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# ECG signal decomposition using Fourier analysis



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# Abstract

This paper explores the Fourier decomposition method to approximate the decomposition of electrocardiogram (ECG) signals into their component waveforms, such as the QRS-complex and T-wave. We compute expansion coefficients using the  $\ell_1$ Fourier transform and the traditional  $\ell_2$  Fourier transform. Numerical examples are presented, and the analysis focuses on ECG signals as a real-world application, comparing the performance of the  $\ell_1$  and  $\ell_2$  Fourier transforms. Our results demonstrate that the  $\ell_1$ Fourier transform significantly enhances the separation of ECG signal components, such as the QRS-complex and T-wave. This improvement is attributed to a notable reduction in the Gibbs phenomenon introduced by the Fourier-series expansion when using the  $\ell_1$ Fourier transform, as opposed to the traditional  $\ell_2$  Fourier transform.

**Keywords:**  $\ell_1$  Fourier analysis,  $\ell_1$  Fourier analysis, Signal decomposition, ECG components

# 1 Introduction

The Fourier-series expansion represents a signal as a linear combination of sinusoidal basis functions [1-3]. In order to know what frequencies are present in the signal, we need to compute the expansion coefficients. The process of computing the expansion coefficients is known as Fourier analysis (transform). The celebrated Fourier transform computes the expansion coefficients by correlating the signal with the corresponding Fourier basis functions. As shown in [4], it is defined based on the  $\ell_2$ -norm minimization of the model error, *i.e.*, the error between the time-series data and the Fourier-series expansion, and the so-called  $\ell_2$  Fourier transform. The traditional  $\ell_2$  Fourier transform has been widely used in many signal and image processing applications [5, 6]. However, its performance is significantly decreased for the application of decomposing a signal into slow and fast components, especially when the fast components contain outliers. In this context, the idea is to filter the signal with a low-pass filter to obtain an estimate of the slow components. The problem is that in the  $\ell_2$  Fourier transform, each component of the error (i.e., fast components) is squared. When a set of error values are squared and then summed together, the sum is most sensitive to the largest error values. It means that the fast components or outliers have more weighting so it can skew results. As a result, some parts of the outliers are mixed with the slow components, which is



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due to the fact that the  $\ell_2$ -norm minimization corresponds to Gaussian distribution while the outliers have non-Gaussian distribution. Therefore, the  $\ell_2$  Fourier transform performance decreases in applications where the Gaussian distribution assumptions of the model errors do not hold in practice. In particular, when dealing with perceptually important signals, such as electrocardiogram (ECG), audio, image, speech, medical, and ocean engineering, we observe non-Gaussian (impulsive and Laplace) noises [7].

The  $\ell_2$  Fourier transform is not the only way to calculate the Fourier coefficients. An alternative approach to compute the Fourier transform has been recently reported in [4] which is based on the calculation of the Fourier coefficients using other norm spaces (*i.e.*,  $\ell_p$ -norm minimization of the model error,  $p = 1, 2, 3, \dots, \infty$ ). It is worth noting that the celebrated  $\ell_2$ -norm minimization which is used in the past for computing the Fourier transform is simple, parameter free, and inexpensive to compute. It satisfies many important properties, such as convexity, symmetry, and differentiability [8]. This might be the reason why other  $\ell_p$ -norm optimization methods were overlooked for such a long time. Among other  $\ell_p$  Fourier transforms, the  $\ell_1$  Fourier transform is an efficient tool for computing the Fourier coefficients which improves the Fourier-series expansion of time-series data such as reducing the effect of Gibbs phenomena in the truncated Fourier expansion of a signal with a jump discontinuity [4]. This method is based on the  $\ell_1$ -norm minimization of the error between the signal and its Fourier-series expansion and so called  $\ell_1$  Fourier transform. Compared to the traditional  $\ell_2$  Fourier transform, the benefits of the  $\ell_1$  Fourier transform include that it reduces the Gibbs effect in truncated Fourier-series expansion and filters the impulsive noises (outliers) in signals and images [4]. In  $\ell_1$ -norm minimization, the absolute value of the error is considered which corresponds to impulsive/Laplace distribution. Therefore,  $\ell_1$  Fourier transform decreases the weights of the outliers in reconstructed signal. That is why the  $\ell_1$  Fourier transform improves the Fourier-series expansion in filtering the impulsive noise from the data. Although other norm spaces can also be considered for computing the Fourier transform, they mostly produce performances in between  $\ell_1$ - and  $\ell_\infty$ -norm minimization [4].

