Revision 1

High-temperature and high-pressure behavior of carbonates in the ternary diagram CaCO3-MgCO3-FeCO3

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Abstract

We report the thermal expansion and the compressibility of carbonates in the ternary compositional diagram CaCO3-MgCO3-FeCO3, determined by *in-situ* X-ray powder and single-crystal diffraction. High-temperature experiments were performed by high-resolution X-ray synchrotron powder diffraction from ambient to decarbonation temperatures (25-850 °C). Single-crystal synchrotron X-ray diffraction experiments were performed in a variable pressure range (0-100 GPa), depending on the stability field of the rhombohedral structure, which is a function of the carbonate composition. The thermal expansion increases from calcite, CaCO₃, α_0 =4.10(7) *10⁻⁵ K⁻¹, to magnesite, MgCO₃, α_0 =7.04(2)*10⁻⁵ K⁻¹. In the magnesite-siderite (FeCO₃) join, the thermal expansion decreases as iron content increases, with an experimental value of $\alpha_0 = 6.44(4)*10^ 5 K⁻¹$ for siderite. The compressibility in the ternary join is higher (*i.e.*, lower bulk modulus) in calcite and Mg-calcite $[K_0 = 77(3)$ GPa for Ca_{0.91}Mg_{0.06}Fe_{0.03}(CO₃)] than in magnesite, $K_0 = 113(1)$ GPa, and siderite, K_0 =125(1) GPa. Single-crystal structure refinements at high pressure illustrate that the main mechanism by which carbonates compress is through the distortion of the octahedral sites. The analysis of thermal expansion and compressibility variation in calcite-magnesite and calcite-iron-magnesite joins clearly shows that the structural changes associated to the order-disorder transitions (*i.e*., *R*-3*c* calcite-type structure *vs. R*-3 $CaMg(CO₃)$ ₂ dolomite-type structure) do not affect significantly the thermal expansion and compressibility of carbonate. On the contrary, the chemical compositions of carbonates play a major role on their thermoelastic properties. Finally, we use our *P-V-T* equation of state data to calculate the unit-cell volume of a natural ternary carbonate, and we compare the calculated volumes to experimental observations, measured *in situ* at elevated pressure and temperatures, using a multianvil device. The experimental and calculated data are in good agreement demonstrating that the equation of state here reported can describe the volume behavior with the accuracy needed, for example, for a direct chemical estimation of carbonates based on experimental unit-cell volume data of carbonates at high pressures and temperatures.

Introduction

The interest in carbonate mineralogy has grown significantly in the last decades. These minerals can in fact provide clues and information about recent and past Earth's climate history. Carbonates are involved in fundamental geological processes related to the global carbon cycle. In particular, a shallow carbon cycle (*i.e.*, involving exchanges between atmosphere and hydrosphere reservoirs) is directly related to the dissolution and precipitation of carbonates, mainly calcite, in the oceans. The carbonate-bearing sediments are involved also in subduction processes. The stability of carbonate in these environments is the key issue to understand the transfer of crustal carbon into the inner Earth or its recycling through volcanism. Carbon is also stored in the Earth's mantle in various forms, as demonstrated by natural occurrence of diamonds. Some of these diamonds also present carbonate, or $CO₂$ as inclusions (e.g., Berg, 1986, Navon, 1999). It is now clear that a thermodynamic modelling of carbonate stability in various environments is fundamental for a deeper understanding of global carbon cycle. Thermodynamic data on the ternary composition diagram calcite-magnesite-siderite are scarce, especially for iron-bearing carbonates. A direct experimental observation of the stable composition as function of pressure, temperature and chemical composition is probably the most successful approach for a detailed understanding of carbonate stability. In recent years, the availability of large-volume high-pressure devices at synchrotron sources has allowed a number of possible experiments to be performed. The accuracy in diffraction and pressure-temperature control can nowadays provide reliable crystallographic information at given *P* and *T* conditions. While planning a series of experiments, we realized that literature data about *P-V-T* equations of state (EoS) of carbonate are not only scarce, but also contradictory. This is especially evident for thermal properties, especially in iron-bearing members. The knowledge of accurate EoS is necessary in order to properly extract chemical information from *in-situ* experiments, where the most direct information on these phases is their unit-cell volume and edges lengths. To fix the discrepancies in *P-V* and *V-T* data and fill the gap for missing compositions, we performed new high-temperature and high-pressure experiments on well-characterized samples.

