

NON-NEGATIVE RICCI CURVATURE AND MINIMAL GRAPHS WITH LINEAR GROWTH

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ABSTRACT. We study minimal graphs with linear growth on complete manifolds M^m with $\text{Ric} \geq 0$. Under the further assumption that the $(m - 2)$ -th Ricci curvature in radial direction is bounded below by $Cr(x)^{-2}$, we prove that any such graph, if non-constant, forces tangent cones at infinity of M to split off a line. Note that M is not required to have Euclidean volume growth. We also show that M may not split off any line. Our result parallels that obtained by Cheeger, Colding and Minicozzi for harmonic functions. The core of the paper is a new refinement of Korevaar's gradient estimate for minimal graphs, together with heat equation techniques.

CONTENTS

1. Introduction	1
Strategy of the proof	6
2. Preliminaries	7
3. Proof of Theorem 5, (i)	9
4. A local gradient estimate	10
5. Uniformly elliptic operators on manifolds with $\text{Ric} \geq 0$	17
6. Proof of Theorem 5, (ii)	23
7. Proof of Theorem 10	25
8. Proof of Corollary 9	26
9. Proof of Proposition 8	27
Appendix A.	29
References	30

1. INTRODUCTION

The theory of entire minimal graphs in Euclidean space \mathbb{R}^m , that is, of functions $u : \mathbb{R}^m \rightarrow \mathbb{R}$ solving the minimal (hyper)surface equation

$$(MSE) \quad \mathcal{M}[u] \doteq \text{div} \left(\frac{Du}{\sqrt{1 + |Du|^2}} \right) = 0$$

is built upon the following foundational results:

Date: December 21, 2021.

2020 Mathematics Subject Classification. Primary 53C21, 53C42; Secondary 53C24, 58J65, 31C12.

Key words and phrases. Bernstein theorem, splitting, minimal graph, Ricci curvature, tangent cone.

- ($\mathcal{B}1$) The *Bernstein theorem*: solutions of (MSE) are all affine if and only if $m \leq 7$.
- ($\mathcal{B}2$) For each $m \geq 2$, positive solutions of (MSE) are constant.
- ($\mathcal{B}3$) For each $m \geq 2$, solutions of (MSE) with at most linear growth on one side are affine (i.e., the Hessian $D^2u \equiv 0$).

Here, u is said to have at most linear growth on one side if, up to changing the sign of u ,

$$u(x) \geq -C(1 + r(x))$$

holds on \mathbb{R}^m for some constant $C > 0$, where r is the distance from a fixed origin. The validity of ($\mathcal{B}1$) is due, as well-known, to the combined effort of S. Bernstein [8] ($m = 2$, see also [53, 43]), W.H. Fleming ([38], still for $m = 2$), E. De Giorgi ([26, 27], $m = 3$), F. Almgren ([2], $m = 4$), J. Simons ([60], $m \leq 7$) and E. Bombieri, De Giorgi and E. Giusti ([13], counterexamples if $m \geq 8$). On the other hand, ($\mathcal{B}2$) and ($\mathcal{B}3$) were both proved in [12] by Bombieri, De Giorgi and M. Miranda for $m \geq 3$; in particular, ($\mathcal{B}3$) refines J. Moser's Theorem [54], which states that u is affine provided that $|Du| \in L^\infty(M)$. Further properties of entire minimal graphs in Euclidean space were obtained in [14] by Bombieri and Giusti: among them, we mention the fact that u is affine whenever $m - 1$ of its partial derivatives are bounded. The result was improved in recent years by A. Farina [34, 35], who showed that u is affine if $m - 7$ partial derivatives of u are bounded on one side. Other extension of Moser' result, proving that $D^2u \equiv 0$ by only assuming that $|Du| = o(r)$ as $r(x) \rightarrow \infty$, were obtained in [17, 32, 59]. We also mention the recent [36], where the rigidity of a minimal graph is obtained by assuming that an upper level set contains, or is contained in, a half-space.

In a Riemannian setting, it is natural to ask the following

Question 1. *For which classes of complete Riemannian manifolds M one could expect results like ($\mathcal{B}1$), ($\mathcal{B}2$), ($\mathcal{B}3$)?*

The problem motivated our previous works [11, 22] as well as the present paper. Recall that a solution of (MSE) on a Riemannian manifold (M^m, σ) gives rise to a graph

$$F : M \rightarrow \mathbb{R} \times M, \quad F(x) = (u(x), x)$$

which is minimal if the ambient space $\mathbb{R} \times M$ is endowed with the product metric $dt^2 + \sigma$. Hereafter, we say that the graph is entire if u is defined on the whole of M .

If M is close to hyperbolic space \mathbb{H}^m , namely, M is a Cartan-Hadamard manifolds with suitably pinched negative curvature, ($\mathcal{B}1$), ($\mathcal{B}2$), ($\mathcal{B}3$) drastically fail, since each continuous function $\phi \in \partial_\infty M$ can be attained as the limit value of an entire minimal graph, which is therefore bounded. An exhaustive literature on the problem can be found in the survey [33], see also the introduction of [9].

Denoting with $g = F^*(dt^2 + \sigma)$ the graph metric and with Δ_g its Laplace-Beltrami operator, equation (MSE) can be written as $\Delta_g u = 0$, suggesting that for ($\mathcal{B}2$), ($\mathcal{B}3$) an answer to the above question may be guided by the global behaviour of harmonic functions on Riemannian manifolds. This is clearly not the case for ($\mathcal{B}1$), since already in \mathbb{R}^m there is no rigidity result for harmonic functions without imposing any growth condition. For ($\mathcal{B}2$) and ($\mathcal{B}3$), it is therefore natural to look at the problem on manifolds satisfying either

$$(1) \quad \text{Sec} \geq 0, \quad \text{or} \quad \text{Ric} \geq 0$$

where Sec, Ric are the sectional and Ricci curvature of (M, σ) . Indeed, if $\text{Ric} \geq 0$, positive harmonic functions are constant, by S.Y. Cheng and S.T. Yau's gradient estimate [62, 20], while a harmonic function with linear growth force any tangent cone at infinity of M to split, by work of J. Cheeger, T. Colding and W. Minicozzi [19]. Indeed, M itself splits off a line whenever $\text{Ric} \geq 0$ is strengthened to $\text{Sec} \geq 0$ (cf. [5] for a complete proof).

In view of the convergence theory developed in the past 50 years for manifolds with $\text{Sec} \geq 0$ or $\text{Ric} \geq 0$, some of the tools used to prove the Bernstein theorem in \mathbb{R}^m are available on manifolds satisfying (1), making these assumptions a natural setting also for the study of $(\mathcal{B}1)$. However, much has to be done and $(\mathcal{B}1)$ seems very challenging to prove even on manifolds with $\text{Sec} \geq 0$. In fact, we are aware of no results in this direction.

The situation is different for $(\mathcal{B}2)$ and $(\mathcal{B}3)$, for which, as we shall detail below, the main difficulty is to prove the results by only requiring $\text{Ric} \geq 0$, arguably the sharp condition for their validity (in this case, however, $(\mathcal{B}3)$ has to be suitably weakened, see later).

Regarding $(\mathcal{B}2)$, after previous work in [57] by H. Rosenberg, F. Schulze and J. Spruck, a complete answer was obtained by the first, third and fourth authors together with M. Magliaro in [22], and independently by Q. Ding in [29] with different methods:

Theorem 2. [22, 29] *Complete manifolds M with $\text{Ric} \geq 0$ satisfy $(\mathcal{B}2)$, that is, entire positive minimal graphs over M are constant.*

In this paper, we address $(\mathcal{B}3)$. In view of the result in [19], it is reasonable to formulate the following

Conjecture 3. *Let M be a complete manifold with $\text{Ric} \geq 0$ and possessing a non-constant entire minimal graph with at most linear growth on one side. Then, every tangent cone at infinity of M splits off a line.*

The problem seems to be considerably harder compared to the case of harmonic functions. We are aware of only two results in the direction of Conjecture 3. The first is [31], where Ding, J. Jost and Y. Xin proved that \mathbb{R}^m is the only manifold satisfying the following assumptions:

$$(2) \quad \begin{cases} (2.\alpha) & \text{Ric} \geq 0, \quad \lim_{r \rightarrow \infty} \frac{|B_r|}{r^m} > 0 \\ (2.\beta) & \text{the curvature tensor decays quadratically} \end{cases}$$

and admitting an entire, non-constant minimal graph with at most linear growth on one side. Very recently, Ding [30] posted on arXiv a paper where he proved Conjecture 3 on manifolds satisfying the assumptions in $(2.\alpha)$. The bulk of his argument is to show the remarkable property that the isoperimetric inequality, satisfied by (M, σ) in view of $(2.\alpha)$, is inherited by the graph of u . This allowed Ding to adapt, in a nontrivial way, tools from [12, 14] and from Cheeger-Colding's theory to reach the goal. We stress that his method heavily depends on the Euclidean volume growth condition in $(2.\alpha)$.

In our work, we address Conjecture 3 without requiring the Euclidean volume growth assumption, but rather a mild further curvature condition. To formulate our main result, we recall the definition of the ℓ -th Ricci curvature:

Definition 4. Let (M, σ) be a manifold of dimension $m \geq 2$. For $\ell \in \{1, \dots, m-1\}$, the ℓ -th (normalized) Ricci curvature is the function

$$v \in T_x M \quad \longmapsto \quad \text{Ric}^{(\ell)}(v) \doteq \inf_{\substack{\mathcal{W} \leq v^\perp \\ \dim \mathcal{W} = \ell}} \left(\frac{1}{\ell} \sum_{j=1}^{\ell} \text{Sec}(v \wedge e_j) \right),$$

where $\{e_j\}$ is an orthonormal basis of \mathcal{W} .

The function $\text{Ric}^{(\ell)}$ interpolates between the sectional and Ricci curvatures, obtained respectively for $\ell = 1$ and, up to the normalization constant $(m-1)$, for $\ell = m-1$. In particular, with our chosen normalization the following implications are immediate:

$$\begin{aligned} \text{Sec} \geq c &\implies \text{Ric}^{(\ell-1)} \geq c \\ &\implies \text{Ric}^{(\ell)} \geq c \implies \text{Ric} \geq (m-1)c. \end{aligned}$$

Hereafter, given $H \in C([0, \infty))$ and denoting with r the distance from a fixed origin $o \in M$, we use the short-hand notation $\text{Ric}^{(\ell)}(\nabla r) \geq -H(r)$ on M to mean the inequality

$$\text{Ric}^{(\ell)}(\nabla r(x)) \geq -H(r(x)) \quad \forall x \in M \setminus (\{o\} \cup \text{cut}(o)),$$

where $\text{cut}(o)$ is the cut-locus of o .

Recall that a manifold M is said to be parabolic if every positive superharmonic function on M is constant. Also recall that a tangent cone at infinity for a complete (non-compact) manifold M is any metric space obtained as a blow-down of M . More precisely, a pointed metric space $(X_\infty, d_\infty, x_\infty)$, $x_\infty \in X_\infty$, is a tangent cone at infinity for (M, σ) if, for some base point $x \in M$ and some sequence $\{\lambda_n\}$ of positive real numbers such that $\lambda_n \rightarrow +\infty$, one has

$$(M, \lambda_n^{-1} \text{dist}_\sigma, x) \rightarrow (X_\infty, d_\infty, x_\infty)$$

in the pointed Gromov-Hausdorff (pGH) sense. If (M, σ) has non-negative Ricci curvature, then tangent cones at infinity exist based at any point $x \in M$, by Gromov's precompactness theorem [41].

We are ready to state

Theorem 5. Let (M, σ) be a complete Riemannian manifold of dimension $m \geq 2$ with

$$\text{Ric} \geq 0,$$

and let $u \in C^\infty(M)$ be a non-constant entire solution of (MSE).

- (i) If M is parabolic, then $M = N \times \mathbb{R}$ with the product metric $\sigma_N + ds^2$, for some complete manifold N with $\text{Ric}_N \geq 0$, and in the variables $(y, s) \in N \times \mathbb{R}$ it holds $u(y, s) = as + b$ for some $a, b \in \mathbb{R}$.
- (ii) If M is non-parabolic and
 - u has at most linear growth on one side;
 - there exists an origin $o \in M$ such that, denoting with r the distance from o ,

$$(3) \quad \text{Ric}^{(m-2)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2} \quad \text{on } M,$$

for some constant $\bar{\kappa} \geq 0$,
then every tangent cone at infinity of M splits off a line.

Remark 6. In Case (ii), since M is non-parabolic then necessarily $m \geq 3$ (see below), so $\text{Ric}^{(m-2)}$ is well-defined. Also, we emphasize in the statement of (ii) that tangent cones at infinity may be based at any point of M , not necessarily at o .

Case (i) in Theorem 5 is easy to obtain, and might be well-known among specialists. We included it for the sake of completeness. Regarding Case (ii), the curvature condition in (3) is only used to infer that u has bounded gradient on M . In other words, as a consequence of our proof we obtain a generalization of Moser's result in [54] to the following

Theorem 7. *Let M be a complete manifold with $\text{Ric} \geq 0$. If u is a non-constant solution of (MSE) and $|Du| \in L^\infty(M)$, then every tangent cone at infinity of M splits off a line.*

It was already observed in [19] that a manifold M with $\text{Ric} \geq 0$ may not split off any line despite each of its tangent cones at infinity splits off a line. A counterexample was constructed in [45], and building on it we get the following result:

Proposition 8. *For $m \geq 4$, there exists a complete manifold M with*

$$\text{Ric}^{(2)} \geq 0, \quad \text{Ric} > 0, \quad |\text{Sec}| \leq \bar{\kappa}^2$$

for some constant $\bar{\kappa} > 0$, which carries a non-constant minimal graph $u : M \rightarrow \mathbb{R}$ with $|Du| \in L^\infty(M)$.

