

Using AS SCIVR to understand Proline vibrational spectrum

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Abstract

Semiclassical dynamics is able to reproduce quantum effects from classical dynamics runs, allowing the vibrational study of very large dimensional systems.[1] Adiabatic switching has already proven capable of improving precision and accuracy of semiclassical results of challenging model potentials and small molecular systems.[2] I extended the technique to larger molecular systems, whose semiclassical spectrum is usually collected by means of a single run evolved with ab initio “on-the-fly” calculations. This application has been benchmarked on small molecules and then tested on glycine, improving the pre-existing SC calculations. Finally, this new approach has permitted a preliminary study of the vibrational spectrum of the 17-atom proline, a still open problem in theoretical and experimental chemistry.[3, 4]

Starting Point: MC SCIVR

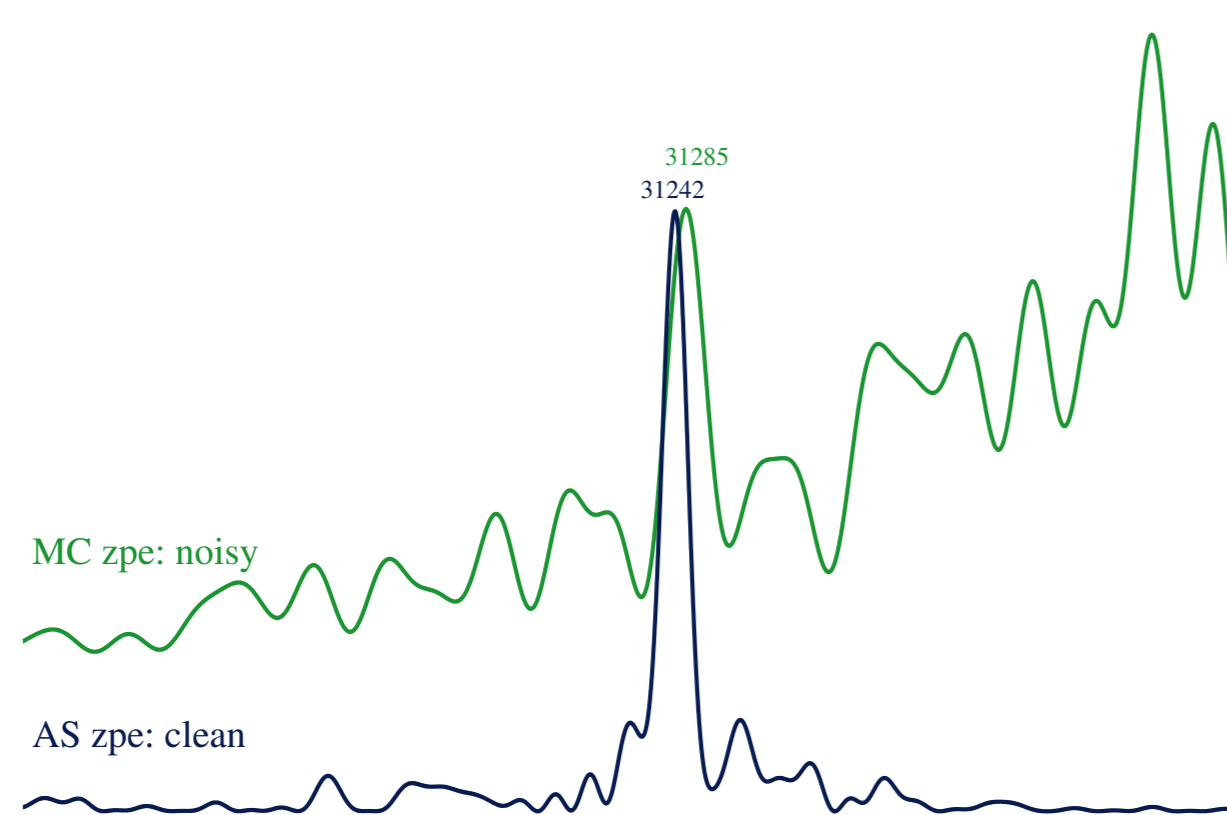
$$I_{\text{qct}}(E) = \frac{1}{2T} \left| \int_0^T dt e^{iEt/\hbar} p(t) \right|^2 \quad (1)$$