We report, in the current paper, the results of high-resolution X-ray powder diffraction measurements at high temperature and single-crystal experiments at high pressure. The *V-T* and *P-V* equations of state are finally compared with *in-situ* experiments at simultaneously high pressure and high temperature.

Materials and methods

Natural and synthetic samples were used in this study. Table 1 reports the samples, their provenance and chemical composition. The chemical analyses were performed with a Jeol Superprobe instrument available at the Earth Sciences Department, University of Milan, Italy (ESD-Mi). The calibration of the instrument was performed with standard mineral samples.

Single-crystal diffraction at ambient conditions was performed with the Oxford Diffraction X'Calibur diffractometer, equipped with Mo-*K*α source and CCD detector, at ESD-Mi. We used standard operating conditions for data collection, performing 0.75° ω-axis scans (Busing and Levy 1967) at variable φ, κ and θ axis configurations. The raw data were integrated with the Crysalis Red software (Oxford Diffraction 2008) and structure refinements were performed with the Jana2006 program (Petricek et al. 2014).

Three samples were synthesized with the end-loaded piston cylinder at ESD-Mi. We employed the standard procedures commonly in use for double stage piston cylinders. The starting material was a mixture of carbonates (calcite:dolomite 5:95%; calcite:Fe-dolomite 80:20%). The samples were finely grounded in ethanol to assure a homogeneous mixing. The starting material was inserted in a welded Pt capsule, and before welding, a tiny amount of water (apx. 2%) was added, to promote growth of homogeneous crystals. The synthesis experiments were run at 2.0 GPa and 1050 °C for 2 days. The synthesis ended up with crystals with variable size, up to $100x100x100 \mu m^3$.

The X-ray high-temperature (room-pressure) powder diffraction experiments were performed at the MCX beamline at the Elettra synchrotron facility (Trieste, Italy), the beamline being described in Rebuffi et al. (2014). The X-ray diffraction patterns were collected with the high-resolution two circle diffractometer available at the beamline. A monochromatic beam of 0.8270 Å was used, and diffraction patterns were collected over the angular range of 10-30 ° 2θ, with a 0.0075° step size. High-temperature conditions were maintained with a hot gas blower. The sample temperature was monitored by the use of quartz as internal standard in five different experiments. The comparison of unit-cell parameters with reference data (Kihara 1990) and the alpha/beta quartz transition demonstrates that the temperature is accurate within 1-2 $^{\circ}$ C and perfectly reproducible in the different runs (Deposit Items). To avoid any possible reaction between the

carbonate and quartz, the quartz powder was inserted in a 0.2 mm quartz capillary, within the larger (0.7mm) capillary containing the sample. The diffraction patterns were analyzed by the Rietveld method, using the GSAS+EXPGUI software (Larson et al. 1994; Toby 2001). Two additional measurements were performed in a CO2 atmosphere using the furnace operating under controlled atmosphere (Riello et al. 2012) equipped with an imaging plate detector.