According to the above discussion, the question naturally arises as to whether the beneficial filtering properties of the  $\ell_1$  Fourier-series expansion would improve decomposition of ECG signals into their depolarization and repolarization component waveforms. Roughly speaking, it is well known that the Fourier-series expansion (or frequency domain filtering) is unable to decompose signals with overlapping spectra. Therefore, it fails to separate the signal components that overlap in frequency domain. The ECG is an example of such signals. ECG is a rich source of information for cardiac diagnoses which makes it an important tool for assessing the cardiac health status. Every ECG beat is composed of different waves, classically labeled as P, Q, R, S and T, which reflect, at the body surface, the electrophysiological activity of the heart. Each ECG component reflects depolarization or repolarization of the heart: P-wave corresponds to depolarization of the atria, while QRS-complex and T-wave correspond to ventricular depolarization and repolarization. Separating ventricular depolarization and repolarization using Fourier-series expansion is a challenging task. It is well known that these ECG components (QRS-complex and T-wave) overlap in frequency domain. It is usually stated in the literature [9] that the content of T-wave lays mostly within a range of [0, 10] Hz. The

content of QRS-complex lays within a range of [8, 50] Hz frequencies. Note that estimation of the frequency content of these ECG components typically employs the  $\ell_2$  Fourier transform. That is why the spectra of the ECG components (QRS-complex and T-wave) overlap with each other and as such cannot be recovered by Fourier decomposition or frequency domain filtering.

In this paper, we show that this frequency overlap is due to the traditional method (*i.e.*,  $\ell_2$  Fourier transform) typically used for computing Fourier-series coefficients. Considering the fact that the changes in T-wave are slower than the changes in QRS-complex, the overlap between the spectra of these two components decreases if we use the  $\ell_1$  Fourier transform to compute the Fourier coefficients. Especially, the QRS-complex is an impulsive component (but not a Gaussian component) when a truncated Fourier-series expansion is used to low-pass filtering the signal to obtain an estimate of the slow component (*i.e.*, T-wave). That is why the traditional  $\ell_2$  Fourier transform is not able to separate T-wave and QRS-complex. We show that the truncated Fourier-series expansion of an ECG signal, when the  $\ell_1$  Fourier transform is used to compute the  $\ell_1$  Fourier transform is used to complex. The provide the sequence of the slow component (*i.e.*, produces a more accurate estimate of the T-wave and rejects the QRS-complex.

## 2 Fourier-series expansion-based signal representation

Fourier-series expansion decomposes a signal into oscillatory components. In this method, a given signal x[n],  $n = \{0, 1, \dots, N-1\}$  is represented as a linear combination of exponential basis functions:

$$x_M[n] = \sum_{k=0}^{M-1} c_k \phi_k[n],$$
(1)

where  $M \le N$ ,  $\phi_k[n] = \exp(i\frac{2\pi k}{N}n)$  and  $i = \sqrt{-1}$ . In vector notation, (1) can be expressed as

$$\boldsymbol{x}_M = \boldsymbol{\Phi} \boldsymbol{c},\tag{2}$$

where  $\mathbf{\Phi}$  is an  $N \times M$  matrix,  $\mathbf{x}_M = (\mathbf{x}_M[0], \cdots, \mathbf{x}_M[N])^T$  is a length-*N* vector, and  $\mathbf{c} = (c_0, \cdots, c_{M-1})^T$  is a length-*M* vector. The Fourier transform is the process of computing the expansion coefficients  $\mathbf{c}$ . It can be calculated by different  $\ell_p$ -norm minimization of the model error, *i.e.*,  $\mathbf{e} = \mathbf{x} - \mathbf{x}_M[4]$ :

$$\ell_p$$
 Fourier transform:  $\operatorname{argmin}_{r} \| \boldsymbol{x} - \boldsymbol{x}_M \|_p$  (3)