Note that $\text{Ric}^{(m-2)} \geq 0$ and that M does not split off any line. Whence, in assumption (ii) of Theorem 5, the conclusion cannot be strengthened to a splitting of M itself.

When $\text{Sec} \geq 0$, however, the above phenomenon does not happen. Leaving aside the case $m = 2$, covered by Case (i) in Theorem 5, we obtain

Corollary 9. *Let (M, σ) be a complete Riemannian manifold of dimension $m \geq 3$ satisfying $\text{Sec} \geq 0$. If there exists a non-constant entire solution $u \in C^\infty(M)$ of (MSE) with at most linear growth on one side, then $M = N \times \mathbb{R}$ with the product metric $\sigma_N + ds^2$, for some complete manifold N with $\text{Sec}_N \geq 0$, and in the variables $(y, s) \in N \times \mathbb{R}$ it holds $u(y, s) = as + b$ for some $a, b \in \mathbb{R}$.*

The corresponding problem for harmonic functions was also studied by A. Kasue [44]. Corollary 9 relates to the results obtained in [49] by P. Li and J. Wang. There, the authors study complete, stable minimal hypersurfaces $\Sigma \rightarrow \bar{M}$ properly immersed into a complete manifold \bar{M} whose sectional curvature is non-negative, and prove that either Σ has only one end or Σ is a totally geodesic cylinder $P \times \mathbb{R}$, for some compact manifold P with non-negative sectional curvature. Our setting falls into their framework, since a minimal graph in $\mathbb{R} \times M$ is stable, properly embedded and $\bar{M} = \mathbb{R} \times M$ has non-negative sectional curvature. However, our conclusion is stronger, since it allows M to have only one end and it also implies a splitting of M itself.

To conclude, we prove the next result for graphs with slower than linear growth on one side, that should be compared to [31, Theorem 3.6] and [30, Theorem 1.4].

Theorem 10. *Let (M, σ) be a complete Riemannian manifold of dimension $m \geq 2$ with $\text{Ric} \geq 0$, and let $u \in C^\infty(M)$ solve (MSE) on M and satisfy*

$$(4) \quad \lim_{r(x) \rightarrow \infty} \frac{u_-(x)}{r(x)} = 0,$$

where $u_-(x) = \max\{-u(x), 0\}$. Assume that either

- M is parabolic, or
- M is non-parabolic and there exists an origin $o \in M$ such that, denoting with r the distance from o ,

$$\text{Ric}^{(m-2)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2} \quad \text{on } M,$$

for some constant $\bar{\kappa} \geq 0$.

Then, u is constant.

Strategy of the proof. Case (i) in Theorem 5 is a direct consequence of the parabolicity of (M, σ) , which in our setting can be transplanted to the graph of u . In particular, since every surface with $\text{Ric} \geq 0$ is parabolic, the result holds if $m = 2$, so we focus on dimension $m \geq 3$. In [12], the authors obtain (B3) in \mathbb{R}^m via the following steps:

- (a) a sharp gradient estimate, implying that a solution $u \in C^\infty(\mathbb{R}^m)$ of (MSE) with at most linear growth on one side satisfies $|Du| \in L^\infty(\mathbb{R}^m)$;
- (b) an argument of Moser [54]: since $|Du| \in L^\infty(\mathbb{R}^m)$, for each coordinate field ∂_j the partial derivative $\partial_j u$ is a bounded solution of a uniformly elliptic PDE. The global Harnack inequality implies that $\partial_j u$ is constant, which implies that u is affine.

Step (b) cannot be implemented on manifolds, which in general lack parallel fields. An alternative idea was proposed in [19] to study harmonic functions, a blowdown argument which exploits the convergence theory of manifolds with $\text{Ric} \geq 0$. Our strategy closely follows the one in [19], and can be split into the following steps:

- (a) we prove that a solution $u \in C^\infty(M)$ of (MSE) with at most linear growth on one side satisfies $|Du| \in L^\infty(M)$;
- (b) for fixed $x_0 \in M$, we show that the functions $|Du|$ and $|D^2u|$ satisfy

$$(5) \quad \lim_{R \rightarrow \infty} \frac{1}{|B_R(x_0)|} \int_{B_R(x_0)} |Du|^2 dx = \sup_M |Du|^2,$$

$$(6) \quad \lim_{R \rightarrow \infty} \frac{R^2}{|B_R(x_0)|} \int_{B_R(x_0)} |D^2u|^2 dx = 0,$$

where $B_R(x_0)$ is the geodesic ball of radius R and center x_0 in (M, σ) .

- (c) we use the blowdown argument to guarantee the splitting of any tangent cone at infinity with base point x_0 .

To be more precise, Step (a) will be achieved by assuming

$$(7) \quad \text{Ric} \geq 0 \quad \text{and} \quad \text{Ric}^{(m-2)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2},$$

while (b) will be shown by requiring

$$(8) \quad \text{Ric} \geq 0 \quad \text{and} \quad |Du| \in L^\infty(M).$$

Though the strategy is the same as that in [19], we emphasize that the techniques in the current literature to prove (a) (respectively, (b)) do not apply under the sole assumptions in (7) (respectively, in (8)). We shall justify this claim in the next sections. Our strategy to obtain (a) is to refine a method due to N. Korevaar [46], see Theorem 13 below, while to get (b) in our needed generality we exploit heat

equation techniques, inspired by works of P. Li [47] and L. Saloff-Coste [58]. In this respect, we underline Theorem 25 below, yielding to (6), that in the stated generality seems to us new and of an independent interest.

2. PRELIMINARIES

We briefly review some formulas for minimal graphs that will be used later on. In local coordinates (x^i) on M , we write

$$\sigma = \sigma_{ij} dx^i \otimes dx^j, \quad g = g_{ij} dx^i \otimes dx^j, \quad du = u_i dx^i.$$

We have $g_{ij} = \sigma_{ij} + u_i u_j$. Denoting as σ^{ij} , g^{ij} the inverse matrices of (σ_{ij}) , (g_{ij}) , respectively, it holds

$$g^{ij} = \sigma^{ij} - \frac{u^i u^j}{W^2},$$

where $W = \sqrt{1 + |Du|^2}$ and $u^i = \sigma^{ij} u_j$. In general, if $\phi \in C^1(M)$ then we denote as $D\phi$ and $\nabla\phi$ the gradients of ϕ with respect to metrics σ and g , respectively, and in local notation we write

$$d\phi = \phi_i dx^i, \quad D\phi = \phi^i \partial_{x_i} \equiv \sigma^{ij} \phi_j \partial_{x_i}, \quad \nabla\phi = g^{ij} \phi_j \partial_{x_i}.$$

Differentiating the upward pointing unit normal vector $\mathbf{n} = W^{-1}(\partial_t - u^i e_i)$, the second fundamental form II in the direction of \mathbf{n} has components

$$(9) \quad \Pi_{ij} = \frac{u_{ij}}{W},$$

where u_{ij} are the components of the Hessian D^2u in the metric σ . Let $H = g^{ij} h_{ij}$ be the mean curvature, which we assume to vanish. Using the relation

$$\Gamma_{ij}^k - \gamma_{ij}^k = \frac{u^k u_{ij}}{W^2}$$

between the Christoffel coefficients Γ_{ij}^k of g and γ_{ij}^k of σ , for every $\phi : M \rightarrow \mathbb{R}$ the Laplace-Beltrami operator Δ_g of g writes as

$$\Delta_g \phi = g^{ij} \phi_{ij} - \phi_k u^k \frac{H}{W} = g^{ij} \phi_{ij},$$

where we used the minimality of Σ . Also, Δ_g has the following local expression:

$$(10) \quad \Delta_g \phi = \frac{1}{\sqrt{|g|}} \operatorname{div} \left(\sqrt{|g|} g^{ij} \phi_i \partial_{x_j} \right) = \frac{1}{W} \operatorname{div} (W g^{ij} \phi_i \partial_{x_j})$$

where div is, as before, the divergence operator in (M, σ) and $|g| = W^2$ is the determinant of g with respect to σ . Next, for every Killing field \bar{X} defined in $\mathbb{R} \times M$, the angle function $\Theta_{\bar{X}} \doteq \langle \mathbf{n}, \bar{X} \rangle$ solves the Jacobi equation

$$(11) \quad \Delta_g \Theta_{\bar{X}} + \left(\|\Pi\|^2 + \overline{\operatorname{Ric}}(\mathbf{n}, \mathbf{n}) \right) \Theta_{\bar{X}},$$

with $\overline{\operatorname{Ric}}$ the Ricci curvature of $\mathbb{R} \times M$. This is the case, for instance, of the angle function $\Theta_{\partial_t} = \langle \mathbf{n}, \partial_t \rangle = W^{-1}$ associated to the Killing field ∂_t . As a consequence, W satisfies

$$(12) \quad \mathcal{L}_W W = \left(\|\Pi\|^2 + \overline{\operatorname{Ric}}(\mathbf{n}, \mathbf{n}) \right) W,$$

where we defined

$$\mathcal{L}_W \phi \doteq W^2 \operatorname{div}_g (W^{-2} \nabla \phi) = \Delta_g \phi - 2 \langle \nabla \log W, \nabla \phi \rangle.$$

Observe that, in terms of the metric σ ,

$$(13) \quad \mathcal{L}_W \phi = W \operatorname{div} (W^{-1} g^{ij} \phi_i \partial_{x_j}).$$

If X is a Killing field in (M, σ) then we can extend it by parallel transport on $\mathbb{R} \times M$ to a Killing field \bar{X} satisfying $\langle \partial_t, \bar{X} \rangle = 0$, with corresponding angle function $\Theta_{\bar{X}} = \langle \mathbf{n}, \bar{X} \rangle = -W^{-1} \sigma(Du, X)$. Since (11) holds for both Θ_{∂_t} and $\Theta_{\bar{X}}$, it can be checked that the quotient

$$v \doteq -\Theta_{\bar{X}} / \Theta_{\partial_y} = \sigma(Du, X)$$

is a solution of

$$(14) \quad \mathcal{L}_W v = 0.$$

We next discuss the implications of ℓ -th Ricci curvature lower bounds. Hereafter, we set $\mathbb{R}^+ = (0, \infty)$ and $\mathbb{R}_0^+ = [0, \infty)$.

Proposition 11. *Let (M, σ) be a complete manifold of dimension $m \geq 2$ satisfying*

$$\operatorname{Ric}^{(\ell)}(\nabla r) \geq -H(r) \quad \text{for } \ell = \max\{1, m-2\},$$

where r is the distance from a fixed origin $o \in M$, and $0 \leq H \in C(\mathbb{R}_0^+)$. Let $h \in C^2(\mathbb{R}_0^+)$ solve

$$(15) \quad \begin{cases} h'' - Hh \geq 0 & \text{on } \mathbb{R}^+, \\ \liminf_{t \rightarrow 0} \left(\frac{h'}{h} - \frac{1}{t} \right) \geq 0. \end{cases}$$

Let $u : M \rightarrow \mathbb{R}$ solve (MSE). Then, denoting with Δ_g the Laplacian on the graph of u ,

$$\Delta_g r \leq m \frac{h'(r)}{h(r)}$$

pointwise on $M \setminus (\{o\} \cup \operatorname{cut}(o))$ and in the barrier sense on $M \setminus \{o\}$.

Proof. Outside of $\{o\} \cup \operatorname{cut}(o)$, denote by $\{\lambda_j(D^2 r)\}$ the eigenvalues of $D^2 r$ in increasing order. The comparison theorem in [52, Prop. 7.4] guarantees that

$$(16) \quad \sum_{j=2}^m \lambda_j(D^2 r) \leq (m-1) \frac{h'(r)}{h(r)}$$

pointwise on $M \setminus (\{o\} \cup \operatorname{cut}(o))$ and in the barrier sense on $M \setminus \{o\}$. Note that the initial assumptions on h therein are $h(0) = 0$, $h'(0) \geq 1$, but the same proof works for the more general (15). In this respect, note that $h' > 0$ on \mathbb{R}^+ follows from $H \geq 0$.