The X-ray high-pressure (room-temperature) single-crystal diffraction experiments were performed at the ID09A beamline (ESRF, Grenoble, France) with the standard beamline setup (Merlini and Hanfland 2013). We used a monochromatic radiation having wavelength λ =0.41414 Å. The diffraction spots were collected with the Mar555 flat panel detector, at a distance of 300 mm from the sample, and images were recorded while integrating over 1° step size during an ω-rotation (Busing and Levy 1967). High-pressure conditions were generated with a membrane-type diamond anvil cell (DAC) using diamonds with Boheler-Almax cut and variable culet size (600, 300, 125 μm diameter), depending on the target pressure. The gaskets were made with a rhenium foil and the *P*-transmitting medium was helium or neon. The pressure was determined by the ruby-fluorescence method (Forman et al. 1972; Mao et al. 1986) or with Sm:borate (Datchi et al. 1997) for the Mbar experiment. The raw data were integrated with the Crysalis Red program and singlecrystal structural refinements were performed with Jana2006 software.

The high-pressure/high-temperature experiments were performed at the ID06 beamline at the ESRF, using the high-resolution setup. High-pressure conditions were generated by a 20MN large-volume device. The experiments were performed using a standard 10/4 mm multi-anvil assembly, with pyrophyllite gaskets and sample directly contained within the graphite furnace. Co-axial to the incident beam direction, a boron:epoxy rod was inserted to avoid X-ray diffraction signal from the spinel and MgO from the Cr-doped octahedron. A 'C'-type thermocouple was included in the assembly for temperature measurements. Pressures and temperatures during the experiments were calculated with cross matching of thermocouple values, heating power, thermal EoS of graphite (Hanfland et al. 1989; Colonna et al. 2011) and MgO (Hazen 1976; Dewaele et al. 2000). Pressure was increased up to 3 GPa at room-*T*, and then the temperature was increased. Monochromatic X-ray diffraction patterns (λ =0.3757 Å) were collected with a Tl:NaI Bicron scintillator detector, scanning on a large-radius (~1350 mm) pseudorotation over the 2θ angular range of 5-25°, with step size 0.001°. Two sets of adjustable slits were used to define the volume of the diffracting sample, and thus shield diffraction signal from the up- and downstream pyrophyllite gaskets and the furnace. Angular calibration of the instrument was performed before the experiment against NIST standard LaB_6 SRM660a. The data were analyzed by the Rietveld method, using the GSAS+EXPGUI software (Larson et al. 1994; Toby 2001).

Results

- **Crystal structure of synthetic and natural samples by single-crystal X-ray diffraction**

Fourteen different carbonates (Table 1) were investigated by single-crystal X-ray diffraction, to characterize the samples used for thermal expansion and compressibility measurements.

All the samples crystallize in calcite-type structure (*R*-3*c*) except dolomite and Fe-dolomite samples (*R*-3). All the structural details are reported in Table 1 and in the Deposit Items.

We may consider different binary compositional joins, in particular the calcite-magnesite; calcite-Femagnesite; dolomite-ankerite and magnesite-siderite. The unit-cell edges and volume vary almost linearly in the four different binary joins considered, and are independent from the symmetry of the sample. These results are summarized in Fig. 1. The variation of the unit-cell parameters in the ternary compositional carbonate system, at a first approximation, can be described with the following simple equations involving the concentration of $CaCO₃$, MgCO₃ and FeCO₃:

 a (Å) = 4.985 X_{CaCO3} + 4.636 X_{MeCO3} + 4.696 X_{FeCO3}

$$
c
$$
 (Å)= 17.064 X_{CaCO3} + 15.033 X_{MgCO3} + 15.414 X_{FeCO3}

$$
Vol (A3) = 366.2 XCaCO3 + 279.2 XMgCO3 + 294.1 XFeCO3
$$

The numerical values are derived from a fit on the experimental data. X_{CaCO3} , X_{MgCO3} , X_{FeCO3} are, respectively, the molar fraction of the calcitic, magnesitic and sideritic components in the considered carbonate. The data agree with literature data (*e.g*., Effenberger et al. 1981; Reeder and Dollase 1989; Boulard et al. 2012).