Among other norm spaces, we consider the Fourier transforms that compute the expansion coefficients by solving one of the following optimization problems:

$$\ell_2$$
 Fourier transform:  $\underset{c}{\operatorname{argmin}} \| \boldsymbol{x} - \boldsymbol{\Phi} \boldsymbol{c} \|_2$  (4a)

$$\ell_1$$
 Fourier transform:  $\underset{c}{\operatorname{argmin}} \| \boldsymbol{x} - \boldsymbol{\Phi} \boldsymbol{c} \|_1$  (4b)

In  $\ell_2$  Fourier transform, the expansion coefficients are analytically computed as

$$\boldsymbol{c} = (\boldsymbol{\Phi}^H \boldsymbol{\Phi})^{-1} \boldsymbol{\Phi}^H \boldsymbol{x},\tag{5}$$

where  $\Phi^H$  is the complex conjugate transpose (Hermitian transpose) of  $\Phi$ . Since the Fourier basis functions are orthogonal, the matrix  $\Phi^H \Phi$  becomes an identity matrix and (5) is simplified to  $c = \Phi^H x$ . Unlike  $\ell_2$  Fourier transform, the solution to  $\ell_1$  Fourier transform cannot be written in explicit form. The solution can be found only by running an iterative numerical algorithm. In [4], a solution was obtained using majorization minimization (MM) approach which is summarized as follows. Using a majorizer for the absolute value and considering an initial value for  $c^{(0)}$ , (4b) is converted to

$$\boldsymbol{c}^{(r+1)} = \operatorname*{argmin}_{\boldsymbol{c}^{(r)}} \left( \boldsymbol{x} - \boldsymbol{x}_M^{(r)} \right)^T \boldsymbol{\Lambda}_r^{-1} \left( \boldsymbol{x} - \boldsymbol{x}_M^{(r)} \right) + \frac{\gamma_r}{2},\tag{6}$$

where  $\gamma_r = \| \pmb{x} - \pmb{x}_M^{(r)} \|_1, \, \pmb{x}_M^{(r)} = \pmb{\Phi} \pmb{c}^{(r)}, \, \pmb{e}_r = \pmb{x} - \pmb{x}_M^{(r)}$  and

$$\mathbf{\Lambda}_{r} = \begin{pmatrix} |\boldsymbol{e}_{r}[0]| & 0 & \dots & 0 \\ 0 & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & |\boldsymbol{e}_{r}[N]| \end{pmatrix},$$
(7)

Taking the derivative of (6) with respect to the coefficients  $c^{(r)}$ , we obtain:

$$\boldsymbol{c}^{(r+1)} = \left(\boldsymbol{\Phi}^{H}\boldsymbol{\Lambda}_{r}^{-1}\boldsymbol{\Phi}\right)^{-1}\boldsymbol{\Phi}^{H}\boldsymbol{\Lambda}_{r}^{-1}\boldsymbol{x}$$
(8)

When M < N, the Fourier-series (1) is known as a truncated Fourier-series expansion. The truncated Fourier-series expansion acts as a low-pass filter as it neglects high-frequency components by setting their coefficients to zero. In the following section, the Fourier-series decomposition methods are used for approximate decomposition of a signal to slow and fast components using numerical examples. We employ  $\ell_1$  and  $\ell_2$  Fourier transform to compute the expansion coefficients. The corresponding Fourier-series expansion is called  $\ell_1$  and  $\ell_2$  Fourier-series expansion. We compare these two methods for signal components separation. As a real application, we compare them for ECG signal analysis and waveform decomposition.

### **3** Numerical examples

In the following section, we present some important functions that are widely used in signal processing literature to illustrate the properties of the Fourier transform.