To estimate $\Delta_g r = g^{ij} r_{ij}$, pick a point x where r is smooth. If $Du(x) = 0$, then $g^{ij} = \delta^{ij}$. In our assumptions, the Ricci curvature satisfies

$$\operatorname{Ric}(\nabla r, \nabla r) \geq -(m-1)H(r),$$

so by the Laplacian comparison theorem and since $h' > 0$,

$$\Delta_g r = \operatorname{Tr}(D^2 r) \leq (m-1) \frac{h'(r)}{h(r)} \leq m \frac{h'(r)}{h(r)}.$$

Assume that $Du(x) \neq 0$, write $\nu = Du/|Du|$ in a neighbourhood of x and complete it to a local orthonormal basis $\{\nu, e_\alpha\}$ with $2 \leq \alpha \leq m$. Note that g^{ij} is diagonalized

with eigenvalues $W^{-2} \leq 1$ in direction ν and 1 in directions $\{e_\alpha\}$. Expressing $\Delta_g r$ in the basis $\{e_\alpha, \nu\}$ we get

$$\begin{aligned}
\Delta_g r &= \frac{1}{W^2} D^2 r(\nu, \nu) + \sum_{\alpha=2}^m D^2 r(e_\alpha, e_\alpha) \\
(17) \quad &= \frac{1}{W^2} \left[\text{Tr}(D^2 r) - \sum_{\alpha=2}^m D^2 r(e_\alpha, e_\alpha) \right] + \sum_{\alpha=2}^m D^2 r(e_\alpha, e_\alpha) \\
&= \frac{1}{W^2} \text{Tr}(D^2 r) + \left[\frac{W^2 - 1}{W^2} \right] \sum_{\alpha=2}^m D^2 r(e_\alpha, e_\alpha).
\end{aligned}$$

By min-max and since the eigenvalues are ordered,

$$\text{Tr}(D^2 r) \leq \frac{m}{m-1} \sum_{\alpha=2}^m \lambda_\alpha(D^2 r), \quad \sum_{\alpha=2}^m D^2 r(e_\alpha, e_\alpha) \leq \sum_{\alpha=2}^m \lambda_\alpha(D^2 r).$$

Therefore,

$$\begin{aligned}
\Delta_g r &\leq \left[\frac{m}{(m-1)W^2} + \frac{W^2 - 1}{W^2} \right] \sum_{\alpha=2}^m \lambda_\alpha(D^2 r) \\
(18) \quad &\leq \left[\frac{1}{(m-1)W^2} + 1 \right] (m-1) \frac{h'(r)}{h(r)} \leq m \frac{h'(r)}{h(r)},
\end{aligned}$$

as claimed. The inequality in the barrier sense on the entire $M \setminus \{o\}$ easily follows from the corresponding one for (16) via the argument in [52, Prop. 7.4]. \square

Remark 12. In particular, letting $\bar{\kappa} \in \mathbb{R}_0^+$, for $\ell = \max\{1, m-2\}$ it holds

$$\begin{aligned}
\text{Ric}^{(\ell)}(\nabla r) \geq -\bar{\kappa}^2 &\implies \Delta_g r \leq m\bar{\kappa} \coth(\bar{\kappa}r) \\
\text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2} &\implies \Delta_g r \leq \frac{m(1+\sqrt{1+4\bar{\kappa}^2})}{2r}
\end{aligned}$$

pointwise outside of $\{o\} \cup \text{cut}(o)$ and in the barrier sense on $M \setminus \{o\}$. Indeed, it is enough to consider for h , respectively, the functions

$$h(t) = \frac{\sinh(\bar{\kappa}t)}{\bar{\kappa}} \quad \text{and} \quad h(t) = t^{\bar{\kappa}'}, \quad \text{where } \bar{\kappa}' = \frac{1 + \sqrt{1 + 4\bar{\kappa}^2}}{2}.$$

3. PROOF OF THEOREM 5, (i)

Since $\text{Ric} \geq 0$, by work of N. Varopoulos [63] and Li and Yau [50], the parabolicity of M is equivalent to

$$(19) \quad \int^\infty \frac{s ds}{|B_s|} = \infty,$$

where B_s is a geodesic ball centered at a fixed origin o . Indeed, (19) is sufficient for the parabolicity of a complete manifold, independently of any curvature requirement. An argument in [48] (cf. also [61]) shows that the volume of the graph $\Sigma = (M, g)$ inside an extrinsic ball $\mathbb{B}_r \subset \mathbb{R} \times M$ centered at a point $(u(o), o)$ satisfies

$$|\Sigma \cap \mathbb{B}_r| \leq |B_r| + \frac{1}{2} |B_{3r} \setminus B_r| \leq 2|B_{3r}|.$$

Hence, the volume of a geodesic ball B_r^g in Σ centered at o is bounded as follows:

$$(20) \quad |B_r^g| \leq |\Sigma \cap \mathbb{B}_r| \leq 2|B_{3r}|,$$

which implies

$$\int^{\infty} \frac{sd s}{|B_s^g|} = \infty$$

and therefore the graph $\Sigma = (M, g)$ is parabolic. By the Jacobi equation

$$\Delta_g \frac{1}{W} = -(\|\Pi\|^2 + \overline{\text{Ric}}(\mathbf{n}, \mathbf{n})) \frac{1}{W},$$

the bounded function $1/W$ is superharmonic on Σ , and it is therefore constant by parabolicity. Hence, $\Pi \equiv 0$ and Σ is totally geodesic in $M \times \mathbb{R}$. Equivalently, by (9), $D^2u \equiv 0$ in M . As a consequence, since u is non-constant, Du is a non-zero parallel vector field, and M splits isometrically as a product $N \times \mathbb{R}$, u being an affine function of the \mathbb{R} -coordinate alone.

4. A LOCAL GRADIENT ESTIMATE

Let $u : B_R \subset M \rightarrow \mathbb{R}$ solve (MSE) on a geodesic ball $B_R = B_R(x)$. The original argument in [12] to prove the gradient estimate in Euclidean setting

$$(21) \quad |Du(x)| \leq c_1 \exp \left\{ c_2 \frac{u(x) - \inf_{B_R} u}{R} \right\},$$

for some constants $c_j = c_j(m)$, makes use of the isoperimetric inequality, which does not hold for minimal graphs over manifolds (M, σ) with $\text{Ric} \geq 0$ unless M has maximal volume growth compatible with the Bishop-Gromov inequality, in the sense that (2.a) is satisfied. Indeed, the isoperimetric inequality forces geodesic balls B_r^g in the graph $\Sigma = (M, g)$ to satisfy $|B_r^g| \geq Cr^m$, which coupled with (20) imply that M has Euclidean volume growth.

In the seminal paper [14], Bombieri and Giusti improved (21) to the following inequality:

$$(22) \quad |Du(x)| \leq c_1 \left\{ 1 + \frac{\sup_{B_R} |u(x)|}{R} \right\}^m,$$

showing that, differently to the case $m = 2$ (cf. [37]), an exponentially growing bound in terms of $|u|$ is not sharp. If M has $\text{Ric} \geq 0$ and Euclidean volume growth, the validity of an isoperimetric inequality on entire minimal graphs was recently shown in [30], see also [15] for the case $\text{Sec} \geq 0$. As a consequence, in [30, Theorems 1.3 and 6.2] the author was able to extend (21) and (22) to such manifolds.

An alternative method to prove (21) in Euclidean setting was given by N. Trudinger [61]. His strategy hinges on a mean value inequality on Σ which, remarkably, is obtained without needing the isoperimetric inequality and is therefore suited to apply to manifolds that do not have a Euclidean volume growth. However, to adapt the proof to minimal graphs over M , it seems that an upper bound on the sectional curvature of M is necessary, see also the related [31]. Later, the striking argument found by N. Korevaar in [46] gave another method to obtain gradient estimates, which only requires lower bounds on the curvatures of M . Exploiting Korevaar's method, in [57] the authors obtained the slightly different estimate

$$(23) \quad |Du(x)| \leq c_1 \exp \left\{ c_2 [1 + \bar{\kappa} R \coth(\bar{\kappa} R)] \frac{(u(x) - \inf_{B_R} u)^2}{R^2} \right\},$$

provided that $\text{Ric} \geq 0$ and $\text{Sec} \geq -\bar{\kappa}^2$. Note that, unless $\bar{\kappa} = 0$, the estimate explodes as $R \rightarrow \infty$ if $u : M \rightarrow \mathbb{R}$ is of linear growth. Extensions to more general

ambient spaces were later given in [18, 23, 24], but they only consider graphs which are bounded on one side or have logarithmic growth.

Inspecting the proofs in [57, 23, 24, 31], to reach the inequality $|Du(x)| \leq C$ for solutions of linear growth the bounds on Sec are instrumental to guarantee that the distance r_x from x satisfies $\Delta_g r_x \leq C_1/r_x$ for some absolute constant C_1 . In view of the arbitrariness of the point x , assumption $\text{Sec} \geq 0$ in [57] seems therefore difficult to replace with a weaker control on Sec from below. For instance, if one considers the inequality $\text{Sec} \geq -\bar{\kappa}^2/(1+r_o^2)$ for some constant $\bar{\kappa} > 0$ and some origin o , comparison theory and standard estimates for ODE would yield to a constant C_1 , hence C , that depends on the distance of x from o and explodes as $r_o(x) \rightarrow \infty$, making the estimate on $\Delta_g r_x$ insufficient to imply the desired gradient bound.

From a different perspective, we mention that a *global* gradient estimate for *positive* entire solutions was obtained in [22] under the sole curvature assumption

$$\text{Ric} \geq -(m-1)\kappa^2, \quad \kappa \in \mathbb{R}_0^+,$$

namely, a positive solution of (MSE) on the entire M shall satisfy

$$\sqrt{1 + |Du(x)|^2} \leq e^{\kappa u(x)\sqrt{m-1}} \quad \forall x \in M.$$

Note that (B2) directly follows if $\kappa = 0$. However, modifying the argument in [22] to allow for linearly growing solutions seems challenging.

Our first main result, Theorem 13, provides an improvement of Korevaar's method that apply to the more general assumption (7).

Theorem 13. *Let (M, σ) be a Riemannian manifold with dimension $m \geq 2$. Let $B_R = B_R(o) \subseteq M$ be a geodesic open ball of radius $R > 0$ centered at $o \in M$ and let $u \in C^3(\overline{B_R})$ be a non-constant solution to (MSE). Assume that*

$$\text{Ric} \geq -(m-1)\kappa^2, \quad \text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2} \quad \text{on } B_R, \quad \ell = \max\{1, m-2\},$$

for some $\kappa, \bar{\kappa} \in \mathbb{R}_0^+$, where r denotes the distance from o . Let $0 < R_1 < R$. Then,

$$\begin{aligned} & \sqrt{1 + |Du(x)|^2} \\ & \leq \max \left\{ \sqrt{1 + a_0^2(\gamma^*)^2}, \sqrt{\frac{a_3}{a_3 - a_2}} \right\} \left(\frac{e^{LR(\sqrt{\varepsilon^2+1}-\varepsilon)} - 1}{e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x)^2/R^2-q\gamma(x)})} - 1} \right) \end{aligned}$$

for every $x \in B_{R_1}(o)$, where

$$\gamma(x) = \frac{u(x) - \inf_{B_R} u}{R}, \quad \gamma^* = \sup_{x \in B_{R_1}} \gamma(x) = \frac{\sup_{B_{R_1}} u - \inf_{B_R} u}{R},$$

$\varepsilon > 0$ and $\tau \in (0, 1)$ are fixed arbitrarily, $q, a_0 \in \mathbb{R}^+$ satisfy

$$\frac{\sqrt{1 + \varepsilon^2} - \sqrt{(R_1/R)^2 + \varepsilon^2}}{\gamma^*} > q > \frac{1}{\sqrt{\tau} a_0 \gamma^*} > 0$$

and $L \in \mathbb{R}^+$ satisfies

$$(1 - \tau) \left(q^2 - \frac{1}{\tau a_0^2(\gamma^*)^2} \right) L^2 - \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} > (m-1)\kappa^2$$

with $\bar{\kappa}_0 = \max\{1, \bar{\kappa}\}$. Finally, a_2, a_3 are defined by

$$a_2 = \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} \quad a_3 = (1 - \tau) \left(q^2 - \frac{1}{\tau a_0^2(\gamma^*)^2} \right) L^2 - (m-1)\kappa^2.$$

Remark 14. *The assumption that u is non-constant ensures that $\gamma_* > 0$ by the maximum principle.*

Before proving the Theorem, we give some applications, starting from the case where $\kappa = 0$.

Corollary 15. *Let (M^m, σ) be a Riemannian manifold with $\text{Ric} \geq 0$ and*

$$\text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2}, \quad \ell = \max\{1, m-2\},$$

for some $\bar{\kappa} \in \mathbb{R}_0^+$, where r denotes the distance from o . Let $u \in C^3(\overline{B_R})$ solve (MSE). Then, for every $\delta \in (0, 1)$ and for every $R_1 \in (0, \delta R)$,

$$(24) \quad \sup_{B_{R_1}} \sqrt{1 + |Du|^2} \leq C_1 \exp \left(C_2 m \bar{\kappa}_0 \frac{[\sup_{B_{R_1}} u - \inf_{B_R} u]^2}{R^2} \right)$$

with $\bar{\kappa}_0 = \max\{1, \bar{\kappa}\}$ and $C_1, C_2 > 0$ only depending on δ .

Proof. The desired inequality is trivial if u is constant, so assume that u is non-constant. It suffices to prove the claim for $\delta \in [1/2, 1)$. Let

$$\gamma^* = \frac{\sup_{B_{R_1}} u - \inf_{B_R} u}{R}.$$

Choose

$$\tau = \frac{1}{2}, \quad \varepsilon = \delta, \quad q = \frac{1-\delta}{2\sqrt{2}\gamma^*}, \quad a_0 = \frac{2}{q\gamma^*}, \quad L = \frac{8(m+1)\bar{\kappa}_0}{\varepsilon R q^2}.$$

With this choice, we have

$$(25) \quad \frac{\sqrt{1+\varepsilon^2} - \sqrt{(R_1/R)^2 + \varepsilon^2}}{\gamma^*} \geq \frac{\sqrt{1+\varepsilon^2} - \sqrt{\delta^2 + \varepsilon^2}}{\gamma^*} \\ = \frac{\sqrt{1+\delta^2} - \sqrt{2}\delta}{\gamma^*} \geq 2q,$$

where, from $\delta < 1$, we used

$$\sqrt{1+\delta^2} - \sqrt{2}\delta = \frac{1-\delta^2}{\sqrt{1+\delta^2} + \sqrt{2}\delta} \geq \frac{1-\delta^2}{\sqrt{2} + \sqrt{2}\delta} = \frac{1-\delta}{\sqrt{2}}.$$

We also have

$$q^2 - \frac{1}{a_0^2(\gamma^*)^2\tau} = q^2 - \frac{q^2}{2} = \frac{q^2}{2}$$

and then

$$a_3 = (1-\tau) \left(q^2 - \frac{1}{a_0^2(\gamma^*)^2\tau} \right) L^2 = \frac{L^2 q^2}{4} = 2 \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} = 2a_2.$$

Hence, all assumptions of Theorem ?? are satisfied and for every $x \in B_{R_1}$ we have

$$\sqrt{1 + |Du(x)|^2} \leq \max \left\{ \sqrt{1 + a_0^2(\gamma^*)^2}, \sqrt{2} \right\} \times \\ \times \frac{e^{LR(\sqrt{1+\varepsilon^2}-\varepsilon)} - 1}{e^{LR(\sqrt{1+\varepsilon^2}-\sqrt{(r(x)/R)^2+\varepsilon^2}-q\gamma(x))} - 1}.$$

