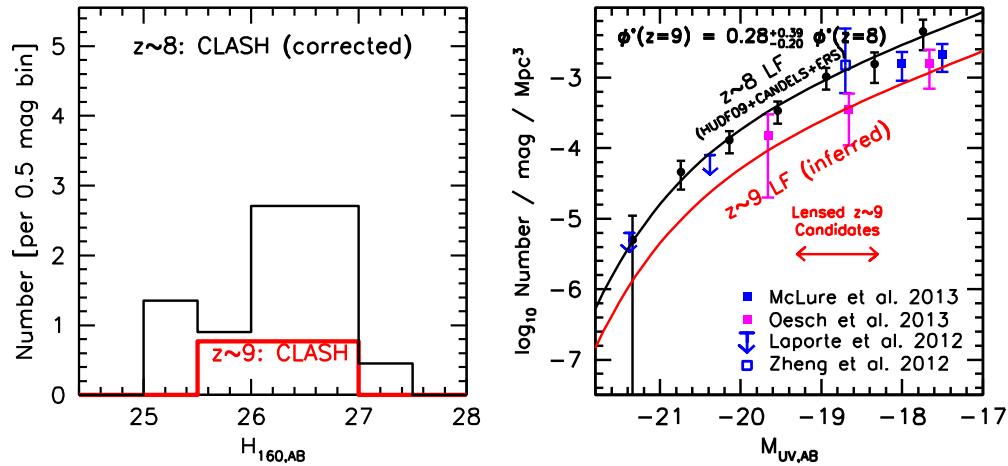


Figure 9. Illustration of our differential approach to deriving the UV LF at $z \sim 9$. Left: the contamination-corrected number of $z \sim 9$ galaxy candidates we find within CLASH (red histogram) vs. the number of $z \sim 8$ galaxy candidates behind the same CLASH clusters, corrected to have the same selection volume as at $z \sim 9$. For simplicity, the contamination rate correction is applied in a magnitude-independent manner (although the contamination rate will clearly be higher near the faint ends of our two samples). The number of $z \sim 9$ galaxy candidates in CLASH, after contamination correction, is just $0.28^{+0.39}_{-0.20} \times$ that at $z \sim 8$. A simple comparison of these surface densities should give us a fairly model-independent measure of the relative normalization of the UV LF at $z \sim 8$ and the UV LF at $z \sim 9$ —assuming that the shape of the LF (i.e., M^* and α) does not change very dramatically from $z \sim 9$ to $z \sim 8$. Right: the observed UV LF at $z \sim 8$ as derived by Oesch et al. (2012b; black points, error bars, and line) based on the HUDF09+CANDELS+ERS data set and our newly inferred UV LF at $z \sim 9$ (red line) based on our differential comparison of our $z \sim 8$ and $z \sim 9$ selections. We infer that the UV LF at $z \sim 9$ has an effective ϕ^* that is just $0.28^{+0.39}_{-0.20} \times$ that at $z \sim 8$. The red horizontal arrow toward the bottom of this panel indicates the approximate luminosities inferred for our three $z \sim 9$ candidates (after correction for lensing magnification: see Section 3.3). For context, we also show recent constraints on the volume density of $z \sim 9$ galaxies from Zheng et al. (2012), Laporte et al. (2012), McLure et al. (2013), and Oesch et al. (2013).

(A color version of this figure is available in the online journal.)

4.4. Inferred UV LF at $z \sim 9$

After correction for possible contamination of our selection by possible low-redshift contaminants (see Section 3.5), the total number of $z \sim 9$ candidates in our $z \sim 9$ selection is 2.3. This number is just $0.28^{+0.39}_{-0.20} \times$ the total number of $z \sim 8$ sources to a similar luminosity limit (corrected for differences in the selection volume: see Figure 9). In calculating the uncertainties on the fraction $0.28^{+0.39}_{-0.20}$, we have accounted for the Poissonian errors on the total number of galaxies in the $z \sim 8$ and $z \sim 9$ samples, as well as the Poissonian uncertainties in the contamination rates.

Assuming that we can approximate the differences between the $z \sim 8$ and $z \sim 9$ LFs as occurring simply through density evolution (i.e., by changing ϕ^*), we infer that the value of ϕ^* at $z \sim 9$ is just $0.28^{+0.39}_{-0.20} \times \phi^*$ at $z \sim 8$ (or $0.28^{+0.84}_{-0.26} \times$ if 2σ errors are used). The present search is inconsistent, with no evolution at $>92\%$ confidence.²⁹

Using the recent Oesch et al. (2012b) determination from HUDF09+CANDELS+ERS field studies (Bouwens et al. 2011b; Grogin et al. 2011; Koekemoer et al. 2011; Windhorst et al. 2011) that ϕ^* at $z \sim 8$ is $5.0^{+7.0}_{-3.3} \times 10^{-4} \text{ Mpc}^{-3}$, we estimate that ϕ^* at $z \sim 9$ is $1.4^{+2.0}_{-1.1} \times 10^{-4} \text{ Mpc}^{-3}$. For the purpose of parameterizing a $z \sim 9$ LF, we will assume that M^* and α at $z \sim 9$ match that derived by Oesch et al. (2012b) at $z \sim 8$. The resultant $z \sim 9$ LF is illustrated in the right panel of Figure 9 and compared with the Oesch et al. (2012b) $z \sim 8$ LF. Our new