A comprehensive discussion on the chemical effects on structural parameters is widely discussed in the literature (*e.g*., Effenberger et al. 1981; Reeder 1983; Reeder and Dollase 1989; Redfern 2000). The *R*-3 double-carbonates analyzed (dolomite and Fe-dolomite) present a fully ordered cation distribution (Ca and

Mg+Fe, respectively) over the two cation sites. The synthetic Mg-calcite samples possess the calcite-type structure (*R*-3*c*) and no evidence of partial ordering is detected by X-rays, within the experimental accuracy. The synthetic sample $Ca_{0.55}Mg_{0.45}(CO_3)$ has also a calcite-type structure, *R-3c*, as demonstrated by structure refinement (Deposit Items). This indicates that a slight Ca excess in dolomite may quench as a fully disordered dolomite structure.

- **Thermal expansion: synchrotron X-ray powder diffraction experiments**

Powdered samples from the batch crystals described in the previous section were used for the hightemperature thermal expansion measurements. Fig. 2 shows an example of an X-ray powder pattern containing calcite and quartz as internal standard. The Full Width at Half Maximum (FWHM) of the diffraction peaks, over the angular range considered, is, on average, 0.03°. All the samples are analyzed by the Rietveld method, which allows the refinement of the unit-cell parameters at variable temperature (Deposit Items).

Calcite presents a negative thermal expansion along the *a-*crystallographic axis (Rao et al. 1968). The expansion parallel to the *c*-axis is positive, like the volume expansion. All the other samples analyzed present a positive expansion along both crystallographic directions except the Mg-Fe-calcite, which has an anomalous thermal behavior parallel to the crystallographic *a* axis, but significantly smoothed compared to pure calcite. This indicate that a minor incorporation of Mg and Fe may affect significantly the thermal behavior. The volume data are fitted with the formalism proposed by Pawley et al. (1996) and Fei (1995) commonly used in mineralogy and petrology (*e.g*., Holland and Powell 1998; 2011; Angel et al. 2014).

Thermal expansion, defined as:

$$
\alpha(T) = \frac{1}{V} * \frac{\partial V}{\partial T} \tag{1}
$$

and can be described with a polynomial expression, for example (Pawley et al. 1996):

$$
\alpha(T) = a_0 + \frac{a_1}{\sqrt{T}} \tag{2}
$$

The equation resulting from integration of (1) using the expression (2) is:

$$
V = V_0 * exp[a_0 * (T - T_0) + 2 * a_1 * (\sqrt{T} - \sqrt{T_0})]
$$
\n(3)

which can be further simplified, by a Taylor expansion to:

$$
V = V_0 * [1 + a_0 * (T - T_0) + 2 * a_1 * (\sqrt{T} - \sqrt{T_0})]
$$
\n(4)

A further simplification results from the empirical observation that for most materials, a_0 and a_1 are correlated such that $a_1 \approx -10a_0$ (Pawley et al. 1996) leading to:

$$
V = V_0 * [1 + a_0 * (T - T_0) - 20 * a_0 * (\sqrt{T} - \sqrt{T_0})]
$$
\n(5)

For comparison, we also used the formalism proposed by Fei (1995), in which the thermal expansion is expressed as follows:

$$
\alpha(T) = a_0 + a_1 * T \tag{6}
$$

It results in the expression:

$$
V = V_0 * exp[a_0 * (T - T_0) + 0.5 * a_1 * (T^2 - T_0^2)]
$$
\n(7)

The experimental data were fitted with the expression (2) and (6) for the thermal expansion and the results are reported in Table 2.

In the magnesite-siderite binary join, the thermal expansion coefficient decreases from the magnesite to the siderite end-member (Fig. 3), and the variation is linear, as a first approximation. The current data provide a clear indication on the degree of the effect of increasing Fe in magnesite, consisting of a sensible decrease of the thermal expansion. In the dolomite-ankerite join, the rate of decrease of thermal expansion as a function

of Fe content is reduced, but Fe-dolomite still presents a lower expansion if compared to pure dolomite. The thermal expansion of dolomite is comparable to literature data (Markgraf and Reeder, 1986).