Note that, for every $x \in B_{R_1}$ and taking into account (25),

$$\begin{aligned} \sqrt{1 + \varepsilon^2} - \sqrt{(r(x)/R)^2 + \varepsilon^2} - q\gamma(x) &\geq \sqrt{1 + \varepsilon^2} - \sqrt{\delta^2 + \varepsilon^2} - q\gamma^* \\ &\geq 2q\gamma^* - q\gamma^* = q\gamma^*, \end{aligned}$$

and also

$$\sqrt{1 + \varepsilon^2} - \varepsilon = \frac{1}{\sqrt{1 + \varepsilon^2} + \varepsilon} \leq \frac{1}{\sqrt{1 + \varepsilon^2}} = \frac{1}{\sqrt{1 + \delta^2}} \leq \frac{1}{\sqrt{2}\delta} = \frac{2}{\delta(1 - \delta)} q\gamma^*.$$

Therefore, we can estimate

$$\frac{e^{LR(\sqrt{1 + \varepsilon^2} - \varepsilon)} - 1}{e^{LR(\sqrt{1 + \varepsilon^2} - \sqrt{(r(x)/R)^2 + \varepsilon^2} - q\gamma(x))} - 1} \leq \frac{e^{LR\frac{2}{\delta(1 - \delta)}q\gamma^*} - 1}{e^{LRq\gamma^*} - 1} \leq C(\delta)e^{LR(\frac{2}{\delta(1 - \delta)} - 1)q\gamma^*}.$$

Here, we have exploited the fact that for every $\alpha \in \mathbb{R}$ one has the validity of an inequality of the form

$$\frac{y^\alpha - 1}{y - 1} \leq C(\alpha)y^{\alpha-1} \quad \forall y > 1$$

for a suitable constant $C(\alpha) > 0$. Recalling that

$$\begin{aligned} a_0^2(\gamma^*)^2 &= \frac{32(\gamma^*)^2}{(1 - \delta)^2}, \\ LRq\gamma^* &= \frac{8(m + 1)\bar{\kappa}_0\gamma^*}{\delta q} = \frac{16\sqrt{2}(m + 1)\bar{\kappa}_0}{\delta(1 - \delta)}(\gamma^*)^2 \end{aligned}$$

we obtain

$$\begin{aligned} \sqrt{1 + |Du(x)|^2} &\leq \max \left\{ \sqrt{1 + \frac{32(\gamma^*)^2}{(1 - \delta)^2}}, \sqrt{2} \right\} \times \\ &\quad \times C(\delta) \exp \left(\frac{16\sqrt{2}(m + 1)\bar{\kappa}_0}{\delta(1 - \delta)} \left(\frac{2}{\delta(1 - \delta)} - 1 \right) (\gamma^*)^2 \right) \end{aligned}$$

and the conclusion follows. \square

Assuming that u has at most linear growth, we get

Corollary 16. *Let (M^m, σ) be a complete Riemannian manifold with $\text{Ric} \geq 0$ and*

$$\text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1 + r^2}, \quad \ell = \max\{1, m - 2\},$$

for some $\bar{\kappa} \in \mathbb{R}_0^+$. If $u \in C^\infty(M)$ solves (MSE) and has at most linear growth on one side, then $|Du| \in L^\infty(M)$.

Proof. Without loss of generality we can assume that the negative part of u has at most linear growth, so that there exists $a > 0$ such that $u(x) \geq -a(1 + r(x))$ for every $x \in M$. Let $R_1 > 0$ be fixed. Choosing $\delta = 1/2$ and letting $R \rightarrow \infty$ in estimate (24) we get

$$\sup_{B_{R_1}} \sqrt{1 + |Du|^2} \leq C_1 \exp(C_2 m \bar{\kappa}_0 a^2)$$

where $C_1, C_2 > 0$ do not depend on R_1 . Since $R_1 > 0$ was arbitrary, we get

$$\sup_M \sqrt{1 + |Du|^2} \leq C_1 \exp(C_2 m \bar{\kappa}_0 a^2).$$

\square

To prove Theorem 13, we need the following

Lemma 17. *Let (M^m, σ) be a complete Riemannian manifold with*

$$\text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2}, \quad \ell = \max\{1, m-2\},$$

for some $\bar{\kappa} \in \mathbb{R}_0^+$, where r is the distance from a fixed origin $o \in M$, and let $u \in C^\infty(B_R)$ solve (MSE) on a geodesic ball B_R centered at o .

For any given $a > 0$, the function $\psi = \sqrt{a^2 + r^2}$ satisfies

$$|D\psi| < 1, \quad \Delta_g \psi \leq (m+1) \frac{\max\{1, \bar{\kappa}\}}{a}$$

in the barrier sense on B_R and pointwise on $B_R \setminus \text{cut}(o)$.

Remark 18. *By its very definition, a solution in the barrier sense is also a solution in the viscosity sense, see [51] for comments.*

Proof. Outside of $\{o\} \cup \text{cut}(o)$, a direct computation yields $|D\psi| = \frac{r}{\psi} |Dr| < 1$ and

$$\Delta_g \psi = \frac{r \Delta_g r}{\sqrt{a^2 + r^2}} + \frac{a^2 \|\nabla r\|^2}{(a^2 + r^2)^{3/2}}.$$

From $\|\nabla r\|^2 = g^{ij} r_i r_j \leq |Dr|^2 = 1$ and Remark 12,

$$\Delta_g \psi \leq \frac{1}{\sqrt{a^2 + r^2}} \left(\frac{m(1 + \sqrt{1 + 4\bar{\kappa}^2})}{2} + \frac{a^2}{a^2 + r^2} \right)$$

and the conclusion follows by observing that $\frac{1}{\sqrt{a^2 + r^2}} \leq \frac{1}{a}$ and that

$$\frac{m(1 + \sqrt{1 + 4\bar{\kappa}^2})}{2} + \frac{a^2}{a^2 + r^2} \leq m(1 + \bar{\kappa}) + 1 \leq (m+1) \max\{1, \bar{\kappa}\}.$$

The validity of the inequality in the barrier sense can be proved by Calabi's trick, see for instance [52, Prop. 7.4]. \square

Proof of Theorem 13. Without loss of generality, we can assume $\inf_{B_R} u = 0$. Then

$$\gamma^* = \sup_{x \in B_{R_1}} \gamma(x) = \frac{\sup_{B_{R_1}} u}{R}.$$

As in the statement of the theorem, fix $\tau \in (0, 1)$ and $\varepsilon > 0$, and choose $q > 0$ and $a_0 > 0$ such that

$$(26) \quad \frac{\sqrt{\varepsilon^2 + 1} - \sqrt{(R_1/R)^2 + \varepsilon^2}}{\gamma^*} > q > \frac{1}{\sqrt{\tau} a_0 \gamma^*}$$

and then $L > 0$ matching

$$(27) \quad (1 - \tau) \left(q^2 - \frac{1}{\tau a_0^2 (\gamma^*)^2} \right) L^2 - \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} > (m-1)\kappa^2,$$

where $\bar{\kappa}_0 = \max\{1, \bar{\kappa}\}$. Set

$$C = qL, \quad \delta = e^{-LR\sqrt{\varepsilon^2 + 1}},$$

define the function

$$\psi = \sqrt{\varepsilon^2 R^2 + r^2}$$

where $r(x) = \text{dist}_\sigma(o, x)$, and let

$$\eta = e^{-Cu - L\psi} - \delta, \quad z = W\eta.$$

By writing

$$\eta = \delta \left(e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+(r/R)^2}-qu/R)} - 1 \right)$$

we see that for every $x \in B_{R_1}$

$$\begin{aligned} \eta(x) &= \delta \left(e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x)^2/R^2}-q\gamma(x))} - 1 \right) \\ &\geq \delta \left(e^{LR(\sqrt{1+\varepsilon^2}-\sqrt{(R_1/R)^2+\varepsilon^2}-q\gamma^*)} - 1 \right) > 0 \end{aligned}$$

as a consequence of (26). Noting that, on ∂B_R ,

$$\eta = \delta (e^{-qLu} - 1) < 0,$$

the set

$$\Omega = \{x \in \overline{B_R} : z(x) > 0\} \equiv \{x \in \overline{B_R} : \eta(x) > 0\}$$

is non-empty and satisfies $B_{R_1} \subseteq \Omega \subseteq B_R$. Therefore, there exists $x_0 \in \Omega$ such that

$$0 < z(x_0) = \max_{\Omega} z.$$

The function z satisfies

$$\Delta_g z - 2\langle \nabla z, \nabla \log W \rangle \geq \left(-(m-1)\kappa^2 \|\nabla u\|^2 + \frac{\Delta_g \eta}{\eta} \right) z \quad \text{on } \Omega.$$

The above inequality has to be interpreted in the viscosity sense, in case x_0 is not a point where r (hence, ψ) is smooth. By the maximum principle, necessarily

$$(28) \quad -(m-1)\kappa^2 \|\nabla u\|^2 + \frac{\Delta_g \eta}{\eta} \leq 0 \quad \text{at } x_0$$

in the viscosity sense. We compute

$$\Delta_g \eta = (\eta + \delta) (-C\Delta_g u - L\Delta_g \psi + \|C\nabla u + L\nabla \psi\|^2).$$

We recall that $W^{-2}(\sigma^{ij})_{i,j} \leq (g^{ij})_{i,j} \leq (\sigma^{ij})_{i,j}$ in the sense of quadratic forms, hence

$$\begin{aligned} \|C\nabla u + L\nabla \psi\|^2 &= g^{ij}(Cu_i + L\psi_i)(Cu_j + L\psi_j) \\ &\geq \frac{1}{W^2} \sigma^{ij}(Cu_i + L\psi_i)(Cu_j + L\psi_j) \\ &\geq \frac{1}{W^2} |CDu + LD\psi|^2. \end{aligned}$$

It follows that

$$\frac{\Delta_g \eta}{\eta + \delta} \geq -C\Delta_g u - L\Delta_g \psi + \frac{1}{W^2} |CDu + LD\psi|^2.$$

Using $\Delta_g u = 0$ and Young's inequality we obtain

$$\frac{\Delta_g \eta}{\eta + \delta} \geq -L\Delta_g \psi + (1-\tau)C^2 \frac{|Du|^2}{W^2} - L^2 \frac{1-\tau}{\tau} \frac{|D\psi|^2}{W^2}.$$

Taking into account Lemma 17 we infer

$$|D\psi| < 1, \quad \Delta_g \psi \leq \frac{(m+1)\bar{\kappa}_0}{\varepsilon R}.$$

Substituting these estimates in the above inequality, we deduce

$$\frac{\Delta_g \eta}{\eta + \delta} \geq (1-\tau)C^2 \frac{|Du|^2}{W^2} - L \left(\frac{(m+1)\bar{\kappa}_0}{\varepsilon R} + \frac{1-\tau}{\tau} \frac{L}{W^2} \right)$$

If $|Du(x_0)| \geq a_0\gamma^*$ then $\frac{|Du(x_0)|^2}{W^2 a_0^2 (\gamma^*)^2} \geq \frac{1}{W^2}$. Thus, we can further estimate

$$\frac{\Delta_g \eta}{\eta + \delta} \geq (1 - \tau) \left(q^2 - \frac{1}{\tau a_0^2 (\gamma^*)^2} \right) L^2 \frac{|Du|^2}{W^2} - \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} \quad \text{at } x_0$$

that is,

$$(29) \quad \frac{\Delta_g \eta}{\eta + \delta} \geq a_1 \|\nabla u\|^2 - a_2$$

with

$$a_1 = (1 - \tau) \left(q^2 - \frac{1}{\tau a_0^2 (\gamma^*)^2} \right) L^2 > 0, \quad a_2 = \frac{(m+1)\bar{\kappa}_0 L}{\varepsilon R} > 0.$$

Since

$$a_3 = a_1 - (m-1)\kappa^2$$

we have $a_1 \geq a_3 > 0$ by condition (27). We claim that

$$\frac{|Du(x_0)|^2}{W^2(x_0)} = \|\nabla u(x_0)\|^2 \leq \frac{a_2}{a_3},$$

that is,

$$(30) \quad W(x_0) \leq \sqrt{\frac{a_3}{a_3 - a_2}}.$$

Indeed, assume by contradiction that

$$(31) \quad \|\nabla u(x_0)\|^2 > \frac{a_2}{a_3}.$$

Then, from (29) it follows

$$\frac{\Delta_g \eta}{\eta + \delta} \geq a_3 \|\nabla u\|^2 - a_2 > 0 \quad \text{at } x_0,$$

hence $\Delta_g \eta > 0$ and, by (28) and (29) again,

$$(m-1)\kappa^2 \|\nabla u\|^2 \geq \frac{\Delta_g \eta}{\eta} \geq \frac{\Delta_g \eta}{\eta + \delta} \geq a_1 \|\nabla u\|^2 - a_2 \quad \text{at } x_0,$$

leading to

$$a_2 \geq (a_1 - (m-1)\kappa^2) \|\nabla u\|^2 = a_3 \|\nabla u\|^2 \quad \text{at } x_0,$$

which contradicts (31) and proves our claim.