$z \sim 9$ LF parameters are also presented in Table 3. Use of other recent $z \sim 8$ LF results (e.g., Bouwens et al. 2014b; Bradley et al. 2012; Schenker et al. 2013; McLure et al. 2013; Schmidt et al. 2014) yield very similar results for M^* ($\lesssim 0.3$ mag), α ($\lesssim 0.2$), and ϕ^* (< 0.1 dex).³⁰

What effect will field-to-field variations (i.e., “cosmic variance”) have on the overall uncertainties here? To estimate the size of these uncertainties, we first considered the case of a single cluster field. We used the Trenti & Stiavelli (2008) cosmic variance calculator, assumed a mean redshift of 8.0 and 9.2 for our two samples (as estimated from our simulations: see Figure 1), took the Δz width for these redshift distributions to be 0.8, and assumed that the relevant area in the source plane was 0.4×0.4 . The latter area in the source plane assumes a factor of ~ 10 dilution of the total search area (consistent with the mean magnification factors found here) and further that only 30% of the total area on our WFC3/IR images is effective for finding $z \sim 9$ galaxies. The fractional uncertainty we estimated in our volume density estimates from field-to-field variations is 0.55 and 0.58 for our $z \sim 8$ and $z \sim 9$ selections, respectively, over a single CLASH cluster field.

Because each of our cluster fields provides an independent sightline on the high-redshift universe, we need to reduce the derived variance by $\sim 19^{0.5} \sim 4.4$ (although we note that the actual reduction will be slightly less than this because all our clusters will not receive equal weight in the total volume calculation and hence the gains from independent sightlines will be less). This results in fractional uncertainties of ~ 0.13 in the total number of sources in the current $z \sim 8$ and $z \sim 9$ samples. Since our $z \sim 9$ LF estimate is based on a differential comparison of the present 19 cluster $z \sim 8$ and $z \sim 9$ samples, we must add both of these uncertainties in quadrature to

²⁹ Given the approximate degeneracy between evolution in M^* and ϕ^* for LFs at $z \sim 7-9$, where a $\Delta M^* = 1$ mag change trades off for a $\Delta \phi^*$ change (e.g., Figure 8 of Oesch et al. 2012b), we could reframe the inferred evolution in ϕ^* from $z \sim 9$ to $z \sim 8$ in terms of an equivalent evolution in M^* (as we have parameterized the LF evolution in the past, e.g., Bouwens et al. 2007; Bouwens et al. 2008; C13). We estimate that the effective M^* at $z \sim 9$ is $0.5^{+0.4}_{-0.3}$ mag fainter than at $z \sim 8$ (keeping ϕ^* fixed). However, in a more recent and comprehensive study of the UV LFs from $z \sim 7$ to $z \sim 4$, Bouwens et al. (2014b) found that the overall evolution can be better represented by an evolution in ϕ^* and α (with a more limited evolution in M^*).

³⁰ We persist in our reliance on the Oesch et al. (2012b) LF results to maintain consistency with our original submission (but note the overall agreement of the Oesch et al. 2012b results with more recent determinations of the LF at $z \sim 8$).