## 3.1 Example 1

Let us consider a sinc function  $x_1[n] = sinc[n]$ ,  $n = 0, \dots, 1000$  plotted via blue color in Fig. 1a. The reconstructed signals provided by a truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1[n]$  using the first 9 harmonics (*i.e.*, M = 9) are plotted via dashed red and dashed green, respectively. The  $\ell_2$  and  $\ell_1$  Fourier transform of the signal is also shown at the bottom of each figure using red and green color, respectively.  $x_1[n]$  is represented perfectly using a truncated  $\ell_2$  or  $\ell_1$  Fourier-series. It means that  $x_1[n]$  is a low-frequency component signal. Let us consider the same sinc function after it has undergone the



#### (a) Truncated $\ell_2$ and $\ell_1$ Fourier-series of $x_1$

(b) Truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_2$ 



**Fig. 1** Truncated  $\ell_2$  and  $\ell_1$ Fourier-series expansion of sinc signals  $x_1[n] = sinc[n]$  and  $x_2[n] = x_1[15(n + 2.5)]$  for M = 9. The  $\ell_2$  and  $\ell_1$ Fourier transform of the signal is also shown at the bottom of each figure

transformation:  $x_2[n] = x_1[15(n + 2.5)]$  shown via blue color in Fig. 1b. The reconstructed signals provided by a truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_2[n]$  using the first 80 harmonics (*i.e.*, M = 80) are plotted via dashed red and dashed green, respectively.  $x_2[n]$  is also represented perfectly using a truncated  $\ell_2$  and  $\ell_1$  Fourier-series which means that it is also a low-frequency component signal. Note that  $x_1[n]$  is slower than  $x_2[n]$ .

Now, consider another signal which is a summation of  $x_1[n]$  and  $x_2[n]$ . The sum is plotted via gray color in Fig. 2a. Suppose that our objective is to estimate  $x_1[n]$  from the sum signal. In this case, the objective is to reconstruct  $x_1[n]$  and reject  $x_2[n]$ . However, the frequency spectrum of  $x_2[n]$  overlaps with  $x_1[n]$  computed by either  $\ell_2$ or  $\ell_1$  Fourier transform (see their Fourier transforms in Fig. 1a, b). In Fig. 2a, we plot the reconstructed signal provided by  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1[n] + x_2[n]$  using the first 9 harmonics (*i.e.*, M = 9) via dashed red and dashed green, respectively. The reconstructed signal using the truncated  $\ell_1$  Fourier-series is more close to  $x_1[n]$ . The  $\ell_1$  Fourier transform of  $x_1[n] + x_2[n]$  is shown via red color at left bottom of Fig. 2a. The  $\ell_1$  Fourier transform of the sum is also shown via green color which is comparable to the  $\ell_1$  Fourier transform of  $x_1[n] + x_2[n]$  and  $x_1[n]$  are different. After subtracting the truncated Fourier-series of  $x_1$  (which was computed in the previous step) from the sum, the truncated  $\ell_2$  and  $\ell_1$  were again used to estimate  $x_2$  using the first 80 harmonics, M = 80. The reconstructed signals provided by the truncated  $\ell_2$  and  $\ell_1$ 



(a) Truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1 + x_2$ 

(b) Truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1 + x_2 - \hat{x}_1$ 



**Fig. 2**  $\ell_2$  and  $\ell_1$ Fourier-series expansion-based separation of the mixture of two sinc signals:  $x_1[n] = sinc[n]$ and  $x_2[n] = x_1[15(n + 2.5)]$ . **a** Truncated  $\ell_2$  and  $\ell_1$ Fourier-series expansion of **a**  $x_1[n] + x_2[n]$  for M = 8 **b**  $x_1[n] + x_2[n] - \hat{x}_1[n]$  for M = 80 where  $\hat{x}_1[n]$  is an estimate of  $x_1[n]$  using the truncated Fourier-series either by  $\ell_2$  or  $\ell_1$ Fourier transform. The  $\ell_2$  and  $\ell_1$ Fourier transform of each signal is also shown via red color at the bottom of each figure. The  $\ell_2$  and  $\ell_1$ Fourier transform of  $x_1[n]$  and  $x_2[n]$  is also plotted via blue color for comparison

Fourier-series of the residual using the first 80 harmonics are plotted via dashed red and dashed green in Fig. 2b. The reconstructed signal using the truncated  $\ell_1$  Fourierseries is more close to the reconstructed signal using the truncated  $\ell_2$  Fourier-series. Comparing to the  $\ell_2$  Fourier transform, the  $\ell_1$  Fourier transform of the residual is closer to the Fourier transform of  $x_2[n]$  as shown at the bottom of Fig. 2b. It means that  $\ell_1$  Fourier-series expansion separates these two sinc signals much better than the  $\ell_2$  Fourier-series expansion.