On the other hand, if $|Du(x_0)| \leq a_0\gamma^*$, then

$$(32) \quad W(x_0) \leq \sqrt{1 + a_0^2 (\gamma^*)^2}$$

Since x_0 is a global maximum point for z in Ω , we have $z(x) \leq z(x_0)$, that is,

$$W(x) \leq W(x_0) \frac{\eta(x_0)}{\eta(x)}$$

for every $x \in B_{R_1} \subseteq \Omega$. Note that

$$\begin{aligned} \frac{\eta(x_0)}{\eta(x)} &= \frac{e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x_0)^2/R^2-qu(x_0)/R})} - 1}{e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x)^2/R^2-q\gamma(x)})} - 1} \\ &\leq \frac{e^{LR(\sqrt{\varepsilon^2+1}-\varepsilon)} - 1}{e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x)^2/R^2-q\gamma(x)})} - 1} \end{aligned}$$

hence

$$W(x) \leq W(x_0) \left(\frac{e^{LR(\sqrt{\varepsilon^2+1}-\varepsilon)} - 1}{e^{LR(\sqrt{\varepsilon^2+1}-\sqrt{\varepsilon^2+r(x)^2/R^2-q\gamma(x)})} - 1} \right).$$

The latter, together with (30) and (32), implies the desired estimate. \square

5. UNIFORMLY ELLIPTIC OPERATORS ON MANIFOLDS WITH $\text{Ric} \geq 0$

Having shown that an entire minimal graph with at most linear growth on one side has globally bounded gradient, we need to show (5) and (6). We shall prove both of them under the only conditions

$$(33) \quad \text{Ric} \geq 0, \quad |Du| \in L^\infty(M).$$

In such generality, it seems difficult to apply the ‘‘elliptic’’ approach in [19], adapted in [31, 30]. To justify the statement, we observe that the method in [19] relies on the construction of a function ϱ satisfying

$$(34) \quad C^{-1}r \leq \varrho \leq Cr, \quad |D\varrho| \leq C, \quad \Delta_g \varrho \leq \frac{C}{\varrho},$$

for some absolute constant C . When considering harmonic functions, the third condition is replaced by $\Delta \varrho \leq C/\varrho$, thus by comparison theory the choice $\varrho = r$ is admissible. On the contrary, to our knowledge, for minimal graphs the existence of ϱ satisfying (34) is currently unknown under the sole assumptions (33). If M has Euclidean volume growth, we mention that in [31, 30] the authors used as ϱ a reparametrization of the Green kernel of the Laplacian on M . Although the inequality $\Delta_g \varrho \leq C/\varrho$ may not hold pointwise, the integral estimates for $|D^2 \varrho|$ provided in [21] suffice to estimate $\Delta_g \varrho$ and apply the method in [19], as done in [30, Lemma 7.1]. However, to our knowledge, estimates like those in [21] are not yet (if ever) available on manifolds with $\text{Ric} \geq 0$ but whose volume growth is less than Euclidean.

For these reasons, inspired by [47, 58] we choose a different approach via the heat equation. Throughout this section, let (M^m, σ) be a complete Riemannian manifold of dimension $m \geq 2$ with $\text{Ric} \geq 0$. Let L be the linear uniformly elliptic operator defined by

$$(35) \quad L\psi = \text{div}(A D\psi)$$

where A is a measurable section of $T^{1,1}M$ satisfying

$$(36) \quad \text{i) } \alpha^{-1}|X|^2 \leq \langle AX, X \rangle \quad \text{and} \quad \text{ii) } |AX| \leq \alpha|X| \quad \forall X \in TM$$

for some constant $\alpha > 0$. Hereafter, we shall assume that A is smooth, the general case being obtainable by approximation.

We denote by $H_L(x, y, t)$ the minimal heat kernel associated to the parabolic operator $\partial_t - L$, that is, the unique continuous function on $M \times M \times \mathbb{R}^+$ such that for every $\psi \in C_0^\infty(M)$ the function u defined by

$$u(t, x) = \int_M H_L(x, y, t) \psi(y) dy \quad \forall (t, x) \in \mathbb{R}^+ \times M$$

is a solution of

$$(37) \quad \partial_t u = Lu$$

on $\mathbb{R}^+ \times M$ satisfying

$$\text{i) } u(t, \cdot) \rightarrow \psi \text{ pointwise on } M \text{ as } t \searrow 0,$$

ii) $u \leq v$ on $(0, T) \times M$ for every $v \in C^2([0, T) \times M)$, $T > 0$, such that

$$\begin{cases} \partial_t v = Lv & \text{on } (0, T) \times M, \\ \psi \leq v(0, \cdot) & \text{on } M. \end{cases}$$

If the endomorphism A is self-adjoint with respect to $\langle \cdot, \cdot \rangle$, the minimal heat kernel H_L is a symmetric function of the space variables, that is,

$$(38) \quad H_L(x, y, t) = H_L(y, x, t) \quad \forall x, y \in M, \forall t > 0.$$

By [58], see Corollary 6.2 and Theorem 6.3, there exist positive constants $C_i > 0$, $1 \leq i \leq 6$, depending only on m and α such that, for every $x, y \in M$ and $t > 0$,

$$(39) \quad C_1 \frac{\exp\left(-C_2 \frac{\text{dist}(x, y)^2}{t}\right)}{\sqrt{|B_{\sqrt{t}}(x)| |B_{\sqrt{t}}(y)|}} \leq H_L(x, y, t) \leq C_3 \frac{\exp\left(-C_4 \frac{\text{dist}(x, y)^2}{t}\right)}{\sqrt{|B_{\sqrt{t}}(x)| |B_{\sqrt{t}}(y)|}}$$

and

$$(40) \quad |\partial_t H_L(x, y, t)| \leq \frac{C_5}{t} \frac{\exp\left(-C_6 \frac{\text{dist}(x, y)^2}{t}\right)}{\sqrt{|B_{\sqrt{t}}(x)| |B_{\sqrt{t}}(y)|}}$$

Remarks on (39) will be given in the Appendix. We first need the following simple estimate on the volume of geodesic balls.

Lemma 19. *Let (M^m, σ) be a complete manifold with $\text{Ric} \geq 0$. For every $x, y \in M$ and for every $R > 0$ it holds*

$$|B_R(x)| \left(1 + \frac{\text{dist}(x, y)}{R}\right)^{-\frac{m}{2}} \leq \sqrt{|B_R(x)| |B_R(y)|} \leq |B_R(x)| \left(1 + \frac{\text{dist}(x, y)}{R}\right)^{\frac{m}{2}}$$

Proof. By Bishop-Gromov's comparison Theorem we have

$$\frac{|B_R(x)|}{|B_r(x)|} \leq \left(\frac{R}{r}\right)^m, \quad 0 < r < R < \infty,$$

thus

$$|B_R(y)| \leq |B_{R+\text{dist}(x, y)}(x)| \leq |B_R(x)| \left(1 + \frac{\text{dist}(x, y)}{R}\right)^m,$$

and the thesis follows. \square

Next, we recall that L generates a diffusion which is stochastically complete (cf. [40]), that is, the following holds:

Lemma 20. *Let M be a complete manifold with $\text{Ric} \geq 0$, and let A, L be as in (35)-(36), with A self-adjoint and smooth. Then*

$$\int_M H_L(x, y, t) dy = 1 \quad \forall (t, x) \in \mathbb{R}^+ \times M.$$

The result is stated with no proof in the discussion following [58, Theorem 7.4]. We here provide an argument for the convenience of the reader.

Proof. Since L is uniformly elliptic and M has polynomial volume growth as a consequence of $\text{Ric} \geq 0$, by Theorem 4.1 of [1] we have that for any $\lambda > 0$ the only entire bounded solution v of $Lv = \lambda v$ on M is $v \equiv 0$. Then the conclusion follows by Theorem 3.11 of [56]. \square

With the above preparation, we are ready to state the following asymptotic mean value theorem. Our method is inspired by the one in [47], where the author considered the case $L = \Delta$, but with a noticeable difference. Indeed, in [47] the author needs to use the Li-Yau's differential Harnack inequality to get rid of a boundary term at infinity. The inequality is available for solutions of the heat equation, but in general it fails for solutions of $\partial_t u = Lu$ unless one has a uniform control on the gradient of A on the entire M . As we will apply our results to $A = W \text{Id} - W^{-1} du \otimes Du$, with $W = \sqrt{1 + |Du|^2}$, in our setting only an L^∞ control on A is available. For this reason, we shall modify the method, exploiting upper level sets of H_L rather than geodesic balls. Note that we do not assume a Euclidean volume growth. We start with the following

Lemma 21. *Let $(M^m, \langle \cdot, \cdot \rangle)$ be a complete, noncompact manifold with $\text{Ric} \geq 0$ and let A, L be as in (35)-(36), with A self-adjoint and smooth. If $f \in C^2(M) \cap L^\infty(M)$ satisfies $Lf \leq 0$ on M then the function $u : \mathbb{R}^+ \times M \rightarrow \mathbb{R}$ given by*

$$(41) \quad u(t, x) = \int_M f(y) H_L(x, y, t) dy \quad \forall (t, x) \in \mathbb{R}^+ \times M$$

satisfies

$$(42) \quad \partial_t u \leq 0 \quad \text{on } \mathbb{R}^+ \times M.$$

Proof. Note that the integral on the RHS of (41) converges for every $(t, x) \in \mathbb{R}^+ \times M$ since $f \in L^\infty(M)$ and because of (39) and Lemma 19. Also note that H_L is smooth as a consequence of the regularity assumptions on A . By Lemma 20, u only varies by an additive constant if so does f , hence without loss of generality we can assume $\inf_M f = 0$. Let $(t, x) \in \mathbb{R}^+ \times M$ be fixed. For notational convenience, for every $a > 0$ we define

$$(43) \quad \varphi^a(y) = H_L(x, y, t) - a \quad \forall y \in M, \quad \Omega_a = \{y \in M : \varphi_a(y) > 0\}.$$

Because of (39) it holds $H_L(x, y, t) \rightarrow 0$ as $y \rightarrow \infty$ in M , hence the collection $\{\Omega_a\}_{a>0}$ is an exhaustion of M by relatively compact open subsets, with $\Omega_a \subseteq \Omega_b$ when $a \geq b$. By (40) and boundedness of f , we can apply Lebesgue's dominated convergence theorem to get

$$(44) \quad \partial_t u(t, x) = \int_M f(y) \partial_t H_L(x, y, t) dy = \lim_{a \rightarrow 0^+} \int_{\Omega_a} f(y) \partial_t H_L(x, y, t) dy.$$

Therefore, since $Lf \leq 0$ and $\varphi_a > 0$ on Ω_a , (42) holds if we prove the inequality

$$(45) \quad \int_{\Omega_a} f(y) \partial_t H_L(x, y, t) dy \leq \int_{\Omega_a} \varphi_a(y) Lf(y) dy.$$

Because of $\partial_t H_L(x, y, t) = L_y H_L(x, y, t) = L\varphi(y) = L\varphi_a(y)$, we have

$$\int_{\Omega_a} f(y) \partial_t H_L(x, y, t) dy = \int_{\Omega_a} f(y) L_y H_L(x, y, t) dy = \int_{\Omega_a} f(y) L\varphi_a(y) dy.$$

Since $H_L \in C^\infty(\mathbb{R}^+ \times M)$, we have $\varphi \in C^\infty(M)$ and for almost every $a > 0$ the set Ω_a has smooth boundary. Let $a > 0$ be a regular value for φ . By Green's identity, since $\varphi_a = 0$ on $\partial\Omega_a$

$$\int_{\Omega_a} f(y) L\varphi_a(t, y) dy = \int_{\Omega_a} \varphi_a(t, y) Lf(y) dy + \int_{\partial\Omega_a} f(y) \langle AD\varphi_a(y), \nu \rangle d\mathcal{H}^{m-1}(y)$$

where $\nu = -D\varphi_a/|D\varphi_a|$ is the outward pointing normal on $\partial\Omega_a$. Noting that $f \geq 0$, that φ_a is non-increasing in the direction of ν and that A is positive definite, we see that $f\langle AD\varphi_a, \nu \rangle \leq 0$ on $\partial\Omega_a$ and therefore the second integral is non-positive, which implies the desired inequality (45). \square

Proposition 22. *Let $(M^m, \langle \cdot, \cdot \rangle)$ be a complete, noncompact manifold with $\text{Ric} \geq 0$ and let A, L be as in (35)-(36), with A self-adjoint and smooth. If $f \in C^2(M) \cap L^\infty(M)$ satisfies $Lf \leq 0$ on M , then for any $x \in M$*

$$(46) \quad \lim_{R \rightarrow \infty} \frac{1}{|B_R(x)|} \int_{B_R(x)} f(y) \, dy = \inf_M f.$$

Proof. Without loss of generality, we assume $\inf_M f = 0$. Let $u : \mathbb{R}^+ \times M \rightarrow \mathbb{R}$ be the function defined by (41). Note that u is the minimal solution of the parabolic equation $\partial_t u = Lu$ on $\mathbb{R}^+ \times M$ corresponding to the initial datum $u(0^+, \cdot) = f$. Hence, by the maximum principle and the monotonicity (42) we have

$$\inf_M f \leq u(t, x) \leq f(x) \quad \forall (t, x) \in \mathbb{R}^+ \times M.$$

In particular, the limit

$$u_\infty(x) = \lim_{t \rightarrow \infty} u(t, x)$$

is well defined for every $x \in M$. The convergence $u(t, \cdot) \rightarrow u_\infty$ is uniform on compact subsets, u_∞ is bounded and $Lu_\infty = 0$. Since M is complete and has non-negative Ricci curvature, the operator L enjoys a Liouville property, see Theorem 7.4 of [58]. In particular, u_∞ must be constant. Since $\inf_M f \leq u_\infty \leq f$, it must be $u_\infty \equiv \inf_M f = 0$, that is,

$$(47) \quad \lim_{t \rightarrow \infty} u(t, x) = 0 \quad \forall x \in M.$$

To conclude the proof of (46), we observe that

$$\begin{aligned} u(t, x) &= \int_M H_L(x, y, t) f(y) \, dy \\ &\geq \frac{C_1}{|B_{\sqrt{t}}(x)|} \int_M \left(1 + \frac{\text{dist}(x, y)}{\sqrt{t}}\right)^{-m/2} \exp\left(-C_2 \frac{\text{dist}(x, y)^2}{t}\right) f(y) \, dy \\ &= \frac{C_1}{|B_{\sqrt{t}}(x)|} \int_0^\infty \left(1 + \frac{r}{\sqrt{t}}\right)^{-m/2} \exp\left(-C_2 \frac{r^2}{t}\right) \int_{\partial B_r(x)} f(y) \, d\mathcal{H}^{m-1}(y) \, dr \\ &\geq \frac{C_1}{|B_{\sqrt{t}}(x)|} \int_0^{\sqrt{t}} \left(1 + \frac{r}{\sqrt{t}}\right)^{-m/2} \exp\left(-C_2 \frac{r^2}{t}\right) \int_{\partial B_r(x)} f(y) \, d\mathcal{H}^{m-1}(y) \, dr \\ &\geq \frac{2^{-m/2} e^{-C_2} C_1}{|B_{\sqrt{t}}(x)|} \int_0^{\sqrt{t}} \int_{\partial B_r(x)} f(y) \, d\mathcal{H}^{m-1}(y) \, dr \\ &= \frac{2^{-m/2} e^{-C_2} C_1}{|B_{\sqrt{t}}(x)|} \int_{B_{\sqrt{t}}(x)} f(y) \, dy. \end{aligned}$$

Since $f \geq 0$, by comparison we have

$$\lim_{t \rightarrow \infty} \frac{1}{|B_{\sqrt{t}}(x)|} \int_{B_{\sqrt{t}}(x)} f(y) \, dy = 0 = \inf_M f$$

as desired. \square

From the above result, we also obtain information on the spherical mean of u . This follows from the next variant of de L'Hôpital's Theorem.