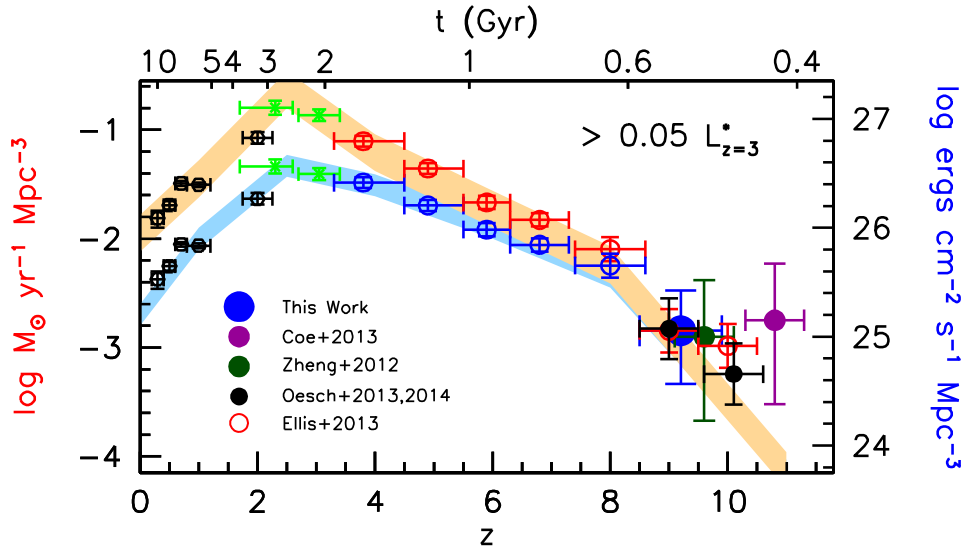


Figure 10. UV luminosity density (right axis) and star formation rate density (left axis) vs. redshift. The UV luminosity and SFR density shown at $z \sim 9$ (large blue solid circle) are from the present work and inferred based on the relative number of $z \sim 8$ and $z \sim 9$ galaxies found within the CLASH cluster program (see Section 4.5). These luminosity densities and SFR densities are only considered down to a limiting luminosity of -17.7 AB mag, which is the approximate limit of both the HUDF09 probe (Bouwens et al. 2011b) and the present search assuming a maximum typical magnification factor of ~ 9 and limiting magnitude of ~ 27.0 mag. The UV luminosity is converted into a star formation rate using the canonical UV -to-SFR conversion factors (Madau et al. 1998; Kennicutt 1998). The upper set of points at every given redshift and orange contour show the dust-corrected SFR densities, while the lower set of points and blue contours show the inferred SFR densities before dust correction. Dust corrections at $z > 3$ are estimated based on the observed UV -continuum slope distribution and are taken from Bouwens et al. (2012b). At $z \leq 3$, the dust corrections are from Schiminovich et al. (2005) and Reddy & Steidel (2009). UV luminosity density and SFR density determinations from the literature are from Schiminovich et al. (2005) at $z < 2$ (black hexagons), Reddy & Steidel (2009) at $z \sim 2-3$ (green crosses, Bouwens et al. (2007) at $z \sim 4-6$ (open red and blue circles), Bouwens et al. (2011b) at $z \sim 7$ (open red and blue circles), Oesch et al. (2012b) at $z \sim 8$ (open red and blue circles), Ellis et al. (2013) at $z \sim 9-10$ (open red circles), Oesch et al. (2013) at $z \sim 9$ (solid black circles), and Oesch et al. (2014) at $z \sim 10$ (solid black circles). Estimates of the SFR density at $z \sim 9.6$ and $z \sim 10.8$ as derived in C13, based on the $z \sim 9.6$ Z12 and $z \sim 10.8$ C13 candidates, are also shown (dark green and magenta solid circles, respectively). Conversion to a Chabrier (2003) IMF would result in a factor of ~ 1.8 (0.25 dex) decrease in the SFR density estimates given here. The present $z \sim 9$ determination is in good agreement with the trend in the SFR density and UV luminosity, as defined by the Oesch et al. (2012a) and Z12 estimates.

(A color version of this figure is available in the online journal.)

derive the approximate fractional uncertainty. The result is 0.19. By comparison, $z \sim 9$ searches using a single 4.4 arcmin² deep field would yield a fractional uncertainty of ~ 0.5 in the volume density of $z \sim 9$ galaxies from large-scale structure (“cosmic variance”). This is much higher than the present uncertainties arising from large-scale structure.

Overall, the uncertainties from large-scale structure only have a fairly marginal impact on our total uncertainty in ϕ^* for the $z \sim 9$ LF, increasing it by just 3% over what one would estimate based on the small numbers in the current $z \sim 9$ selection.

4.5. UV Luminosity and Star Formation Rate Density at $z \sim 9$

We can utilize our newly estimated $z \sim 9$ LF to determine the approximate UV luminosity density and SFR density at $z \sim 9-10$. We compute these luminosity densities to a limiting luminosity $0.05 L_{z=3}^*$, which is the effective limit of the Oesch et al. (2012b) $z \sim 8$ LF we used as a reference point for inferring the $z \sim 9$ LF. This limiting luminosity is also what one would expect for $z \sim 8-9$ searches in the CLASH program to ~ 27 AB mag assuming a $\sim 9\times$ magnification factor, which is equivalent to the average magnification factor for $z \sim 9$ galaxy candidates uncovered in the present search. We can convert the UV luminosities we estimate to SFR densities using the canonical UV luminosity-to-SFR conversion factor (Madau et al. 1998; see also Kennicutt 1998).

The $z \leq 8$ SFR density determinations are corrected for dust extinction based on the values Bouwens et al. (2014a) estimate based on the observed UV -continuum slopes β and the Meurer et al. (1999) IRX- β . Given the observed trends toward bluer UV -continuum slopes β at very high redshifts

Table 3

Estimated Schechter Parameters for the UV LF at $z \sim 9$ and a Comparison with UV LF Determinations at other Redshifts $z \sim 4-10$ (see Section 4.4)

Dropout Sample	Redshift	M_{UV}^* ^a	ϕ^* (10^{-3} Mpc ⁻³)	α
$J_{110} + J_{125}$	9.2	-20.04 (fixed)	$0.14^{+0.20}_{-0.11}$	-2.06 (fixed)
B	3.8	-21.07 ± 0.08	$1.41^{+0.23}_{-0.20}$	-1.64 ± 0.04
V	4.9	-21.19 ± 0.11	$0.64^{+0.14}_{-0.12}$	-1.78 ± 0.05
i	5.9	-21.16 ± 0.20	$0.33^{+0.15}_{-0.10}$	-1.91 ± 0.09
z	6.8	-21.04 ± 0.26	$0.22^{+0.14}_{-0.09}$	-2.06 ± 0.12
Y	8.0	$-20.04^{+0.44}_{-0.48}$	$0.50^{+0.70}_{-0.33}$	$-2.06^{+0.35}_{-0.28}$

Notes. ^a Values of M_{UV}^* are at 1600 \AA for the Bouwens et al. (2014b) $z \sim 4$, $z \sim 5$, $z \sim 6$, and $z \sim 7$ LFs and at $\sim 1750 \text{ \AA}$ for the Oesch et al. (2012b) constraints on the $z \sim 8$ LF. While the $z \sim 8$ LF results from Bouwens et al. (2014b) likely represent an improvement on those from Oesch et al. (2012b), we quote the Oesch et al. (2012b) results here because they represent our baseline for extending the LF results to $z \sim 9.2$ (to maintain consistency with our earlier submission and because of excellent agreement between the Oesch et al. 2012b LF results and subsequent work at $z \sim 8$).