$$I_{\text{ta}}(E) = \left(\frac{1}{2\pi\hbar} \right)^{N_v} \iint d\mathbf{p}_0 d\mathbf{q}_0 \frac{1}{2\pi\hbar T} \left| \int_0^T dt e^{i[S_t(\mathbf{p}_t, \mathbf{q}_t) + Et + \phi_t(\mathbf{p}_t, \mathbf{q}_t)]} \langle \psi | g_t(\mathbf{p}_t, \mathbf{q}_t) \rangle \right|^2 \quad (2)$$

$$I_{\text{mc}}(E; \mathbf{p}_{\text{mc}}, \mathbf{q}_{\text{mc}}) = \left(\frac{1}{2\pi\hbar} \right)^{N_v} \frac{1}{2\pi\hbar T} \left| \int_0^T dt e^{i[S_t(\mathbf{p}'_t, \mathbf{q}'_t) + Et + \phi_t(\mathbf{p}'_t, \mathbf{q}'_t)]} \langle \psi(\mathbf{p}_{\text{eq}}, \mathbf{q}_{\text{eq}}) | g_t(\mathbf{p}'_t, \mathbf{q}'_t) \rangle \right|^2 \quad (3)$$

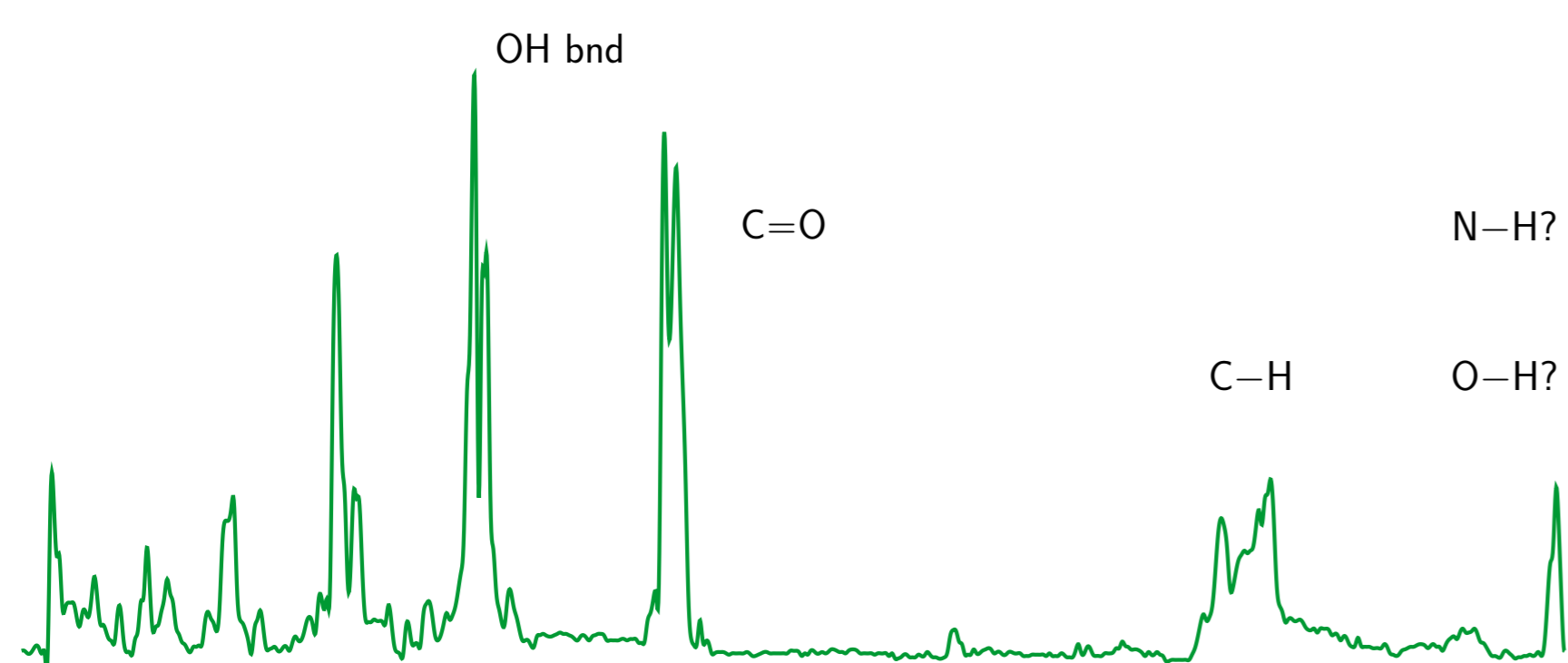
The need for AS SCIVR