Lemma 23. *Let $h, g \in L^\infty_{\text{loc}}(\mathbb{R}^+)$ satisfy $h \geq 0$, $g > 0$ a.e. and $g \notin L^1(\infty)$. Then,*

$$(48) \quad \text{ess lim inf}_{r \rightarrow \infty} \frac{h(r)}{g(r)} \leq \liminf_{r \rightarrow \infty} \frac{\int_0^r h(t) dt}{\int_0^r g(t) dt}.$$

Proof. Denote by A and B , respectively, the left-hand side and right-hand side of (48). For $A' < A$, fix R_0 such that $h \geq A'g$ a.e. on (R_0, ∞) . Then, for each $r > R_0$,

$$\frac{\int_0^r h(t) dt}{\int_0^r g(t) dt} = \frac{\int_0^{R_0} h(t) dt + \int_{R_0}^r h(t) dt}{\int_0^{R_0} g(t) dt + \int_{R_0}^r g(t) dt} \geq \frac{\int_0^{R_0} h(t) dt + A' \int_{R_0}^r g(t) dt}{\int_0^{R_0} g(t) dt + \int_{R_0}^r g(t) dt}.$$

Since $g \notin L^1(\infty)$, letting $r \rightarrow \infty$ along a sequence realizing B we get $B \geq A'$, and the thesis follows by letting $A' \uparrow A$. \square

Corollary 24. *Let $(M^m, \langle \cdot, \cdot \rangle)$ be a complete (connected) Riemannian manifold with infinite volume. Let $0 \leq f \in L^1_{\text{loc}}(M)$ and $x \in M$ and assume that*

$$\liminf_{R \rightarrow \infty} \frac{1}{|B_R(x)|} \int_{B_R(x)} f(y) dy = \inf_M f.$$

Then

$$\text{ess lim inf}_{R \rightarrow \infty} \frac{1}{|\partial B_R(x)|} \int_{\partial B_R(x)} f(y) d\mathcal{H}^{m-1}(y) dy = \inf_M f.$$

Proof. The functions h and g defined by

$$h(t) = \int_{\partial B_t(x)} f(y) d\mathcal{H}^{m-1}(y) \quad \text{and} \quad g(t) = |\partial B_t(x)| \quad \forall t > 0$$

satisfy the assumptions of the previous Lemma (note that $1/g \in L^\infty_{\text{loc}}(\mathbb{R}^+)$ by [10, Prop. 1.6], since M is non-compact). The thesis follows from the next chain of inequalities:

$$\begin{aligned} \inf_M f &\leq \text{ess lim inf}_{R \rightarrow \infty} \frac{1}{|\partial B_R(x)|} \int_{\partial B_R(x)} f(y) d\mathcal{H}^{m-1}(y) dy \\ &\leq \liminf_{R \rightarrow \infty} \frac{1}{|B_R(x)|} \int_{B_R(x)} f(y) dy \leq \inf_M f. \end{aligned}$$

\square

We are ready to state our second main result of the section, which will enable us to prove the Hessian estimate (6). The argument below seems to be new.

Theorem 25. *Let $(M^m, \langle \cdot, \cdot \rangle)$ be a complete manifold with $\text{Ric} \geq 0$ and let A, L be as in (35)-(36), with A self-adjoint. If $f \in L^\infty(M)$ satisfies $Lf \leq 0$ on M , then for any $x \in M$*

$$(49) \quad \lim_{R \rightarrow \infty} \frac{R^2}{|B_R(x)|} \int_{B_R(x)} Lf(y) dy = 0.$$

Proof. Without loss of generality, we assume $\inf_M f = 0$. Fix $x \in M$. We refer to the proof of Lemma 21 for notation, and in particular, for $t > 0$ and $a > 0$ we define $\varphi_a(y)$ and Ω_a as in (43). As already observed, $\{\Omega_a\}$ is an exhaustion of M , increasing as a decreases. Furthermore, for almost every $a > 0$ the boundary $\partial\Omega_a$ is smooth. From the proof of Lemma 21 we get

$$(50) \quad \int_{\Omega_a} f(y) \partial_t H_L(x, y, t) \, dy \leq \int_{\Omega_a} \varphi_a(y) Lf(y) \, dy.$$

On the other hand, since $f \geq 0$, by (40) and Lemma 19 we can estimate

$$\begin{aligned} & \int_{\Omega_a} f(y) \partial_t H_L(x, y, t) \, dy \\ & \geq -\frac{C_5}{t} \frac{1}{|B_{\sqrt{t}}(x)|} \int_{\Omega_a} f(y) \left(1 + \frac{\text{dist}(x, y)}{\sqrt{t}}\right)^{m/2} \exp\left(-C_6 \frac{\text{dist}(x, y)^2}{t}\right) \, dy. \end{aligned}$$

By (39) and Lemma 19 we also have the bounds

$$\begin{aligned} & \frac{C_1}{|B_{\sqrt{t}}(x)|} \left(1 + \frac{\text{dist}(x, y)}{\sqrt{t}}\right)^{-m/2} \exp\left(-C_2 \frac{\text{dist}(x, y)^2}{t}\right) \\ & \leq H_L(x, y, t) \leq \frac{C_3}{|B_{\sqrt{t}}(x)|} \left(1 + \frac{\text{dist}(x, y)}{\sqrt{t}}\right)^{m/2} \exp\left(-C_4 \frac{\text{dist}(x, y)^2}{t}\right). \end{aligned}$$

Now, fix $k > 1$ large enough so that

$$C_3(1+s)^{m/2} e^{-C_4 s^2} \leq \frac{1}{2} C_1 2^{-m/2} e^{-C_2} \quad \forall s \geq k$$

and pick

$$a = \frac{C_1 2^{-m/2} e^{-C_2}}{2|B_{\sqrt{t}}(x)|}.$$

With this choice, we have

$$\begin{cases} \varphi \leq a & \text{on } M \setminus B_{k\sqrt{t}}(x), \\ \varphi \geq 2a & \text{on } B_{\sqrt{t}}(x), \end{cases}$$

hence $B_{\sqrt{t}}(x) \subseteq \Omega_a \subseteq B_{k\sqrt{t}}(x)$ and $\varphi_a \geq a$ on $B_{\sqrt{t}}(x)$. Thus, using also (50) we can estimate

$$\begin{aligned} 0 & \geq \frac{C_1 2^{-m/2} e^{-C_2}}{2|B_{\sqrt{t}}(x)|} \int_{B_{\sqrt{t}}(x)} Lf(y) \, dy = a \int_{B_{\sqrt{t}}(x)} Lf(y) \, dy \\ & \geq \int_{B_{\sqrt{t}}(x)} \varphi_a(y) Lf(y) \, dy \geq \int_{\Omega_a} \varphi_a(y) Lf(y) \, dy \\ & \geq \int_{\Omega_a} f(y) \partial_t H_L(x, y, t) \, dy \\ & \geq -\frac{C_5}{t|B_{\sqrt{t}}(x)|} \int_{\Omega_a} f(y) \left(1 + \frac{\text{dist}(x, y)}{\sqrt{t}}\right)^{m/2} \exp\left(-C_6 \frac{\text{dist}(x, y)^2}{t}\right) \, dy \\ & \geq -\frac{C_7}{t|B_{\sqrt{t}}(x)|} \int_{B_{k\sqrt{t}}(x)} f(y) \, dy \end{aligned}$$

where

$$C_7 = C_5 \sup \left\{ (1+s)^{m/2} e^{-C_6 s^2} : s > 0 \right\} < \infty.$$

Summing up, there exists a constant $C > 0$, depending only on C_i , $1 \leq i \leq 7$, such that

$$0 \geq \frac{t}{|B_{\sqrt{t}}(x)|} \int_{B_{\sqrt{t}}(x)} Lf(y) \, dy \geq -\frac{C}{|B_{\sqrt{t}}(x)|} \int_{B_{k\sqrt{t}}(x)} f(y) \, dy.$$

Since $f \geq 0$, by Bishop-Gromov theorem we also have

$$0 \geq \frac{t}{|B_{\sqrt{t}}(x)|} \int_{B_{\sqrt{t}}(x)} Lf(y) \, dy \geq -\frac{Ck^m}{|B_{k\sqrt{t}}(x)|} \int_{B_{k\sqrt{t}}(x)} f(y) \, dy.$$

By Proposition 22 we have that the RHS of this inequality converges to $\inf_M f = 0$ as $t \rightarrow \infty$, and the conclusion follows. \square

6. PROOF OF THEOREM 5, (ii)

Combining Corollary 16, Proposition 22 and Theorem 25, we get

Proposition 26. *Let (M^m, σ) be a complete Riemannian manifold with $\text{Ric} \geq 0$ and*

$$\text{Ric}^{(\ell)}(\nabla r) \geq -\frac{\bar{\kappa}^2}{1+r^2}, \quad \ell = \max\{1, m-2\},$$

for some $\bar{\kappa} \in \mathbb{R}_0^+$ and where r is the distance from a fixed origin. Let $u \in C^\infty(M)$ be a non-constant solution of (MSE) which grows at most linearly on one side. Then, for each $x \in M$,

$$(51) \quad \lim_{R \rightarrow \infty} \frac{1}{|B_R(x)|} \int_{B_R(x)} |Du|^2 dx = \sup_M |Du|^2,$$

$$(52) \quad \lim_{R \rightarrow \infty} \frac{R^2}{|B_R(x)|} \int_{B_R(x)} |D^2u|^2 dx = 0.$$

Proof. Because of Corollary 16, in our assumptions $|Du| \in L^\infty(M)$, hence by (10) the operator

$$L\phi \doteq W\Delta_g\phi = \text{div}(Wg^{ij}\phi_i\partial_{x_j})$$

is uniformly elliptic on M . By the Jacobi equation, $f = 1/W$ is a non-negative solution of $Lf \leq -\|\text{II}\|^2 \leq 0$, and therefore $-W^2 \in L^\infty(M)$ satisfies $L(-W^2) \leq 0$. Applying Proposition 22 to $-W^2$ and Theorem 25 to f we deduce

$$(53) \quad \lim_{R \rightarrow \infty} \frac{1}{|B_R(x)|} \int_{B_R(x)} W^2 dx = \sup_M W^2$$

$$(54) \quad \limsup_{R \rightarrow \infty} \frac{R^2}{|B_R(x)|} \int_{B_R(x)} \|\text{II}\|^2 dx \leq -\lim_{R \rightarrow \infty} \frac{R^2}{|B_R(x)|} \int_{B_R(x)} Lf dx = 0.$$

From (53) we readily deduce (51). On the other hand, note that

$$\begin{aligned} \|\text{II}\|^2 &= W^{-2}g^{ik}u_{kj}g^{jl}u_{li} \\ &= W^{-2} \left\{ |D^2u|^2 - 2 \left| D^2u \left(\frac{Du}{W}, \cdot \right) \right|^2 + \left[D^2u \left(\frac{Du}{W}, \frac{Du}{W} \right) \right]^2 \right\} \end{aligned}$$

If $du(x) = 0$, then $\|\text{II}\|^2 \geq W^{-2}|D^2u|^2$. Otherwise, let $e_1 = Du/|Du|$ and choose a local orthonormal frame $\{e_\alpha\}$ for e_1^\perp around x , where $2 \leq \alpha \leq m$. Then,

$$\begin{aligned} &|D^2u|^2 - 2 \left| D^2u \left(\frac{Du}{W}, \cdot \right) \right|^2 + \left[D^2u \left(\frac{Du}{W}, \frac{Du}{W} \right) \right]^2 \\ &= \sum_{\alpha, \beta} u_{\alpha\beta}^2 + 2 \sum_{\alpha} u_{1\alpha}^2 + u_{11}^2 - 2 \frac{W^2-1}{W^2} \sum_j u_{1j}^2 + \frac{(W^2-1)^2}{W^4} u_{11}^2 \\ &= \sum_{\alpha, \beta} u_{\alpha\beta}^2 + \frac{2}{W^2} \sum_{\alpha} u_{1\alpha}^2 + \frac{1}{W^4} u_{11}^2 \geq W^{-4} |D^2u|^2 \end{aligned}$$

Summarizing, we have $\|\mathbb{II}\|^2 \geq W^{-4}|D^2u|^2$, thus from the boundedness of W and from (54) we conclude (52). \square

We now conclude the proof of Theorem 5 with a blow-down procedure, for which we use some basic convergence results in the theory of limit spaces and nonsmooth spaces with Ricci curvature bounded below. All the tools needed herein can be found in [42, 3, 4].