(e.g., Stanway et al. 2005; Bouwens et al. 2006, 2009, 2012b, 2014a; Finkelstein et al. 2012; Wilkins et al. 2011), we would expect the dust extinction at $z \sim 9-10$ to be zero, and therefore apply no dust correction to the SFR density determinations there.

We present the UV luminosity and SFR densities that we estimate at $z \sim 9$ in Figure 10 and also in Table 4. For context, we also provide the SFR and UV luminosity densities of several noteworthy determinations in the literature over the redshift range $z \sim 0$ to $z \sim 10$ (Bouwens et al. 2007, 2010; Ellis et al.

Table 4
UV Luminosity Densities and Star Formation Rate Densities to -17.7
AB mag (0.05 $L_{z=3}^*$: see Section 4.5)^{a,b}

Dropout Sample	$\langle z \rangle$	$\log_{10} \mathcal{L}$ (ergs s^{-1} $\text{Hz}^{-1} \text{Mpc}^{-3}$)	\log_{10} SFR density ($M_{\odot} \text{Mpc}^{-3} \text{yr}^{-1}$) Uncorrected	Corrected ^c
<i>J</i>	9.2	$25.03^{+0.37}_{-0.49}$	$-2.87^{+0.37}_{-0.49}$	$-2.87^{+0.37}_{-0.49}$
<i>B</i>	3.8	26.42 ± 0.05	-1.48 ± 0.05	-1.10 ± 0.05
<i>V</i>	5.0	26.20 ± 0.06	-1.70 ± 0.06	-1.36 ± 0.06
<i>i</i>	5.9	25.98 ± 0.08	-1.92 ± 0.08	-1.67 ± 0.08
<i>z</i>	6.8	25.84 ± 0.10	-2.06 ± 0.10	-1.83 ± 0.10
<i>Y</i>	8.0	25.58 ± 0.11	-2.32 ± 0.11	-2.17 ± 0.11
<i>J</i> ^d	10.0	24.45 ± 0.36	-3.45 ± 0.36	-3.45 ± 0.36

Notes.

^a Integrated down to 0.05 $L_{z=3}^*$. Based upon the $z \sim 9$ inferred here (Table 3; Section 4.4) and the LF parameters in Oesch et al. (2012a, 2012b) and Table 4 of Bouwens et al. (2014b) (see Section 4.5). The SFR density estimates assume $\gtrsim 100$ Myr constant SFR and a Salpeter IMF (e.g., Madau et al. 1998). Conversion to a Chabrier (2003) IMF would result in a factor of ~ 1.8 (0.25 dex) decrease in the SFR density estimates given here.

^b Uncertainties on the luminosity densities and star formation rate densities at $z \sim 9.2$ are calculated by adding in quadrature the logarithmic uncertainties on both the $z \sim 8$ densities and the differential evolutionary factors from $z \sim 9.2$ to $z \sim 8$. Uncertainties on the luminosity densities and star formation rate densities at $z \sim 4-8$ are computed by marginalizing over the likelihood contours for Schechter fits to the $z \sim 4-8$ LFs (from Bouwens et al. 2014b).

^c Dust corrections are from Bouwens et al. (2014a) and are based on the observed UV -continuum slopes. No dust correction is assumed at $z \gtrsim 9$.

^d $z \sim 10$ determinations and limits are from Bouwens et al. (2014b).

2013; Oesch et al. 2013, 2014). We also show the SFR density estimates at $z \sim 9.6$ and $z \sim 10.8$ from the $z \sim 9.6$ Z12 and $z \sim 10.8$ C13 candidates, as estimated by C13.

4.6. Implications of the Present $z \sim 9-10$ Search for the Evolution of the LF at $z > 6$

One of our primary motivations for obtaining constraints on the UV LF at $z \sim 9$ was to characterize the evolution of the UV LF at $z > 8$ and to test whether the UV LF at $z > 8$ really evolves more rapidly as a function of redshift—as recently found by Oesch et al. (2012a; see also Bouwens et al. 2012a)—or if the evolution is more consistent with a simple extrapolation of the UV LF trends found by Bouwens et al. (2011b; see also Bouwens et al. 2008) over the redshift range $z \sim 4-8$. Several theoretical models (Trenti et al. 2010; Lacey et al. 2011) support the idea that the UV LF might indeed evolve faster at $z > 8$ as a function of redshift than at $z \sim 4-8$ (e.g., Figure 8 of Oesch et al. 2012a or Figure 10 of Oesch et al. 2014), and we want to test this hypothesis using our current results.