The effect of calcium on thermal expansion can be evaluated in the calcite-magnesite and calcite-Femagnesite compositional joins (Fig. 4). Calcite has the lowest expansion. A small incorporation of Mg and Fe in calcites drives the expansion values close to the ones of dolomite and Fe-dolomite. The calcite-Femagnesite join (after deriving the value for $Mg_{0.6}Fe_{0.4}CO_3$ from a linear fit on the magnesite-siderite mixture) indicates, further, that there is no evident effect of the actual symmetry (*R*-3*c vs. R*-3) on the thermal expansion of carbonates, which seems mostly determined by their composition.

All the samples here analyzed show an incipient decomposition at temperatures higher than 550-650 °C. In order to extend the experimental temperature range, a few measurements were performed under controlled atmosphere (P_{CO2} = 2 bar) and with a fast detector. Pure calcite undergoes two high temperature transitions, above 500 °C and 1000 °C respectively, to the CaCO₃-IV (R -3 c) and CaCO₃-V (R -3 m) polymorphs. These transitions are marked by a change in elasticity and symmetry, in agreement with previous observation (Mirwald 1976; Dove and Powell 1989; Redfern et al. 1989; Dove et al. 2005; Antao et al. 2009; Ishizawa et al. 2013). A detailed description of high temperature polymorphism is to be reported elsewhere. Dolomiteankerite and magnesite-siderite do not extend their stability significantly, and decarbonation is detected above 650 °C. Mg-Fe-calcite has an enlarged stability domain, persisting as a single *R*-3*c* phase up to 800 $^{\circ}$ C. Above this temperature, a decomposition into CaCO₃, Fe₃O₄ and MgO is observed.

- **Compressibility: synchrotron single crystal and powder X-ray diffraction experiments**

Single-crystal high-pressure diffraction experiments were performed on nine different samples in the ternary compositional diagram (Table 3). The maximum pressure investigated depends on the stability field of rhombohedral carbonates (at ambient temperature). The experimental volume data are fitted with the formalism of the Birch-Murnagham (BM) Equation of State (EoS), truncated to the third order:

$$
P = 1.5 K \sigma * \left[\left(\frac{V}{V_0} \right)^{-\frac{7}{3}} - \left(\frac{V}{V_0} \right)^{-\frac{5}{3}} \right] * \left[1 + \frac{3}{4} (K' - 4) * \left(\left(\frac{V}{V_0} \right)^{-\frac{2}{3}} - 1 \right) \right]
$$

where *K*0 is the isothermal bulk modulus and *K'* its pressure derivative.

Pure calcite is stable in rhombohedral *R*-3*c* symmetry up to a maximum pressure of 1.7 GPa (Merrill and Bassett 1975) at ambient temperatures, and was not analyzed further in this study. Single-crystal data for calcite at high pressure are reported in Redfern and Angel (1999). Dolomite and ankerite are stable in rhombohedral symmetry up to 15-20 GPa (Santillan et al. 2006; Mao et al. 2011; Merlini et al. 2012), with a transition pressure that is a function of composition and degree of order (Zucchini et al. 2014). Magnesite and siderite (Fiquet et al. 2002; Merlini and Hanfland 2013) are observed in rhombohedral symmetry above the Mbar. Iron, in rhombohedral carbonate, undergoes a high- to low-spin state transition above 45 GPa (Mattila et al. 2007; Lavina et al. 2010a; Lin et al. 2012). This electronic transition induces a volume collapse easily detectable by X-ray diffraction experiments, even when Fe is present as minor element (Lavina et al. 2012; Lin et al. 2012).