AS trajectory collects more signal than the MC one

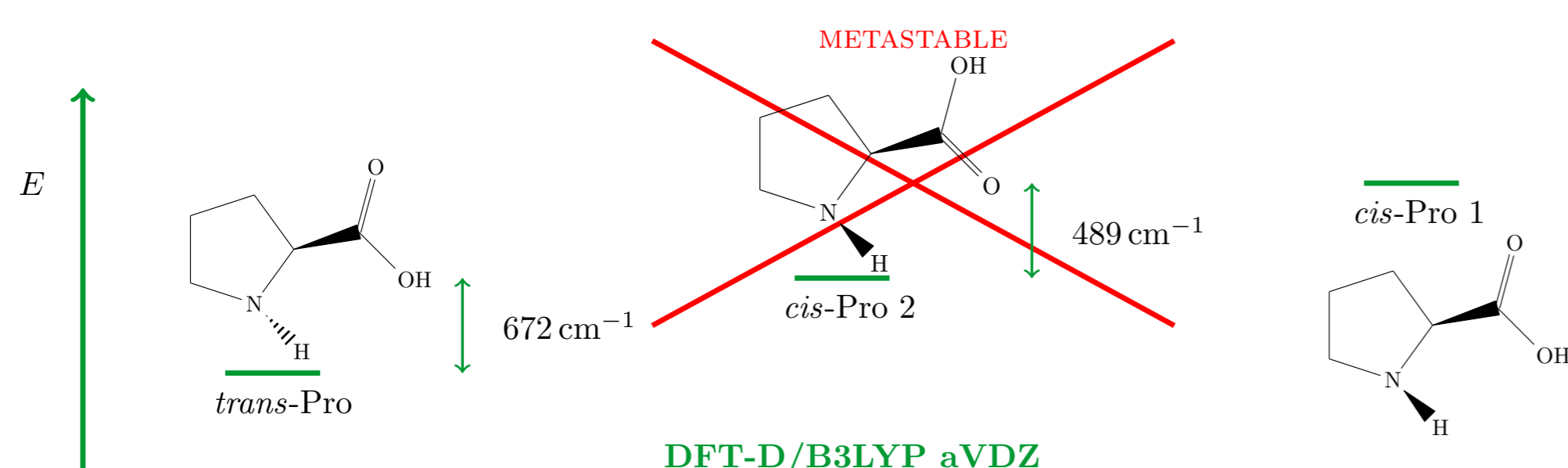


Proline: the experiment

The experimental spectrum of Proline from [5]: still an open problem



Proline has a lot of similar conformers: which ones are in the spectrum?



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The novelty: of AS SCIVR

AS dynamics

- Velocity Verlet integrator
- Harmonic potential in normal modes
- *ab initio* energy and gradient at each step
- the smoother, the better

$$H_{\text{as}} = [1 - \lambda(t)] H_{\text{harm}} + \lambda(t) H_{\text{anh}} \quad (4)$$

$$\dot{\mathbf{p}} = (\lambda - 1) \mathbf{g}_{\text{harm}} - \lambda \mathbf{C}^T \mathbf{g}_{\text{anh}} \quad (5)$$