To determine which of these two scenarios the present observations favor, we first compute the change in ϕ^* that each would predict. Using the Bouwens et al. (2014b) fitting formula for the evolution of the UV LF, we estimate an expected change of $\Delta \log_{10} \phi^* = 0.23$ and $\Delta M_{UV}^* \sim 0.15$ in the UV LF from $z \sim 8$ to $z \sim 9.2$ (the mean redshift of our sample). Taking advantage of the approximate degeneracy between M^* and ϕ^* at $z \sim 7$ and $z \sim 8$ ($\Delta M^* = 1$ is nearly degenerate with $\Delta \log_{10} \phi^* = 1$; see Figure 8 of Oesch et al. 2012b), we can convert this to a change in ϕ^* over the redshift interval $z \sim 8$ to $z \sim 9.2$ (i.e., $\Delta \log_{10} \phi^* \sim 0.4$ dex) so that $\phi^*(z = 9.2) = 0.4\phi^*(z = 8)$. Oesch et al. (2012a) also estimated the rate of evolution from $z \sim 10$ to $z \sim 8$ based

on their $z \sim 10$ HUDF09+ERS+CANDELS search results, which is more rapid than implied by the Bouwens et al. (2011b) fitting formula. The best-fit evolution in ϕ^* that Oesch et al. (2012a) find is a $0.54^{+0.36}_{-0.19}$ dex change per unit redshift, so that $\phi^*(z = 9.2) = 0.23^{+0.15}_{-0.15}\phi^*(z = 8)$.

The evolution we measure from $z \sim 9.2$ to $z \sim 8$ (Section 4.4) is such that $\phi^*(z = 9.2) = 0.28^{+0.39}_{-0.20}\phi^*(z = 8)$ (fixed M^* and α). As compared with the two different evolutionary scenarios, we can see that the observed evolution may suggest marginally more rapid evolution than seen at lower-redshifts $z \sim 4-8$, but is nonetheless consistent (versus what one would expect using the Bouwens et al. 2011b LF fitting formula where $dM^*/dz \sim 0.33$, which is consistent with the new Bouwens et al. 2014b results if one excludes constraints at the bright end from wide-area searches).

The new ultra-deep WFC3/IR data over the HUDF/XDF field from the HUDF12 program also tentatively support a more rapid evolution. Using a sample of four $z \sim 9$ sources, Ellis et al. (2013) find $\phi^*(z = 9) = 0.25^{+0.15}_{-0.09}\phi^*(z = 8)$, while Oesch et al. (2013) find $\phi^*(z = 9) = 0.26^{+0.15}_{-0.12}\phi^*(z = 8)$ using similar samples. Searches to $z \sim 10$ (Oesch et al. 2013, 2014; Bouwens et al. 2014b) are again consistent with a slightly more rapid evolution. However, the results are not at all definitive, and indeed lensed candidate galaxies at redshifts as high as $z \sim 10.8$ identified by C13 would appear more consistent with a slower evolution. In any case, it seems clear that more observations, such as available with the Frontier Fields program or further study of the CANDELS fields (GO 13792: PI Bouwens), will be required to resolve this situation.

5. SUMMARY

In this paper, we explore the use of a two-color Lyman-break selection to search for $z \sim 9-10$ galaxies in the first 19 clusters observed with the CLASH program. Building on the important exploratory studies of Z12 and C13, we extend the CLASH $z \sim 9-10$ selections even deeper to the approximate magnitude limit of the CLASH program (~ 27 mag). Such a search is possible by making full use of the noteworthy *Spitzer*/IRAC observations over the CLASH clusters (Egami et al. 2008; Bouwens et al. 2011c), which allows us to determine which $z \sim 9-10$ galaxy candidates have a blue spectral slope redward of the break (and therefore strongly favor a $z \sim 9-10$ solution) and which candidates do not.

In total we find three plausible $z \sim 9-10$ galaxy candidates from the CLASH program that satisfy a two-color Lyman-break-like selection criteria (i.e., $(J_{110} + J_{125})/2 - H_{160} > 0.7$ and $JH_{140} - H_{160} < 0.5$) and have a combined $JH_{140} + H_{160}$ S/N of ≥ 6.0 . The $H_{160,AB}$ magnitudes for sources in our selection range from ~ 25.7 AB mag to 26.9 AB mag. The candidates are found behind the galaxy clusters MACSJ1149.6+2223, MACSJ1115.9+0129, and MACSJ1720.3+3536. The highest S/N source in our selection is the $z \sim 9.6$ Z12 candidate (here $z_{ph} \sim 9.7$). All three of our candidates have reasonably blue $H_{160,AB} - \text{IRAC}$ colors, which strongly favors the $z \sim 9-10$ solution for all three sources in our selection.

As in other $z \sim 9$ and $z \sim 10$ selections we have performed (Bouwens et al. 2011a; Z12; C13; Oesch et al. 2013, 2014; Bouwens et al. 2014b), we carefully considered the possibility of contamination. We find that the only significant source of contamination is from the “photometric scatter” of lower-redshift galaxies into our selection and that this likely contributes only ~ 0.7 source to our $z \sim 9-10$ sample (Section 3.5),

with MACS1720-JD1 or MACS1115-JD1 being the most probable contaminant. However, we emphasize that we cannot completely exclude the possibility that the contamination rate may be somewhat higher.

To determine the implications of the present search results for the UV LF, UV luminosity density, and SFR density at $z \sim 9$ –10, we introduce a novel differential approach for deriving these quantities. Our procedure is to simply compare the number of candidate $z \sim 9$ galaxies found in the CLASH fields with the number of $z \sim 8$ galaxies found in the CLASH fields and then correct this ratio for the relative selection volume at $z \sim 8$ and $z \sim 9$. This procedure takes advantage of the fact that the ratio of selection volumes at $z \sim 8$ and $z \sim 9$ for a given cluster is not greatly dependent on the details of the gravitational lensing model one is using (e.g., see Figure 8). This procedure therefore provides us with a very robust technique for measuring the evolution of the UV LF to $z > 9$ using searches over lensing cluster fields. The $z \sim 8$ and $z \sim 9$ selection volumes we derive are from detailed simulations where artificial sources are added to the real imaging data and then reselected using the same criteria as applied to the real data (Section 4.3).