## 3.2 Example 2

Sum of Gaussian kernels are common in signal modeling, especially they become popular in ECG signal modeling, due to their morphological similarity with the ECG components waveform [10-12]. A sum of Gaussian model is defined as follows:

$$x[n] = \sum_{i=1}^{N} \alpha_i \exp\left[-\frac{(nT_s - \mu_i)^2}{2b_i^2}\right],$$
(9)

where  $\alpha_i$ ,  $b_i$ , and  $\mu_i$  are the amplitude, angular spread, and location of the Gaussian functions. Let us consider a sum of two Gaussian functions:

$$x[n] = x_1[n] + x_2[n] = \sum_{i=1}^{2} \alpha_i \exp\left[-\frac{(nT_s - \mu_i)^2}{2b_i^2}\right],$$
(10)

where  $n = 0, \dots, 1000, \alpha_1 = 7, 5, b_1 = 0.4, \mu_1 = 750, \alpha_2 = 30, b_2 = 0.05$  and  $\mu_2 = 500$ .  $x_1[n]$  and its truncated  $\ell_2$  and  $\ell_1$  Fourier-series using the first 7 harmonics (*i.e.*, M = 7) are plotted via blue color, dashed red and dashed green, respectively in Fig. 3a.  $x_1[n]$ is reconstructed perfectly using a truncated  $\ell_2$  or  $\ell_1$  Fourier-series which means it is a low-frequency component signal. In Fig. 3b, we plotted  $x_2[n]$  and its truncated  $\ell_2$  and  $\ell_1$ Fourier-series using the first 51 harmonics (*i.e.*, M = 51) via blue color, dashed red and dashed green, respectively.  $x_2[n]$  is also reconstructed perfectly using a truncated  $\ell_2$  or  $\ell_1$  Fourier-series which means it is also a low-frequency component signal. The  $\ell_2$  and  $\ell_1$ Fourier transform is also shown at the bottom of each figure. Figure 4a shows the truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1[n] + x_2[n]$  using the first 7 harmonics. The Fourier transforms are also shown at the bottom of the figure. The results show that the  $\ell_1$  Fourier transform of  $x_1[n] + x_2[n]$  is close to the  $\ell_1$  Fourier transform of  $x_1[n]$ . However, the  $\ell_2$  Fourier transform of  $x_1[n]$  is far from the  $\ell_2$  Fourier transform of  $x_1[n] + x_2[n]$ . The estimated  $x_2$  using truncated  $\ell_2$  and  $\ell_1$  Fourier-series is also plotted in Fig. 4b. The results of Figs. 3 and 4 show that the Fourier-series expansion of a mixture of Gaussian signals, when the  $\ell_1$  Fourier transform is used to identify the expansion coefficients, produces a more accurate estimate of its components.





 $\ell_2$  inverse DFT  $\ell_1$  inverse DFT Amp<sup>2</sup> Amp<sup>2</sup> 0 100 200 300 400 500 600 700 800 900 1000 100 200 300 400 500 600 700 800 900 1000 Samples Samples 400 പ് ് 30 10 30 40 50 k

(b) Truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_2$  for M = 51

**Fig. 3**  $\ell_2$  and  $\ell_1$  Fourier-series expansion of two mixed Gaussian signals. **a**  $x_1[n]$  **b**  $x_2[n]$ 



(a) Truncated  $\ell_2$  and  $\ell_1$  Fourier-series of  $x_1 + x_2$  for M = 7

 $\ell_2 \text{ iDFT}(x_1 + x_2 - \hat{x}_1)$  $\ell_1 \text{ iDFT}(x_1 + x_2 - \hat{x}_1)$ 20 Amp 10 200 700 800 900 100 200 400 600 700 800 900 1000 Samples Samples 60  $\ell_2 \operatorname{DFT}(x_2)$  $\ell_1 \operatorname{DFT}(x_2)$  $\ell_2 DFT(x_1 + x_2 - \hat{x}_1)$  $\ell_1 \text{ DFT}(x_1 + x_2 - \hat{x}_1)$ 400 പ് 20 k k