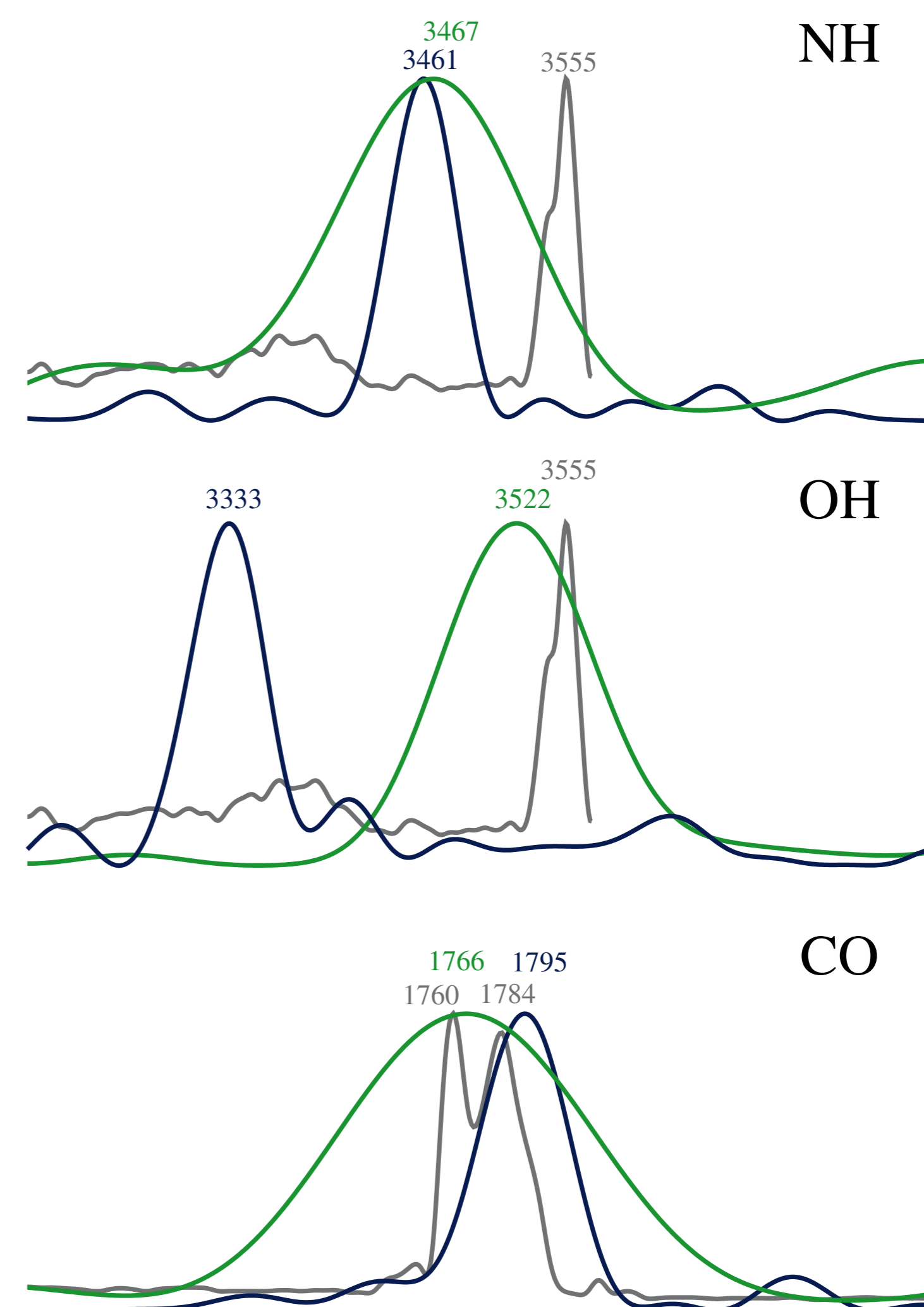
$$\lambda(t) = \frac{t}{T_{\text{AS}}} - \frac{1}{2\pi} \sin\left(\frac{2\pi t}{T_{\text{AS}}}\right) \quad (6)$$

The new trajectory, starting from “switching conditions” ($\mathbf{p}_{\text{as}}, \mathbf{q}_{\text{as}}$) is more stable and descriptive, allowing a better sampling of the phase space:

$$I_{\text{as-ta}}(E; \mathbf{p}_{\text{as}}, \mathbf{q}_{\text{as}}) = \left(\frac{1}{2\pi\hbar} \right)^{N_v} \frac{1}{2\pi\hbar T} \left| \int_0^T dt e^{i[S_t(\mathbf{p}'_t, \mathbf{q}'_t) + Et + \phi_t(\mathbf{p}'_t, \mathbf{q}'_t)]} \langle \psi(\mathbf{p}_{\text{eq}}, \mathbf{q}_{\text{eq}}) | g_t(\mathbf{p}'_t, \mathbf{q}'_t) \rangle \right|^2 \quad (7)$$

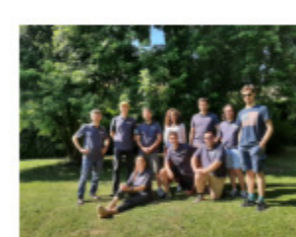
The spectra

trans-Pro and cis-Pro are in the experimental spectrum



References

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