Comparing our sample of three candidate $z \sim 9$ galaxies with a sample of 19 $z \sim 8$ galaxies found to similar 6σ detection significance over the same CLASH cluster fields (and correcting for the expected 23% contamination in our $z \sim 9$ selection), we derive the approximate evolution in the UV LF to $z \sim 9$. One strength of the present evolutionary estimate is that we are particularly insensitive to large-scale structure uncertainties, due to our many independent lines of sight on the high-redshift universe (Section 4.4).

We find that ϕ^* for the $z \sim 9$ LF is just $0.28^{+0.39}_{-0.20} \times$ the equivalent ϕ^* at $z \sim 8$ (Section 4.4: keeping M^* and α fixed). We would expect the normalization ϕ^* of the LF at $z \sim 9$ to be just $0.4 \times$ that at $z \sim 8$ if the evolution in the UV LF proceeded at the same rate as $z \sim 4$ –8. While the present result is consistent with there being no significant change in the rate of evolution of the UV LF from $z \sim 9$ to $z \sim 4$, our result does favor slightly more rapid evolution of the LF at $z > 8$, as suggested by Oesch et al. (2012a) based on their early search results for $z \sim 10$ galaxies. Using the best-fit evolutionary trend from Oesch et al. (2012a; see also Oesch et al. 2014), we would have predicted the normalization of our $z \sim 9$ LF to be $0.23^{+0.15}_{-0.15} \times$ that at $z \sim 8$. Several theoretical models (Trenti et al. 2010; Lacey et al. 2011) support the idea that the UV LF may evolve faster at $z > 8$ as a function of redshift than at $z \sim 4$ –8 (see Figure 8 of Oesch et al. 2012a and Figure 10 of Oesch et al. 2014). Despite the excellent agreement between the present evolutionary result and new findings from Oesch et al. (2012a), Ellis et al. (2013), and Oesch et al. (2014), the uncertainties on the evolution of the LF at $z > 8$ are still somewhat large.

In the future, we expect further advances in our constraints on the UV LF at $z \geq 9$ from the Frontier Fields program (Lotz et al. 2014),^{31,32} pointed follow-up observations of $z \sim 9$ –10 candidates over CANDELS (GO 13792: PI Bouwens), and the new wide-area BoRG program (GO 13767: PI Trenti). Substantially deeper *Spitzer* observations over the CLASH clusters, as part of the Surf’s Up program (Bradač et al. 2012) and other programs, should allow us both to obtain better constraints on the nature of current $z \sim 9$ candidates and to provide initial estimates of the stellar mass density at $z \sim 9$.

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APPENDIX A

DESCRIPTION OF THE REDSHIFT PRIORS

In Section 3.4, we present redshift likelihood distributions for the three candidate $z \sim 9$ galaxies in our selection. This allows us to estimate the relative probability that sources in this sample correspond to higher- or lower-redshift galaxies (Figure 5). However, in doing so we must utilize a prior. We consider three different redshift priors: (1) a flat redshift-independent prior, (2) a prior calibrated to published LF or LF trends, and (3) a prior tuned to reproduce the results from our photometric scattering experiments (Section 3.5). This section discusses the latter two priors in detail.

LF-calibrated Prior. The second prior we consider is calibrated according to published LFs or LF trends. For this prior, we give special attention to two galaxy populations: star forming galaxies at $z \sim 9$ and faint red galaxies at $z \sim 1.3$ –2. These are the only two galaxy populations that can at least provide approximate fits to the sources in our selection and therefore have nominal χ^2 ’s that are not especially large. For the faint red $z \sim 1.3$ –2 galaxy case, we calibrate our priors based on the LF results of Giallongo et al. (2005) for red galaxies using deep near-IR observations available over the HDF-North and HDF-South fields (Williams et al. 1996; Casertano et al. 2000) and the K20 spectroscopic sample (Cimatti et al. 2002). At $z \sim 2$, their $< m^*/m$ (bimodal) LF results correspond to $M_{B,0}^* = -21.90$ mag, $\phi^* = 2 \times 10^{-4}$ Mpc⁻³ mag⁻¹, and $\alpha = -0.53$. The basic validity of these LF results has been verified with very improved statistics based on new results for red galaxies over the UKIDSS Ultra Deep Survey field (Lawrence et al. 2007), where fits yield $M_V = -21.9$, $\phi^* = 2 \times 10^{-4}$ Mpc⁻³, and $\alpha = 0.07$ (R. Quadri et al. 2012, private communication). Meanwhile, at $z \sim 1.3$, the Giallongo et al. (2005) $< m^*/m$ (bimodal) LF results correspond to $M_{B,0}^* = -21.49$ mag, $\phi^* = 5 \times 10^{-4}$ Mpc⁻³ mag⁻¹, and $\alpha = -0.53$. Finally, for the $z \sim 9$ star forming galaxy case, our priors use the Bouwens et al. (2011b) LF fitting formula as a guide (which is a parameterization of the evolution of the UV LF from $z \sim 8$ to $z \sim 4$; see Section 7.5 of that paper).