**Fig. 4**  $\ell_2$  and  $\ell_1$ Fourier-series expansion-based separation of two mixed Gaussian signals. **a**  $x_1[n] + x_2[n]$ **b**  $x_1[n] + x_2[n] - \hat{x}_1[n]$  where  $\hat{x}_1[n]$  is the estimated  $x_1[n]$  using the truncated Fourier-series either by  $\ell_2$  or  $\ell_1$  Fourier transform. The Fourier transform obtained by  $\ell_2$  and  $\ell_1$ minimization is also shown at the bottom of each figure

## 4 ECG signal representation using Fourier-series expansion

As mentioned before, the  $\ell_2$ -norm corresponds to the Gaussian distribution and its performance decreases in non-Gaussian model errors. It is notable that the ECG is a sparse band-limited signal and the changes in T-waves are much slower than the changes in QRS-complex. In other words, for low-frequency (slow) components like T-wave, the distribution of the fast component (QRS-complex) is non-Gaussian. Therefore, the  $\ell_1$ Fourier transform (which is based on the  $\ell_1$ -norm minimization) would provide better ECG components separation (*i.e.*, ventricular depolarization and repolarization) than the  $\ell_2$  Fourier transform. In this section, we analyze the ECG signal using  $\ell_2$  and  $\ell_1$  Fourier transform. We show that the Fourier-series expansion when the coefficients are computed using  $\ell_1$  Fourier transform, produces a more accurate estimate of its slow components (T-waves) and rejects the fast component (QRS-complex). Especially, the Gibbs phenomenon introduced by the truncated Fourier-series expansion is significantly decreased when the expansion coefficients are computed using  $\ell_1$  Fourier transform compared to the traditional  $\ell_2$  Fourier transform.

As a first example of the application of Fourier-series expansion, we consider the truncated Fourier-series expansion of a specific case (record s0017lrem from the PhysioNet PTB Diagnostic ECG Database (ptbdb) [13]) shown in Fig. 5. Each record of the database sampled at 1 kHz ( $f_s = 1000$ ). In this example, we consider a portion of the ECG record with 4000 samples, *i.e.*, N = 4000. We represent the ECG signal using a truncated



Fig. 5 Real ECG record s0017Irem from the PhysioNet PTB Diagnostic ECG Database (ptbdb)

(a) Truncated  $\ell_2$  Fourier-series expansion



(b) Truncated  $\ell_2$  Fourier-series expansion (Blackman window)



(c) Truncated  $\ell_2$  Fourier-series expansion (Hanning window)





Fig. 6 T-wave and QRS-complex detection using truncated Fourier-series expansion (using rectangular, Blackman and Hanning window) and zero-phase Butterworth filter

Fourier-series expansion with M < N. In this case, the truncated Fourier-series expansion acts as a low-pass filter that passes the low-frequency components with frequencies less than  $M/N \times f_s$  Hz and rejects the high-frequency components with frequencies greater than it. We compare the Fourier-series expansion of the ECG signal when the coefficients are computed using  $\ell_2$  and  $\ell_1$  Fourier transform. The question is which Fourier transform type better separates ventricular repolarization (T-wave) and ventricular depolarization (QRS-complex). The result of ECG representation using the truncated  $\ell_2$  Fourier-series expansion, *i.e.*,  $x_{32}[n]$ , is shown in Fig. 6a via red curve. Selecting M = 32



Fig. 7 T-wave and QRS-complex detection using truncated  $\ell_1$ Fourier-series expansion



Fig. 8 Abnormal ECG record 8378m from MIT-BIH Atrial Fibrillation Database

means that the cutoff frequency is 8 Hz. The residual is QRS-complex component which can be computed using Fourier-series expansion when the frequencies are greater than 8 Hz. The extracted QRS-complex is shown via yellow curve. The truncated  $\ell_2$  Fourierseries expansion is unable to accurately extract the T-wave and reject the QRS-complex. The Gibbs phenomenon is evident in the truncated  $\ell_2$  Fourier-series expansion as shown by arrows in Fig. 6a. A method to reduce the Gibbs phenomena is multiplication by a tapered window (*e.g.*, Blackman, Hanning, Kaiser), versus a rectangular window. We also employed the Blackman and Hanning window for computing the truncated  $\ell_2$  Fourierseries expansion. The results of  $\ell_2$  Fourier analysis using Blackman and Hanning window are shown in Fig. 6b, c, respectively. Although the Gibbs phenomena is reduced using these tapered windows, it causes that some parts of T-waves are wrongly extracted and mixed with QRS-complexes.