We performed high-pressure experiments on calcite and Mg-Fe-calcite samples and observed the calcite to calcite-II transition at 1.7 GPa for pure calcite, at similar pressure for $(Ca_{0.96}Mg_{0.02}Fe_{0.02})CO_3$ and at 2.5 GPa for $(Ca_{0.91}Mg_{0.06}Fe_{0.03})CO_3$. We collected enough data points for the determination of the equation of state of

 $(Ca_{0.91}Mg_{0.06}Fe_{0.03})CO₃$. The evolution of the Eulerian finite strain (*f*) as function of the normalized pressure (*F*), the *f-F* plot (Angel 2000), for this sample is consistent with an extremely low pressure derivative of bulk modulus (Deposit Item) and a fit to a 3rd order BM-EoS results in a bulk modulus K_0 =77(3) GPa with negative $K' = -4(1)$. This anomaly is probably related to the structural behavior, which rapidly proceeds to the second-order transition of calcite to calcite-II structure. Such uncommon elastic anomalies, resulting in a negative *K'*, are sometimes observed and it is the case, for example, of cordierite (Miletich et al. 2014). A fit with a 2nd order BM EoS results in K_0 = 65(1) GPa, significantly lower than pure calcite, as determined by Redfern and Angel (1999), *i.e.*, *K*₀=73.5(3) GPa.

The compressibilities were determined for three samples close to dolomite-ankerite compositions. Pure dolomite presents the dolomite to dolomite-II transition at 14 GPa, in agreement with recent results obtained by Zucchini et al. (2014). The *f-F* plot (Angel 2000) for this sample (Deposit Items), indicates that a 2nd order BM EoS is appropriate to describe the elastic behavior. The bulk modulus, 94(1) GPa, is similar to the reported data and also to Ross and Reeder (1992). A sample close to dolomite in its disordered form, $(Ca_{0.55}Mg_{0.45})CO₃$, has a quenched high-temperature disordered cation configuration. The *f-F* plot indicates that for this sample the pressure derivative of bulk modulus is not 4, and a $3rd$ order BM should be used. The refined values are $K_0 = 89(2)$ GPa and $K' = 3.0(2)$. The K' value lower than 4 may also relate to anomalous high-pressure behavior, as discussed in Zucchini et al. (2014). A detailed discussion is beyond the aim of this work; however, we also noticed, and in agreement with Zucchini et al. (2014), a lack or a shift towards high pressure of the dolomite to dolomite-II transition in disordered dolomite when compared to *R*-3 dolomite. Fe-dolomite presents a similar bulk modulus (*i.e.*, $K_0=93(1)$ GPa) to dolomite.

Samples in the magnesite-siderite join have a compressional behavior that is described by a $2nd$ order BM EoS. The bulk modulus is reasonably higher than Ca-bearing carbonates, and increase from magnesite to siderite (Fig. 5a). The measured values agree with literature data (*i.e.*, Fiquet et al. 1994; Ross et al. 1997; Lavina et al. 2010b; Lin et al. 2013; Litasov et al. 2013). A linear regression curve indicates a bulk modulus K_0 =122.7 (GPa) for pure siderite. The scattering of these experimental data are likely related to the deviation for pure magnesite-siderite join, due do the presence of minor Mn and Ca content. From these data we may extrapolate the bulk modulus of 118 GPa for a sample with a composition $Mg_{0.6}Fe_{0.4}CO_3$. This value is useful to compare the bulk modulus variation over calcite-magnesite and calcite-Fe-magnesite joins (Fig. 5b). We observe that the bulk modulus increases from calcite to magnesite, and from calcite to Fe-magnesite, with an almost linear trend. The presence of Fe increases the bulk modulus; this is most noticeable along the magnesite-siderite join.