Assuming a deep blank search at ~ 28.5 mag (the approximate intrinsic magnitude of our candidates after correction for magnification) with a $\Delta z \sim 1$, $\Delta \text{mag} \sim 1$ selection window, we find that these LFs predict ~ 1.2 $z \sim 9.2$ galaxies per arcmin², but 0.14 faint red galaxies per arcmin² over the redshift range $z \sim 1.3$ –2.5. Surprisingly, these results suggest that we would be much more likely (i.e., by $\sim 9 \times$) to find a blue galaxy at $z \sim 9$ with our selection than a faint red galaxy at $z \sim 1.3$ –2. Even correcting these predictions based on the present search results for $z \sim 9$ galaxies (where we find just $\sim 70^{+98}_{-50}$ % as many

³¹ <http://www.stsci.edu/hst/campaigns/frontier-fields/>

³² http://www.stsci.edu/hst/campaigns/frontier-fields/HDFI_SWGReport2012.pdf

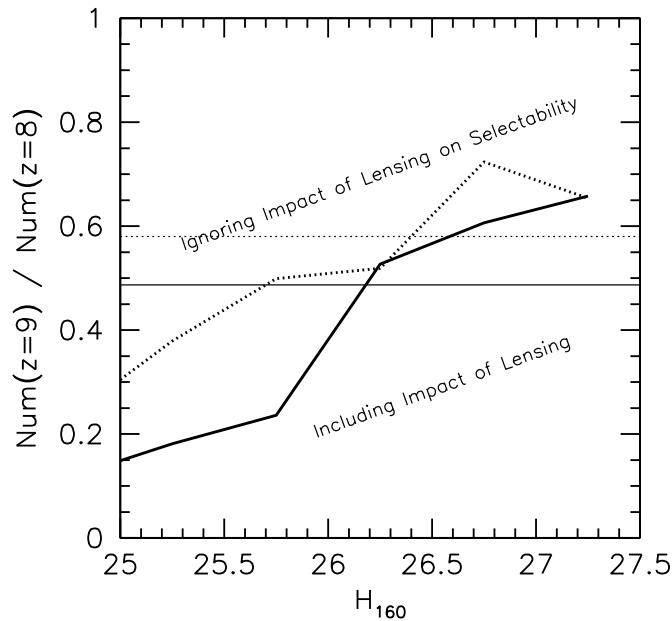


Figure 11. Relative number of galaxies we expect to be present in our $z \sim 9.2$ and $z \sim 8$ samples, assuming no evolution including the impact of lensing on surface brightness profiles and the selectability of individual galaxies in simulations behind galaxy clusters (thick solid line), and ignoring the impact of lensing in the simulations (thin dashed line), as a function of H_{160} -band magnitude (see Appendix B). The thin lines show the average relative numbers averaging over the full magnitude range. In the simulations where we ignore lensing, a factor-of-nine magnification boost to the LF is nonetheless considered (see Section 4.2). The slight magnitude dependence on the relative numbers is due to slight differences in the distance modulus to the two samples, and due to the slightly smaller sizes of $z \sim 9$ galaxies in our simulations. In both the case where the effects of lensing are fully included in estimating the selectability of individual sources, and where this is ignored, we expect approximately the same ratio of galaxies in the two samples.

galaxies as expected from the Bouwens et al. 2011b fitting formula), $z \sim 9$ galaxies would still be ~ 7 times more abundant on the sky at ~ 28.5 mag than red (old and/or dusty) galaxies at $z \sim 1.3$ – 2.5 . For the purposes of our “LF calibrated” prior, we assume that $z \sim 9$ galaxies have a five times higher surface density on the sky than $z \sim 1.3$ – 2.5 red (old and/or dusty) galaxies to be conservative.

Contamination-tuned Prior. Of course, it is not simply the faint red (old and/or dusty) galaxies at $z \sim 1.3$ – 2 that can contaminate $z > 8$ selections. Other galaxies can scatter into $z \sim 9$ selections through noise. This makes the low-redshift solution more likely than what we would calculate based on observationally based LFs. Considered alone, each photometrically scattered source would be unlikely to look like a probable $z \sim 9$ candidate, but one must account for the fact that there are some $\sim 4 \times 10^4$ sources in our fields that noise could conspire to make look like such a $z \sim 9$ – 10 candidate. We account for this possibility with our third “contamination-tuned” prior. With this prior, we adjust the relative likelihood of the high- and low-redshift peaks for our entire three source $z \sim 9$ – 10 galaxy sample so that it matches the 23% contamination rate estimated in the photometric scattering experiments described in Section 3.5.³³ However, it is worth noting that results based on the third prior likely overweight the

probability that sources are low-redshift contaminants. This is because our photometric scatter simulations do not account for the fact that red (old and/or dusty) galaxies are even rarer at ~ 27 – 28 mag than the ~ 24 – 25.5 mag sources we use as inputs to our photometric scattering simulations (e.g., Figure 11 from Oesch et al. 2012a).

APPENDIX B

ESTIMATING THE RATIO OF EFFECTIVE VOLUMES FOR DIFFERENT SELECTIONS BEHIND LENSING CLUSTERS USING SIMILAR VOLUMES IN BLANK FIELD SEARCHES

In deriving the differential evolution in the UV LF from $z \sim 9.2$ to $z \sim 8$, one particularly significant assumption we made in Section 4.2 was that the relative numbers of galaxies expected to be present in our $z \sim 9.2$ and $z \sim 8$ samples would remain roughly the same whether or not we included lensing in the simulations.