It is mentionable that the zero-phase LTI filters are the standard and simple choice for ECG signal preprocessing, as they impose little assumptions on the signals and preserve their phase contents. Therefore, a naive third order zero-phase low-pass Butterworth filter with cutoff frequency 8 Hz was also used to extract the T-waves in the ECG signal. The T-waves and QRS-complexes provided by zero-phase low-pass Butterworth filter are, respectively, shown via red and yellow curve in Fig. 6d. As shown by arrows, we observe that the zero-phase Butterworth filter which is commonly used in the literature does not perfectly reject the QRS-complex. Finally, we model the same ECG record using a truncated  $\ell_1$  Fourier-series expansion. In order to compute the expansion coefficients, we set M = 32 (cutoff frequency is 8 Hz), and the number of iterations to 100, *i.e.*,  $x_{32}^{(100)}[n]$ . Figure 7 shows the results of T-waves detection using  $\ell_1$  Fourier-series expansion. The truncated  $\ell_1$  Fourier-series expansion extracts the T-waves and rejects the QRS-complexes much better than the truncated  $\ell_2$  Fourier-series expansion and zero-phase Butterworth filter.

In the second example, we consider the truncated Fourier-series expansion of an abnormal case (record 08378m from MIT-BIH Atrial Fibrillation Database [13]). This record is shown in Fig. 8. We consider a portion of the ECG record with 2500 samples,

*i.e.*, N = 2500. We model the ECG signal using a truncated Fourier-series with M = 60. Figure 9a–d shows the results of ECG representation using the truncated  $\ell_2$  Fourier-series expansion and  $\ell_1$  Fourier-series expansion, respectively.  $\ell_2$  Fourier-series expansion leads to the distortion of T-waves before premature ventricular contractions (PVC), while  $\ell_1$  Fourier-series expansion does not, as highlighted using arrows. The Gibbs phenomenon is also evident in the truncated  $\ell_2$  Fourier-series expansion. A PVC is a heartbeat which is autonomously triggered in the ventricles, and not in the sinus node. PVCs are common events which do not necessarily imply a negative heart condition [14].

## **5** Simulation results

For evaluating the performance of the  $\ell_1$  Fourier analysis, we applied it on simulated data, which permit to quantify the decomposition error directly. The single-channel synthetic data were obtained using the ECG dynamical model (EDM) proposed in [10]. The original ECG components are also needed for quantifying the decomposition error. To this purpose, we used the extended version of EDM proposed by



(b) P and QRS-complex detection using truncated  $\ell_2$  Fourier-series expansion



(c) T-wave detection using truncated  $\ell_1$  Fourier-series expansion



#### (d) P and QRS-complex detection using truncated $\ell_1$ Fourier-series expansion



**Fig. 9** ECG components separation using truncated Fourier-series expansion **a** T-wave detection using truncated  $\ell_2$  Fourier-series expansion **b** P and QRS-complex detection using truncated  $\ell_2$  Fourier-series expansion **c** T-wave detection using truncated  $\ell_1$ 2 Fourier-series expansion **d** P and QRS-complex detection using truncated  $\ell_1$ Fourier-series expansion

Sayadi *et al.* for generating ECG characteristic waveforms (CWs) [15]. Seven Gaussian kernels were employed to model ECG beats, corresponding to each of the ECG components (P-wave, QRS-complex, and T-wave), and for modeling asymmetries two Gaussian kernels were used for P- or T-waves (indicated by  $^+$  and  $^-$  superscripts), leading to:

$$\begin{cases} \dot{\theta} = \omega \\ \dot{P} = -\sum_{i \in \{P^-, P^+\}} \alpha_i \omega \frac{\theta - \theta_i}{b_i^2} \exp\left[-\frac{(\theta - \theta_i)^2}{2b_i^2}\right] \\ Q\dot{R}S = -\sum_{i \in \{Q, R, S\}} \alpha_i \omega \frac{\theta - \theta_i}{b_i^2} \exp\left[-\frac{(\theta - \theta_i)^2}{2b_i^2}\right], \\ \dot{T} = -\sum_{i \in \{T^-, T^+\}} \alpha_i \omega \frac{\theta - \theta_i}{b_i^2} \exp\left[-\frac{(\theta - \theta_i)^2}{2b_i^2}\right] \\ x = P + QRS + T \end{cases}$$
(11)