Discussion and implications

Despite the large amount of experimental data acquired and published in the recent years, the effect of Fe in thermal and elastic properties of carbonates was unconstrained, with the exception of a recent paper on siderite (Litasov et al. 2013) based on *in-situ P-V-T* data from energy-dispersive diffraction experiment. The high-temperature data show that thermal expansion decreases as Fe content increases in carbonates. This behavior is also observed in pyroxenes, *e.g.* Hugh-Jones (1997). The high-pressure behavior shows that the Fe-bearing samples are less compressible (*i.e*., higher bulk moduli) than the Fe-free corresponding carbonates. These data are important for the thermodynamic modelling of Fe-carbonates at high temperatures (*e.g*., Kang et al. 2015). The data reported for Fe-free carbonates are comparable to the most accurate data in literature. The *V-P* and *T-V* equation of state here reported are useful in modelling the volume behavior of carbonates at variable pressures and temperature. We here compare the calculated volume of dolomite and Fe-dolomite as a function of pressures and temperatures, with experimental data, based on two *in-situ* experiments. The computed volumes are based on *P-V* and *V-T* EoS, assuming the derivative of bulk modulus with temperature, ∂*K*/∂*T* = -0.02, a value which can approximate the behavior of mantle minerals (*e.g*., Shim et al. 2000; Angel et al. 2014). The experimental volume of dolomite and Fe-dolomite is based on monochromatic X-ray powder diffraction experiments performed with a multianvil press and synchrotron radiation. An example of the high-resolution X-ray powder pattern is reported in Fig. 6. The data are fitted by the Rietveld method. About twenty X-ray diffraction patterns were collected for dolomite and Fedolomite at variable pressures and temperatures. The measured volume is compared with the calculated one at the corresponding pressure and temperature and the difference is plotted in Fig. 7 (as function of temperature). The numerical data are reported in the Deposit Items. For comparison, the volume difference

calculated using the EoS parameters of calcite and magnesite are also reported. The ordered/disordered configuration in dolomite does not affect significantly the volume behavior, at least as a first approximation and within the experimental accuracy achieved in these experiments. We noticed the disappearance of superstructure peaks of the ordered configuration at high temperature, without volume discontinuity, in agreement with previous determined temperatures (Hammouda et al. 2011; Franzolin et al. 2012). The minimum discrepancies between measured and calculated volume data indicates that the reported EoS parameters and the current experimental accuracy in monochromatic *in-situ* experiments, may provide the needed accuracy and precision for planning *in-situ* kinetic experiments, where change of composition of carbonates is expected, especially concerning the variation of $Ca/(Mg + Fe)$ ratio. This is the case, for example, during exsolution and incongruent melting of carbonates (*e.g*., Franzolin et al., 2011), whose relations are still not known completely, and especially, in the subsolidus relations of calcite-Fe-magnesite system, with direct application to the understanding of carbonate stability in subduction environments. The quantitative knowledge of such phenomena may, in fact, provide fundamental insights to understand the major geological processes involving carbon exchange between the different reservoirs in the Earth's upper mantle.