Fix $o \in M$, and write $B_R = B_R(o)$. Because of Corollary 16 and Proposition 26,

$$(55) \quad \lim_{R \rightarrow \infty} \frac{1}{|B_R|} \int_{B_R} |Du|^2 dx = \sup_M |Du|^2,$$

$$(56) \quad \lim_{R \rightarrow \infty} \frac{R^2}{|B_R|} \int_{B_R} |D^2u|^2 dx = 0.$$

Consider a tangent cone at infinity X_∞ for M based at o . By statement (2.1) in [28], the limit space X_∞ also supports a Borel measure \mathfrak{m}_∞ such that, up to a subsequence,

$$(57) \quad (M, \lambda_n^{-1} \text{dist}_\sigma, \lambda_n^{-m} dx, o) \xrightarrow{\text{pmGH}} (M_\infty, d_\infty, \mathfrak{m}_\infty, o_\infty)$$

in the pointed-measured-Gromov-Hausdorff (pmGH) sense. For the precise definition of pGH and pmGH convergence we refer to [39]. Here, $\{\lambda_n\} \subset \mathbb{R}^+$, $\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$, and $\lambda_n^{-1} \text{dist}_\sigma$ is the distance function induced by the rescaled metric $\sigma_n \doteq \lambda_n^{-2} \sigma$. Denote with D_n and dx_n the induced connection and volume measure, and B_R^n the metric balls centered at o in (M, σ_n) . Therefore, $B_R^n = B_{\lambda_n R}$. Define $u_n = u/\lambda_n$. Then,

$$(58) \quad |D_n u_n|_{\sigma_n} = |Du|, \quad |D_n^2 u_n|_{\sigma_n} = \lambda_n |D^2 u|$$

and therefore, by Arzelá-Ascoli Theorem, up to subsequences $u_n \rightarrow u_\infty \in \text{Lip}(M_\infty)$ locally uniformly, hence $u_n \rightarrow u_\infty$ strongly in L^2 on $B_R^\infty = B_R^{\text{d}_\infty}(x_\infty)$, that is,

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_{B_R^n} |u_n|^2 dx_n &= \int_{B_R^\infty} |u_\infty|^2 d\mathfrak{m}_\infty, \\ \lim_{n \rightarrow \infty} \int_{B_R^n} u_n \varphi dx_n &= \int_{B_R^\infty} u_\infty \varphi d\mathfrak{m}_\infty \end{aligned}$$

for each φ bounded and continuous on a metric space Z in which (B_R^n, d_n) and (B_R^∞, d_∞) are isometrically embedded and converge in Hausdorff sense, with $o_n \rightarrow o_\infty$ and o_n the center of B_R^n . From

$$W_n \doteq \sqrt{1 + |D_n u_n|_{\sigma_n}^2} = \sqrt{1 + |Du|^2} = W$$

Scaling (55) and (56) we therefore get, for each fixed $R > 0$

$$(59) \quad \lim_{n \rightarrow \infty} \frac{1}{|B_R^n|_{\sigma_n}} \int_{B_R^n} |D_n u_n|_{\sigma_n}^2 dx_n = \sup_M |Du|^2,$$

$$(60) \quad \lim_{n \rightarrow \infty} \frac{R^2}{|B_R^n|_{\sigma_n}} \int_{B_R^n} |D_n^2 u_n|_{\sigma_n}^2 dx_n = 0.$$

In particular, from Newton's inequality $|\Delta_n u_n|^2 \leq n |D_n^2 u_n|_{\sigma_n}^2$ and Bishop-Gromov's Theorem, $|B_R^n|_{\sigma_n} \leq \omega_{m-1} R^m / m$ we deduce

$$(61) \quad \int_{B_R^n} |\Delta_n u_n|^2 dx_n \leq \frac{\omega_{m-1} R^m}{|B_R^n|_{\sigma_n}} \int_{B_R^n} |D_n^2 u_n|_{\sigma_n}^2 dx_n \rightarrow 0$$

as $n \rightarrow \infty$, and therefore

$$(62) \quad \begin{aligned} \int_{B_R^n} \varphi \Delta_n u_n dx_n &\leq \left(\int_{B_R^n} \varphi^2 dx_n \right)^{\frac{1}{2}} \left(\int_{B_R^n} |\Delta_n u_n|^2 dx_n \right)^{\frac{1}{2}} \\ &\leq \max |\varphi| \left[\frac{\omega_{m-1} R^m}{m} \right]^{\frac{1}{2}} \left(\int_{B_R^n} |\Delta_n u_n|^2 dx_n \right)^{\frac{1}{2}} \rightarrow 0. \end{aligned}$$

By (61) and (62), $\Delta_n u_n \rightarrow 0$ strongly in L^2 . Combining $u_n \rightarrow u_\infty$ strongly in L^2 with

$$\sup_n \left(\int_{B_R^n} [|u_n|^2 + |D_n u_n|_{\sigma_n}^2 + (\Delta_n u_n)^2] dx_n \right) < \infty$$

we infer by [3, Thm. 4.4] that

- (i) $u_\infty \in \mathcal{D}(\Delta, B_R^\infty)$, the domain of the Laplacian on B_R^∞ ;
- (ii) $\Delta_n u_n \rightarrow \Delta u_\infty$ on B_r^n weakly in L^2 , so in particular $\Delta u_\infty = 0$;
- (iii) $|D_n u_n|_{\sigma_n}^2 \rightarrow |D_\infty u_\infty|_\infty^2$ in L^1 -strongly in B_r^n , for each $r < R$;

In particular, setting $P \doteq \sup_M |Du|^2$, from (58) and (59) we get

$$\lim_{n \rightarrow \infty} \int_{B_R^n} ||D_n u_n|_{\sigma_n}^2 - P| dx_n \leq \frac{\omega_{m-1} R^m}{|B_R^n|_{\sigma_n}} \int_{B_R^n} (P - |D_n u_n|_{\sigma_n}^2) dx_n = 0$$

Using (iii) and [16, Prop. 1.27 (i)] (cf. also [4]), we therefore deduce $|D_n u_n|_{\sigma_n}^2 - P \rightarrow |D_\infty u_\infty|_\infty^2 - P$ strongly in L^1 on B_r^∞ for each $r < R$, and thus

$$0 = \lim_{n \rightarrow \infty} \int_{B_r^n} ||D_n u_n|_{\sigma_n}^2 - P| dx_n = \int_{B_r^\infty} ||D_\infty u_\infty|_\infty^2 - P| d\mathbf{m}_\infty.$$

Concluding, u_∞ solves

$$\Delta u_\infty = 0, \quad |D_\infty u_\infty|^2 = P \neq 0$$

on the RCD(0, m) space $(M_\infty, d_\infty, \mathbf{m}_\infty, x_\infty)$. Bochner inequality (see [42, Thm. 1.4]) guarantees that $|D^2 u_\infty| \equiv 0$ on M_∞ . One concludes that $M_\infty = N \times \mathbb{R}$ by using [6, Lem. 1.21].

7. PROOF OF THEOREM 10

If M is parabolic, clearly the result follows from Theorem 5. If M is non-parabolic, the argument goes as in [31, Theorem 3.6], so we only sketch the main steps. In our assumptions, by Corollary 16, $|Du| \in L^\infty(M)$, hence $L = W\Delta_g$ is uniformly elliptic. The Harnack inequality in [58] together with (4) imply that $|u(x)| = o(r(x))$ as x diverges. By a standard cutoff argument using $Lu = 0$, the next Caccioppoli inequality holds: for each $\varphi \in \text{Lip}_c(M)$

$$(63) \quad \int_M \varphi^2 |Du|^2 dx \leq 4\alpha^2 \int_M u^2 |D\varphi|^2 dx.$$

In particular, having fixed $\varepsilon > 0$, by condition $u = o(r)$ we can also fix $R_0 = R_0(\varepsilon) > 0$ such that for every $R \geq R_0$ we have $u^2 \leq \varepsilon R^2$ on B_{2R} . Considering

the Lipschitz cutoff function φ which is 1 on B_R , 0 outside of B_{2R} and satisfies $|D\varphi| \leq 1/R$, we get

$$\int_{B_R} |Du|^2 dx \leq \frac{4\alpha^2}{R^2} \int_{B_{2R} \setminus B_R} u^2 dx \leq \varepsilon |B_{2R}| \leq C\varepsilon |B_R|$$

for every $R \geq R_0$, where we used the doubling property on M coming from condition $\text{Ric} \geq 0$. From (51) we finally infer

$$\sup_M |Du|^2 = \lim_{R \rightarrow \infty} \frac{1}{|B_R|} \int_{B_R} |Du|^2 dx \leq C\varepsilon,$$

and the thesis follows by letting $\varepsilon \rightarrow 0$.

8. PROOF OF COROLLARY 9

By Theorem 5, $|Du| \in L^\infty(M)$ and any tangent cone at infinity of M splits off a line. It is a general fact that, if $\text{Sec} \geq 0$, a tangent cone splits if and only if M itself splits. A proof of this result can be found in [5, Thm. 4.6]. Therefore, it remains to prove that u only depends on the coordinate of a split line. Write $M = N^{m-1} \times \mathbb{R}$ with coordinates (y_1, s_1) , for some complete manifold N^{m-1} with $\text{Sec} \geq 0$, and consider the function $v_1 = \sigma(Du, \partial_{s_1})$, which by (14) satisfies $Lv_1 = 0$ on M , where we set

$$L\phi \doteq W^{-1} \mathcal{L}_W \phi = \text{div} (W^{-1} g^{ij} \phi_i \partial_{x_j}).$$

Our gradient estimate guarantees that v_1 is bounded and that L is uniformly elliptic on M , and therefore, by [58, Theorem 7.4] we deduce that v_1 is constant on M . Hence,

$$u(y_1, s_1) = a_1 s_1 + b_1 + u_2(y_1) \sqrt{1 + a_1^2},$$

for some smooth function $u_2 : N^{m-1} \rightarrow \mathbb{R}$ and some $a_1, b_1 \in \mathbb{R}$. One easily checks that u_2 solves (MSE) on N^{m-1} . Since v_2 has at most linear growth on one side, and N^{m-1} has non-negative sectional curvature, by the first part of the proof we deduce that either u_2 is constant or that $N^{m-1} = N^{m-2} \times \mathbb{R}$ and $u_2(y_2, s_2) = a_2 s_2 + b_2 + u_3(y_2) \sqrt{1 + a_2^2}$. Iterating, we can write $M = N^{m-k} \times \mathbb{R}^k$ for some $k \in \{1, \dots, m-2\}$ and for some complete manifold N^{m-k} with $\text{Sec} \geq 0$, and

$$u(z, (s_1, \dots, s_k)) = \sum_{j=1}^k a_j s_j + b + v_{k+1}(z) \sqrt{1 + a_k^2}$$

for some $a_i, b \in \mathbb{R}$ and $v_{k+1} : N^{m-k} \rightarrow \mathbb{R}$. Indeed, we can continue the iteration procedure up until either v_{k+1} is constant, or $k = m-2$ and v_{m-1} is non-constant. In the latter case, observe that N^2 is a complete surface with $\text{Sec} \geq 0$, hence N^2 is parabolic. Being v_{m-1} non-constant, both N^2 and v_{m-1} split as indicated in Theorem 5, (i). Summarizing, in each case we can conclude that $M = N^{m-k} \times \mathbb{R}^k$ for some $k \in \{1, \dots, m-1\}$, and that

$$(64) \quad u(z, (s_1, \dots, s_k)) = \sum_{j=1}^k a_j s_j + b$$

for some $b \in \mathbb{R}$, as required. It is therefore sufficient to consider the splitting $\mathbb{R}^k = \mathbb{R}^{k-1} \times \mathbb{R}$ along a line in direction (a_1, \dots, a_k) to get the desired splitting $M = N \times \mathbb{R}$ of M in such a way that $u(y, s) = as + b$.

9. PROOF OF PROPOSITION 8

The following example is essentially that in [45, p. 913]. Let $m \geq 4$. We consider a manifold (P^{m-2}, h) and smooth functions $f, \eta \in C^\infty(\mathbb{R}^+)$ to be chosen later, and define the following metric on $M \doteq \mathbb{R} \times \mathbb{R}^+ \times P$:

$$\sigma = f(r)^2 dt^2 + dr^2 + \eta(r)^2 h.$$

To compute the curvatures of M , we use the index agreement $1 \leq a, b, c, l \leq m$, $3 \leq \alpha, \beta, \gamma, \delta \leq m$. Let $\{\theta^\alpha\}$ be a local orthonormal coframe on P , with associated connection forms ω_β^α obeying the structure equations

$$\begin{cases} d\theta^\alpha = -\omega_\beta^\alpha \wedge \theta^\beta \\ \omega_\beta^\alpha = -\omega_\alpha^\beta \end{cases}$$

and related curvature forms $\Theta_\beta^\alpha = d\omega_\beta^\alpha + \omega_\gamma^\alpha \wedge \omega_\beta^\gamma$. Then, a local orthonormal coframe $\{\bar{\theta}^a\}$ on M is given by

$$\bar{\theta}^1 = f dt, \quad \bar{\theta}^2 = dr, \quad \bar{\theta}^\alpha = \eta \theta^\alpha,$$

where, as usual, pull-backs to M via the canonical projections onto \mathbb{R}, \mathbb{R}^+ and P are implicit. Differentiating, one checks that the forms

$$\bar{\omega}_1^\alpha = 0, \quad \bar{\omega}_2^\alpha = \frac{\eta'}{\eta} \bar{\theta}^\alpha, \quad \bar{\omega}_\beta^\alpha = \omega_\beta^\alpha, \quad \bar{\omega}_1^2 = -\frac{f'}{f} \bar{\theta}^1.$$

satisfy the structure equations on M for the coframe $\{\bar{\theta}^a\}$, hence they are the connection forms of $\{\bar{\theta}^a\}$. The associated curvature forms $\bar{\Theta}_b^a = d\bar{\omega}_b^a + \bar{\omega}_c^a \wedge \bar{\omega}_b^c$ are therefore

$$(65) \quad \begin{aligned} \bar{\Theta}_1^\alpha &= -\frac{\eta' f'}{\eta f} \bar{\theta}^\alpha \wedge \bar{\theta}^1, & \bar{\Theta}_2^\alpha &= \frac{\eta''}{\eta} \bar{\theta}^2 \wedge \bar{\theta}^\alpha, \\ \bar{\Theta}_\beta^\alpha &= \Theta_\beta^\alpha - \left(\frac{\eta'}{\eta}\right)^2 \bar{\theta}^\alpha \wedge \bar{\theta}^\beta & \bar{\Theta}_1^2 &= -\frac{f''}{f} \bar{\theta}^2 \wedge \bar{\theta}^1, \end{aligned}$$