We set the sampling rate to 250 Hz and generated 1000 synthetic series with 10s long. We allowed the RR interval durations to have a random fluctuation of up to 5% in each beat to make the synthetic ECGs more realistic. We added the random noise to synthetic ECGs. The signal-to-noise ratio (SNR) was modulated from 10 to 50 dB.

The  $\ell_2$  and  $\ell_1$  Fourier transforms were used to analyze the ECG signals. We compared the efficiency of the  $\ell_2$  and  $\ell_1$  Fourier transform in extracting the ECG components. To quantify the performances of the methods, we used the measures of improvement given by the normalized mean absolute error (MAE):

$$MAE = \frac{\sum_{k} |x_k - \hat{x}_k|}{\sum_{k} |x_k|},$$

where  $x_k$  and  $\hat{x}_k$  denote the original and the estimated components (either using  $\ell_2$  or  $\ell_1$ Fourier transforms). The mean of MAE at different input SNRs is plotted in Fig. 10.  $\ell_1$ Fourier analysis outperforms zero-phase Butterworth filter and  $\ell_2$  Fourier analysis.



Fig. 10 Mean values of MAE as a function of SNR



Fig. 11 Abnormal ECG record 8455m from MIT-BIH Atrial Fibrillation Database

## 6 Conclusion

The  $\ell_1$  Fourier transform improves the Fourier-series expansion of time-series in reducing the effect of Gibbs phenomena and filtering the impulsive noise from the data. In this method, the Fourier coefficients are computed by minimizing the  $\ell_1$ -norm of the error between the time-series and its Fourier-series expansion. This paper presented the application of Fourier transform to decompose an ECG signal to its components waveform as a real application. We showed that the Fourier-series expansion of an ECG signal, when the





(c) T-wave detection using truncated  $\ell_1$  Fourier-series expansion  $\underbrace{\mathsf{ECG}_{\mathsf{P}} \mathsf{and}_\mathsf{T-wave}}_{\mathsf{and}}}}}}}}}}}}}}}}}}}}}}}}}}}}}}}} abus but be a but a$ 





**Fig. 12** ECG components separation using truncated Fourier-series expansion **a** T-wave detection using truncated  $\ell_2$  Fourier-series expansion **b** P and QRS-complex detection using truncated  $\ell_2$  Fourier-series expansion **c** T-wave detection using truncated  $\ell_1$ 2 Fourier-series expansion **d** P and QRS-complex detection using truncated  $\ell_1$ Fourier-series expansion

 $\ell_1$  Fourier transform is used to identify the expansion coefficients, produced a much accurate estimate of its components waveform (*e.g.*, QRS-complex and T-wave). Especially, the Gibbs phenomenon introduced by the Fourier-series expansion is significantly decreased when the expansion coefficients are computed using  $\ell_1$  Fourier transform compared to the traditional  $\ell_2$  Fourier transform. The efficiency of  $\ell_1$  Fourier analysis was compared with the traditional  $\ell_2$  Fourier transform and zero-phase Butterworth filter. The  $\ell_1$  Fourier transform significantly improves the Fourier-series expansion to decompose a signal to slow and fast components. Based on our assessments, the P-wave is indistinguishable with either the T-wave or QRS-complex in the frequency spectrum, as demonstrated in Figs. 9, 11 and 12, even when employing the  $\ell_1$  Fourier transform. Our future research will examine the utilization of the  $\ell_p$  Fourier transform for values of p between 0 and 1, aiming at further reducing the Gibbs effect and thus being able to better separate signals which do not strongly overlap in frequency.

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#### Author contributions

Arman Kheirati Roonizi wrote the main manuscript and all authors reviewed the manuscript.

#### Data availability

No datasets were generated or analyzed during the current study.

#### Competing interests

The authors declare no competing interests.

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