In this section, we test the accuracy of this assumption by making use of four different gravitational lensing models and galaxy clusters from the CLASH program (i.e., MACS1149, MACS1115, MACS1720, and MACS0416) when creating mock galaxy fields. We then create mock galaxy fields over the same cluster ignoring the lensing deflection fields. By comparing the relative number of $z \sim 8$ and $z \sim 9.2$ galaxies that we select from the two simulations, we test the assumption made in Section 4.2 of this paper.

In simulating the mock galaxy fields subject to lensing, we use exactly the same set of assumptions that we used for the simulations described in Section 4.2. Starting with the same (and non-evolving) LF of galaxies at $z \sim 8$ and $z \sim 9$, we construct a mock catalog of sources over multiple cluster fields. We then create artificial images of each source by artificially redshifting similar luminosity $z \sim 4$ galaxies from the HUDF to higher redshift. While redshifting the sources, we scale their sizes as $(1+z)^{-1}$ at fixed luminosity and take the UV -continuum slope β of galaxies in our simulations to have a mean value and 1σ scatter of -2.3 and 0.45 , respectively, to match the observed trends extrapolated to $z \sim 8$ – 9 . We remap the simulated images of galaxies in the source plane onto the image plane using the lensing models we have available for these clusters (Zitrin et al. 2012). We then add these simulated fields to the actual CLASH observations and attempt to recover the mock sources using the same $z \sim 8$ and $z \sim 9$ selection criteria as given in this paper.

In Figure 11, we present the relative surface density of the $z \sim 8$ and $z \sim 9.2$ galaxies we recover from non-evolving LF from our CLASH observations and compare with the surface density of galaxies we find without including lensing (but assuming a uniform factor-of-nine magnification in the luminosity of all sources; see Section 4.2). Overall, we find that the ratio of galaxies we select in the two cases is similar, but not exactly the same.

In fact, the relative numbers of $z \sim 9.2$ galaxies to $z \sim 8$ galaxies is slightly ($\sim 16\%$) lower in simulations where we include the effect of lensing. We find similar results for all four cluster lensing models that we have run simulations, but note that the relative numbers can show a dependence on the LFs or average magnification factors we assume ($\sim 10\%$ – 20%). However, for all reasonable LF or magnification factors we consider, the relative numbers do not differ by $\gtrsim 20\%$ from the ratios we give here.

³³ Admittedly, a more accurate approach would be to determine the actual redshift distribution of the intermediate-magnitude sources scattering into our selection and present it in Figure 5.

Table 5
Candidate $z \sim 8$ Galaxies Selected from CLASH to Compare
with a Similar Selection at $z \sim 9$

Right Ascension	Declination	H_{160}
17:22:25.76	32:08:58.2	$26.7 \pm 0.2^{a,d}$
13:47:30.47	-11:45:27.6	26.5 ± 0.2^b
13:47:33.89	-11:45:09.3	26.1 ± 0.2^a
13:47:31.82	-11:44:13.20	25.4 ± 0.1
21:29:24.92	-7:42:04.2	26.8 ± 0.2
03:29:40.34	-2:12:44.9	26.3 ± 0.2
06:47:59.14	70:14:15.2	26.4 ± 0.1^b
06:47:40.49	70:14:15.3	26.5 ± 0.2
06:47:40.90	70:14:41.5	25.8 ± 0.1^a
06:47:40.67	70:14:56.5	26.0 ± 0.2^b
06:47:44.76	70:15:37.9	27.1 ± 0.2^b
11:15:52.85	1:28:56.7	25.5 ± 0.1
11:15:51.09	1:30:34.9	25.5 ± 0.1^a
15:32:58.81	30:21:02.2	25.3 ± 0.1^b
19:31:45.22	-26:34:24.6	25.9 ± 0.2^d
21:29:43.60	0:05:31.9	26.4 ± 0.2^c
01:31:48.72	-13:36:49.1	26.3 ± 0.2^c
01:31:54.51	-13:36:00.5	25.4 ± 0.2^c

Notes.

^a Present in the $z \sim 8$ sample of Bradley et al. (2014).

^b Present in the $z \sim 7$ sample of Bradley et al. (2014).

^c Search field not considered in Bradley et al. (2014).

^d This $z \sim 8$ candidate is sufficiently compact (i.e., the SExtractor stellarity parameter is >0.9 in at least two of the five near-IR bands probed where 1 and 0 corresponds to a point source and extended source, respectively) that it may correspond to a low-mass star.

APPENDIX C

$Z \sim 8$ COMPARISON SAMPLE

In deriving the UV LF at $z \sim 9$ from our CLASH search results (Section 4), we make use of a baseline sample of $z \sim 8$ candidate galaxies in CLASH that we contrast with a similar $z \sim 9$ sample to establish the evolution from $z \sim 8$ to $z \sim 9$.

We tabulate this sample of $z \sim 8$ candidate galaxies in Table 5. Four of the 18 candidates from our selection are reported as $z \sim 8$ candidates by Bradley et al. (2014), five candidates are reported as $z \sim 7$ candidates by Bradley et al. (2014), four candidates do not appear in the Bradley et al. (2014) $z \sim 6-8$ compilation, and three candidates are found over CLASH clusters not considered by Bradley et al. (2014).

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