The components $R_{\beta\gamma\delta}^\alpha$ and \bar{R}_{bcl}^a of the (3, 1) curvature tensors of, respectively, P and M , are given by the identitis

$$\Theta_\beta^\alpha = \frac{1}{2} R_{\beta\gamma\delta}^\alpha \theta^\gamma \wedge \theta^\delta, \quad \bar{\Theta}_b^a = \frac{1}{2} \bar{R}_{bcl}^a \bar{\theta}^c \wedge \bar{\theta}^l,$$

and thus, from (65), we deduce

$$(66) \quad \begin{aligned} 0 &= \bar{R}_{12\alpha}^2 = \bar{R}_{1\alpha 1}^2 = \bar{R}_{1\alpha\beta}^2 = \bar{R}_{12\beta}^\alpha = \bar{R}_{1\gamma\delta}^\alpha = \bar{R}_{2\gamma\delta}^\alpha \\ \bar{R}_{121}^2 &= -\frac{f''}{f}, \quad \bar{R}_{1\beta 1}^\alpha = -\frac{\eta' f'}{\eta f} \delta_\beta^\alpha, \quad \bar{R}_{2\beta 2}^\alpha = -\frac{\eta''}{\eta} \delta_\beta^\alpha, \\ \bar{R}_{\beta\gamma\delta}^\alpha &= \frac{1}{\eta^2} R_{\beta\gamma\delta}^\alpha - \left(\frac{\eta'}{\eta}\right)^2 [\delta_\gamma^\alpha \delta_{\beta\delta} - \delta_\delta^\alpha \delta_{\beta\gamma}]. \end{aligned}$$

Assume that (P, h) is the round sphere with curvature 1, and let $\{e_\alpha\}$ and $\{\bar{e}_a\}$ be, respectively, the dual frames of $\{\theta^\alpha\}$ and $\{\bar{\theta}^a\}$. From (66) we deduce that the curvature operator is diagonalized by the simple planes $\{\bar{e}_a \wedge \bar{e}_b\}$, so for $m \geq 4$ we get

$$|\overline{\text{Sec}}(\pi)| \leq \max \left\{ \left| \frac{f''}{f} \right|, \left| \frac{\eta' f'}{\eta f} \right|, \left| \frac{1 - (\eta')^2}{\eta^2} \right| \right\}.$$

In [45], the authors chose the following functions f, η : given $\alpha, \beta \in (0, 1)$ such that $m - 1 - \beta > 2 + \alpha$, let $0 < \zeta_1, \zeta_2 \in C^\infty(\mathbb{R}^+)$ satisfy

$$\zeta_1(t) = \begin{cases} t & \text{if } t \in (0, 1] \\ t^{-1-\alpha} & \text{if } t \in [2, \infty), \end{cases} \quad \zeta_2(t) = \int_t^\infty \zeta_1(s) ds.$$

Then, for $b, c \in \mathbb{R}^+$ they defined

$$\eta(r) = \frac{1}{2}r + \frac{1}{2\zeta_2(0)} \int_0^r \zeta_2(s) ds, \quad f(r) = (b + r^2)^{\frac{\beta+3-m}{2}} + c.$$

Note that with such a choice the metric extends in a C^2 way at $r = 0$, giving rise to a complete manifold. Since the curvature operator is diagonalized by $\{\bar{e}_a \wedge \bar{e}_b\}$,

$$(67) \quad \begin{aligned} \overline{\text{Ric}}^{(2)} \geq & \min \left\{ -\frac{f''}{f} + \frac{1 - (\eta')^2}{\eta^2}, -\frac{f''}{f} - \frac{\eta' f'}{\eta f}, -\frac{f''}{f} - \frac{\eta''}{\eta}, \right. \\ & \frac{1 - (\eta')^2}{\eta^2} - \frac{\eta' f'}{\eta f}, \frac{1 - (\eta')^2}{\eta^2} - \frac{\eta''}{\eta}, -\frac{\eta''}{\eta} - \frac{\eta' f'}{\eta f}, -2\frac{\eta' f'}{\eta f}, \\ & \left. +2\frac{1 - (\eta')^2}{\eta^2}, -2\frac{\eta''}{\eta} \right\}. \end{aligned}$$

By the expression of η, f , the four terms in the second line of (67) are positive, and it is easy to see that, when b, c are large enough, the three terms in the first line are positive as well. The two terms in the third line are positive except at $r = 0$. Whence, $\overline{\text{Ric}}^{(2)} \geq 0$, and moreover $|\overline{\text{Sec}}| \leq \bar{\kappa}^2$ holds for a suitable $\bar{\kappa} > 0$. Moreover, from the fact that $\overline{\text{Ric}}$ is diagonal in the basis $\{\bar{e}_a\}$ with

$$\begin{aligned} \overline{\text{Ric}}_{11} &= -\frac{f''}{f} - (m-3)\frac{\eta' f'}{\eta f}, & \overline{\text{Ric}}_{22} &= -\frac{f''}{f} - (m-3)\frac{\eta''}{\eta}, \\ \overline{\text{Ric}}_{\alpha\beta} &= \left[-\frac{\eta' f'}{\eta f} - \frac{\eta''}{\eta} + (m-3)\frac{1 - (\eta')^2}{\eta^2} \right] \delta_{\alpha\beta} \end{aligned}$$

we deduce that $\text{Ric} > 0$ if b, c are chosen large enough. To construct linearly growing minimal graphs, consider a function $u : M \rightarrow \mathbb{R}$ of the coordinate t alone. It follows that $du = u_a \bar{\theta}^a$ with $u_1 = (\partial_t u)/f$ and $u_a = 0$ for $a \geq 2$. The components of the Hessian D^2u obey the relation

$$u_{ab} \bar{\theta}^b = du_a - u_c \bar{\omega}_a^c,$$

and from the expression of $\bar{\omega}_a^c$ we get

$$\begin{aligned} u_{11} &= \frac{\partial_t^2 u}{f^2}, & u_{21} &= -\frac{f'}{f^2} \partial_t u, \\ u_{1\alpha} &= u_{22} = u_{2\alpha} = u_{\alpha\beta} = 0, \end{aligned}$$

In particular, setting $W = \sqrt{1 + |Du|^2} = \sqrt{1 + u_1^2}$,

$$\text{div} \left(\frac{Du}{\sqrt{1 + |Du|^2}} \right) = \frac{\Delta u}{W} - \frac{D^2u(Du, Du)}{W^3} = \frac{\partial_t^2 u}{f^2 W} - \frac{(\partial_t^2 u) u_1^2}{f^2 W^3} = \frac{\partial_t^2 u}{f^2 W^3}.$$

It follows that any affine function $u(t) = at + b$ gives rise to a minimal graph. Furthermore, $|Du| = a/f$ is bounded on M since f is bounded below by a positive constant, thus u has at most linear growth.

APPENDIX A.

Let M be a complete Riemannian manifold with non-negative Ricci curvature, $\dim M = m$, and let A, L, H_L be as in section 5. In this Appendix, we discuss the two-sided bound in (39) for H_L . While the upper bound is shown in [58], the argument for the lower bound is merely indicated with no proof. The approach relies on the following parabolic Harnack inequality in [58, Corollary 5.4]: given $p \in M$, $R > 0$, $T > 0$ and $\delta \in (0, 1)$, if u is a positive solution of $\partial_t u = Lu$ on $B_R(p) \times (0, T)$, then

$$(68) \quad \log \left(\frac{u(t, y)}{u(s, x)} \right) \leq C \left(\frac{\text{dist}(x, y)^2}{s - t} + \left(\frac{1}{R^2} + \frac{1}{t} \right) (s - t) + 1 \right)$$

for every $x, y \in B_{\delta R}(p)$ and $0 < t < s < T$, with $C = C(m, \delta, \alpha) > 0$. A note of warning: in [58, Corollary 5.4], the final $+1$ in brackets in (68) is missing. However, necessity of this correction becomes apparent by direct inspection of Moser's original proof, [55, pages 110–112], in Euclidean setting (the analogue of (68) is [55, Formula (1.5)]). For the reader's convenience, we give a proof that the lower bound in (39) follows from the upper one coupled with (68), along the lines of the argument developed by Aronson and Serrin [7] in the Euclidean case. A few observations are in order.

First, in view of Lemma 19 the upper bound in (39) implies

$$(69) \quad H_L(x, y, t) \leq \frac{C'_3}{|B_{\sqrt{t}}(x)|} \exp \left(-C'_4 \frac{\text{dist}(x, y)^2}{t} \right) \quad \forall x, y \in M, t > 0$$

with $C'_3, C'_4 > 0$ depending only on m and α (the ellipticity constant of A). Secondly, the differential Harnack inequality (68) applied to $u = H_L(x, \cdot, \cdot)$ yields

$$(70) \quad H_L(x, y_1, t_1) \leq H_L(x, y_2, t_2) \exp \left(C \frac{\text{dist}(y_1, y_2)^2}{t_2 - t_1} + C \frac{t_2}{t_1} \right)$$

for every $y_1, y_2 \in M$ and $0 < t_1 < t_2 < \infty$, with $C = C(m, \alpha) > 0$. Lastly, note that if we have the validity of a lower bound of the form

$$(71) \quad H_L(x, y, t) \geq \frac{C'_1}{|B_{\sqrt{t}}(x)|} \exp \left(-C'_2 \frac{\text{dist}(x, y)^2}{t} \right) \quad \forall x, y \in M, t > 0$$

with $C'_1, C'_2 > 0$ depending only on m and α , then, again by Lemma 19, a lower bound as that in (39) holds for suitable constants $C_1 \in (0, C'_1)$ and $C_2 > C'_2$ depending only on C'_1, C'_2 and m . Hence, we limit ourselves to the proof that (71) follows from (69) and (70) under the assumption $\text{Ric} \geq 0$.

Fix a constant $c_0 > 2$ such that

$$(72) \quad \gamma \doteq mC'_3 \int_{\sqrt{c_0}}^{+\infty} s^{m-1} e^{-C'_4 s^2} ds < 1.$$

Let $(x, y, t) \in M \times M \times \mathbb{R}^+$ be given. By (70) we have

$$(73) \quad H_L(x, x, t/2) \leq H_L(x, y, t) \exp \left(2C \frac{\text{dist}(x, y)^2}{t} + 2C \right)$$

and also

$$H_L(x, z, t/c_0) \leq H_L(x, x, t/2) \exp \left(c_0^* C \frac{\text{dist}(x, z)^2}{t} + \frac{c_0}{2} C \right)$$

for every $z \in M$, with $c_0^* = \frac{2c_0}{c_0-2} = \left(\frac{1}{2} - \frac{1}{c_0}\right)^{-1}$. Integrating on $B_{\sqrt{t}}(x)$ we get

$$(74) \quad \int_{B_{\sqrt{t}}(x)} H_L(x, z, t/c_0) dz \leq e^{(c_0^*+c_0/2)C} |B_{\sqrt{t}}(x)| H_L(x, x, t/2).$$

Putting together (73) and (74) we obtain

$$(75) \quad H_L(x, y, t) \geq \frac{e^{-(2+c_0/2+c_0^*)C}}{|B_{\sqrt{t}}(x)|} \exp\left(-2C \frac{\text{dist}(x, y)^2}{t}\right) \int_{B_{\sqrt{t}}(x)} H_L(x, z, t/c_0) dz.$$

From the upper bound (69) and the co-area formula we have

$$\begin{aligned} \int_{M \setminus B_{\sqrt{t}}(x)} H_L(x, z, t/c_0) dz &\leq C'_3 \int_{\sqrt{t}}^{\infty} \frac{|\partial B_r(x)|}{|B_{\sqrt{t/c_0}}(x)|} \exp\left(-c_0 C'_4 \frac{r^2}{t}\right) dr \\ &= C'_3 \int_{\sqrt{c_0}}^{\infty} \frac{\sqrt{t/c_0} |\partial B_{s\sqrt{t/c_0}}(x)|}{|B_{\sqrt{t/c_0}}(x)|} e^{-C'_4 s^2} ds \end{aligned}$$

where we have changed variable $s = r\sqrt{c_0}/t$. Since $\text{Ric} \geq 0$ we have

$$\frac{\sqrt{t/c_0} |\partial B_{s\sqrt{t/c_0}}(x)|}{|B_{\sqrt{t/c_0}}(x)|} \leq s^{m-1} \frac{\sqrt{t/c_0} |\partial B_{\sqrt{t/c_0}}(x)|}{|B_{\sqrt{t/c_0}}(x)|} \leq m s^{m-1},$$

where the first inequality follows by Bishop-Gromov theorem and the second from the inequality $R|\partial B_R(x)| \leq m|B_R(x)|$, holding for every $R > 0$ and for any base point x on a Riemannian manifold with $\text{Ric} \geq 0$, see for instance [47, Formula (19)]. Substituting in the above estimate and recalling (72) and Lemma 20 we get

$$\int_{B_{\sqrt{t}}(x)} H_L(x, z, t/c_0) dz = 1 - \int_{M \setminus B_{\sqrt{t}}(x)} H_L(x, z, t/c_0) dz \geq 1 - \gamma > 0$$

and from (75) we obtain

$$H_L(x, y, t) \geq \frac{C'_1}{|B_{\sqrt{t}}(x)|} \exp\left(-2C \frac{\text{dist}(x, y)^2}{t}\right)$$

where $C'_1 = (1 - \gamma)e^{-(2+c_0/2+c_0^*)C} > 0$ only depends on m and α . This proves (71).

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