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Tables

Sample	Locality	Composition	Crystallographic data	Experiment
Calcite	Esztramos Hill (Hungary)	CaCO ₃	$a=4.9902(2), c=17.0631(9), V=367.98(3)$	SC.HT
Mg-Fe-calcite	Synthetic	$(Ca_{0.96}Mg_{0.02}Fe_{0.01})CO_3$	$a=4.9755(2)$, $c=17.0053(17)$, $V=364.58(7)$	SC.
Mg-Fe-Calcite	Synthetic	$(Ca_{0.91}Mg_{0.06}Fe_{0.03})CO_3$	$a=4.9462(7)$, $c=16.881(3)$, $V=357.66(16)$	SC, HT, HP
Ca-Dolomite (disordered)	Synthetic	$(Ca_{0.55}Mg_{0.45})CO3$	$a=4.8337(3)$, $c=16.2329(12)$, $V=328.46(3)$	SC, HP
Dolomite	Bazena (Italy)	CaMg(CO ₃) ₂	$a=4.8091(1), c=16.020(1), V=320.86(11)$	SC, HT, HP, HT/HP
Fe-dolomite	Col du Petit San Bernard (France/Italy)	$Ca_{1.01}Mg_{0.9}Fe_{0.08}Mn_{0.01}(CO_3)$	$a=4.8108(3)$, $c=16.0412(8)$, $V=321.51(10)$	SC.
Fe-dolomite	La Mure (France)	$Ca(Mg_{0.6}Fe_{0.4})(CO_3)$	$a=4.8189(2), c=16.0912(7), V=323.60(2)$	SC, HT, HP, HT/HP
Magnesite	Val Solda	$(Mg_{0.99}Fe_{0.01})CO3$	$a=4.6348(19), c=15.0235(8), V=279.49(3)$	SC, HP
Magnesite	Czechoslovakia	$(Mg_{0.99}Fe_{0.01})CO3$	$a=4.6371(2)$, $c=15.0365(13)$, $V=280.01(5)$	SC, HT
Fe-Magnesite	Tyrol (Austria)	$Mg_{0.89}Fe_{0.11}(CO3)$	$a=4.64475(17), c=15.0805(6),$ $V = 281.754(14)$	SC, HT, HP
Siderite	Puits de Villaret (France)	$(Fe_{0.70}Mg_{0.26}Mn_{0.025}Ca_{0.015})C$	$a=4.6836(2), c=15.3177(14), V=290.99(2)$	SC. HT. HP
Siderite	Foppolo (Italy)	$(Fe_{0.74}Mg_{0.21}Mn_{0.04}Ca_{0.01})CO_3$	$a=4.6875(3), c=15.3547(9), V=292.19(3)$	SC, HP
Siderite	Ivigtut (Greenland)	$(Fe_{0.95}Mn_{0.045}Mg_{0.005})CO_3$	$a=4.6924(2), c=15.3865(9), V=293.39(2)$	SC, HT, HP

Table 1 – Samples used in the current study for the single crystal (SC) diffraction experiments at ambient conditions, high-temperature (H*T*) X-ray powder diffraction experiments and high-pressure (H*P*) single crystal X-ray diffraction. The compositions were determined by electron microprobe analysis. Axial unit-cell parameters in Å, unit-cell volume in \mathring{A}^3 .

Table 2 – Results of fitting procedure on high-temperature volume data, using the Pawley et al. (1996) formalism with one thermal expansion coefficient (eq. 5), and Fei (1995) formalism (eq. 7) with two thermal expansion coefficients. Additional results using the Pawley expression with two coefficients and a linear expression are in the Deposit Items.

sample	$V_0(\AA^3)$	K_0 (GPa)	
(Ca0.91Mg0.06Fe0.03)CO3 syn	360.5(1)	77(3)	$-3.8(17)$
$(Ca0.55Mg0.45)CO3$	329.1(4)	89(2)	3.0(2)
dolomite Bazena	320.7(3)	94(1)	4
Fe-dolomite La Mure	321.1(4)	93(1)	4
magnesite Czech	279.0(5)	113(1)	4
Fe-magnesite Tyrol	280.9(2)	115.9(9)	4
siderite Puits de Villaret	288.0(4)	123(1)	4
siderite Foppolo	291.0(8)	125(2)	4
siderite Ivigtut	293.4(2)	117.7(9)	

Table 3 – Elastic parameters of carbonates (by BM EoS fit) in the ternary diagram.

Figures

Fig. 5a – Bulk moduli along the siderite-magnesite join.

Figure 5b. Bulk moduli along the calcite-magnesite and calcite-Fe-magnesite binary joins.

Fig. 6 – X-ray powder pattern of Fe-dolomite (cross) at high pressure and temperature, Rietveld full-profile fit (black line) and difference curve (gray line). The diffraction peaks of graphite and MgO are marked with asterisk.

Fig 7a

Fig. 7b

Figure 7 – Difference in calculated and experimental volume of dolomite (a) and Fe-dolomite (b). The experimental and computed data refer to a series of *P* and *T* points (Deposit Items). All the points are plotted using only temperature as variable. For comparison, also the volume difference with pure calcite and magnesite are reported, with the data of calcite and magnesite computed from the EoS of this paper at the corresponding *P* and *T*.